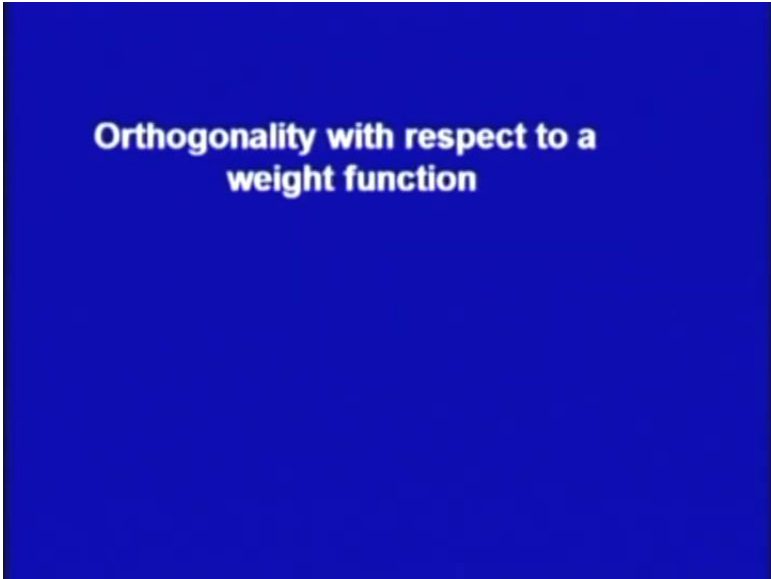


**Mathematics - III**  
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**Lecture - 7**  
**Bessel Functions and their Properties (Contd.)**

Dear viewers, in my last lecture we had discussed the Bessel's equation and its solutions, which are known as Bessel's function of first kind and Bessel's function of second kind. We had also discussed the transformation of Bessel's equation, where we obtained a fairly general linear differential equation of second order, whose solutions could be written in terms of Bessel functions. Now, today we shall discuss some properties of the Bessel functions.

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**Orthogonality with respect to a  
weight function**

We begin with an important property of Bessel's functions known as orthogonality of Bessel functions. The orthogonality of Bessel functions is significant, it is used in the problems on the vibration of circular membranes.

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**Orthogonality with respect to a weight function**

A set of real functions  $g_1, g_2, \dots$  is said to be orthogonal with respect to the *weight function*  $p(x)$  on  $[a, b]$  if

$$\int_a^b p(x)g_m(x)g_n(x)dx = 0 \quad \text{when } m \neq n.$$

If we set  $h_m(x) = (\sqrt{p})g_m$

Let us first define the orthogonality with respect to a weight function, a set of real functions  $g_1, g_2$  and so on, is said to be orthogonal with respect to the weight function  $p(x)$  on the interval  $a, b$ . If  $\int_a^b p(x)g_m(x)g_n(x)dx$  is equal to 0, whenever  $m$  is not equal to  $n$ , if we set  $h_m(x)$  equal to  $\sqrt{p(x)}g_m(x)$ .

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then the above equation becomes

$$\int_a^b h_m(x)h_n(x)dx = 0 \quad \text{when } m \neq n,$$

i.e. the functions  $h_m$  form an orthogonal set in the usual sense.

Example. Chebyshev polynomials of the first kind defined by

$$T_n(x) = \cos(ncos^{-1}x), \quad n = 0, 1, 2, \dots$$

are orthogonal on the interval  $[-1, 1]$  w.r.t the weight function

$$p(x) = (1-x^2)^{-1/2}.$$

Then, we see that the above equation reduces to  $\int_a^b h_m(x)h_n(x)dx$  equal to 0, whenever  $m$  is not equal to  $n$  which implies that the functions  $h_m(x)$  are orthogonal in the usual sense. That is the functions  $h_m(x)$  satisfy the orthogonal property  $\int_a^b h_m(x)h_n(x)dx = 0$ .

to  $b^h m^x$  into  $h n^x d^x$  is equal to 0, so they form an orthogonal set in the usual sense. For example, let us consider Chebyshev polynomials of the first kind, which are defined by  $T_n(x) = \cos n \cos^{-1} x$ , where  $n$  is equal to 0, 1, 2, 3 and so on.

Now, from here we can see that  $T_0(x)$  is equal to 1,  $T_1(x)$  is equal to  $\cos$  of  $\cos^{-1} x$ , so it is  $x$ .  $T_2(x)$  is  $\cos 2 \cos^{-1} x$ , if you put  $\cos^{-1} x$  as  $\theta$ , then you can see that  $\cos 2\theta$  is  $2 \cos^2 \theta - 1$ . So,  $T_2(x)$  will be equal to  $2x^2 - 1$  and similarly, we can see that  $T_3(x)$  which is  $\cos 3 \cos^{-1} x$  turns out to be  $4x^3 - 3x$ . Because,  $\cos^{-1} x$  when we put as  $\theta$  we have got  $\cos 3\theta$  and  $\cos 3\theta$  is  $4 \cos^3 \theta - 3 \cos \theta$ .

So,  $T_3(x)$  turns out to be  $4x^3 - 3x$ , so in general  $T_n(x)$  is polynomial  $n^x$  of degree  $n$ . And we note that these functions  $T_n(x)$  are orthogonal on the interval  $-1$  into  $1$  with respect to the weight function  $p(x) = 1 - x^2$  raised to the power  $-1/2$ , this one can easily show by taking  $\cos^{-1} x$  equal to  $\theta$ . And then, we can show that the integral over  $-1$  to  $1$   $T_m(x) T_n(x) dx$  is equal to 0 whenever  $m$  is not equal to  $n$ .

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**Theorem: For each fixed non negative integer  $n$ , the sequence of Bessel functions of the first kind  $J_n(k_{1n}x), J_n(k_{2n}x), J_n(k_{3n}x), \dots$  where  $k_{mn}$ ,  $m = 1, 2, 3, \dots$  are defined as  $J_n(k_{mn}R) = 0$ , forms an orthogonal set on the interval  $0 \leq x \leq R$  with respect to the weight function  $p(x) = x$  i.e.**

$$\int_0^R x J_n(k_{in}x) J_n(k_{jn}x) dx = 0, \quad i \neq j.$$

Let us study this theorem on orthogonality Bessel functions, for each fixed a non negative integer  $n$ , the sequence of Bessel functions of the first kind  $J_n(k_{1n}x)$ ,  $J_n(k_{2n}x)$ ,  $J_n(k_{3n}x)$  and so on. Where  $k_{mn}$ ,  $m$  is equal to 1, 2, 3 are defined as  $J_n(k_{mn}R) = 0$  that is  $k_{mn}R$  are the roots of the Bessel function  $J_n(x)$  is equal to 0. Form

an orthogonal set on the interval  $0 < x \leq R$  with respect to the weight function  $p(x) = x$ , that is we have the result  $\int_0^R J_n(k_i x) J_n(k_j x) dx = 0$ , whenever  $i \neq j$ .

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**Proof.** We know that  $u = J_n(k_i x)$  and  $v = J_n(k_j x)$ ,  $i \neq j$  are solutions of the differential equations

(1)  $x^2 u'' + xu' + (k_i^2 x^2 - n^2)u = 0$   
and  
(2)  $x^2 v'' + xv' + (k_j^2 x^2 - n^2)v = 0$   
respectively.

Multiplying (1) by  $v/x$  and (2)  $u/x$  and then subtracting, we get

$$x(u''v - v''u) + (u'v - v'u) + (k_i^2 - k_j^2)xuv = 0$$

Let us look at the proof of this theorem, we know that  $u = J_n(k_i x)$  and  $v = J_n(k_j x)$ ,  $i \neq j$  are solutions of the differential equations.  $x^2 u'' + xu' + (k_i^2 x^2 - n^2)u = 0$  and  $x^2 v'' + xv' + (k_j^2 x^2 - n^2)v = 0$ .

We had obtained such a differential equation in the article on transformation of Bessel's equation, when we had made the first substitution, we change the independent variable from  $T = \lambda x$  there. So, from that it comes that  $u = J_n(k_i x)$  is a solution of this, and  $v = J_n(k_j x)$  is the solution of this equation number 2. Now, let us multiply the equation 1 by  $v/x$  and equation 2 by  $u/x$  and then, subtract we shall get  $x(u''v - v''u) + (u'v - v'u) + (k_i^2 - k_j^2)xuv = 0$ .

