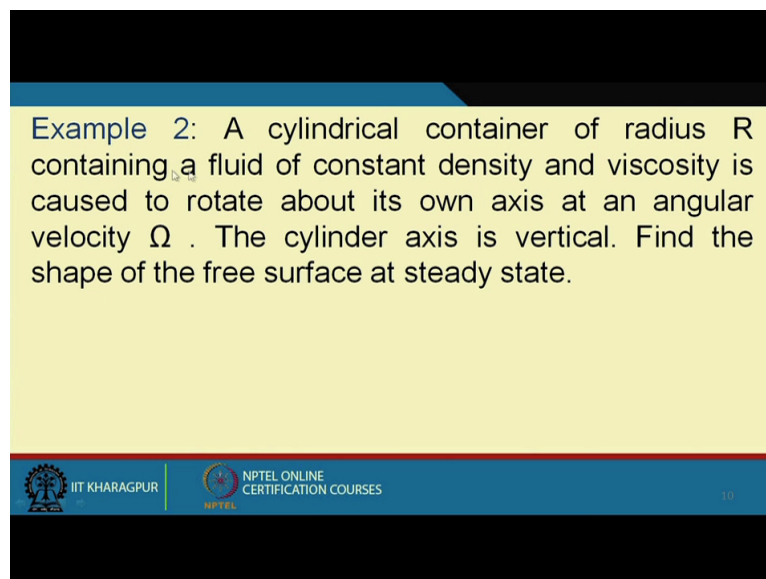


Course on Momentum Transfer in Process Engineering
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Lecture 08
Module 2
Navier stokes equations utilization for solving problems

Now I hope you have tried with several other problems in solving different problems using the Navier Stokes Equations. So if you if you have any problem, any difficulty in understanding obviously you can come back to us we will try or our best to solve those, but as we said earlier that to to understand the utilities the the implications the the the places over it can be utilized, the areas where it can be utilized the equation of motions you need to find out you need to solve some more problems and that some more problems let us take one more problem today and that problem will let us come to that right so if we if we if we think that any problem which you can do or which you can that you bring to us and we will solve them, ok.

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Example 2: A cylindrical container of radius R containing a fluid of constant density and viscosity is caused to rotate about its own axis at an angular velocity Ω . The cylinder axis is vertical. Find the shape of the free surface at steady state.

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So let us look into the problem second that is a cylindrical container of radius R containing a fluid of constant density and viscosity is caused to rotate about its own axis at an angular velocity ω . The the cylinder axis is vertical. Find the shape of the free surface at steady state, right.

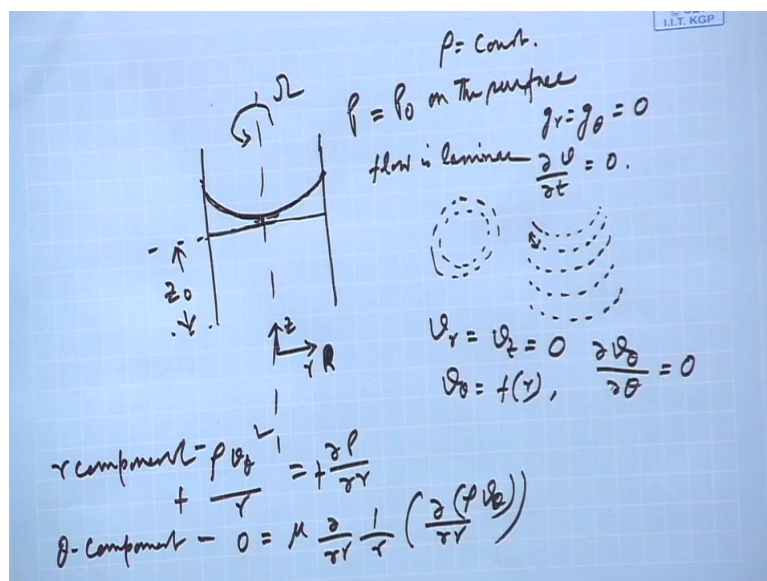
I hope in secondary we have done a locus of equation or locus of some surface etc those things we have done in a secondary. So here also a similar problem we will do in understanding the problem here you just imagine I do not know whether during your childhood you had played it or you had that option or playing with that is also true because nowadays it is being so much what should I say urbanization that those openness is also not available all the places.

