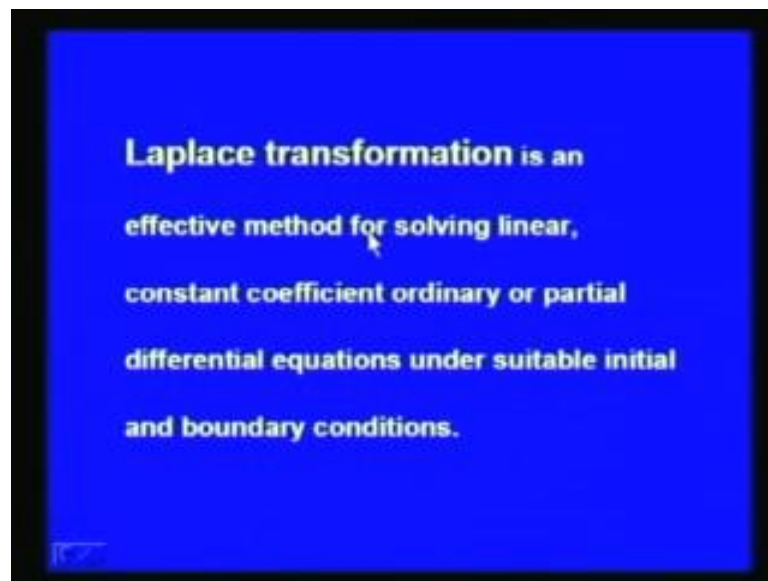


**Mathematics III**  
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**Department of Mathematics**  
**Indian Institute of Technology, Roorkee**

**Lecture - 8**  
**Laplace Transformation**

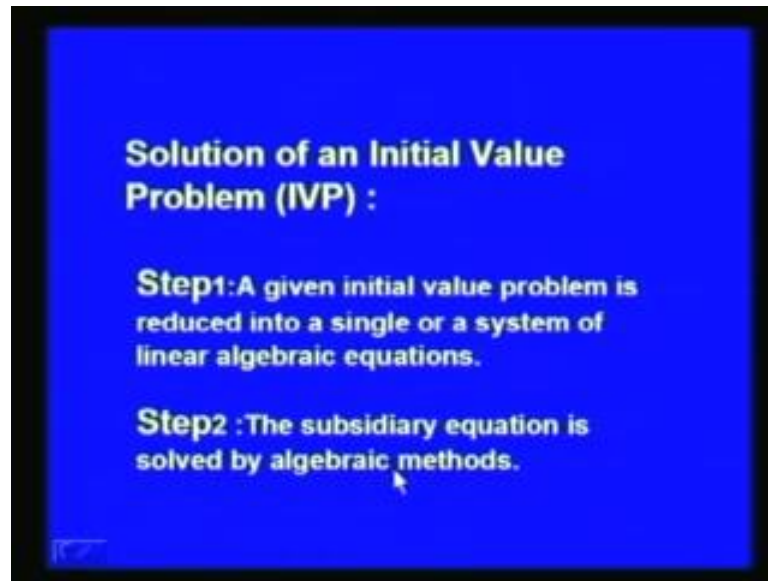
Dear viewers, in our talk today we shall be discussing the Laplace Transformation. Laplace transformation is a method for solving ordinary linear differential equations with constant coefficients. Partial differential equations can also be treated by using the Laplace transformation, it has got a tremendous application in the problems on in mechanical and electrical engineering and it is used extensively by the scientist and engineers.

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So, Laplace transformation is an effective method for solving linear constant coefficient ordinary or partial differential equations, under the suitable initial and boundary conditions.

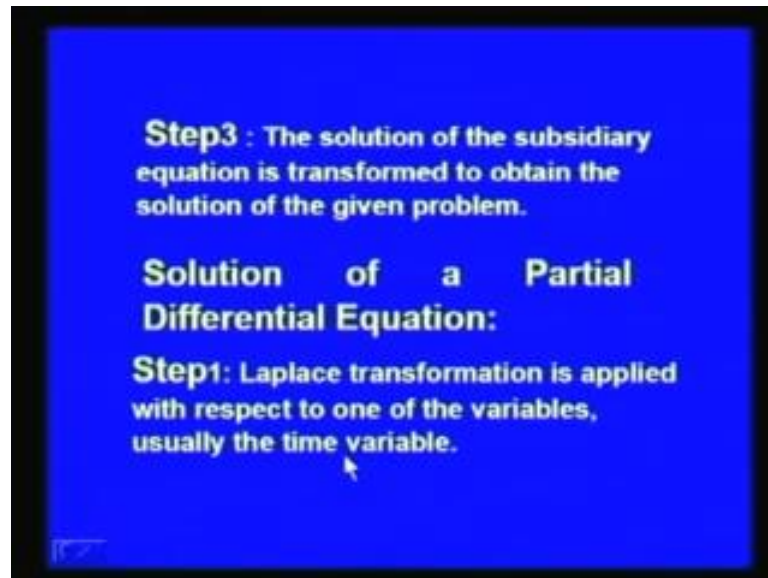
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Let us first discuss how do, we find the solution of an initial value problem using the Laplace transformation method. So, the step 1 is first we apply the Laplace transform to be given initial value problem, it will be reduced into a single or a system of linear algebraic equations. If we have a single equation in the initial value problem it will be giving us a single linear algebraic equation and we have a system of differential equations in the initial value problem.

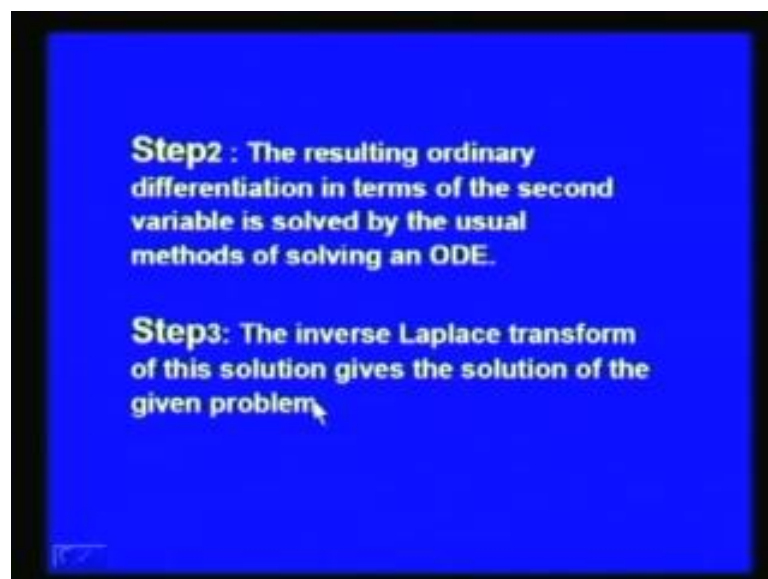
Then, after using the Laplace transform, we will get a system of linear algebraic equations, the algebraic equations are called as the subsidiary equation or subsidiary equations. In the step 2, the subsidiary equation is then solved by purely algebraic methods.

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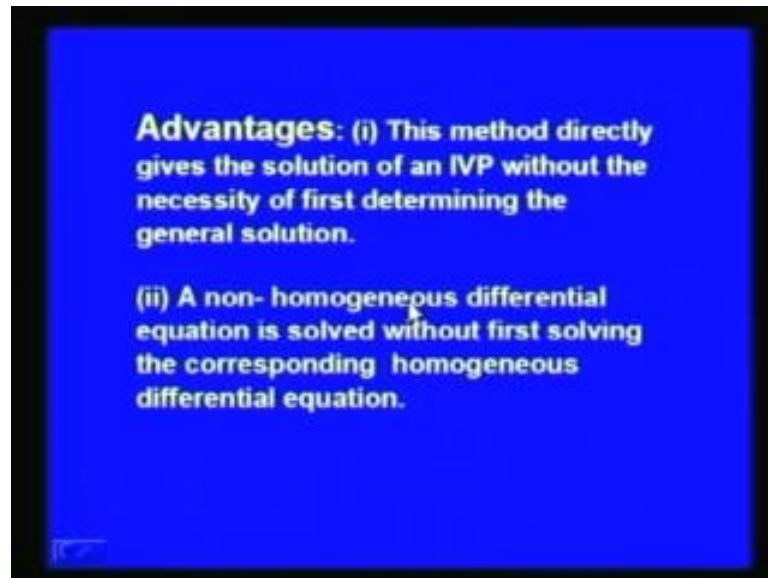
Then, the solution, in the step 3 the solution of the subsidiary equation is transformed to obtain the solution of the given initial value problem, that is we take the inverse Laplace transform of the solution of the subsidiary equation, in order to obtain the solution of the given problem. Let us see how do, we obtain the solution of a partial differential equation. In partial differential equations, first step is we apply the Laplace transformation on the given partial differential equation. It will be applied with respect to one of the two variables usually we take the variable as the time variable.

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And then the resulting ordinary differentiation in terms of the other variable is solved by using the methods of solving an ordinary differential equation. After getting this solution of the ordinary differential equation so obtained, in terms of the other variable we then take the inverse Laplace transform of the solution, which it gives us the solution of the given boundary value problem.

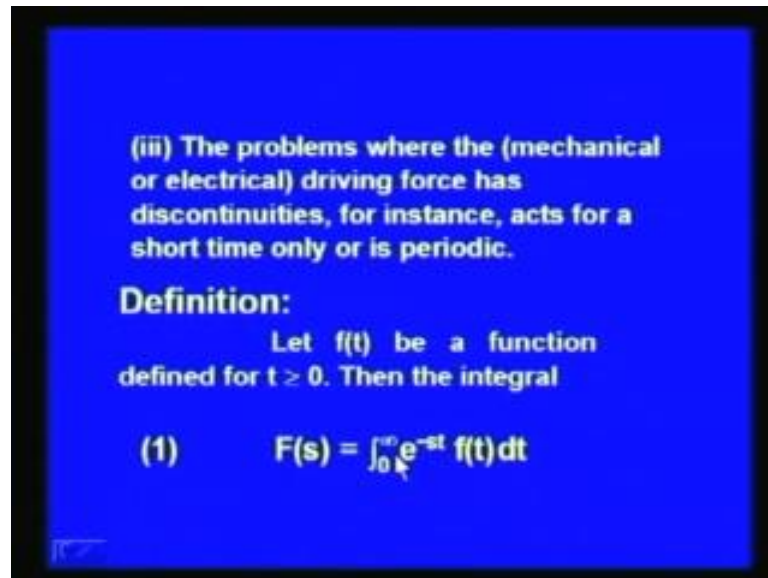
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Let us now see, the advantages of the Laplace transformation method, the Laplace transformation method gives us the solution of an initial value problem without the necessity of first determining the general solution. See the when we solve an ordinary differential equation with constant coefficients in the theory of ordinary differential equation, what we do is, we first find the general solution of the given differential equation.

And then use the initial conditions in order to arrive at the solution of the initial value problem. But, in the case of the Laplace transformation method we directly obtain the solution of the initial value problem, we do not have to find the general solution of the problem first. Non homogeneous differential equation can also be solved using the Laplace transformation method. Here, we do not have to solve the corresponding homogeneous differential equation which we do when we solve a differential equation otherwise, a non homogeneous differential equation otherwise.