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$$\begin{aligned}
 &\text{or, } \frac{d}{dx} \{x(u'v - v'u)\} + (k_{in}^2 - k_{jn}^2)xuv = 0. \\
 &\text{Now, integrating with respect to } x \text{ over } (0,R) \text{ we obtain} \\
 &\quad (k_{jn}^2 - k_{in}^2) \int_0^R xuv dx = [x(u'v - v'u)]_0^R \\
 &= \left[ x \left\{ k_{in} J_n'(k_{in}x) J_n(k_{jn}x) - k_{jn} J_n(k_{in}x) J_n'(k_{jn}x) \right\} \right]_0^R \\
 &= 0 \\
 &\Rightarrow \int_0^R x J_n(k_{in}x) J_n(k_{jn}x) dx = 0, \quad i \neq j.
 \end{aligned}$$

Now, this first two term in the last equation can be combined and we were write them as the differential of  $x$  into  $u$  dash  $v$  minus  $v$  dash  $u$  and then, the third term is as it is  $k_{in}$  square minus  $k_{jn}$  square  $x$  into  $u$  into  $v$  equal to 0. Let us now integrate with respect to  $x$  over the interval  $(0, R)$ , then we shall have  $k_{jn}$  square minus  $k_{in}$  square into 0 to  $R$   $x$  into  $u$  into  $v$   $dx$  equal to  $x$  into  $u$  dash  $v$  minus  $v$  dash  $u$  evaluated 0  $R$ .

Let us substitute the values of  $u$  and  $v$  here, the right hand side becomes  $x$  into the derivative of  $u$  is  $J_n(k_{in}x)$ , when we differentiate that with respect to  $x$  we get  $k_{in}$  into  $J_n$  dash  $(k_{in}x)$ ,  $v$  was  $J_n(k_{jn}x)$ , then minus  $v$  dash is  $k_{jn}$  into  $J_n$  dash  $(k_{jn}x)$  and  $u$  is  $J_n(k_{in}x)$ . So, when you put the limits  $R$  and 0 here, then because of the effect that  $J_n(k_{in}x) J_n(k_{in}R)$  is equal to 0 and  $J_n(k_{jn}R)$  equal to 0.

When we put  $x$  equal to  $R$  here, this becomes 0 and also this becomes 0 and when you put the lower limit  $x$  is equal to 0, again because of  $x$  here this whole thing becomes 0, so we have the right hand side equal to 0. And this therefore, implies that if  $i$  is not equal to  $j$ , we can divide by  $k_{jn}$  square minus  $k_{in}$  square in this equation and that will give us integral to 0 to  $R$   $x$  into  $u$  into  $v$   $dx$  equal to 0 replacing the values of  $u$  and  $v$ , we get integral 0 to  $R$   $x$  into  $J_n(k_{in}x)$  into  $J_n(k_{jn}x)$   $dx$  equal to 0.

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case  $i = j$ : Let  $k_{in} \in R$  be a zero of  $J_n(x) = 0$  and  $k_{jn} \rightarrow k_{in}$ , then

$$\int_0^R x J_n^2(k_{in} x) dx$$

$$= \lim_{k_{jn} \rightarrow k_{in}} \int_0^R x J_n(k_{in} x) J_n(k_{jn} x) dx$$

$$= R \lim_{k_{jn} \rightarrow k_{in}} \left[ \frac{k_{in} J_n'(k_{in} R) J_n(k_{jn} R)}{k_{jn}^2 - k_{in}^2} \right]$$

$$= R k_{in} J_n'(k_{in} R) \lim_{k_{jn} \rightarrow k_{in}} \frac{R J_n(k_{jn} R)}{2k_{jn}}$$

$$= \frac{1}{2} R^2 \{J_n'(k_{in} R)\}^2 = \frac{R^2}{2} J_{n+1}^2(k_{in} R),$$

as  $J_n'(k_{in} R) = -J_{n+1}(k_{in} R)$ .

Next let us study the case  $i$  is equal to  $j$ , so let us assume that  $k_{in} \in R$  is a zero of  $J_n(x) = 0$  that is we assume that  $J_n(k_{in} R) = 0$ . And let us assume that  $k_{jn} \rightarrow k_{in}$ , then integral  $\int_0^R x J_n^2(k_{in} x) dx$  can be expressed as the limit of integral  $\int_0^R x J_n(k_{in} x) J_n(k_{jn} x) dx$  as  $k_{jn} \rightarrow k_{in}$ . Now, from our last slide we see that integral  $\int_0^R x J_n(k_{in} x) J_n(k_{jn} x) dx$  can be written as  $k_{in} \int_0^R J_n'(k_{in} x) J_n(k_{jn} x) dx$  over  $k_{jn}^2 - k_{in}^2$ , because  $J_n(k_{in} R) = 0$ .

So, this limit of this multiplied by  $R$  we have and then, this becomes  $R$  times  $k_{in}$  into  $J_n'(k_{in} R)$  and  $\lim_{k_{jn} \rightarrow k_{in}} \frac{R J_n(k_{jn} R)}{2k_{jn}}$ , this we get by using the L'Hospital's rule, because this is of the form  $0/0$ , then  $k_{jn} \rightarrow k_{in}$ . So, differentiating with respect to  $k_{jn}$  we get this expression as  $\frac{R J_n'(k_{jn} R)}{2}$  and which is equal to  $\frac{1}{2} R J_n'(k_{in} R)$  whole square. And that is further equal to  $\frac{R^2}{2} J_{n+1}^2(k_{in} R)$ , as we know that  $J_n'(k_{in} R) = -J_{n+1}(k_{in} R)$  in view of the fact that  $J_n(k_{in} R) = 0$ .

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**Fourier-Bessel Expansion of  $f(x)$ :**  
If  $f$  is a continuous function having finite no. of oscillations in  $(0,a)$ , then we can write

$$f(x) = c_1 J_n(\lambda_1 x) + c_2 J_n(\lambda_2 x) + \dots$$

(3) 
$$= \sum_{i=1}^{\infty} c_i J_n(\lambda_i x),$$

Next we study the Fourier Bessel expansion of  $f(x)$ , if we have a continuous function which has finite number of oscillations in an interval  $0, a$ , then it can be expressed as in terms of Bessel's functions of a given order  $n$ . So, let us say  $f(x)$ , we want to write  $f(x)$  as an infinite series in terms of Bessel functions of order  $n$ , then we assume that  $f(x)$  is equal to  $c_1 J_n(\lambda_1 x) + c_2 J_n(\lambda_2 x) + \dots$  and so on. Or in short we can write the function  $f(x)$  as  $\sum_{i=1}^{\infty} c_i J_n(\lambda_i x)$ .

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where  $\lambda_1, \lambda_2, \dots$  are the roots of  $J_n(\lambda a) = 0$ .

To determine  $c_i$ , multiplying both sides of (3) by  $x J_n(\lambda_i x)$  and integrating from 0 to  $a$ ,

Where  $\lambda_1, \lambda_2, \lambda_3$  and so on, are the roots of the equation  $J_n(\lambda a) = 0$  to determine the unknown coefficient  $c_i$ , let us multiply both sides of the equation (3) by  $J_n(\lambda_i x)$  and integrate from 0 to  $a$ .

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$$\begin{aligned}
 \int_0^a x f(x) J_n(\lambda_i x) dx &= \int_0^a \left( \sum_{j=1}^{\infty} c_j J_n(\lambda_j x) \right) x J_n(\lambda_i x) dx \\
 &= \sum_{j=1}^{\infty} c_j \int_0^a x J_n(\lambda_j x) J_n(\lambda_i x) dx \\
 &= c_i \int_0^a x J_n^2(\lambda_i x) dx,
 \end{aligned}$$

by the orthogonality property of Bessel functions.

Then, we shall have  $\int_0^a x f(x) J_n(\lambda_i x) dx = \int_0^a \sum_{j=1}^{\infty} c_j J_n(\lambda_j x) x J_n(\lambda_i x) dx$ , which can be also written as  $\int_0^a x f(x) J_n(\lambda_i x) dx = \sum_{j=1}^{\infty} c_j \int_0^a x J_n(\lambda_j x) J_n(\lambda_i x) dx$ . Now, making use of the orthogonality of Bessel functions, we ignore that  $\int_0^a x J_n(\lambda_j x) J_n(\lambda_i x) dx = 0$ , whenever  $j$  is not equal to  $i$ . And therefore, this equation further reduces to  $c_i \int_0^a x J_n^2(\lambda_i x) dx$  by the orthogonal property of Bessel functions.