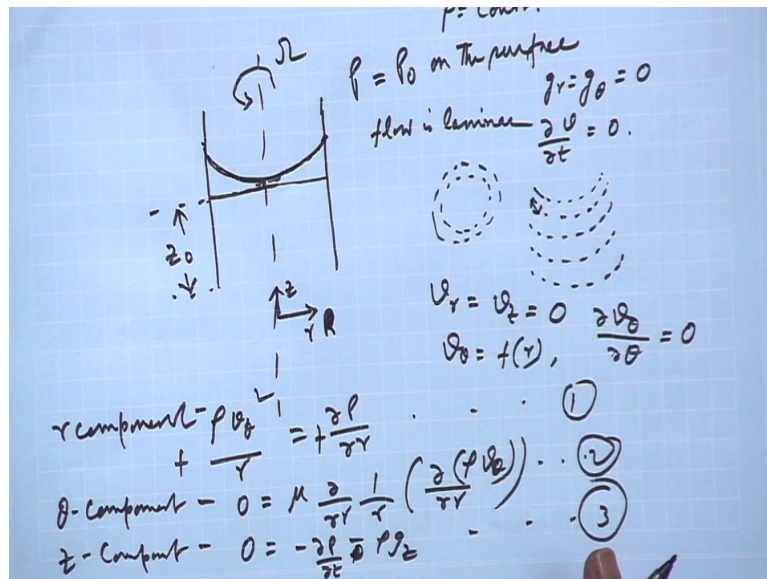
So if you if you think if you if you if you go back to your child ages and their if you had ever a bucket full of water right and you take that full of water and start making a round right, start rotating like this right bucket full of water and you are rotating on around you, right with that bucket. If you had remember if you remember that the surface of the bucket water that looks like this right that looks like this that the surface is like that when you are rotating with that bucket, right.

If you remember and know we will find out this surface of the bucket which is like this and what is the locus of that surface this you can easily do with the help of Navier Stokes Equations. So this is another kind of of application of the Navier Stokes to solve your time and again different problems, ok. So let us look into that.

So we say again this problem a cylindrical container of radius R containing a fluid of constant density and viscosity is caused to rotate about its own axis at an angular velocity ω . The cylinder axis is vertical. Find the shape of the free surface at steady state, right.

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So let us let us look into this that we have a cylinder right and this cylinder has the axis like this ok and originally originally we had a a layer of liquid like this right. Now suddenly we started moving this with an omega in the axis in the own axis omega right and if we look at the this is the R or say let us write small r and this is z right and third one is the theta of course right. So initially if we say this is the height z₀ initial height that is where we started with right and after some time when the steady state has arrived we will have a surface like this right we will have a surface like this.

Now we have to find out what is the locus of this surface right that is this equation of this surface that if we can solve then we can say that the understanding of Navier Stokes Equation has become little bit more easier or we have we have understood a little more application of the Navier Stokes Equation.

Here one more thing has to be said that P is equal to P₀ on the surface right, P is equals to P₀ on the surface if that be true again we bring the again we bring the condition the Physical understanding of the problem right. We said that we have a liquid in a container and this container is rotating on its own axis like this and with a velocity angular velocity of omega right.

So we have this axis at R and z these are the two initial height where we started with it was z₀ and after some time when steady state arrived we got this surface with an angular velocity omega like this and P is equals to P₀ obviously of that pressure on all the surface is same, right. So from the understanding of the problem we can say that since this is a vertical gr is

equals to $g \sin \theta$ is equals to 0. We also said flow is equals to constant, we also said at steady state, so $\frac{d}{dt}$ of any v_r, v_z, v_θ all are is equals to 0.

Fourth thing which we have not said let us say again the flow or it is such a way that terminance is not occurring so flow is laminar right, again laminar means this layer, this layer they are all not mixing with each other, right a free layers is like that again when it became a surface like this the layers in the in the bucket were like this right. So that means there is no in mixing of the layers interaction of v_r, v_z, v_θ these were not there, right it is a streamline laminar flow, ok.

So this indicates that v_r is equals to v_z is equals to 0, right and v_θ is also a function of r right and since it is laminar we can also say $\frac{d v_\theta}{d \theta}$ is equals to 0, right. That means this layer is moving right this layer is moving but there is no interaction between this v_θ and this v_θ , right. So otherwise there would have been a rippling there would have been a mixing of the layers which is not there.