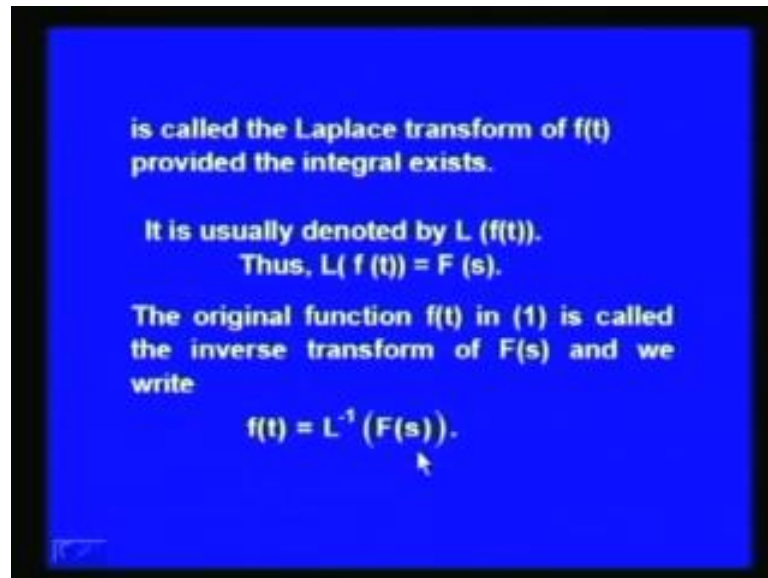
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The problems where the driving force has discontinuities say for example, the problems in a simple electric circuit, where when we give the voltage to the circuit from battery or a say it takes for a shorted short time say for the time  $t$ , and then it ceases or in the problems where the function is periodic. So, in the problems in mechanical and electrical engineering the Laplace transform is quite useful.

Let us define, what to be mean be a Laplace transform, let us say let  $f(t)$  be a function which is defined for all  $t$  greater than or equal to 0; then the integral given by  $F(s)$  equal to  $\int_0^{\infty} e^{-st} f(t) dt$  is called the Laplace transform of  $f(t)$  provided the integral in the equation 1 exists.

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is called the Laplace transform of  $f(t)$  provided the integral exists.

It is usually denoted by  $L (f(t))$ .

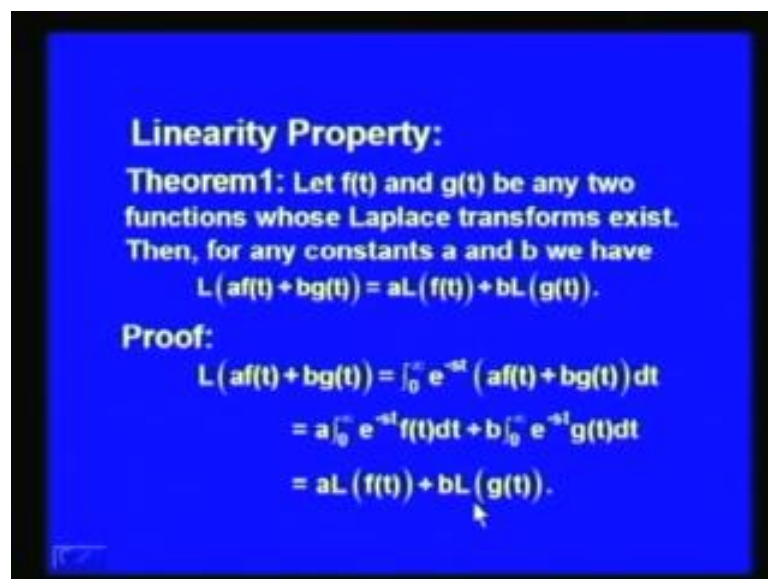
Thus,  $L( f (t)) = F (s)$ .

The original function  $f(t)$  in (1) is called the inverse transform of  $F(s)$  and we write

$$f(t) = L^{-1} (F(s)).$$

Now, we usually denote this integral by  $L f t$ , so we get  $L$  of  $f t$  equal to  $F s$  this  $L$  means the Laplace transforms of the function  $f t$ . The original function  $f t$  in the equation 1 is called the inverse transform of the function  $F s$  and we write  $f t$  equal to  $L$  inverse of  $F s$ .

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**Linearity Property:**

**Theorem 1:** Let  $f(t)$  and  $g(t)$  be any two functions whose Laplace transforms exist. Then, for any constants  $a$  and  $b$  we have

$$L (af(t) + bg(t)) = aL (f(t)) + bL (g(t)).$$

**Proof:**

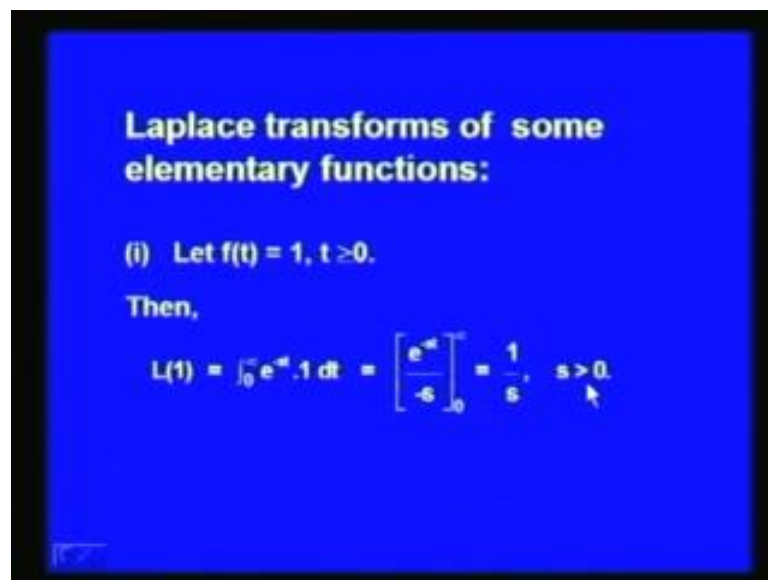
$$\begin{aligned} L (af(t) + bg(t)) &= \int_0^{\infty} e^{-st} (af(t) + bg(t)) dt \\ &= a \int_0^{\infty} e^{-st} f(t) dt + b \int_0^{\infty} e^{-st} g(t) dt \\ &= aL (f(t)) + bL (g(t)). \end{aligned}$$

Now, let us look at the linearity property of the Laplace transform, it says that let  $f t$  and  $g t$  be any two functions whose Laplace transforms exist. Then for any constants  $a$  and  $b$  could be real or complex constants, we have  $L$  of  $a f t$  plus  $b g t$  equal to  $a$  times  $L f t$  plus  $b$  times  $L g t$ .

Let us, look at the proof of this theorem, we can write  $L$  of  $f(t) + b g(t)$  by definition of the Laplace transform as  $\int_0^{\infty} e^{-st} (f(t) + b g(t)) dt$ ; which will be equal to  $a$  times  $\int_0^{\infty} e^{-st} f(t) dt$  plus  $b$  times  $\int_0^{\infty} e^{-st} g(t) dt$ , and which is equal to  $a$  times Laplace transform of  $f(t)$  plus  $b$  times Laplace transform of  $g(t)$ .

We have already assumed that the Laplace transform of  $f(t)$  and  $g(t)$  exists, so the integral that are occurring on the right side here  $\int_0^{\infty} e^{-st} f(t) dt$  exists and  $\int_0^{\infty} e^{-st} g(t) dt$  exists and thus we have the right hand side as  $a$  times  $L f(t)$  plus  $b$  times  $L g(t)$ .

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Now, let us discuss the Laplace transforms of some elementary function, let us assume that  $f(t)$  is the constant function 1 for all  $t$  greater than or equal to 0. Then Laplace transform of 1 by definition will be equal to  $\int_0^{\infty} e^{-st} \cdot 1 dt$ , the integral of  $e^{-st}$  is  $\frac{e^{-st}}{-s}$ ; when  $t$  goes to infinity is the upper limit here, when  $t$  goes to infinity  $e^{-st}$  goes to 0, provided we assume that  $s$  is greater than 0. So, here the condition on  $s$  comes into picture in order to make  $e^{-st}$  go to 0, we have to assume that  $s$  is greater than 0. So, when  $t$  goes to infinity this goes to 0 and at the lower limit its value is  $-\frac{1}{s}$  and therefore, when we substitute the limits we get the result as  $\frac{1}{s}$ , so Laplace transform of 1 is  $\frac{1}{s}$  for all  $s$  greater than 0.

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$$\begin{aligned} \text{(ii) } L(t^a) &= \int_0^{\infty} e^{-st} t^a dt \\ &= \frac{1}{s^{a+1}} \int_0^{\infty} e^{-x} x^a dx, \text{ on putting } st = x \\ &= \frac{\Gamma(a+1)}{s^{a+1}}, \text{ provided } a > -1 \text{ and } s > 0. \end{aligned}$$

In particular, if  $a$  is a positive integer say  $n$  then

$$L(t^n) = \frac{n!}{s^{n+1}}$$

Next let us discuss, the Laplace transform of  $t$  to the power  $a$  let  $f(t)$  be the  $t$  to the power  $a$  here, so the Laplace transform of this is  $\int_0^{\infty} e^{-st} t^a dt$ . Let us now put  $st$  equal to  $x$  then  $e^{-st}$  will become  $e^{-x}$  and  $t^a$  will be  $x^a$  over  $s^a$ . And therefore, we will have the integral we have we have the right hand side as  $\frac{1}{s^{a+1}} \int_0^{\infty} e^{-x} x^a dx$  where we have assumed that  $a > -1$  and  $s > 0$ . Because, when  $t$  goes to infinity the upper limit when we have change the variable from  $t$  to  $x$  will become infinity only when we assume  $s$  to be greater than 0. So, here the assumption is that  $s$  is greater than 0.

Now, this is equal to  $\frac{\Gamma(a+1)}{s^{a+1}}$  if you recall the definition of gamma function, then  $\int_0^{\infty} e^{-x} x^a dx$  becomes  $\Gamma(a+1)$  and therefore, we have the right hand side as  $\frac{\Gamma(a+1)}{s^{a+1}}$ .

And, we know that the gamma function or the integral  $\int_0^{\infty} e^{-x} x^a dx$  is  $\Gamma(a+1)$  provided  $a > -1$ . So, we the condition on  $a$  that it must be greater than minus 1 and we have the condition on  $s$



that it must be positive this can be said that  $s$  is positive we had used here in order to have the limit upper limit of  $s$  as infinity.

Now, let us discuss a particular case of this example let us take  $a$  to be a positive integer say  $a$  equal to  $n$ . Then Laplace transform of  $t$  to the power  $n$  will be  $\Gamma(n+1)$  over  $s$  to the power  $n+1$ , and below that  $\Gamma(n+1)$  is  $n$  factorial if  $n$  is a positive integer. So,  $L$  of  $t$  to the power  $n$  is  $n$  factorial over  $s$  to the power  $n+1$ .