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in view of

$$\int_0^R x J_n^2(k_{in} x) dx = \frac{R^2}{2} J_{n+1}^2(k_{in} R),$$

$$\int_0^a x J_n^2(\lambda_i x) dx = \frac{a^2}{2} J_{n+1}^2(\lambda_i a)$$

$$\Rightarrow \int_0^a x f(x) J_n(\lambda_i x) dx = c_i \frac{a^2}{2} J_{n+1}^2(\lambda_i a),$$

$$\Rightarrow c_i = \frac{2}{a^2 J_{n+1}^2(\lambda_i a)} \int_0^a x f(x) J_n(\lambda_i x) dx.$$

Now, we have seen that when  $i$  is equal to  $J$  integral 0 to  $R$   $x J_n$  square  $k_{in} x dx$  is equal to  $R$  square by 2  $J_n$  plus 1 square  $k_{in} R$ , therefore integral 0 to  $a$   $x J_n$  square  $\lambda_i x dx$  will be equal to  $a$  square by 2  $J_n$  plus 1 square  $\lambda_i a$ . And hence, integral 0 to  $a$   $x f(x) J_n \lambda_i x dx$  is equal to  $c_i$  times  $a$  square by 2  $J_n$  plus 1 square  $\lambda_i a$ . From where we get the value of  $c_i$  as  $2$  over  $a$  square into  $J_n$  plus 1 square  $\lambda_i a$  multiplied by integral 0 to  $a$   $x f(x) J_n \lambda_i x dx$ , and thus we get the values of the unknown coefficients  $c_i$  which occur in the Fourier Bessel expansion of the function  $f(x)$ .

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Putting  $i=1,2,3,\dots$  we can find  $c_1, c_2, c_3, \dots$  and hence the function  $f(x)$ .

**Example.** If  $\alpha_1, \alpha_2, \dots, \alpha_n$  are the positive roots of  $J_0(x) = 0$ , prove that

$$\frac{1}{2} = \sum_{i=1}^{\infty} \frac{J_0(\alpha_i x)}{\alpha_i J_1(\alpha_i)}$$

**Solution.** If  $f(x) = \sum_{i=1}^{\infty} c_i J_0(\lambda_i x)$

So, putting  $i$  equal to 1, 2, 3, we can find the values of  $c_1, c_2, c_3$  and so on, and hence the function  $f(x)$  can be expanded in the desired finite series of Bessel functions. Now, let us study an example based on this article if  $\alpha_1, \alpha_2$  and so on,  $\alpha_n$  are the positive roots of the equation  $J_n(\lambda x) = 0$ ,  $J_n$  is the Bessel function of order  $n$ , so  $J_n(\lambda x)$  is equal to 0.

Let us show that  $\frac{1}{2} = \sum_{i=1}^{\infty} \frac{J_1(\alpha_i)}{\alpha_i}$  or we may say that  $1 = \sum_{i=1}^{\infty} 2 \frac{J_1(\alpha_i)}{\alpha_i}$ . So, we wish to expand the function  $f(x) = 1$ , in terms of infinite series of Bessel functions of order 0, now we know that if  $f(x)$  is equal to  $\sum_{i=1}^{\infty} c_i J_0(\alpha_i x)$ .

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then

$$c_i = \frac{2}{a^2 J_{n+1}^2(\lambda_i a)} \int_0^a x f(x) J_n(\lambda_i x) dx.$$

Taking  $f(x)=1$ ,  $a=1$  and  $n=0$ , we get

$$c_i = \frac{2}{J_1^2(\alpha_i)} \int_0^1 x J_0(\alpha_i x) dx$$

$$= \frac{2}{J_1^2(\alpha_i)} \left[ \frac{x J_1(\alpha_i x)}{\alpha_i} \right]_0^1 = \frac{2}{\alpha_i J_1(\alpha_i)}$$

Then, we have seen that  $c_i$  are given by  $\frac{2}{a^2 J_{n+1}^2(\lambda_i a)} \int_0^a x f(x) J_n(\lambda_i x) dx$ , where  $\lambda_1, \lambda_2$  and  $\lambda_3$  and so on, are the roots of the equation  $J_n(\lambda a) = 0$ . So, in this let us take  $f(x) = 1$ ,  $a = 1$  and  $n = 0$ , because we wish to expand the function  $f(x) = 1$ , in terms of Bessel functions of order 0.

So, we take  $f(x) = 1$  and moreover we are given that  $\alpha_1, \alpha_2, \alpha_n$  are the roots of the equation,  $J_n(\lambda x) = 0$ , so we take  $a = 1$ , then we will get  $c_i = \frac{2}{J_1^2(\alpha_i)} \int_0^1 x J_0(\alpha_i x) dx$ . Now, we know that from the differential formulae for Bessel function, we know that integral of  $x J_n$

naught x is  $x J_1 x$ . So, integral of  $x J_1 \alpha_i x$  over the interval 0 to 1 will be equal to  $x J_1 \alpha_i x$  over  $\alpha_i$  evaluated at over 0, 1. And therefore,  $c_i$  is will be equal to  $2$  over  $J_1^2 \alpha_i$  into  $x J_1 \alpha_i x$  over  $\alpha_i$ , when you put the limits 0 and 1 here, we get this  $c_i$  is to be equal to  $2$  over  $\alpha_i J_1 \alpha_i$ .

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Therefore,

$$f(x) = \sum_{i=1}^{\infty} c_i J_n(\lambda_i x)$$

$$\Rightarrow 1 = \sum_{i=1}^{\infty} \frac{2}{\alpha_i J_1(\alpha_i)} J_0(\alpha_i x)$$

or,

$$\frac{1}{2} = \sum_{i=1}^{\infty} \frac{J_0(\alpha_i x)}{\alpha_i J_1(\alpha_i)}$$

Therefore,  $f(x)$  is equal to  $\sum_{i=1}^{\infty} c_i J_n(\lambda_i x)$  implies that 1 is equal to  $\sum_{i=1}^{\infty} \frac{2}{\alpha_i J_1(\alpha_i)} J_0(\alpha_i x)$ . And then, dividing by 2 we see that, half is equal to  $\sum_{i=1}^{\infty} \frac{J_0(\alpha_i x)}{\alpha_i J_1(\alpha_i)}$ , which gives us the record proof.

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**GENERATING FUNCTION:**

It follows that

$$\exp\left\{x(t - t^{-1})/2\right\} = \sum_{n=-\infty}^{\infty} J_n(x)t^n.$$

The function on the left hand side of the above equation is called the generating function of the Bessel's function  $J_n(x)$ .

Let us now study the generating function for Bessel functions, we are going to show that exponential of  $x t$  minus  $t$  to the power of minus 1 over 2 can be written as sigma  $n$  equal to minus infinity to plus infinity  $J_n x$  into  $t$  to the power  $n$ . So, the Bessel functions of various orders can be derived from the coefficients of the powers of  $t$ , and that is why we call this exponential of  $x t$  minus  $t$  to the power minus 1 over 2 as the generating function for the Bessel's functions. The function on the left hand side of the above equation is called the Bessel generating function of the Bessel's function  $J_n x$ .

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**We have**

$$e^{\frac{1}{2}x\left(t - \frac{1}{t}\right)} = e^{(xt/2)} e^{-(x/2t)}$$
$$(4) = \left[ 1 + \frac{xt}{2} + \frac{1}{2!} \left(\frac{xt}{2}\right)^2 + \dots \right] \left[ 1 - \frac{x}{2t} + \frac{1}{2!} \left(\frac{x}{2t}\right)^2 - \dots \right].$$

The terms that are independent of  $t$  are given by

We can write  $e^{x/2}$  into  $t$  minus  $t$  to the power minus 1 equal to  $e^{-t}$  to the power  $x/2$  into  $e^{-t}$  to the power  $x/2$ . Let us then put the Maclaurin series expansion of  $e^{x/2}$  and  $e^{-t}$  to have the right hand side as  $1 + x/2 + 1/2 \text{ factorial } (x/2)^2$  and so on. And then,  $1 - x/2 + 1/2 \text{ factorial } (x/2)^2$  and so on, then let us collect the coefficients of  $t$ .

First we look at the coefficients of  $t$  to the power 0, that is the terms that are independent of  $t$ , such terms will come when we will multiply  $1$  by  $1$  had been multiplied  $t$  by  $1$  by  $t$  or  $t$  square by  $1$  by  $t$  square. That is the  $t$  to the power and coefficient of  $t$  power and is multiplied by the coefficient of  $t$  to the power minus and from here, we will just gives us the terms that are independent of  $t$ .