So $\frac{d v_\theta}{d \theta}$ is also 0, right then we can say that r component r component of the velocity that is equals to ρv_θ^2 by r this is $\frac{d P}{d r}$, right you remember in the previous thing we have shown it was minus minus versus it is on the both site so we can write ρv_θ^2 is equals to $\frac{d P}{d r}$. Similarly theta component we can write for theta component we can write that this was 0 equals to $\mu \frac{d}{d r} \left(\frac{1}{r} \frac{d}{d r} (r v_\theta) \right)$ if you remember the same thing we have written.

And z component we can write this is 0 is equals to minus $\frac{d P}{d z}$ plus or here it was minus minus $\rho g z$ since it is there. So $\rho g z$ vertical thing is there so $\rho g z$ is also acting or $g z$ is also there, that these are the three equations we have three equations from the Navier Stokes Equation which originally started with and putting the boundaries putting the putting the understanding of this solution understanding of the problem we can we have identified different individual components and said that yes this is the equations which come after putting the after putting the right situations involved in the problem, right.

So the terms which are which are which can be neglected or which can be written as 0 is already shown also in the previous class right so we are not repeating their otherwise again and again the same thing we will be doing, right. So we understand that again if we have any any understanding problem you can come back to ask and we will I will definitely try to solve your doubt if there be any, ok.

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$$0 = \mu \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial (r v_\theta)}{\partial r} \right)$$

$$\frac{\partial}{\partial r} (r v_\theta) = A$$

$$\frac{\partial (r v_\theta)}{\partial r} = A r$$

$$r v_\theta = A r^2 / 2 + B$$

$$v_\theta = \frac{A r}{2} + \frac{B}{r}$$

$r=0$ since $v_\theta = \infty$, $B=0$
 $r=R$, $v_\theta = \omega R$
 $\omega R = A R / 2$, or $A = 2 \omega$
 $v_\theta = \omega r$

So if this is these are the three equations then we can say that the integrating the theta component that is 0 is equals to theta component towards 0 is equals to mu del del r of 1 by r del del r of 0 sorry r rv theta rv theta right. So if this is there this was that on first integration we can write del del r 1 by r del del r of rv theta is equals to A right on simplification we can write del del r of rv theta is equals to Ar right.

So this was first integration of with constant A and the second integration we can say rv theta is equals to Ar square by 2 plus B, where B is the second integration constant, right which we have also seen earlier. So we can write v theta is equals to A by A by sorry A by 2 into r Ar by 2, so v theta is equals to Ar by 2 r r so it goes out plus B by r right.

Now at r is equals to 0 since at r is equals to 0 since v theta is not equals to infinity at r is equals to 0 since v theta is not equals to infinity we can write B to be equals to 0, otherwise if r is 0 v theta has to be infinity then only you will have some value otherwise B is equals to 0, since this is applicable we can write B to be equals to 0 right the second (())(15:58) at r is equals to capital R, v theta is equals to capital omega R, right.

So this omega R is the second equation so we can write omega R is equals to AR by 2 or A is equals to 2 omega, right. Therefore v theta is equals to omega r right so this is v theta is equals to v is 0 right and A is equals to 2 omega, so A is 2 omega so v theta is equals to omega r this is one situation from the solution of one component that is theta component.

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$$\frac{\rho \theta_0^2}{r} = \frac{\partial p}{\partial r}$$

$$\frac{\partial p}{\partial r} = \frac{\rho \omega^2 r}{r} = \rho \omega^2 r$$

$$\frac{\partial p}{\partial z} = -\rho g \quad p = f(r, z)$$

$$dp = \frac{\partial p}{\partial r} dr + \frac{\partial p}{\partial z} dz$$

$$dp = \frac{\rho \omega^2 r}{r} dr - \rho g dz$$

$$p = \frac{\rho \omega^2 r^2}{2} - \rho g z + C$$

$$p = p_0 \text{ at } z = z_0, r = 0 \quad p_0 = -\rho g z_0 + C \quad \therefore C = p_0 + \rho g z_0$$