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$$\begin{aligned}
 \text{(iii) } L(e^{at}) &= \int_0^{\infty} e^{-st} e^{at} dt \\
 &= \int_0^{\infty} e^{-(s-a)t} dt \\
 &= \left[ \frac{e^{-(s-a)t}}{-(s-a)} \right]_0^{\infty} = \frac{1}{s-a}, \quad s > a
 \end{aligned}$$

Let us discuss. now the Laplace transform of  $e^{at}$  by definition  $e^{at}$  we may write as integral 0 to infinity  $e^{-st} e^{at} dt$ . And this is equal to integral 0 to infinity  $e^{-(s-a)t} dt$ . When we integrate this we get  $e^{-(s-a)t} / (s-a)$ , integral the limits are 0 and infinity. This equal to now when  $s$  is greater than  $a$   $e^{-(s-a)t}$  will go to 0 as  $t$  goes to infinity. So, this will become 0 when  $t$  goes to infinity and at the lower limit the value is  $-1/(s-a)$ . So, the value of the right hand side is  $1/(s-a)$  provided  $s$  is greater than  $a$ , the condition on  $s$  is required to make this exponential of  $-(s-a)t$  0 as  $t$  goes to infinity.

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(iv)  $L(\cos at) = \frac{s}{s^2 + a^2}$  and  $L(\sin at) = \frac{a}{s^2 + a^2}$   
We have  
$$L(e^{iat}) = \int_0^{\infty} e^{-st} e^{iat} dt$$
$$= \int_0^{\infty} e^{-(s-ia)t} dt$$
$$= \left[ \frac{e^{-(s-ia)t}}{s-ia} \right]_0^{\infty}$$
$$= \frac{1}{s-ia}, \quad s > 0.$$

Let us, now discuss the Laplace transform of  $\cos a t$  and  $\sin a t$ , we will see that Laplace transform of  $\cos a t$  is  $s$  over  $s$  square plus  $a$  square and Laplace transform of  $\sin a t$  is  $a$  over  $s$  square plus  $a$  square. So, in order to find the Laplace transform of  $\cos at$  and  $\sin a t$  we consider Laplace transform of  $e$  to the power  $i a t$ .

We know that by the Euler's formula  $e$  to the power  $i a t$  is  $\cos a t$  plus  $i \sin a t$ , and then we will make use of the linearity property to write the Laplace transformer of  $e$  to the power  $i a t$  as Laplace transform of  $\cos a t$  plus  $i$  times Laplace transform of  $\sin a t$ . And then we will find the right hand side, equate the real and imaginary parts on both sides to get the values of Laplace transform of  $\cos a t$  and  $\sin a t$ .

So, both Laplace transforms can be found just by considering the Laplace transform of  $e$  to the power  $i a t$  and using the linearity property of the Laplace transforms. Now, Laplace transform  $e$  to the power of  $i a t$  can be written as integral  $0$  to infinity  $e$  to the power minus  $s t$  into  $e$  to the power  $i a t$   $d t$  which we can write as integral  $0$  to infinity  $e$  to the power minus  $s$  minus  $a s$  minus  $i a$  into  $t d t$ .

When we take, when we find the integral of this we get  $e$  to the power minus  $s$  minus  $i a$  into  $t$  over  $s$  minus  $i a$  the limits are  $0$ , and infinity. Then  $t$  goes to infinity  $e$  to the power minus  $s$  minus  $i a$  into  $t$ , what is the limit of this, you see we can write  $e$  to the power minus  $s$  minus  $i a$  into  $t$  as  $e$  to the power minus  $s t$  into  $e$  to the power  $i a t$ .

And then if you take the modulus of  $e$  to the power  $st$  into  $e$  to the power  $iat$ , then what you get is  $e$  to the power  $st$ , because modulus of  $e$  to the power  $iat$  is equal to 1. Now,  $e$  to the power  $ist$  minus  $st$ ,  $e$  to the power  $st$  then goes to 0 as  $t$  goes to infinity, so we get the limit of  $e$  to the power  $st$  minus  $iat$  into  $t$ ,  $st$  goes to infinity equal to 0. And therefore, the right hand side is equal to  $1$  over  $s$  minus  $ia$  provided  $s$  is greater than 0.

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Or,  $L(\cos at + i \sin at) = \frac{s + ia}{s^2 + a^2}, s > 0.$

Using linearity property of  $L$  and equating real and imaginary parts we obtain

$$L(\cos at) = \frac{s}{s^2 + a^2}$$

and

$$L(\sin at) = \frac{a}{s^2 + a^2}, s > 0.$$

So, Laplace transform of  $\cos at$  plus  $i \sin at$  is equal to  $s$  plus  $ia$  over  $s$  square plus  $a$  square, we have multiplied  $s$  plus  $ia$  in the numerator and denominator on the right side, in order to bracket into real and imaginary in order to bracket into its real and imaginary part. Now, we will use the linearity property of  $L$ , and then equate the real and imaginary parts to be able to get the values of Laplace transform of  $\cos at$  and Laplace transform of  $\sin at$ . So, when we use the linearity property here, the left hand side becomes Laplace transform of  $\cos at$  plus  $i$  times Laplace transform of  $\sin at$  right hand side is  $s$  over  $s$  square plus  $a$  square plus  $ia$  over  $s$  square plus  $a$  square. So, when we equate the real and imaginary parts on both sides we get the Laplace transform of  $\cos at$  as  $s$  over  $s$  square plus  $a$  square and Laplace transform of  $\sin at$  as  $a$  over  $s$  square plus  $a$  square, where the condition on that  $s$  is greater than 0.

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$$(v) \quad L(\cosh at) = \frac{a}{s^2 - a^2}, \quad L(\sinh at) = \frac{a}{s^2 - a^2}, \quad s > |a|$$
$$L(\cosh at) = L\left(\frac{e^{at} + e^{-at}}{2}\right)$$
$$= \frac{1}{2} \left[ \frac{1}{s-a} + \frac{1}{s+a} \right] = \frac{s}{s^2 - a^2}, \quad s > |a|$$

Similarly, we can show

$$L(\sinh at) = \frac{a}{s^2 - a^2}, \quad s > |a|$$

Now, we will find the Laplace transform of hyperbolic functions the cos hyperbolic  $a t$  and sin hyperbolic  $a t$ , and we shall see that Laplace transform of cos hyperbolic  $a t$  is a over  $s$  square minus  $a$  square, where the Laplace transform of sin hyperbolic  $a t$  is  $a$  over  $s$  square minus  $a$  square the condition on  $s$  is that it must be greater than the absolute value of  $a$ . We know the definition of cos hyperbolic  $a t$  it is equal to  $e$  to the power  $a t$  plus  $e$  to the power minus  $a t$  by 2. So, Laplace transform of cos hyperbolic  $a t$  is Laplace transform of  $e$  to the power  $a t$  plus  $e$  to the power minus  $a t$  by 2.

Now, let us use the linearity property of  $L$  here, so that we can write it as half of Laplace transform of  $e$  to the power  $a t$  plus half of Laplace transform of  $e$  to the power minus  $a t$ . And then we call the Laplace transform of  $e$  to the power  $a t$  it is  $1$  over  $s$  minus  $a$  provided  $s$  is greater than  $a$ , similarly Laplace transform of  $e$  to the power minus  $a t$  is  $1$  over  $s$  plus  $a$  provided  $s$  is greater than minus  $a$ .

Thus, the right hand side of Laplace transform of cos hyperbolic  $a t$  is equal to half of  $1$  over  $s$  minus  $a$  plus  $1$  over  $s$  plus  $a$  provided  $s$  is greater than  $a$  as well as  $s$  is greater than minus  $a$ , which means that  $s$  must be greater than the absolute value of  $a$ . Hence, we get the Laplace transform of cos hyperbolic  $a t$  as  $s$  over  $s$  square minus  $a$  square whenever  $s$  is greater than minus  $a$ . Similarly, we can show that Laplace transform of sin hyperbolic  $a t$  is equal to  $a$  over  $s$  square minus  $a$  square whenever  $s$  is greater than mod of  $a$ .

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**Change of Scale Property:**

**Theorem 2.** If  $L(f(t)) = F(s)$  then

$$L(f(at)) = (1/a) F(s/a), \quad a > 0.$$

**Proof.** By definition, we have

$$L(f(at)) = \int_0^{\infty} e^{-st} f(at) dt$$

Now, let us discuss the change of scale property, let us assume that Laplace transform of  $f(t)$  is equal to  $F(s)$ . Then we shall show that Laplace transform of  $f(at)$  is equal to  $1/a$  times  $F(s/a)$  whenever  $a$  is greater than 0. Now, by definition of the Laplace transform, we can write Laplace transform of  $f(at)$  as integral from 0 to infinity of  $e^{-st} f(at) dt$ .

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$$= \frac{1}{a} \int_0^{\infty} e^{-\left(\frac{s}{a}\right)z} f(z) dz, \quad \text{where } z = at$$
$$= \frac{1}{a} F\left(\frac{s}{a}\right).$$

**Example.** We know that

$$L(\sin t) = \frac{1}{s^2 + 1}$$

Hence,

$$L(\sin 3t) = \frac{1}{3} \frac{1}{(s/3)^2 + 1} = \frac{3}{s^2 + 9}$$

Let us, put  $z$  equal to  $at$  now, so then  $e^{-st}$  will become  $e^{-s(z/a)}$  or  $e^{-\frac{s}{a}z}$ ,  $f(at)$  will become  $f(z)$ ,  $dt$  will become  $dz/a$  and the

limits will become 0, the lower limit will be 0, because when  $t$  is 0,  $z$  is 0. The upper limit will be infinity because we have assumed  $a$  to be positive, so whenever  $t$  goes to infinity  $z$  goes to infinity. And thus the Laplace transform of  $f(at)$  is equal to  $\frac{1}{a}$  over an integral over 0 to infinity of  $e^{-sz}$  into  $f(z) dz$ .

And, which is equal to  $\frac{1}{a}$  into  $f(s)$  by a if  $F(s)$  is equal to we know that  $F(s)$  is equal to integral 0 to infinity of  $e^{-st}$  into  $f(t) dt$ . So, in place of the variability here we have the variable  $z$ , and  $s$  in  $F(s)$  is replaced by  $s/a$  here, so we get the right hand side here as  $\frac{1}{a}$  into  $f(s/a)$ .

Let us, look at an example which is based on this change of scale property, we know that Laplace transform of  $\sin t$  is equal to  $\frac{1}{s^2 + 1}$ , because Laplace transform we have shown that Laplace transform of  $\sin at$  is  $\frac{1}{s^2 + a^2}$ . So, taking  $a$  equal to 1 we get the Laplace transform of  $\sin t$  as  $\frac{1}{s^2 + 1}$ .

Hence, Laplace transform of  $\sin 3t$ , so here we are taking  $f(t)$  is equal to  $\sin t$   $f(at)$  is equal to  $\sin 3t$ . So,  $a$  is equal to 3 here, so  $a$  is clearly positive. So, Laplace transform of  $\sin 3t$  by change of scale property will be equal to  $\frac{1}{3}$  into  $\frac{1}{(s/3)^2 + 1}$  which is equal to  $\frac{1}{3}$  into  $\frac{1}{s^2/9 + 1}$ .  $F(s)$  here  $F(s)$  is this,  $F(s)$  is  $\frac{1}{s^2 + 1}$ , so we replace  $F(s)$  by  $f(s)$  by 3; that means, we replace  $s$  by  $s/3$ , so we get  $\frac{1}{3}$  into  $\frac{1}{s^2/9 + 1}$  which is equal to  $\frac{3}{s^2 + 9}$ . Now, this Laplace transform of  $\sin 3t$  we can get directly also by using the formula of Laplace transform of  $\sin at$  equal to  $\frac{a}{s^2 + a^2}$  in that you put  $a$  equal to 3, you get Laplace transform of  $\sin 3t$  as  $\frac{3}{s^2 + 9}$ .