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$$1 - \left(\frac{x}{2}\right)^2 + \frac{1}{(2!)^2} \left(\frac{x}{2}\right)^4 - \frac{1}{(3!)^2} \left(\frac{x}{2}\right)^6 + \dots$$

$$= \sum_{k=0}^{\infty} \frac{(-1)^k}{(k!)^2} \left(\frac{x}{2}\right)^{2k}$$

$$= 1 - \frac{x^2}{1^2 \cdot 2^2} + \frac{x^4}{2^2 \cdot 4^2} - \frac{x^6}{2^2 \cdot 4^2 \cdot 6^2} + \dots$$

$$= J_0(x).$$

Collecting the coefficient of  $t^n$  in the product on the r.h.s of (4), we obtain

So, we collect those terms that are independent of  $t$  and we see that, that are come out to be  $1 - x/2$  whole square plus  $1/2$  factorial whole square into  $x/2$  raise to the power 4 minus  $1/3$  factorial whole square  $x/2$  raise to the power 6 and so on, which can be written in the form of sum sigma  $k$  equal to 0 to infinity minus 1 to the power  $k$  over  $k$  factorial square multiplied  $x/2$  raise to the power  $2k$ .

And then, we can put at in this convenient form  $1 - x$  square by  $1$  square into  $2$  square plus  $x$  to the power 4 over  $2$  square  $4$  square minus  $x$  to the power 6 over  $2$  square  $4$  square  $6$  square and so on, which we know is nothing but  $J_0(x)$ . So, the coefficient

of  $t$  to the power 0 is nothing but  $J_n(x)$ , now let us collect the coefficient of  $t$  to the power  $n$  in the product on the right hand side of 4.

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$$\frac{x^n}{2^n n!} - \frac{x^{n+1}}{2^{n+1} (n+1)!} \left(\frac{x}{2}\right) + \frac{x^{n+2}}{2^{n+2} (n+2)! 2!} \left(\frac{x}{2}\right)^2 - \dots$$

$$= \sum_{k=0}^{\infty} (-1)^k \frac{x^{n+2k}}{2^{n+2k} k!(n+k)!} = J_n(x),$$

**The coefficient of  $t^n$  is**

$$\frac{(-1)^n x^n}{2^n n!} \left[ 1 - \frac{x^2}{2 \cdot (2n+2)} + \frac{x^4}{2 \cdot 4 (2n+2)(2n+4)} - \dots \right]$$

$$= (-1)^n J_n(x) = J_{-n}(x).$$

So, that means in the first series, we will multiply the term with that power  $n$  by 1, then we will multiply  $t$  to the power  $n+1$  by  $t$ , then  $t$  to the power  $n+2$  by  $1$  by  $t^2$  and so on, and to have the coefficients of  $t$  to the power  $n$ . And they come out to be  $x$  to the power  $n$  over  $2^n$  to the power  $n$  over  $n$  factorial minus  $x$  to the power  $n+1$  over  $2^{n+1}$  to the power  $n+1$  into  $n+1$  factorial into  $x$  by  $2$ .

Or you can say the second term is  $x$  by  $2$  raise to the power  $n+2$  and then,  $x$  by  $2$  raise to the power  $n+2$  is into  $x$  by  $2$  whole square is  $x$  by  $2$  raise to the power  $n+4$  here divided by  $n+2$  factorial into  $2$  factorial and so on, which can be written as  $\sum_{k=0}^{\infty} (-1)^k \frac{x^{n+2k}}{2^{n+2k} k!(n+k)!}$  and which is equal to  $J_n(x)$ .

Next we collect the coefficient of  $t$  to the power  $-n$ , so for that we may multiply the coefficient of  $t$  to the power  $-n$  by  $1$ , then  $t$  to the power  $-n-1$  by  $t$  and then  $t$  to the power  $-n-2$  by  $t^2$  and so on. And we see that we get the coefficients of  $t$  to the power  $-n$  as  $(-1)^n \frac{x^n}{2^n n!}$  multiplied by this infinite series  $1 - \frac{x^2}{2 \cdot (2n+2)} + \frac{x^4}{2 \cdot 4 (2n+2)(2n+4)} - \dots$  and so on, which

is equal to minus 1 to the power n into J n x. And we know that when n is an integer, minus 1 to the power n into J n x is J minus n x.

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**Hence**

$$\begin{aligned}
 e^{x(t-1)/2} &= J_0(x) + J_1(x)t + J_2(x)t^2 + \dots J_n(x)t^n \\
 &+ \dots + J_{-1}(x)t^{-1} + J_{-2}(x)t^{-2} + \dots \\
 &+ J_{-n}(x)t^{-n} + \dots \\
 &= \sum_{n=-\infty}^{\infty} J_n(x)t^n,
 \end{aligned}$$

which proves the required result.

Hence, we can say that e to the power x into t minus 1 over t by 2 is equal to the term that contains that are independent of t is J naught x, the terms which give us the terms in t, t square, t cube, t to the power n and so on, their coefficients are J 1 x, J 2 x, J n x and so on. Because, we saw that the coefficient of t to the power n is J n x and then, the coefficient of t to the power minus n and we have seen is J minus n x, so the coefficients of t to the power minus 1 is J minus 1 x.

The coefficient of t to the power minus 2 is J minus 2 x and so on, the coefficient of t to the power minus n is J minus n x and so on. So, we can write this in the form of the summation sigma n equal to minus infinity to infinity J n x into t to the power n, which gives us the required result.

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**Example.**  $J_{n-1}(x) + J_{n+1}(x) = \frac{2n}{x} J_n(x).$

**Solution.** Differentiating both sides of

$$\exp\left\{x(t - t^{-1})/2\right\} = \sum_{n=-\infty}^{\infty} J_n(x)t^n$$

w.r.t.  $t$ , we obtain

Now, let us study an example based on the generating function, let us establish this formula for Bessel functions  $J_{n-1}x + J_{n+1}x$  is equal to  $2n$  over  $x$   $J_nx$ , we had earlier shown this. Now, we are going to prove this, using the article on generating functions, so we know that  $e$  to the power  $x$  into  $t$  minus  $1$  by  $t$  divided by  $2$  is equal to  $\sum_{n=-\infty}^{\infty} J_n(x)t^n$ .

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or

$$\frac{x}{2}(1 + t^{-2})\exp\left(\frac{x}{2}(t - t^{-1})\right) = \sum_{n=-\infty}^{\infty} nJ_n(x)t^{n-1}$$

$$\frac{x}{2} \sum_{n=-\infty}^{\infty} J_n(x)t^n + \frac{x}{2} \sum_{n=-\infty}^{\infty} J_n(x)t^{n-2} = \sum_{n=-\infty}^{\infty} nJ_n(x)t^{n-1}$$

Now equating the coefficient of  $t^{n-1}$  on both sides, the required result follows.

When we differentiate this equation with respect to  $t$  what we get is this,  $x$  by  $2$   $1 + t$  plus  $t$  to the power minus  $2$  into exponential of  $x$  by  $2$   $t$  minus  $1$  by  $2$ ,  $t$  is equal to  $\sum_{n=-\infty}^{\infty} nJ_n(x)t^{n-1}$ .



equal to minus infinity to infinity  $\int_0^1 t^{n-1} dt$  to the power  $n-1$ . Now, we know that the left hand side can be written as  $x^{n-1} \int_0^1 t^{n-1} e^{-xt} dt$  to the power  $n-1$ . And the generating function  $e^{-xt} = \sum_{n=0}^{\infty} \frac{(-x)^n t^n}{n!}$  to the power  $n-1$  is equal to  $\sum_{n=0}^{\infty} \frac{(-x)^n}{n!} \int_0^1 t^{n-1} dt$ .

So, we get  $x^{n-1} \int_0^1 t^{n-1} dt = \sum_{n=0}^{\infty} \frac{(-x)^n}{n!} \int_0^1 t^{n-1} dt$  and then, plus  $x^{n-1} \int_0^1 t^{n-2} dt$  multiply to  $\sum_{n=0}^{\infty} \frac{(-x)^n}{n!} \int_0^1 t^{n-1} dt$  gives you  $x^{n-1} \int_0^1 t^{n-2} dt = \sum_{n=0}^{\infty} \frac{(-x)^n}{n!} \int_0^1 t^{n-1} dt$ . So, left hand side becomes the sum of these two terms and then, it is equal to  $\sum_{n=0}^{\infty} \frac{(-x)^n}{n!} \int_0^1 t^{n-1} dt$  and  $\int_0^1 t^{n-1} dt$  to the power  $n-1$ .

So, in this equation what we have done is simply we have put the value of  $e^{-xt} = \sum_{n=0}^{\infty} \frac{(-x)^n t^n}{n!}$  to the power  $n-1$  as  $\sum_{n=0}^{\infty} \frac{(-x)^n}{n!} \int_0^1 t^{n-1} dt$  to the power  $n$ . Now, we equate the coefficient of  $t^{n-1}$  on both sides, coefficient  $n-1$  here will be  $x^{n-1} \int_0^1 t^{n-1} dt$ , the coefficient of  $t^{n-1}$  here will be  $x^{n-1} \int_0^1 t^{n-1} dt$  and the coefficient of  $t^{n-1}$  will be  $n \int_0^1 t^{n-1} dt$  and that will give us the required result.