$$\therefore p = \frac{\rho \omega^2 r^2}{2} + p_0 + \rho g z_0 - \rho g z$$

$$\frac{\partial p}{\partial r} = \frac{\rho \omega^2 r}{r} = \rho \omega^2 r$$

$$\frac{\partial p}{\partial z} = -\rho g \quad p = f(r, z)$$

$$dp = \frac{\partial p}{\partial r} dr + \frac{\partial p}{\partial z} dz$$

$$dp = \frac{\rho \omega^2 r}{r} dr - \rho g dz$$

$$p = \frac{\rho \omega^2 r^2}{2} - \rho g z + C$$

$$p = p_0 \text{ at } z = z_0, r = 0 \quad p_0 = -\rho g z_0 + C \quad \therefore C = p_0 + \rho g z_0$$

$$\therefore p = \frac{\rho \omega^2 r^2}{2} + p_0 + \rho g z_0 - \rho g z$$

$(p - p_0) = 0$ at all points on the surface. $g(z - z_0) = \frac{\omega^2 r^2}{2}$

Now if we see the other one $\rho \omega^2 r$ is equals to $\frac{\partial p}{\partial r}$, $\rho \omega^2 r$ square by r is equals to $\frac{\partial p}{\partial r}$ this was from the r component right we did. So this on this on simplification we can also right $\frac{\partial p}{\partial r}$ is equals to already we have found out what is the value of ω so that is $\rho \omega^2 r$ square by r this is $\omega^2 r$ square right r into r square $\omega^2 r$ it was $(\omega r)^2$ was ωr so this is equals to this so that means $\rho \omega^2 r$ is that $\frac{\partial p}{\partial r}$, right.

And the third equation was $\frac{\partial p}{\partial z}$ this is equals to minus ρg right if you remember that it was 0 is equals to $\frac{\partial^2 p}{\partial z^2}$ (18:09) $\frac{\partial p}{\partial z}$ minus ρg so we can write $\frac{\partial p}{\partial z}$ $\frac{\partial p}{\partial z}$ minus ρg so no it was not minus it was plus $\frac{\partial^2 p}{\partial z^2}$ minus ρg so that on simplification it was it comes $\frac{\partial p}{\partial z}$ is equals to minus ρg right.

So now if that be true we can write since P is a function of r and z right since P is a function of r and z right we can write that $\text{del } P$ or dP is equals to $\text{del } P \text{ del } r \text{ dr}$ plus $\text{del } P \text{ del } z \text{ dz}$. So dP is equals to $\text{del } P \text{ del } r \text{ dr}$ plus $\text{del } P \text{ del } z \text{ dz}$. This we can write right. So now we have already found out the value of $\text{del } P \text{ del } r$ this is $\rho v \theta^2$ by r and also we have found out $\text{del } P \text{ del } z$ is $\rho g z$.

So we can write that dP is equals to this is $\text{del } P \text{ del } r$ is $0 v \theta^2$ by r right. So dr (plus) minus so this is minus this plus goes out $\rho g \text{ del } P \text{ del } z$ is we have found out minus $\rho g z$ so minus $\rho g z \text{ dz}$, right. So this is then dP is equals to $\rho v \theta^2$ by $r \text{ dr}$ minus $\rho g z \text{ dz}$, right.

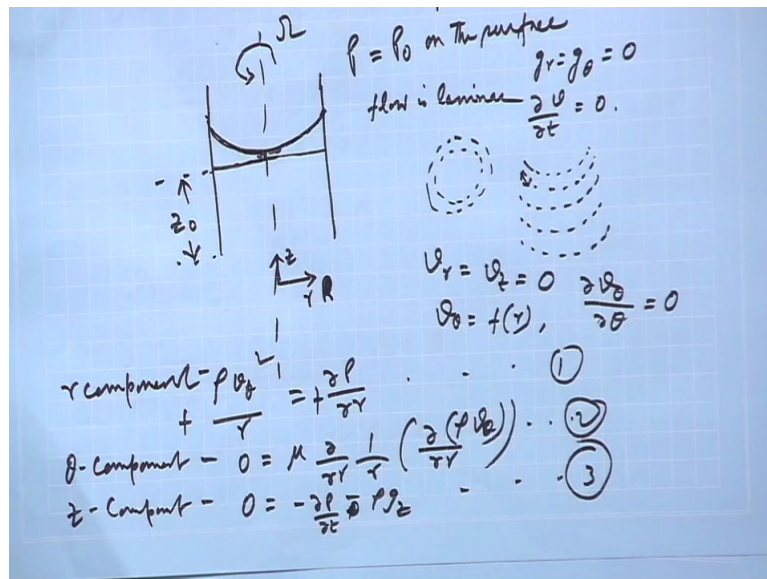
So on integration of this we can write P is equals to ρ right $v \theta^2$ we have already found out that was equals to $\omega^2 r$ so we can we we we can write this is $\omega^2 r^2$ right by $\rho v \theta^2$ by r is $\text{del } P \text{ del } r$ ok. So this is P is equals to $\rho \omega^2 r^2$ by 2 because this dr this becomes 1 by $r \text{ dr}$ then that becomes r^2 by 2 and on integration minus $\rho g z$, right this plus C .