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**Example.** Find the Laplace transform of the Bessel function  $J_0$  and hence find  $L(J_0(at))$ ,  $a > 0$ .

**Solution.** We know that

$$J_0(t) = \left( 1 - \frac{t^2}{2^2} + \frac{t^4}{2^2 \cdot 4^2} - \frac{t^6}{2^2 \cdot 4^2 \cdot 6^2} + \dots \right)$$

Taking Laplace transform of both sides, we have

$$L(J_0(t)) = \frac{1}{s} - \frac{1}{2^2} \frac{2!}{s^3} + \frac{1}{2^2 \cdot 4^2} \frac{4!}{s^5} - \frac{1}{2^2 \cdot 4^2 \cdot 6^2} \frac{6!}{s^7} + \dots$$

$$= \frac{1}{s} \left[ 1 - \frac{1}{2} \left( \frac{1}{s^2} \right) + \frac{1 \cdot 3}{2 \cdot 4} \left( \frac{1}{s^4} \right) - \frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6} \left( \frac{1}{s^6} \right) + \dots \right]$$

Let us, now find the Laplace transform of the Bessel function of order 0 which we denoted by  $J_0(t)$ . And then find the Laplace transform of  $J_0(at)$  where  $a$  is greater than 0, we shall make use of the change of scale property here to find the Laplace transform of  $J_0(at)$  once we have the Laplace transform of  $J_0(t)$ .

Now, we know that the Bessel function of order 0 that is  $J_0(t)$  is given by  $1 - \frac{t^2}{2^2} + \frac{t^4}{2^2 \cdot 4^2} - \frac{t^6}{2^2 \cdot 4^2 \cdot 6^2} + \dots$  and so on.

So, when we take the Laplace transform on both sides we get Laplace transform of  $J_0(t)$  as Laplace transform of 1 which is equal to  $\frac{1}{s}$  minus Laplace transform of  $\frac{t^2}{2^2}$  which is  $\frac{1}{2^2} \frac{2!}{s^3}$  plus Laplace transform of  $\frac{t^4}{2^2 \cdot 4^2}$  which is  $\frac{1}{2^2 \cdot 4^2} \frac{4!}{s^5}$ , and then minus  $\frac{1}{2^2 \cdot 4^2 \cdot 6^2} \frac{6!}{s^7}$ . Laplace transform of 6 similarly is  $\frac{6!}{s^7}$ .

Now, let us take  $\frac{1}{s}$  common from all the terms, then we get  $\frac{1}{s}$  times  $1 - \frac{1}{2} \left( \frac{1}{s^2} \right) + \frac{1 \cdot 3}{2 \cdot 4} \left( \frac{1}{s^4} \right) - \frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6} \left( \frac{1}{s^6} \right) + \dots$ . So, we have written the right hand side in a proper form, so that we can write it as follows.

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$$\begin{aligned}
 &= \frac{1}{s} \left[ 1 + \left(-\frac{1}{2}\right) \left(\frac{1}{s^2}\right) + \frac{\left(-\frac{1}{2}\right) \left(-\frac{3}{2}\right)}{2!} \left(\frac{1}{s^2}\right)^2 \right. \\
 &\quad \left. + \frac{\left(-\frac{1}{2}\right) \left(-\frac{3}{2}\right) \left(-\frac{5}{2}\right)}{3!} \left(\frac{1}{s^2}\right)^3 \dots \right] \\
 &= \frac{1}{s} \left(1 + \frac{1}{s^2}\right)^{-1/2} \quad \text{(By Binomial theorem)} \\
 &= \frac{1}{\sqrt{s^2 + 1}} \quad \leftarrow
 \end{aligned}$$

This is further equal to 1 by s times 1 plus minus half into 1 by s square plus minus half into minus 3 by 2 over 2 factorial into 1 by s square over raise to the power 2 plus minus half into minus 3 by 2 into minus 5 by 2 over 3 factorial into 1 by s square raise to the power 3 and so on.

And, this expression inside the brackets we can write as 1 plus 1 by s square raise to the power minus half. So, which follows by binomial theorem we know that 1 plus 1 by s square raise to the power minus half can be written as infinite series given here, inside the bracket. So, this is through right hand side is equal to 1 by s into 1 plus 1 by s square raise to the power minus half provided s is greater than 1. And, this after simplification gives us 1 over under root s square plus 1. So, Laplace transform of the Bessel function of order 0 is equal to s square plus 1 raise to the power minus half.



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We know that

$$L(f(at)) = \frac{1}{a} F\left(\frac{s}{a}\right)$$

if  $L(f(t)) = F(s)$ .

Therefore,

$$L(J_0(at)) = \frac{1}{a} \frac{1}{\sqrt{\frac{s^2}{a^2} + 1}} = \frac{1}{\sqrt{s^2 + a^2}}.$$

Now, we know that by change of scale property Laplace transform of  $f$  of  $t$  is  $1$  over  $a$  times  $F$  of  $s$  by  $a$  provided  $a$  is greater than  $0$ . So, let us apply this  $F$   $s$  here, we let us we call that  $F$   $s$  here is  $L$  of  $f$   $t$ . So, making use of this change of scale property we can write the Laplace transform of  $J_0$   $at$  as  $1$  over  $a$  times  $f$  of  $s$  by  $a$  and  $F$   $s$ , we have seen is  $1$  over square root  $s$  square plus  $1$ . So, in that  $s$  is replaced by  $s$  over  $a$  to have the Laplace transform of  $J_0$   $at$  which is equal to  $1$  over  $a$  into  $1$  by under root  $s$  square by  $a$  square plus  $1$  which after simplification gives us  $1$  over square root  $s$  square plus  $a$  square.

(Refer Slide Time: 25:08)

Example. Show that  $\int_0^{\infty} \cos x^2 dx = \frac{1}{2} \sqrt{\pi/2}$ .

Solution. Let  $F(t) = \int_0^{\infty} \cos(tx^2) dx$ .

Then  $L(F(t)) = \int_0^{\infty} e^{-st} F(t) dt$

$$= \int_0^{\infty} e^{-st} \left[ \int_0^{\infty} \cos(tx^2) dx \right] dt$$
$$= \int_0^{\infty} \left[ \int_0^{\infty} e^{-st} \cos(tx^2) dt \right] dx$$

Now, let us take an example where we have to evaluate the integral of 0 to infinity  $\cos x$  square  $d x$ , we know that the integral of  $\cos x$  square  $d x$  over 0 infinity cannot be evaluated by the known methods of integration. So, we will see that this can be obtained using the Laplace transform technique, let us assume that  $f t$  is equal to integral 0 to infinity  $\cos$  of  $t$  into  $x$  square  $d x$ .

Then, Laplace transform of  $f t$  which is equal to integral 0 to infinity  $e$  to the power minus  $s t$  into  $f t d t$  by definition will become equal to integral 0 to infinity  $e$  to the power minus  $s t$  into integral 0 to infinity  $\cos$  of  $t$  into  $x$  square  $d x$  into  $d t$  changing the order of integration here, we have integral 0 to infinity into integral 0 to infinity, then integral 0 to infinity  $e$  to the power minus  $s t \cos t x$  square  $d t$  into  $d x$ . Now, this inner integral which is integral 0 to infinity  $e$  to the power minus  $s t \cos t x$  square  $d t$  this is nothing but Laplace transform of  $\cos t x$  square.

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$$= \int_0^{\infty} L(\cos(x^2)) dx = \int_0^{\infty} \frac{s}{s^2 + x^4} dx.$$

Putting  $x = \sqrt{s \tan \theta}$  so that  $dx = \frac{s \sec^2 \theta d\theta}{2\sqrt{s \tan \theta}}$ ,

we get  $L(F(t)) = \int_0^{\pi/2} \frac{1}{2\sqrt{s \tan \theta}} d\theta$

$$= \frac{1}{2\sqrt{s}} \int_0^{\pi/2} (\sin \theta)^{-1/2} (\cos \theta)^{1/2} d\theta$$

$$= \frac{1}{2\sqrt{s}} \frac{\Gamma(1/4)\Gamma(3/4)}{2\Gamma(1)} = \frac{\Gamma(1/4)\Gamma(1-(1/4))}{4\sqrt{s}}$$

So, we get the right hand side as integral 0 to infinity  $\cos$  Laplace transform of  $\cos t x$  square into  $d x$  and Laplace transform of  $\cos a t$  we know, it is  $s$  over  $s$  square plus  $a$  square. So here,  $a$  is equal to  $x$  square, and therefore Laplace transform of  $\cos$  of  $t x$  square is  $s$  over  $s$  square plus  $x$  to the power 4. So, right hand side becomes integral 0 to infinity  $s$  over  $s$  square plus  $x$  to the power 4  $d x$ .

Now, let us put  $x$  equal to square root of  $s \tan \theta$  here, so then we shall have  $d x$  equal to  $1$  over  $2$  root  $s \tan \theta$  into  $s$  into  $\sec$  square  $\theta d \theta$ . And therefore, we will get

the Laplace transform of  $f(t)$  as  $\int_0^{\pi/2} \frac{1}{\sqrt{s}} \sin(\theta) d\theta$ , now here we see that when  $x$  is  $0$   $\theta$  is  $0$ .

And then  $x$  goes to infinity,  $\theta$  goes to  $\pi/2$ , so the limit lower limit of  $\theta$  is  $0$  while the upper limit of  $\theta$  is  $\pi/2$ , and thus we get the Laplace transform of  $f(t)$  as  $\int_0^{\pi/2} \frac{1}{\sqrt{s}} \sin(\theta) d\theta$ , which is equal to  $\int_0^{\pi/2} \frac{1}{\sqrt{s}} \sin(\theta) d\theta$ ,  $\sin(\theta)$  raise to the power minus half into  $\cos(\theta)$  raise to the power half  $d\theta$ .

And, the value of this integral we can obtain by using the gamma function  $\int_0^{\pi/2} \sin^m(\theta) \cos^n(\theta) d\theta = \frac{\Gamma(m+1) \Gamma(n+1)}{2 \Gamma(m+n+2)}$  is as it is here, and then we this we know that  $\int_0^{\pi/2} \sin^m(\theta) \cos^n(\theta) d\theta = \frac{\Gamma(m+1) \Gamma(n+1)}{2 \Gamma(m+n+2)}$  is  $\frac{\Gamma(m+1) \Gamma(n+1)}{2 \Gamma(m+n+2)}$  divided by  $2$  times  $\Gamma(m+n+2)$ .