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INTEGRAL REPRESENTATION OF  $J_n(x)$

Let  $t = \cos\theta + i \sin\theta$ . Then,

(5)  $e^{\frac{x}{2} \left( t - \frac{1}{t} \right)} = e^{ix \sin\theta} = \cos(x \sin\theta) + i \sin(x \sin\theta)$ .

Now,

$$t^n + \frac{(-1)^n}{t^n} = (\cos n\theta + i \sin n\theta) + (-1)^n (\cos n\theta - i \sin n\theta)$$

$$= \begin{cases} 2 \cos n\theta, & \text{when } n \text{ is even} \\ 2i \sin n\theta, & \text{when } n \text{ is odd} \end{cases}$$

Now, let us study the integral representation of Bessel functions of first kind that is  $J_n(x)$ , we are going to show that the Bessel function of first kind of order  $n$  which is an in finite

series can be expressed as an integral, which we call as Bessel integral, if n is an integer. So, let us put t equal to cos theta plus i sin theta, let t be equal to cos theta plus i sin theta, then the generating function for the Bessel functions e to the power half x t minus 1 by t will be equal to e to the power i x sin theta.

Because, t is cos theta plus i sin theta, so 1 by t will be cos theta minus i sin theta, so t minus 1 by t will be 2 i sin theta and therefore, e to the power half x into 2 i sin theta will give us e to the power i x sin theta. And by the Euler's formula we know that e to the power i x sin theta is cos x sin theta plus i sin x sin theta. Now, t to the power n plus minus 1 to the power n over t to the power n will be equal to t to the power n by De Moivre's theorem will be cos theta plus i sin theta to the power n.

So, n is an integer, therefore t to the power n will be cos n theta plus i sin n theta, then minus 1 to the power n, 1 over t to the power n will be t to the power minus n, so cos theta plus i sin theta to the power minus n will be cos n theta minus i sin n theta. If n is an even integer, then the right hand side, therefore is 2 cos n theta and when n is an odd integer the right hand side becomes 2 i sin n theta. So, the value of t to the power n plus minus 1 to the power n over t to the power n is 2 cos n theta, when n is an even integer and it is 2 i sin n theta when n is an odd integer.

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**Hence**

$$e^{\frac{1}{2}x\left(\frac{t-1}{t}\right)} = J_0(x) + \sum_{n=1}^{\infty} \left( t^n + \frac{(-1)^n}{t^n} \right) J_n(x)$$

$$(6) = \left[ J_0(x) + 2 \{ J_2(x)\cos 2\theta + J_4(x)\cos 4\theta + \dots \} \right] + 2i \left[ J_1(x)\sin \theta + J_3(x)\sin 3\theta + \dots \right]$$

**comparing the real and imaginary parts of equations (5) and (6), we obtain**

Hence, the generating function e to the power half x t minus 1 over t, which is equal to J naught x plus sigma n equal to 1 to infinity t to the power n J n x plus sigma n equal to 1

to infinity minus 1 to the power n over t to the power n  $J_n x$  will be equal to  $J_0 x$  plus 2 times  $J_2 x \cos 2\theta$   $J_4 x \cos 4\theta$  and so on. Because, when n is even here, t to the power n plus minus 1 to the power n over t to the power n becomes  $2 \cos n\theta$ , so  $2 \cos n\theta$  into  $J_n x$  when n is an even integer.

And when n is an odd integer, the value of t to the power n plus minus 1 to the power n over t to the power n is  $2i \sin n\theta$ , so for odd integral values of n, this sum will give us the terms  $2i J_1 x \sin \theta$  plus  $J_3 x \sin 3\theta$  and so on. Now, let us compare the real and imaginary parts of equations 5 and 6, in the equation 5 we have seen that  $e^{ix} = \cos x + i \sin x$ , so  $\cos x + i \sin x$  is equal to this.

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(7)  $\cos(x \sin \theta) = J_0(x) + 2\{J_2(x)\cos 2\theta + J_4(x)\cos 4\theta + \dots\}$

and

(8)  $\sin(x \sin \theta) = 2\{J_1(x)\sin \theta + J_3(x)\sin 3\theta + \dots\}$ .

These series are known as Jacobi series.  
From the above series, it follows that

So, when we equate real and imaginary parts on the two sides, we get  $\cos x \sin \theta$  equal to  $J_0 x$  plus 2 times  $J_2 x \cos 2\theta$  plus  $J_4 x \cos 4\theta$  and so on. And  $\sin x \sin \theta$  equal to 2 times  $J_1 x \sin \theta$  plus  $J_3 x \sin 3\theta$  and so on, these series are known as the Jacobi series.

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$$\int_0^\pi \cos(x \sin \theta) \cos n \theta d\theta = \begin{cases} \pi J_n(x), & \text{when } n \text{ is even or } 0 \\ 0, & \text{when } n \text{ is odd} \end{cases}$$

$$\int_0^\pi \sin(x \sin \theta) \sin n \theta d\theta = \begin{cases} 0, & \text{when } n \text{ is even} \\ \pi J_n(x), & \text{when } n \text{ is odd} \end{cases}$$

On adding these two relations, we get

$$\int_0^\pi \{ \cos(x \sin \theta) \cos n \theta + \sin(x \sin \theta) \sin n \theta \} d\theta = \pi J_n(x)$$

Now, from these series it follows that, if you multiply the Jacobi series for  $\cos x \sin \theta$  by  $\cos n \theta$  and integrate over 0 to  $\pi$ , then we know that the cosine functions  $\cos n \theta$  are orthogonal over the interval 0 to  $\pi$  that is whenever  $m$  is not equal to  $n$ , that is  $\int_0^\pi \cos m \theta \cos n \theta d\theta = 0$  whenever  $m$  is not equal to  $n$ . And in the case when  $m$  is equal to  $n$ , we know that  $\int_0^\pi \cos^2 n \theta d\theta = \frac{\pi}{2}$ , so making use of that  $\int_0^\pi \cos x \sin \theta \cos n \theta d\theta = \frac{\pi}{2} J_n(x)$  when  $n$  is an even integer or  $n$  is equal to 0.

And it will be 0 when  $n$  is an odd integer, because when  $n$  is an odd integer  $\cos n \theta$  into  $\cos n \theta$  integral of that will be 0, because  $m$  will be never equal to  $n$ , on the right hand side of  $\cos x \sin \theta$  we have only even integral multiples of  $\theta$  in cosine terms. And then, in the Jacobi series for  $\sin x \sin \theta$ , we multiply  $\sin x \sin \theta$  by  $\sin n \theta$  and similarly, we see that the value of  $\int_0^\pi \sin x \sin \theta \sin n \theta d\theta$  is equal to 0 when  $n$  is even.

Because, in the  $\sin x \sin \theta$  expansion of that we have  $\sin$  of odd multiples of  $\theta$  on the right side, so when we multiply by  $\sin n \theta$  where  $n$  is even, then  $\int_0^\pi \sin m \theta \sin n \theta d\theta = 0$ , because  $m$  will be not equal to  $n$ . And then,  $\int_0^\pi \sin^2 n \theta d\theta = \frac{\pi}{2}$ , when  $n$  is odd will give us  $\frac{\pi}{2} J_n(x)$ , so now let us add these two equations, when we add these two equations, we find that  $\int_0^\pi$

$\int_0^\pi (\cos x \sin \theta \cos n \theta + \sin x \sin \theta \sin n \theta) d\theta$  is equal to  $\pi J_n(x)$ .

Because, if  $n$  is 0 then this is 0, obviously this is 0 and when  $n$  is 0,  $\int_0^\pi \cos x \sin \theta \cos n \theta d\theta$  is equal to  $\pi J_n(x)$ . And when  $n$  is an even integer, then  $\int_0^\pi \sin x \sin \theta \sin n \theta d\theta$  is 0, we have seen here an integral  $\int_0^\pi \cos x \sin \theta \cos n \theta d\theta$  is  $\pi J_n(x)$ , so again we get the right side. And when  $n$  is an odd integer here, then  $\int_0^\pi \sin x \sin \theta \sin n \theta d\theta$  is  $\pi J_n(x)$ , while the other integral  $\int_0^\pi \cos x \sin \theta \cos n \theta d\theta$  is 0. So, when we add these two equations for all integral values of  $n$ , 0, 1, 2, 3,  $n$  equal to 0, 1, 2, 3 and so on, we see that  $\int_0^\pi (\cos x \sin \theta \cos n \theta + \sin x \sin \theta \sin n \theta) d\theta$  is  $\pi J_n(x)$ .