So dP become P 1 by $r \text{ dr}$ this becomes $\rho \omega^2 r^2$ is $v \theta^2$ was $\omega^2 r^2$ right. So $\omega^2 r^2$ that became r so on $r \text{ dr}$ integration r^2 by 2 $\rho g z$ plus C , right so this is that and from the given boundary we know P is equals to P_0 right that means at z is equals to z_0 , r is equals to 0 , this is true at P is equals to 0 when this is at z is equals to z_0 and r is equals to 0 right. Therefore we can write P_0 is equals to minus ρ minus $\rho g z$ right or or minus $\rho g z, z_0$ plus C right or C is equals to P_0 plus $\rho g z_0$, right.

Therefore we can write P is equals to $\rho \omega^2 r^2$ by 2 plus P_0 plus $\rho g z_0$ minus $\rho g z$ right or we can write P minus P_0 P minus P_0 this is equals to 0 at all the points on the all the points on the surface, right. Therefore we can write g into z minus z_0 is equals to $\omega^2 r^2$ by 2 this is $\omega^2 r^2$ by 2 right or we can write that z minus z_0 this is equals to $\omega^2 r^2$ by $2g$, right.

So the locus of this equation right locus of this curve we can write this is nothing but z minus z_0 is equals to $\omega^2 r^2$ by $2g$, right. So we have solved this problem by using the equation of motion or by using the Navier Stokes Equation, obviously we have not done some of the things which we have done earlier that showing initially that how the conditions r applicable right.

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Now since we have little time so what we can do we can what we can do oh this was our problem right we can go back a little and show that from the original equations we have how arrived. The problem was given that we have a bucket and that bucket was rotating that I gave the example we have a bucket which we were rotating like this but in this case say we have some bucket which is rotating at its own axis with an angular velocity of omega right and now we were said what is the steady state flow this curve there will be there will be there will be a circular this thing so that will be how where is the locus of this right this we were asked that what is the equation right.

And we said which starts from the Navier Stokes Equation we have already shown you what is the what is the locus of this line that we have shown you but again for recapitulation or a reunderstanding that let us look into the r component, theta component and z component.



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R-component

$$\rho \left[\frac{\partial v_r}{\partial t} + v_r \frac{\partial v_r}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_r}{\partial \theta} - \frac{v_\theta^2}{r} + v_z \frac{\partial v_r}{\partial z} \right]$$

$$= \mu \left[\frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} (r v_r) \right) + \frac{1}{r^2} \frac{\partial^2 v_r}{\partial \theta^2} - \frac{2}{r^2} \frac{\partial v_\theta}{\partial \theta} - \frac{2}{r^2} \frac{\partial v_\theta}{\partial \theta} + \frac{\partial^2 v_r}{\partial z^2} \right]$$

$$- \frac{\partial p}{\partial r} + \rho g_r$$


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$\rho = \text{const.}$

$p = p_0$ on the surface

flow is laminar $\frac{\partial v}{\partial t} = 0$

$v_r = v_z = 0$

$\frac{\partial v_\theta}{\partial \theta} = 0$

$\frac{\partial v_\theta}{\partial r} = f(r)$

1) r-component $- \rho v_\theta = + \frac{\partial p}{\partial r}$

2) θ -Component $0 = \mu \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial (r v_\theta)}{\partial r} \right)$

3) z-Component $0 = - \frac{\partial p}{\partial z}$

So here we have said that the given condition P is P_0 on this surface, g_r is g_θ is equals to 0 in earlier we had earlier we also had v_r is equals to v_θ is equals to 0 but here we have v_r is v_z is equals to 0 and v_θ is a function of r right and in this case no in that v_r , v_z was 0 and v_θ was also a function of r here also v_θ is a function of r and $\frac{\partial v_\theta}{\partial \theta}$ is 0. These r the three from the (1)(27:22) v_r is v_θ v_z 0 because have there be any any r component right so there would have been a mixing, (1)(27:34) there be any z component then there would have been a mixing, which is not there it is under laminar condition right.