So, making use of that formula we have here  $m$  equal to minus half  $n$  equal to half, so we get  $\Gamma(m+1) \Gamma(n+1)$  that is  $\Gamma(1/2) \Gamma(3/2)$ . Then  $\Gamma(m+n+2)$  that is  $\Gamma(2)$  over  $2$  times  $\Gamma(m+n+2)$ ,  $m$  is minus half  $n$  is half, so  $m+n+2$  gives you  $2$ . And therefore, we have the right hand side equal to  $\frac{\Gamma(1/2) \Gamma(3/2)}{2 \Gamma(2)}$  which is  $\frac{\Gamma(1/2) \Gamma(3/2)}{2}$  divided by  $2$  into square root  $s$  into  $\Gamma(1/2)$  which is equal to  $1$ .

(Refer Slide Time: 29:06)

$$= \frac{1}{4\sqrt{s}} \cdot \frac{\pi}{\sin(\pi/4)} \quad \left( \because \Gamma(n)\Gamma(1-n) = \frac{\pi}{\sin n\pi}, 0 < n < 1. \right)$$

$$= \frac{\pi}{2\sqrt{2}} \cdot \frac{1}{\sqrt{s}}$$

Therefore,

$$F(t) = \frac{\pi}{2\sqrt{2}} L^{-1} \left\{ \frac{1}{\sqrt{s}} \right\} = \frac{\pi}{2\sqrt{2}} \frac{1}{\sqrt{\pi t}} = \frac{1}{2} \left( \frac{\pi}{2t} \right)^{1/2}$$

or

$$F(t) = \int_0^{\infty} \cos(tx^2) dx = \frac{1}{2} \left( \frac{\pi}{2t} \right)^{1/2}$$

Putting  $t=1$ , we get

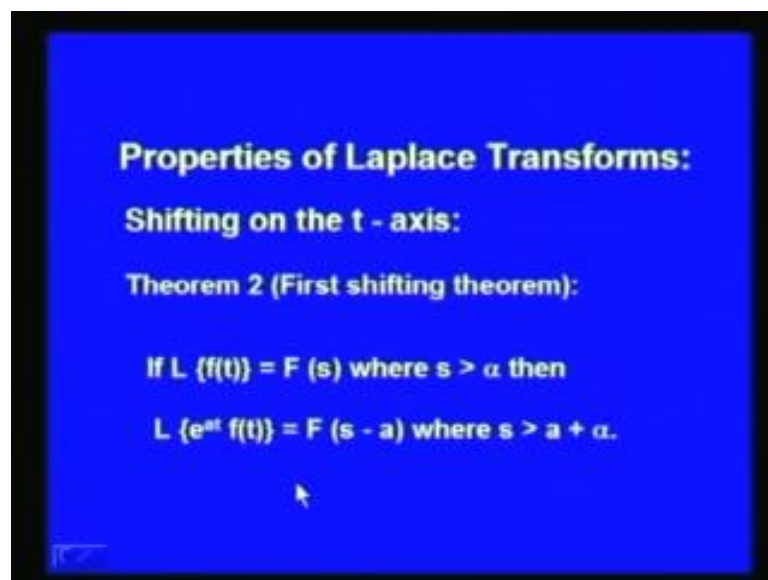
$$\int_0^{\infty} \cos x^2 dx = \frac{1}{2} \sqrt{\pi/2}$$

And then the right hand side becomes equal to  $\frac{1}{2} \frac{\Gamma(4) \Gamma(s)}{\Gamma(s+4)} \sin \frac{\pi}{4}$ , because we know that  $\Gamma(n) \Gamma(1-n)$  is equal to  $\frac{\pi}{\sin n\pi}$  whenever  $0 < n < 1$ . So, this is equal to  $\frac{\pi}{2} \frac{2 \cdot 3 \cdot 2 \cdot 1}{\Gamma(s+4)}$ . And therefore,  $f(t)$  is equal to  $\frac{\pi}{2} \frac{2 \cdot 3 \cdot 2 \cdot 1}{\Gamma(s+4)}$  inverse Laplace transform of  $\frac{1}{\Gamma(s+4)}$ .

Now, we know that Laplace transform of  $t^a$  is equal to  $\frac{\Gamma(a+1)}{s^{a+1}}$  whenever  $a > -1$ . So, taking  $a$  equal to  $-\frac{1}{2}$  there, we get the Laplace transform of  $t^{-\frac{1}{2}}$  is  $\frac{\Gamma(\frac{1}{2})}{s^{\frac{1}{2}}}$ . And therefore, Laplace transform of  $\frac{1}{\Gamma(s+4)}$  is  $\frac{1}{\Gamma(s+4)}$ , so we get the function  $f(t) = \frac{\pi}{2} \frac{2 \cdot 3 \cdot 2 \cdot 1}{\Gamma(s+4)}$ .

And, which is equal to  $\frac{1}{2} \frac{\pi}{2} t^{\frac{1}{2}}$ , and thus  $f(t)$  becomes equal to  $\int_0^\infty \cos t x^2 dx = \frac{1}{2} \frac{\pi}{2} t^{\frac{1}{2}}$ . Now, let us take  $t$  equal to 1 to have the value of the desired integral, we get the desired integral  $\int_0^\infty \cos x^2 dx = \frac{1}{2} \frac{\pi}{2}$ , that is square root of  $\frac{\pi}{2}$ .

(Refer Slide Time: 30:50)



Now, we are going to discuss an important property of the Laplace transform, we call it first shifting theorem. If  $L \{f(t)\} = F (s)$  where  $s$  is greater than  $\alpha$  then  $L \{e^{at} f(t)\} = F (s - a)$ , where  $s$  is greater than  $a + \alpha$ .

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**Proof: We have**

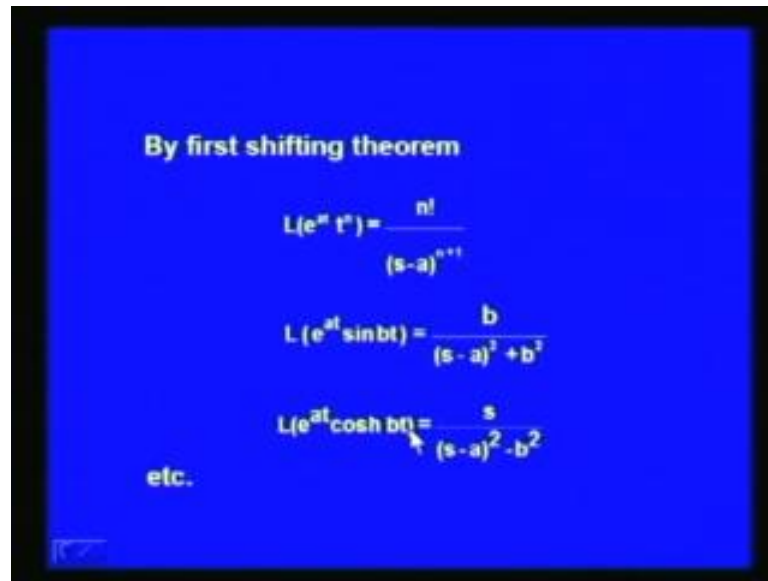
$$\begin{aligned}(2) \quad F(s - a) &= \int_0^{\infty} e^{-(s-a)t} f(t) dt \\ &= \int_0^{\infty} e^{-st} \{e^{at} f(t)\} dt \\ &= L\{e^{at} f(t)\}.\end{aligned}$$

Since  $F(s)$  exists for  $s > \alpha$ , it follows that the integral in (2) exists for  $s > a + \alpha$ .

This theorem tells us that  $s$  is replaced by  $s$  minus  $a$  in the Laplace transform of the function  $f(t)$ , then  $f(t)$  is multiplied by  $e$  to the power  $a t$ . So, when we shift from  $s$  to  $s$  minus  $a$  we get that  $f(t)$  is multiplied by  $e$  to the power  $a t$ . Now, let us prove this theorem by definition of the Laplace transform we can write  $F(s - a)$  equal to integral 0 to infinity  $e^{-(s-a)t} f(t) dt$  replacing  $s$  by  $s - a$  in the definition of the Laplace transform of the function  $f(t)$ .

And, this  $e^{-(s-a)t}$ , we can write as  $e^{-st} e^{at}$  into  $e^{at}$  we can combine with the function  $f(t)$ . So, that we get the right hand side as integral 0 to infinity  $e^{-st} e^{at} f(t) dt$ , and we can then write it as the Laplace transform of  $e^{at} f(t)$ . Since,  $F(s)$  exists for  $s$  greater than  $\alpha$  by our assumption, therefore the integral here the integral here exists for  $s - a$  greater than  $\alpha$ , that is  $s$  greater than  $a + \alpha$ . And therefore, we may say that  $L\{e^{at} f(t)\}$  is equal to  $F(s - a)$  provided  $s$  is greater than  $a + \alpha$ .

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By first shifting theorem

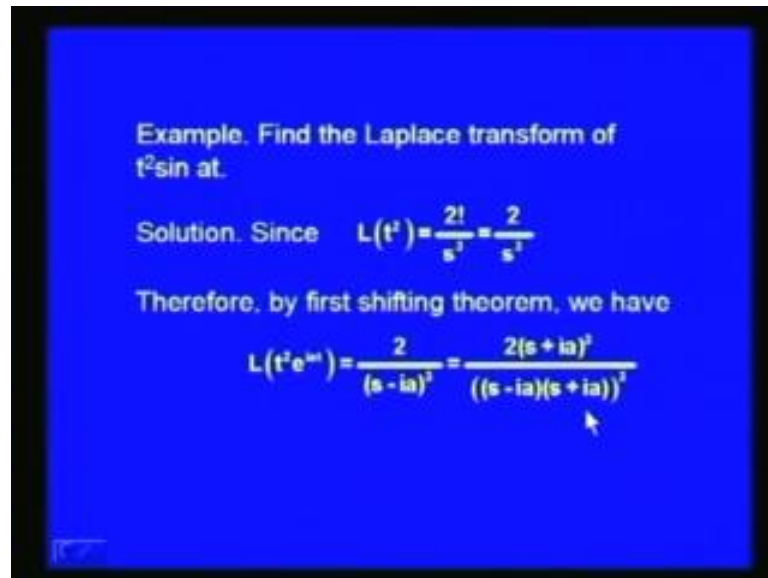
$$L(e^{at} t^n) = \frac{n!}{(s-a)^{n+1}}$$
$$L(e^{at} \sin bt) = \frac{b}{(s-a)^2 + b^2}$$
$$L(e^{at} \cosh bt) = \frac{s}{(s-a)^2 - b^2}$$

etc.

Now, by the first shifting theorem, we know that Laplace transform of  $t$  to the power  $n$  is  $n$  factorial over  $s$  to the power  $n$  plus 1. So, if you multiply  $t$  to the power  $n$  by  $e$  to the power  $a$   $t$  by the first shifting theorem we can say that Laplace transform of  $e$  to the power  $a$   $t$  into  $t$  to the power  $n$  will be equal to  $n$  factorial over  $s$  minus  $a$  to the power  $n$  plus 1 replacing  $s$  by  $s$  minus  $a$ , here we are assuming that  $n$  is a positive integer.