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Therefore, for integral  $n$ ,

$$J_n(x) = \frac{1}{\pi} \int_0^\pi \cos(n\theta - x \sin \theta) d\theta.$$

For  $n = 0$ , we get the integral representation of  $J_0(x)$  as

$$J_0(x) = \frac{1}{\pi} \int_0^\pi \cos(x \sin \theta) d\theta.$$

$$= \frac{1}{\pi} \int_{-\pi/2}^{\pi/2} \cos(x \cos \phi) d\phi$$

And which can be put in an alternate form as  $J_n(x) = \frac{1}{\pi} \int_0^\pi \cos(n\theta - x \sin \theta) d\theta$ , because we know that  $\cos(a - b) = \cos a \cos b + \sin a \sin b$ . For  $n$  equal to 0, from here we can see for  $n$  equal to 0 we get that  $J_0(x)$  is equal to  $\frac{1}{\pi} \int_0^\pi \cos x \sin \theta d\theta$ , now using the property of definite integrals we can see that, if you put here  $\theta = \pi/2 - \phi$ ,

If you make this substitution  $\theta = \pi/2 - \phi$ , then this integral representation of  $J_0(x)$  changes into  $\frac{1}{\pi} \int_{-\pi/2}^{\pi/2} \cos(x \cos \phi) d\phi$ , here we have put  $\theta = \pi/2 - \phi$ . Now, since  $\cos x$

cos phi is an even function of phi may be use of the property of definite integrals, we can write the integral over minus pi by 2 to pi by 2 cos x cos phi d phi as 2 times integral 0 to pi by 2 cos x cos phi d phi.

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$$= \frac{2}{\pi} \int_0^{\pi/2} \cos(x \cos \phi) \, d\phi$$

$$= \frac{2}{\pi} \int_0^{\pi/2} \cos(x \sin \phi) \, d\phi$$

$$= \frac{1}{\pi} \int_0^{\pi/2} \cos(x \sin \phi) \, d\phi.$$

We note that

$$\int_0^{\pi/2} \cos(x \cos \phi) \, d\phi = \frac{1}{2} \int_0^{\pi/2} \cos(x \cos \phi) \, d\phi$$

And so we will get 2 over pi 0 to pi by 2 cos x cos phi d phi, which is equal to 2 over pi integral 0 to pi by 2 cos x sin phi d phi, because we have the property of the definite integrals that integral 0 to a, f (x) d x is equal to integral 0 to a f a minus x d x; so when you replace phi by pi by 2 minus phi here, you get cos x sin phi. Now, if you look at integral 0 to pi cos x sin phi d phi, then we see that 0 to pi cos x sin phi d phi is equal to 2 times 0 to pi by 2 cos x sin phi d phi, again by using the property of definite integrals.

Because, we have the property of definite integrals they says that integral 0 to 2 a f (x) d x is equal to 2 times 0 to a f (x) d x provided f 2 a minus x is equal to f x, so making use of that property these two are equal. And therefore, we have J naught x also equal to 1 over pi integral 0 to pi cos x sin phi d phi, now we also note that integral 0 to pi by 2 cos x cos phi d phi is equal to half of integral 0 to pi cos x cos phi d phi.

Because, again when you replace phi by pi minus phi here, you get cos of minus x cos phi which cos of minus theta is cos theta, so you get cos of x cos phi. And therefore, again making use of the property of definite and definite integrals 0 to pi cos x cos phi d phi is nothing but it is equal to 2 times 0 to pi by 2 cos x cos phi d phi.

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hence

$$J_0(x) = \frac{1}{\pi} \int_0^{\pi} \cos(x \cos \Phi) d\Phi$$
$$= \frac{1}{\pi} \int_0^{\pi} \cos(x \sin \Phi) d\Phi.$$

Aliter:  
Changing  $\theta$  to  $\pi/2 - \Phi$  in (7), we get

$$\cos(x \cos \Phi) = J_0(x) - 2J_2(x) \cos 2\Phi + 2J_4(x) \cos 4\Phi - \dots$$

So, the integral representation of  $J_0(x)$  has another form  $J_0(x)$  can be also written as  $\frac{1}{\pi} \int_0^{\pi} \cos(x \cos \phi) d\phi$ , so we have two integral representations, when the limits are 0 and  $\pi$  of the Bessel functions of order 0. One is  $J_0(x) = \frac{1}{\pi} \int_0^{\pi} \cos(x \cos \phi) d\phi$ , another one is  $J_0(x) = \frac{1}{\pi} \int_0^{\pi} \cos(x \sin \phi) d\phi$ .

Now, we can arrive at this integral representation of  $J_0(x)$  in another way from Jacobi series, also we can get this integral representation of  $J_0(x)$  directly. If in the Jacobi series equation number 7 you change the  $\theta$  to  $\pi/2 - \phi$ , if you put  $\theta = \pi/2 - \phi$  there in the Jacobi series given in equation 6. Then, we get that  $\cos x \cos \phi$ ,  $\sin \phi$  becomes  $\cos \phi$  there and then,  $\cos 2\theta$ ,  $\cos 4\theta$  becomes  $\cos 2\phi$  minus  $\cos 2\phi$  and then, plus  $\cos 4\phi$ . So, right hand side of that equation becomes  $J_0(x) - 2J_2(x) \cos 2\phi + 2J_4(x) \cos 4\phi$  and so on.

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Integrating both sides w.r.t.  $\Phi$  from 0 to  $\pi$ , we get

$$\int_0^{\pi} \cos(x \cos \Phi) d\Phi$$
$$= \int_0^{\pi} [J_0(x) - 2J_2(x)\cos 2\Phi + 2J_4(x)\cos 4\Phi - \dots] d\Phi$$
$$= \left[ J_0(x)\Phi - 2J_2(x)\frac{1}{2}\sin 2\Phi + 2J_4(x)\frac{1}{4}\sin 4\Phi - \dots \right]_0^{\pi}$$
$$= J_0(x)\pi.$$

Therefore, we have

And when we integrate this equation with respect to  $\phi$  from 0 to  $\pi$ , what we get is  $\int_0^{\pi} \cos x \cos \phi d\phi$  is equal to  $\int_0^{\pi} J_0(x) - 2J_2(x)\cos 2\phi + 2J_4(x)\cos 4\phi$  and so on,  $d\phi$ . Now, let us integrate it we get integral of  $J_0(x) d\phi$  with respect to  $\phi$  is  $J_0(x) \phi$ , then  $2J_2(x)$  by  $\sin$  integral of  $\cos 2\phi$  is  $\sin 2\phi$  by 2 and then,  $2J_4(x)$  integral of  $\cos 4\phi$  is  $\sin 4\phi$  by 4.

So, now let us put the limits 0 in to  $\pi$ , we will get the value of this expression as  $J_0(x) \pi$ , because  $\sin n\phi$  is always equal to 0. So, we get  $\sin 2\pi$ ,  $\sin 4\pi$  all are 0's and at  $\phi$  equal to 0 again this whole expression is 0, so the value of this expression is  $J_0(x) \pi$ .



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$$J_0(x) = \frac{1}{\pi} \int_0^{\pi} \cos(x \cos \Phi) d\Phi.$$

Example.

$$J_0(x) = 1 - \frac{x^2}{2^2} + \frac{x^4}{2^2 \cdot 4^2} - \dots = \sum_0^{\infty} \frac{(-1)^r x^{2r}}{2^{2r} (r!)^2}.$$

Solution. We know that

$$J_0(x) = \frac{1}{\pi} \int_0^{\pi} \cos(x \cos \Phi) d\Phi.$$

And therefore, we have  $J_0(x)$  equal to  $\frac{1}{\pi} \int_0^{\pi} \cos(x \cos \phi) d\phi$ , so the value of  $J_0(x)$  can be evaluated in an independent manner directly from the Jacobi series equation number 7, rather than obtaining yet as a particular case of  $J_n(x)$  for  $n$  equal to 0. Now, let us take an example which is based on the integral representation of Bessel functions, let us show that reading the integral representation of  $J_0(x)$ , the value of  $J_0(x)$  is equal to  $1 - \frac{x^2}{2^2} + \frac{x^4}{2^2 \cdot 4^2} - \dots$  and so on.

That is  $\sum_{r=0}^{\infty} \frac{(-1)^r x^{2r}}{2^{2r} (r!)^2}$ , this expansion infinite series expansion of  $J_0(x)$  we shall derive using the integral representation of  $J_0(x)$ ; we know that  $J_0(x)$  is equal to  $\frac{1}{\pi} \int_0^{\pi} \cos(x \cos \phi) d\phi$ .