If that is true then v_r is v_z is equals to 0 and $\frac{\partial v_\theta}{\partial \theta}$ is equals to 0 and also v_r is a function of r right. So if we go back to that r component, right since v_r is 0 and $\frac{\partial v_r}{\partial t}$ that is it is steady state, ρ is constant ok but this is 0 this component is 0, this component is

0, right but v_θ^2 by r with a negative minus ρv_θ^2 by r that remains v_z 0, this is 0 right since v_r is 0 this 0, v_r is 0 this is 0 $\frac{\partial v_\theta}{\partial \theta}$ is 0, so that $\frac{\partial v_\theta}{\partial \theta}$ 0, $\frac{\partial v_\theta}{\partial \theta}$ also becomes 0 or this should be $\frac{\partial^2 v_\theta}{\partial \theta^2}$ so this is 0, this is 0 and v_r is 0 this is 0 minus $\frac{\partial p}{\partial r}$ plus ρg_r .

So we got minus ρv_θ^2 by r minus is equals to minus $\frac{\partial p}{\partial r}$ plus ρg_r , right. But here also g_r is 0 so we get $\frac{\partial p}{\partial r}$ so this we said. And since there was two negatives we can make it positive. Now for that theta component we have said right the first $v_\theta \frac{\partial v_\theta}{\partial \theta}$ $\frac{\partial v_\theta}{\partial \theta}$ $\frac{\partial v_\theta}{\partial t}$ is 0 which is steady state, v_r is 0 this is out $\frac{\partial v_\theta}{\partial \theta}$ is 0, v_r is 0 this is out v_z is 0 this is out right $\mu \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} (r v_\theta) \right)$ this is not 0 this remains, but $\frac{\partial^2 v_\theta}{\partial \theta^2}$ is out then $\frac{\partial v_r}{\partial \theta}$ is out v_r is 0, $\frac{\partial v_\theta}{\partial z^2}$ is also 0.

So $\frac{\partial^2 v_\theta}{\partial z^2}$ is also 0. So this is minus $\frac{1}{r} \frac{\partial p}{\partial \theta}$ plus ρg_θ , right g_θ is also and this should be g_θ , right. So we know g_θ is we said that g_r is equals to g_θ is equals to 0 so $g_r g_\theta$ is 0. So this becomes 0, but $\frac{\partial p}{\partial \theta}$ that is also we said this $\frac{\partial p}{\partial \theta}$ cannot be there so that is also 0 because this p on this on this surface this p is same at every place. So got the second equation 0 is equals to $\mu \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} (r v_\theta) \right)$ of ρv_θ .

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Z-component

$$\rho \left[\frac{\partial v_z}{\partial t} + v_r \frac{\partial v_z}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta} + v_z \frac{\partial v_z}{\partial z} \right]$$

$$= \mu \left[\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 v_z}{\partial \theta^2} + \frac{\partial^2 v_z}{\partial z^2} \right] - \frac{\partial p}{\partial z} + \rho g_r$$

And the third z component we got this $\frac{\partial v_z}{\partial t}$ is 0 v_r is 0 this $\frac{\partial v_z}{\partial \theta}$ v_z is 0 v_z is 0 this all terms goes out then v_z is 0 this comes out this v_z is 0 this comes out $\frac{\partial^2 v_z}{\partial z^2}$ goes out what remains $\frac{\partial p}{\partial z}$ minus plus ρg_z right.

So we wrote 0 equal to minus $\text{del } p \text{ del } z$ of course this minus $\rho g z$, there was also a negative ok . Now this was $\rho g z$ that was since it is acting vertical so here we r saying that gravity is acting on that so means $\rho g z$ plus this has become minus and this on solution we have found out that ultimately we got this $v \text{ theta}$ was ωr and we got the locus of the equation as $z \text{ minus } z_0$ is equals to $\omega^2 r^2$ by 2 right, $z \text{ minus } z_0$, $\omega^2 r^2$ by $2g$, this is the solution, right.

If you remember that we said that we will do some problems for the application of the Navier Stokes Equation, right so this we say that ok we have done some problem on Navier Stokes Equation, thank you.