Now, similarly we know that Laplace transform of  $\sin b t$  is  $b$  over  $s$  square plus  $b$  square. So, when we multiply  $\sin b t$  by  $e$  to the power  $a$   $t$  by first shifting theorem, Laplace transform of  $e$  to the power  $a$   $t$  into  $\sin b t$  will be equal to  $b$  over  $s$  minus  $a$  whole square plus  $b$  square. And, similarly Laplace transform of  $e$  to the power  $a$   $t$  into  $\cosh b t$  is equal to  $s$  over  $s$  minus  $a$  whole square minus  $b$  square, because we know that Laplace transform of  $\cosh b t$  is  $s$  over  $s$  square minus  $b$  square. So, by using first shifting theorem, the Laplace transform of  $e$  to the power  $a$   $t$   $\cosh b t$  will be obtained by replacing  $s$  by  $s$  minus  $a$  in the Laplace transform of  $\cosh b t$ .

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Example. Find the Laplace transform of  $t^2 \sin at$ .

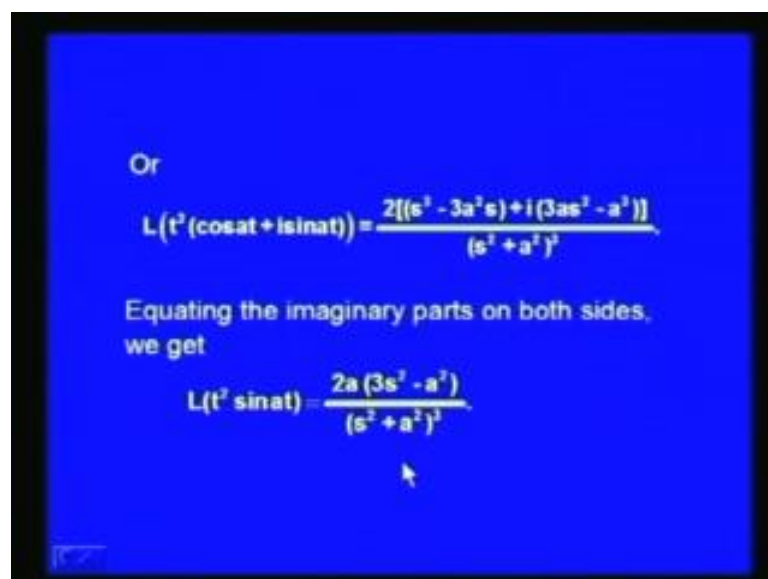
Solution. Since  $L(t^2) = \frac{2!}{s^3} = \frac{2}{s^3}$

Therefore, by first shifting theorem, we have

$$L(t^2 e^{iat}) = \frac{2}{(s-ia)^3} = \frac{2(s+ia)^3}{((s-ia)(s+ia))^3}$$

Now, let us find the Laplace transform of  $t^2 \sin at$ . We know that the Laplace transform of  $t^2$  is  $\frac{2!}{s^3}$ , that is, it is equal to  $\frac{2}{s^3}$ . And therefore, by the first shifting theorem, we can write the Laplace transform of  $t^2 \sin at$  as equal to  $\frac{2}{(s-ia)^3}$ , replacing  $s$  by  $s-ia$  and which is equal to  $\frac{2(s+ia)^3}{(s^2+a^2)^3}$ .

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Or

$$L(t^2 (\cos at + i \sin at)) = \frac{2[(s^2 - 3a^2 s) + i(3as^2 - a^3)]}{(s^2 + a^2)^3}$$

Equating the imaginary parts on both sides, we get

$$L(t^2 \sin at) = \frac{2a(3s^2 - a^2)}{(s^2 + a^2)^3}$$

Now, we can write it as  $L\{t^2 e^{iat}\} = \cos at + i \sin at$  and then right hand side is  $\frac{2s^3 - 3a^2s + i(3as^2 - a^3)}{s^2 + a^2}$ . So, equating the imaginary parts both sides we get Laplace transform of  $t^2 \sin at$  as equal to  $\frac{2a(3s^2 - a^2)}{s^2 + a^2}$ .

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**Example:**  
 Let us determine  $L^{-1}\left\{\frac{1}{s(s^2+9)}\right\}$

**Solution :** We have  

$$F(s) = \frac{1}{s(s^2+9)} = \frac{1}{9}\left(\frac{1}{s} - \frac{s}{s^2+9}\right)$$

then  

$$f(t) = L^{-1}(F(s)) = \frac{1}{9}L^{-1}\left(\frac{1}{s} - \frac{s}{s^2+9}\right)$$

$$= \frac{1}{9}\left[L^{-1}\left(\frac{1}{s}\right) - L^{-1}\left(\frac{s}{s^2+9}\right)\right]$$

$$= \frac{1}{9}(1 - \cos 3t).$$

Let us now determine the inverse Laplace transform of  $\frac{1}{s(s^2+9)}$ . Let us denote by  $F(s)$  the function  $\frac{1}{s(s^2+9)}$ . So,  $F(s)$  is equal to  $\frac{1}{s(s^2+9)}$ . Breaking  $\frac{1}{s(s^2+9)}$  into its partial fractions we get  $\frac{1}{9}\left(\frac{1}{s} - \frac{s}{s^2+9}\right)$ . And then  $f(t)$  is equal to inverse Laplace transform of  $F(s)$ , so  $L^{-1}$  of  $F(s)$  is equal to  $\frac{1}{9}$  using linearity property we have  $\frac{1}{9}L^{-1}\left(\frac{1}{s} - \frac{s}{s^2+9}\right)$  which is equal to  $\frac{1}{9}\left(L^{-1}\left(\frac{1}{s}\right) - L^{-1}\left(\frac{s}{s^2+9}\right)\right)$ . And, we know that inverse Laplace transform of  $\frac{1}{s}$  is 1 and inverse Laplace transform of  $\frac{s}{s^2+9}$  is  $\cos 3t$ . So, inverse Laplace transform of  $F(s)$  which is  $f(t)$  is equal to  $\frac{1}{9}(1 - \cos 3t)$ .



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**Piecewise Continuous function.**

A function  $f(t)$  is called piecewise continuous on a finite interval  $[a, b]$  if  $f(t)$  is defined on  $[a, b]$  and is such that the interval can be subdivided into finitely many intervals, in each of which

- (i)  $f(t)$  is continuous and
- (ii) has finite limits as  $t$  approaches either endpoint of the interval of subdivision from the interior.

Now, let us we are going to discuss the existence theorem for the Laplace transform, so before that we define what do we mean by a piecewise continuous function. A function  $f(t)$  defined over an interval  $a$  to  $b$  is said to be piecewise continuous. If it is if the interval  $a$  to  $b$  can be subdivided into finitely many sub intervals in each of which the function  $f(t)$  is continuous and has finite limits as  $t$  approaches either end point of the interval of subdivision from the interior that is the function  $f(t)$  has only ordinary discontinuities inside the interval  $a$  to  $b$ .

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**Example. The function**

$$f(t) = \begin{cases} t, & 0 < t < \frac{1}{2} \\ t-1, & \frac{1}{2} < t < 1. \end{cases}$$

is piecewise continuous on the interval  $[0, 1]$ .

Let us, look at the function  $f(t)$  equal to  $t$  where  $t$  takes values in the interval  $0 \leq t < \frac{1}{2}$  or  $t \geq \frac{1}{2}$ . And  $f(t)$  is defined as  $t - 1$  when  $\frac{1}{2} < t \leq 1$  or  $t > 1$ . So, you can see that this function is piecewise continuous, because the left hand limit of this function here is half, while the right hand limit is minus half. The function is defined over the whole interval  $[0, 1]$  and the interval  $[0, 1]$  has been divided into 2 sub intervals  $[0, \frac{1}{2}]$  and  $[\frac{1}{2}, 1]$  in each of which the function  $f(t)$  is continuous.

And, as  $t$  approaches to the point of discontinuity that is half, the left hand limit is half while the right hand limit is minus half. So, the function  $f(t)$  has an ordinary discontinuity at the point half, by the function  $f$  is said to be have said to have an ordinary discontinuity at a point in the interval  $[a, b]$  if it has finite on either side. So, of the point of discontinuity, so here the function  $f$  has ordinary discontinuity at the point half.

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Example. The function

$$f(t) = \begin{cases} t, & 0 \leq t < \frac{1}{2} \\ t-1, & \frac{1}{2} < t \leq 1 \\ 0, & t > 1. \end{cases}$$

We have  $L(f(t)) = \int_0^{\infty} e^{-st} f(t) dt$

$$= \int_0^{\frac{1}{2}} t e^{-st} dt + \int_{\frac{1}{2}}^1 (t-1) e^{-st} dt + 0$$

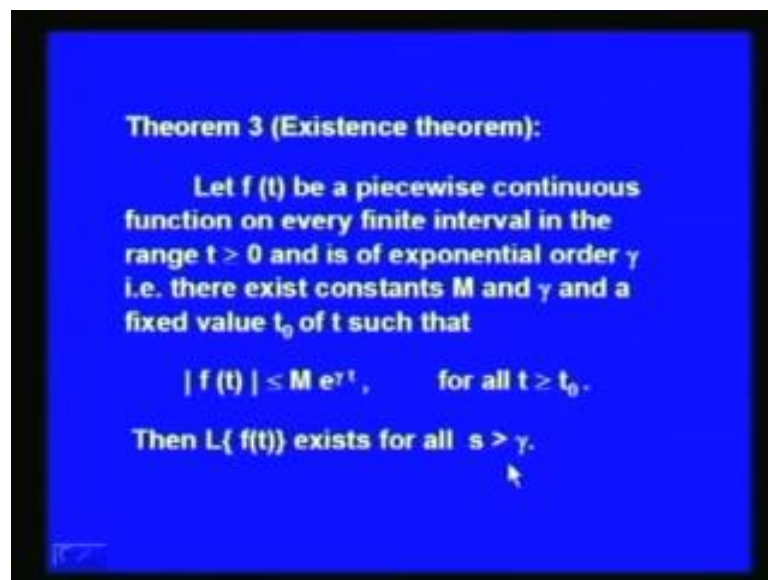
$$= \frac{1}{s^2} - \frac{1}{s} e^{-s/2} - \frac{1}{s^2} e^{-s}.$$

And, we note here that if a function is continuous it is almost piecewise continuous, but the converse is not true. Now, let us take an example of the function  $f(t)$  defined as  $f(t)$  equal to  $t$  when  $0 \leq t < \frac{1}{2}$  or  $t \geq \frac{1}{2}$  or  $t > 1$  and we define it equal to  $t - 1$  when  $\frac{1}{2} < t \leq 1$ . So, then let us see what is the Laplace transform of this by definition Laplace transform  $f(t)$  is  $\int_0^{\infty} e^{-st} f(t) dt$ , which is which is equal to  $\int_0^{\frac{1}{2}} t e^{-st} dt + \int_{\frac{1}{2}}^1 (t-1) e^{-st} dt + \int_1^{\infty} 0 \cdot e^{-st} dt$

$\int_0^{\infty} e^{-st} f(t) dt$ . The third integral which is integral over half 1 to infinity, because 0 because the function  $f$  is defined as 0 there.

Now, let us see value of the values of these two integrals, the integral first integral and the second integral here, when we evaluate their values and put the limits it turns out that Laplace transform of  $f(t)$  is equal to  $\frac{1}{s^2} - \frac{1}{s} + \int_0^{\infty} e^{-st} f(t) dt$ .