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$$J_0(x) = \frac{1}{\pi} \int_0^\pi \left( 1 - \frac{x^2 \cos^2 \Phi}{2!} + \frac{x^4 \cos^4 \Phi}{4!} - \frac{x^6 \cos^6 \Phi}{6!} + \dots \right) d\Phi.$$

Since

$$\frac{1}{\pi} \int_0^\pi \cos^{2r} \Phi \, d\Phi = \frac{1 \cdot 3 \cdot 5 \dots (2r-1)}{2 \cdot 4 \cdot 6 \dots (2r)}$$

we have

$$\frac{1}{\pi} \int_0^\pi \frac{(-1)^r x^{2r} \cos^{2r} \Phi}{(2r)!} d\Phi = \frac{(-1)^r x^{2r} 1 \cdot 3 \cdot 5 \dots (2r-1)}{(2r)! 2 \cdot 4 \cdot 6 \dots (2r)}$$

Now, let us put the Maclaurin series expansion of  $\cos x$ , then we shall get  $J_0(x)$  equal to  $\frac{1}{\pi} \int_0^\pi \left( 1 - \frac{x^2 \cos^2 \Phi}{2!} + \frac{x^4 \cos^4 \Phi}{4!} - \frac{x^6 \cos^6 \Phi}{6!} + \dots \right) d\Phi$ . Now, since  $\frac{1}{\pi} \int_0^\pi \cos^{2r} \Phi \, d\Phi$  is equal to  $\frac{1 \cdot 3 \cdot 5 \dots (2r-1)}{2 \cdot 4 \cdot 6 \dots (2r)}$ .

This result can be derived by using the reduction formula for  $\cos^{2r} \Phi$  or it can also be obtained by the gamma function, because when you replace  $\Phi$  by  $\pi - \Phi$ , because of the even power of  $\cos \Phi$  the sign will not change that is  $\cos^{2r} \Phi$  will remain  $\cos^{2r} \Phi$ , when we replace  $\Phi$  by  $\pi - \Phi$ . So, the limits will change from 0 to  $\pi$  by 2 and we will get  $\frac{2}{\pi} \int_0^{\pi/2} \cos^{2r} \Phi \, d\Phi$ , from functions we know that  $\int_0^{\pi/2} \cos^{2r} \Phi \, d\Phi = \frac{1}{2} \frac{1 \cdot 3 \cdot 5 \dots (2r-1)}{2 \cdot 4 \cdot 6 \dots (2r)}$ .

$\int_0^{\pi/2} \sin^m \theta \cos^n \theta \, d\theta$  is equal to  $\frac{\Gamma(m+1) \Gamma(n+1)}{2 \Gamma(m+n+2)}$ . So, making use of that, we can see that the value of this integral is  $\frac{1 \cdot 3 \cdot 5 \dots (2r-1)}{2 \cdot 4 \cdot 6 \dots (2r)}$ , so then we shall have  $\frac{1}{\pi} \int_0^\pi \frac{(-1)^r x^{2r} \cos^{2r} \Phi}{(2r)!} d\Phi$ .

Now, there is nothing but the general term of this expansion, we have taken the general term from this expansion and evaluated the integral of that with respect to  $\Phi$  over the interval 0 to  $\pi$  and multiply by  $\frac{1}{\pi}$ . So, we see that making use of this result the

value of this integral comes out to be minus 1 to the power r x to the power 2 r 1, 3, 5 and so on, 2 r minus 1 over 2 r factorial 2 into 4 into 6 and so on, up to 2 r.

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and  $\int_0^\pi \cos^{2r} \theta d\theta = \frac{(-1)^r x^{2r}}{2^2 \cdot 4^2 \cdot 6^2 \dots (2r)^2} \cdot \pi/2,$

Consequently, we get

$$J_0(x) = 1 - \frac{x^2}{2^2} + \frac{x^4}{2^2 \cdot 4^2} - \dots = \sum_0^\infty \frac{(-1)^r x^{2r}}{2^{2r} (r!)^2}.$$

Squaring (7) and (8) and integrating w.r.t.θ from 0 to π and in view of the following results

$$\int_0^\pi \cos m\theta \cos n\theta d\theta = \int_0^\pi \sin m\theta \sin n\theta d\theta = 0, \quad (m \neq n)$$

And which is equal to minus 1 to the power r x to the power 2 r over 2 square into 4 square into 6 square and so on, 2 r whole square. And consequently the value of J naught x is equal to sigma r equal to 0 to infinity minus 1 to the power r x to the power 2 r over 2 to the power 2 r into r factorial square, r in the expanded form; it has 1 minus x square by 2 square plus x to the power 4 over 2 square into 4 square and so on.

Now, let us square the equation number 7 and 8 and integrate with respect to theta over the interval 0 to pi, then in view of the following results. Integral 0 to pi cos m theta into cos n theta d theta and 0 to pi sin m theta sin n theta d theta being 0, whenever m is not equal to n.

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and

$$\int_0^{\pi} \cos^2 m\theta d\theta = \int_0^{\pi} \sin^2 m\theta d\theta = \pi/2,$$

we obtain

$$[J_0(x)]^2 \pi + 4[J_2(x)]^2 \frac{\pi}{2} + \dots = \int_0^{\pi} \cos^2(x \sin \theta) d\theta$$

and

$$4[J_1(x)]^2 \frac{\pi}{2} + 4[J_3(x)]^2 \frac{\pi}{2} + \dots = \int_0^{\pi} \sin^2(x \sin \theta) d\theta$$

And integral 0 to pi cos square m theta d theta and 0 to pi sin square m theta d theta being equal to pi by 2, we obtain J naught x whole square into pi plus 4 J 2 x whole square into pi by 2 plus and so on, equal to integral 0 to pi cos square x sin theta d theta. And 4 times J 1 x whole square into pi by 2 plus J 3 x whole square into pi by 2 and so on, equal to integral 0 to pi sin square x sin theta d theta.

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Adding these two equations

$$\pi[J_0^2 + 2J_1^2 + 2J_2^2 + 2J_3^2 + \dots] = \int_0^{\pi} d\theta = \pi$$

hence  $J_0^2 + 2J_1^2 + 2J_2^2 + 2J_3^2 + \dots = 1.$

Consequently,  $|J_0| \leq 1$

and  $|J_n| \leq 2^{-1/2}, n = 1, 2, 3..$

On adding these two equations we get pi into J naught square plus 2 times J 1 square plus 2 times J 2 square plus 2 times J 3 square and so on, equal to 0 to pi integral 0 to pi d

theta. Because,  $\cos^2 \theta + \sin^2 \theta = 1$ , so the right hand side becomes  $\int_0^{\pi} d\theta$  which is equal to  $\pi$ , and hence  $J_0^2 + 2J_1^2 + 2J_2^2 + 2J_3^2 + \dots$  is equal to 1.

Now, from here we can easily see that we are adding positive quantity  $2J_1^2, 2J_2^2, 2J_3^2$  are positive quantities which are added to  $J_0^2$ . And then, we get this as equal to 1, so  $J_0^2$  is less than or equal to 1, which implies that  $|J_0|$  is less than or equal to 1 and  $2J_n^2$  where  $n$  is equal to 1, 2, 3 and so on is less than or equal to 1.

So,  $|J_n|$  or you can say  $J_n^2$  is less than or equal to half, so  $|J_n|$  is less than or equal to  $\frac{1}{\sqrt{2}}$ , this whole square  $n$  equal to 1, 2, 3 and so on. So, these are the properties of Bessel functions of order 0 and Bessel functions of order  $n$ , where  $n$  is equal to 1, 2, 3 and so on.

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**Modified Bessel's Equation**  
 The differential equation  
 (9)  $x^2 y'' + xy' - (x^2 + \nu^2)y = 0$   
 is known as modified Bessel's equation of order  $\nu$  since it can be written as  
 $x^2 y'' + xy' + (i^2 x^2 - \nu^2)y = 0$   
 which is Bessel's equation of order  $\nu$  with the imaginary parameter  $\lambda = i$ .

Now, we shall study a modification of Bessel's equation, the differential equation  $x^2 y'' + xy' - (x^2 + \nu^2)y = 0$  is called as the modified Bessel's equation of order  $\nu$ . Because, we were write it as  $x^2 y'' + xy' + i^2 x^2 - \nu^2 y = 0$ , where  $i$  is the  $\sqrt{-1}$  and so on.

And we know that this is Bessel's equation of order  $\nu$  with the imaginary parameter  $\lambda$  equal to  $i$ , the parameter  $\lambda$  here is the imaginary  $\lambda$  equal to  $i$ .