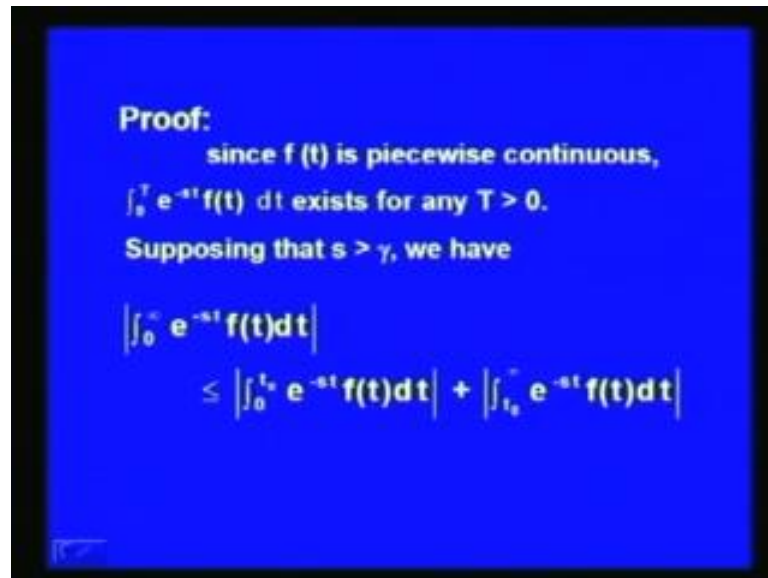
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Now, we discuss the existence theorem for the Laplace transform, that is we are going to see the conditions under which the Laplace transform of a given function  $f$  exists. So, let us assume that  $f(t)$  is a piecewise continuous function on every finite interval in the range  $t$  greater than or equal to 0 and is of exponential order  $\gamma$ , that is we assume that there exist constants  $M$  and  $\gamma$  and a fixed value  $t_0$  of  $t$  such that  $\text{mod of } f(t)$  is less than or equal to  $M$  times  $e^{\gamma t}$  for all  $t$  greater than or equal to 0.

In simple terms, we can say that  $f$  function  $f$  is said to be of exponential order  $\gamma$  if  $\lim_{t \rightarrow \infty} \frac{e^{-\gamma t} f(t)}{e^{-\gamma t}}$  exists and is finite as  $t$  goes to infinity. So, then Laplace transform of the function  $f(t)$  exists for all  $s$  greater than  $\gamma$ .

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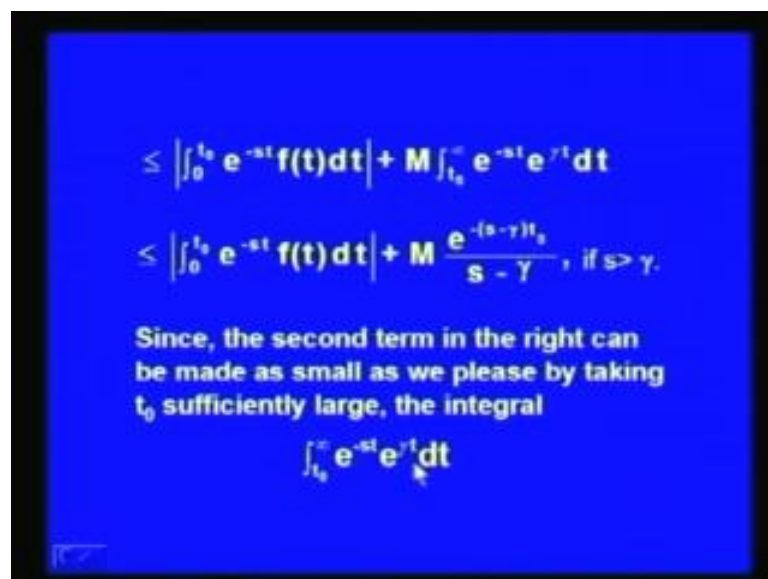


**Proof:**  
since  $f(t)$  is piecewise continuous,  
 $\int_0^T e^{-st} f(t) dt$  exists for any  $T > 0$ .  
Supposing that  $s > \gamma$ , we have

$$\left| \int_0^\infty e^{-st} f(t) dt \right| \leq \left| \int_0^{t_0} e^{-st} f(t) dt \right| + \left| \int_{t_0}^\infty e^{-st} f(t) dt \right|$$

Let us, look at the proof of this theorem we have assume that  $f(t)$  is the piecewise continuous function, therefore integral over 0 to  $t_0$  of  $e^{-st} f(t) dt$  exists for any  $T$  greater than 0. Now, supposing that  $s$  is greater than  $\gamma$ , we can write the modulus of integral 0 to infinity  $e^{-st} f(t) dt$  less than or equal to modulus of integral 0 to  $t_0$   $e^{-st} f(t) dt$  plus modulus of integral  $t_0$  to infinity  $e^{-st} f(t) dt$ . We bracket into 2 parts integral over  $t_0$  0 to  $t_0$ .

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$$\leq \left| \int_0^{t_0} e^{-st} f(t) dt \right| + M \int_{t_0}^\infty e^{-st} e^{\gamma t} dt$$
$$\leq \left| \int_0^{t_0} e^{-st} f(t) dt \right| + M \frac{e^{-(s-\gamma)t_0}}{s-\gamma}, \text{ if } s > \gamma.$$

Since, the second term in the right can be made as small as we please by taking  $t_0$  sufficiently large, the integral

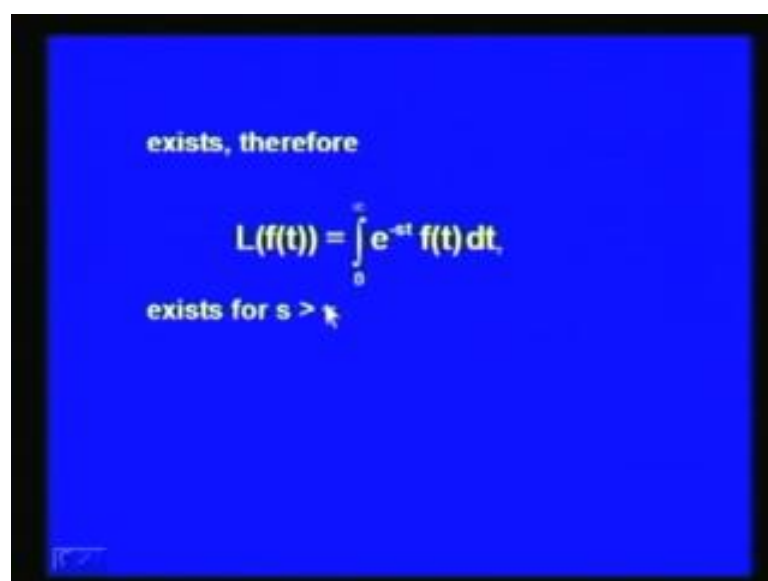
$$\int_{t_0}^\infty e^{-st} e^{\gamma t} dt$$

And then  $t$  naught to infinity,  $t$  naught is the value after which all values of  $t$  for all values of  $t$  mod of  $f(t)$  less than or equal to  $M$  times  $e^{-\gamma t}$ , which is less than or equal to modulus of integral 0 to  $t$  naught,  $e^{-s t}$  into  $f(t) dt$  plus  $M$  times integral  $t$  naught to infinity  $e^{-s t}$  into  $e^{-\gamma t}$   $dt$  and which is further less than or equal to modulus of integral 0 to  $t$  naught  $e^{-s t}$  into  $f(t) dt$ , plus  $M$  times  $e^{-s t}$  into  $e^{-\gamma t}$   $dt$  over  $s - \gamma$ . The integral of  $e^{-s t}$  into  $e^{-\gamma t}$   $dt$  is  $e^{-s t}$  into  $e^{-\gamma t}$  over  $s - \gamma$  if we assume  $s$  to be greater than  $\gamma$  then  $s - \gamma$  goes to infinity.

The value of  $e^{-s t}$  into  $e^{-\gamma t}$  is 0, so at the lower limit its value is  $e^{-s t}$  into  $e^{-\gamma t}$  at  $t$  naught. And therefore, the value of the integral is integral  $t$  naught to infinity  $e^{-s t}$  into  $e^{-\gamma t}$   $dt$  is  $e^{-s t}$  into  $e^{-\gamma t}$  over  $s - \gamma$  wherever  $s$  is greater than  $\gamma$ .

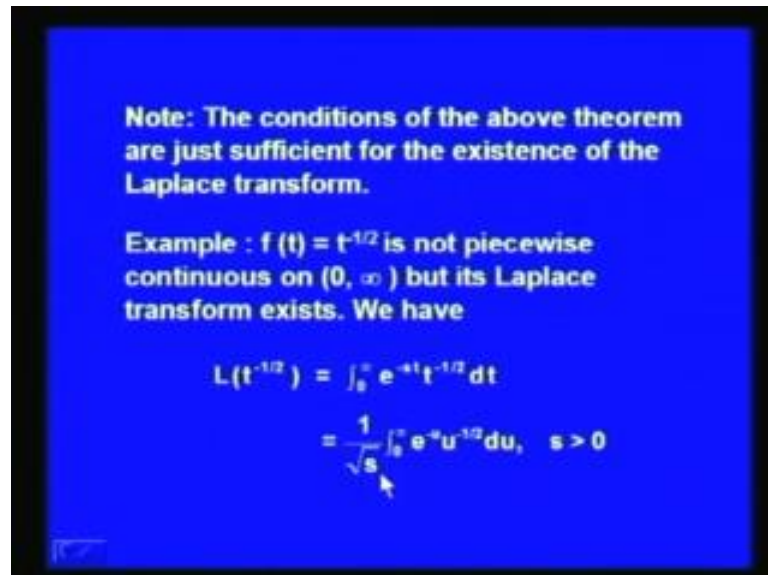
And, the second term here that is  $e^{-s t}$  into  $e^{-\gamma t}$  over  $s - \gamma$  can be made as small as we please by taking  $t$  naught to be sufficiently large. And therefore, the integral  $t$  naught to infinity  $e^{-s t}$  into  $e^{-\gamma t}$   $dt$  exists.

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And, thus Laplace transform of  $f(t)$  which is  $t^{-1/2}$  exists for all values of  $s$  greater than  $\gamma$ .

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Now, we note here that the conditions of the above theorem are just sufficient for the existence of the Laplace transform.

For example, let us consider the function  $f(t)$  equal to  $t$  to the power minus half we can see that this function is not piecewise continuous over the interval  $0$  to infinity, because  $s$  goes to  $0$  plus that is  $s$  goes to  $0$  from the right  $t$  to the power minus half goes to plus infinity. And therefore, it is not a piecewise continuous function.

But, we see that its Laplace transform exists, its Laplace transform is  $\frac{\sqrt{\pi}}{s}$  which follows from Laplace transform of  $t$  to the power  $a$  by taking  $a$  equal to minus half there. So, this we have independently shown here also the Laplace transform of  $t$  to the power minus half is  $\frac{1}{\sqrt{s}}$  integral  $0$  to infinity  $e^{-u}$  to the power minus  $u$  to the power minus half  $du$  whenever  $s$  is greater than  $0$ .

Here, we have taken  $s$   $t$  equal to  $u$ , so the limits of integration for  $u$  are remain become  $0$  to infinity if we assume  $s$  to be greater than  $0$ . And which is equal to using the definition of gamma function,  $\frac{\Gamma(1/2)}{\sqrt{s}}$  that is  $\frac{\sqrt{\pi}}{\sqrt{s}}$  whenever  $s$  is greater than  $0$ .

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$$= \frac{1}{\sqrt{s}} \sqrt{1/2} = \sqrt{\frac{x}{s}}, \quad s > 0.$$

Example.  $\sin t, \cos t, \cosh t, t^n$   
( $n = 0, 1, 2, \dots$ ) all satisfy the conditions  
of the existence theorem because  
 $\sin t$  and  $\cos t$  are bounded,  
 $\cosh t \leq e^t, t^n \leq n! e^t, (n = 0, 1, 2, \dots)$   
 $\forall t \geq 0$  hence their Laplace transforms exist.

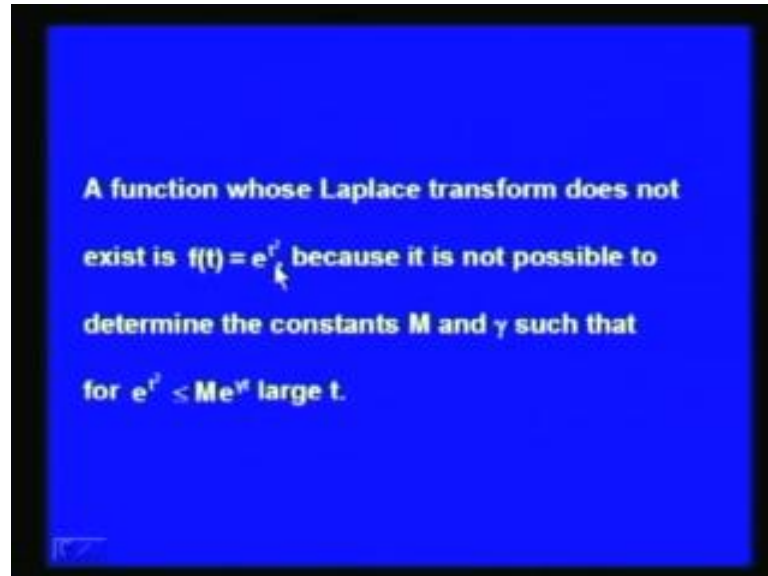
Now, let us look at some of the functions which satisfy the conditions of the existence theorem say let us see take the function  $f(t)$  equal to  $\sin t$   $f(t)$  equal to  $\cos t$ , we know that  $f(t)$  equal to  $\sin t$  and  $\cos t$  are both bounded functions modulus of  $\sin t$  and modulus of  $\cos t$  is less than or equal to 1 for all values of  $t$ . So, they are both bounded function for all  $t$  greater than or equal to 0 in particular. And,  $\cos$  hyperbolic  $t$  is  $e$  to the power  $t$  plus  $e$  to the power minus  $t$  by 2, so  $\cos$  hyperbolic  $t$  is less than or equal to  $e$  to the power  $t$  for all  $t$  greater than or equal to 0 and  $t$  to the power  $n$ . Then  $n$  takes the non negative integral values 0 1 2 and so 1 is less than or equal to  $n$  factorial into  $e$  to the power  $t$  for all  $t$  greater than or equal to 0.

So,  $\sin t$  and  $\cos t$  being both bounded functions that is mod of  $\sin t$  and mod of  $\cos t$  are both less than or equal to 1 means they are both of exponential order. I mean you can write in mod of  $\sin t$  less than or equal to 1 which is further less than or equal to  $e$  to the power  $t$ , and mod of  $\cos t$  less than or equal to 1 can be written less than or equal to  $e$  to the power  $t$ . So, they are both of exponential order.

And, similarly  $\cos$  hyperbolic  $t$  is of exponential order, because it is less than or equal to  $e$  to the power  $t$ ,  $t$  to the power  $n$  for a given  $n$  is less than or equal to  $n$  factorial  $e$  to the power  $t$ , so it is also of exponential order. And they are all continuous functions for all  $t$  greater than or equal to 0, so in particular they are all piecewise continuous functions.

And therefore, they satisfy all the conditions of the existence theorem, and hence they are Laplace transforms.

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Let us take now an example of a function whose Laplace transform does not exist, when such a function is  $f(t) = e^{t^2}$ .  $f(t) = e^{t^2}$  is a continuous function for all  $t \geq 0$  being in the exponential function of  $t^2$ . So, it is piecewise continuous, but it is not of exponential order because for any value of  $a$  greater than 0,  $e^{-at} e^{t^2}$  goes to infinity as  $t$  goes to infinity. It is always finite, it is always infinite, it does not exist. So, therefore, we can say that it is not possible to determine the constants  $M$  and  $\gamma$  such that  $e^{t^2} \leq M e^{\gamma t}$  for large  $t$  and therefore, its Laplace transform does not exist.



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Example . Let

$$f(t) = 2te^{t^2} \cos(e^{t^2}),$$

then  $f(t)$  is continuous in  $[0, \infty)$  but not of exponential order. However, its Laplace transform exists, since integration by parts yields

$$L(f(t)) = \int_0^{\infty} e^{-st} 2te^{t^2} \cos(e^{t^2}) dt$$
$$= e^{-st} \sin(e^{t^2}) \Big|_0^{\infty} + s \int_0^{\infty} e^{-st} \sin(e^{t^2}) dt$$

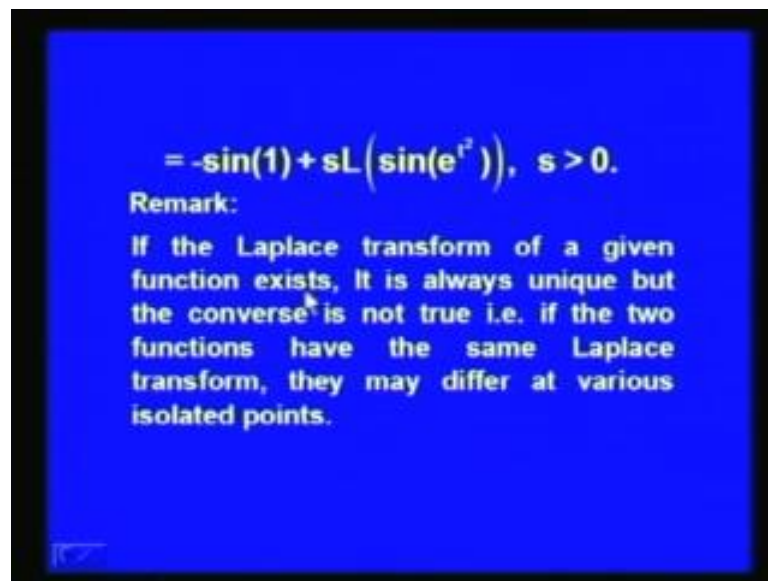
Now, let us take an example of a function which is not of exponential order, but still its Laplace transform exists  $f(t) = 2te^{t^2} \cos(e^{t^2})$ . You can see that  $t e^{t^2} \cos(e^{t^2})$  is not of exponential order. Now, we can see that  $t e^{t^2} \cos(e^{t^2})$  are all continuous functions, so their product is a continuous function in  $[0, \infty)$ . And therefore, it is a piecewise continuous function in  $[0, \infty)$ , but it is not of exponential order as we can see. Because  $e^{-st} f(t)$  for any  $s > 0$  as  $t$  goes to infinity will not go to a finite limit. So; however, its Laplace transform still exists.

Since, integration by parts yields Laplace transform of  $f(t)$  is  $\int_0^{\infty} e^{-st} 2te^{t^2} \cos(e^{t^2}) dt$ . When we integrate by parts here  $2te^{t^2} \cos(e^{t^2})$  is nothing but the differential for  $\sin(e^{t^2})$ .

So, let us take  $e^{-st}$  as first function and  $2te^{t^2} \cos(e^{t^2})$  as second function. We have by integration by parts using integration like parts we have first function that is  $e^{-st}$  and second function that is  $\sin(e^{t^2})$  evaluated at  $0$  and  $\infty$  and then derivative  $e^{-st}$  is  $-s e^{-st}$ . So, we get  $+s \int_0^{\infty} e^{-st} \sin(e^{t^2}) dt$ .

Now, when  $t$  goes to infinity modulus of  $e$  to the power minus  $s t$  into  $\sin e$  to the power  $t$  square is less than or equal to  $e$  to the power minus  $s e$  to the power minus  $s t$ . Because, mod of  $\sin e$  to the power  $t$  square is less than or equal to 1, so if we assume that  $s$  is greater than 0, then  $e$  to the power minus  $s t$ ,  $s t$  goes to infinity will go to 0. And for therefore, the limit of  $e$  to the power minus  $s t$  into  $\sin e$  to the power  $t$  square is 0,  $s t$  approaches infinity. And at the lower limit  $t$  equal to 0 its value is 1, this is 1 and  $\sin e$  to the power 0 is  $e \sin e$  to the power 0 is 1. So, with the lower limit its value is  $\sin 1$ .

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$= -\sin(1) + sL(\sin(e^{t^2})), s > 0.$

**Remark:**  
 If the Laplace transform of a given function exists, it is always unique but the converse<sup>is</sup> is not true i.e. if the two functions have the same Laplace transform, they may differ at various isolated points.

So, we will get the right hand side as minus  $\sin 1$  plus  $s$  times Laplace transform of  $\sin e$  to the power  $t$  square, whenever  $s$  is greater than 0. And,  $\sin$  of  $e$  to the power  $t$  square being a bounded function mod of  $\sin e$  to the power  $t$  square is less than or equal to 1. So, it is of exponential order and it is a continuous function for all  $t$  greater than or equal to 0, so it is piecewise continuous. For all  $t$  greater than or equal to 0, and therefore it satisfies all the conditions of the existence theorem, so its Laplace transform exists.

And therefore, the Laplace transform of the given function  $f(t)$  exists for all  $s$  greater than 0. Now, we may remark here that if the Laplace transform of a given function exist, it is always unique, but the converse is not true, if the two functions have the same Laplace transform they may differ at various isolated points.

This is our lecture on Laplace transformation today we have seen some of the properties of the Laplace transform like the first shifting theorem which refers to the shifting on the

x axis, and then we discussed the change of a scale property. Then we have discussed the existence theorem which gives us the sufficient conditions under which the Laplace transform of a given function exist. The conditions are that the function must be piecewise continuous for all  $t$  greater than or equal to 0, and then it should be of exponential order.

In our next lecture, we shall be discussing some other properties of the Laplace transform that is when you differentiate the function  $f$ , then if you take the Laplace transform what happens will see that the transform of the function gets multiplied by  $s$ . And then we shall see when we take the Laplace transform of the integral of a function it is then what we get is the Laplace transform of the function gets divided by  $s$ .

And then we shall also find see that when we take the derivative of the Laplace transform of the function. Then the function gets multiplied by minus  $t$  and similarly when we take the integral of the Laplace transform will get that the function  $f(t)$  gets divided by  $t$ . So, we will be seeing all those properties of the Laplace transform and then we shall see how to apply those properties of the Laplace transforms to the various problems. The problems that we shall take up will be how to find the solution of an ordinary differential equation with constant coefficients and some more problems.

Thank you.