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The solution of (9) is

$$J_{\nu}(ix) = \sum_{k=0}^{\infty} \frac{(-1)^k (ix)^{\nu+2k}}{2^{\nu+2k} k! \Gamma(\nu+k+1)}$$

$$= i^{\nu} \sum_{k=0}^{\infty} \frac{x^{\nu+2k}}{2^{\nu+2k} k! \Gamma(\nu+k+1)}$$

On multiplying both sides by  $i^{-\nu}$ , we obtain

So, the solution of 9 is  $J_{\nu}(ix) = \sum_{k=0}^{\infty} \frac{(-1)^k i^{\nu+2k} x^{\nu+2k}}{2^{\nu+2k} k! \Gamma(\nu+k+1)}$  which below also write as  $i^{\nu} \sum_{k=0}^{\infty} \frac{x^{\nu+2k}}{2^{\nu+2k} k! \Gamma(\nu+k+1)}$ .

Because,  $i$  to the power  $2k$  is  $i^2$  raised to the power  $k$ , which is  $-1$  to the power  $k$  and  $-1$  to the power  $k$  here also we have, so  $-1$  to the power  $2k$  will give us  $+1$ . Now, let us multiply both sides of this equation by  $i^{-\nu}$ .

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$$i^{-\nu} J_{\nu}(ix) = \sum_{k=0}^{\infty} \frac{x^{\nu+2k}}{2^{\nu+2k} k! \Gamma(\nu+k+1)}$$

which is also a solution of (9) since  $i^{-\nu}$  is a constant.

This new function  $i^{-\nu} J_{\nu}(ix)$  is called modified Bessel function of the first kind of order  $\nu$  and is denoted by  $I_{\nu}(x)$ .

If  $\nu$  is not an integer, the function  $I_{-\nu}(x)$  obtained from  $I_{\nu}(x)$  by replacing  $\nu$  by  $-\nu$  is a second independent solution of (9).

Then, we will obtain  $i^{-\nu} J_{\nu}(ix) = \sum_{k=0}^{\infty} \frac{x^{\nu+2k}}{2^{\nu+2k} k! \Gamma(\nu+k+1)}$  which is also a solution of (9). Since  $i^{-\nu}$  is a constant, then  $i^{-\nu} J_{\nu}(ix)$  is called modified Bessel function of first kind of order  $\nu$  and we denoted by  $I_{\nu}(x)$ . Now, if  $\nu$  is not an integer, the function  $i^{-\nu} J_{-\nu}(ix)$ , which is obtained from  $i^{-\nu} J_{\nu}(ix)$  by replacing  $\nu$  by  $-\nu$  is then a second independent solution of the equation (9).

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Thus

$$I_{-\nu}(x) = i^{-\nu} J_{-\nu}(ix) = \sum_{k=0}^{\infty} \frac{x^{-\nu+2k}}{2^{-\nu+2k} k! \Gamma(k-\nu+1)}$$

Therefore, the complete solution of (9) is

$$y(x) = A I_{\nu}(x) + B I_{-\nu}(x).$$

As in the case of Bessel function of second kind we may take the second solution of modified Bessel equation to be the linear combination

And thus  $x^{-\nu}$  is equal to  $x^i$  to the power  $\nu$  into  $x^{-i}$ , which is equal to  $\sum_{k=0}^{\infty} x^{\nu+2k} / 2^k \Gamma(k-\nu+1)$ . And thus the complete solution of equation 9 we may write as  $y = a x^{\nu} + b x^{-\nu}$ .

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$$K_{\nu}(x) = \frac{\pi I_{-\nu}(x) - I_{\nu}(x)}{2 \sin \nu \pi}$$

and define for  $\nu = n$  (an integer)

$$K_n(x) = \lim_{\nu \rightarrow n} K_{\nu}(x) = \lim_{\nu \rightarrow n} \frac{\pi I_{-\nu}(x) - I_{\nu}(x)}{2 \sin \nu \pi}$$

It can be shown that the limit exists. The function  $K_{\nu}(x)$  is known as modified Bessel function of the second kind.

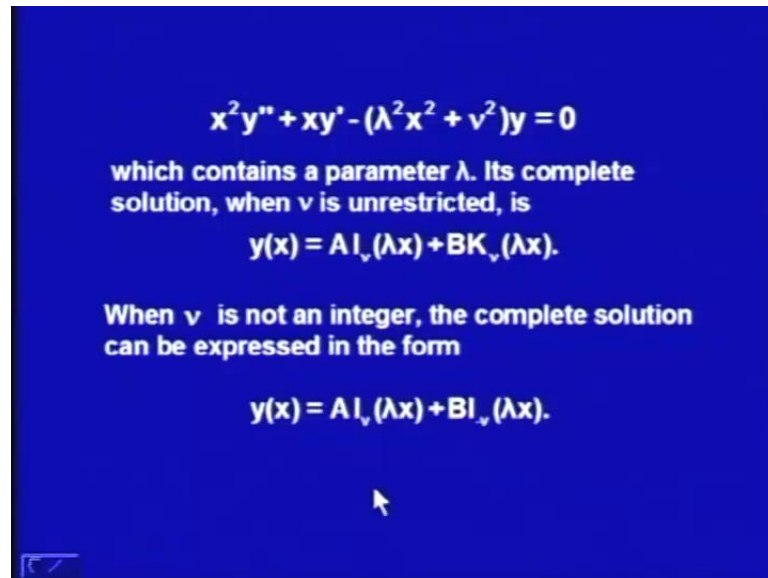
Equation (9) is a special case ( $\lambda=1$ ) of the equation

Now, as in the case of Bessel function of second kind, we may take the second solution of modified Bessel equation to be the linear combination of  $x^{\nu}$  and  $x^{-\nu}$  when  $\nu$  is not an integer, that is we can take  $K_{\nu}(x) = \frac{\pi}{2 \sin \nu \pi} (x^{-\nu} - I_{\nu}(x))$ . And for  $\nu$  equal to  $n$  an integer we define  $K_n(x) = \lim_{\nu \rightarrow n} K_{\nu}(x) = \lim_{\nu \rightarrow n} \frac{\pi}{2 \sin \nu \pi} (x^{-\nu} - I_{\nu}(x))$ .

Now, we can show that this limit exists, so when  $\nu$  is not only integer  $K_{\nu}(x)$  being the linear combination of  $x^{\nu}$  and  $x^{-\nu}$  is also a solution of the equation 9. And when  $\nu$  is the integer,  $K_n(x)$  being the limit of  $K_{\nu}(x)$  and since the limit exists is also a solution of the equation 9. And therefore, we call the function  $K_{\nu}(x)$  as the modified Bessel function of the second kind, now equation 9 is a special case of the following equation.



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$x^2 y'' + xy' - (\lambda^2 x^2 + \nu^2)y = 0$

which contains a parameter  $\lambda$ . Its complete solution, when  $\nu$  is unrestricted, is

$$y(x) = A I_\nu(\lambda x) + B K_\nu(\lambda x).$$

When  $\nu$  is not an integer, the complete solution can be expressed in the form

$$y(x) = A I_\nu(\lambda x) + B I_{-\nu}(\lambda x).$$

If you put lambda equal to 1 in this equation, we get  $x^2 y'' + xy' - (\lambda^2 x^2 + \nu^2)y = 0$  reduces to  $x^2 y'' + xy' - x^2 y = 0$ , which contains a parameter lambda. Its complete solution is, therefore  $y(x) = A I_\nu(\lambda x) + B K_\nu(\lambda x)$ , where  $\nu$  can take integer or non integral value, so there is no restriction on  $\nu$ .

Now, when  $\nu$  is not an integer the complete solution may also be expressed as  $y(x) = A I_\nu(\lambda x) + B I_{-\nu}(\lambda x)$ , because we know that  $I_{-\nu}(\lambda x)$  is independent of  $I_\nu(\lambda x)$ . So, modified Bessel's equation is the particular case of this equation of lambda equal to 1, whose solution are given by for all values of  $\nu$  by  $y(x) = A I_\nu(\lambda x) + B K_\nu(\lambda x)$ . And when  $\nu$  is not an integer  $y(x)$  is equal to  $A I_\nu(\lambda x) + B I_{-\nu}(\lambda x)$ .

In our next lecture we shall discuss the Laplace transformation, Laplace transformation is very useful in engineering mathematics, the linear differential equation with constant coefficients are solved by using Laplace transformation method. When we solve the linear differential equation with constant coefficients by finding the complementary function and particular integral, then the constants that occur there, their values are obtained after putting the conditions given.

While in the case Laplace transformation it takes in to account the initial conditions while taking the Laplace transformation itself. So, we do not have to find the values of the arbitrary constants after we have found the general solution as the case while solving the linear differential equation with constant coefficient. So, it is very useful for solving linear differential equation with constant coefficients, we can also solve simultaneous system of simultaneous linear equations by using Laplace transformation method. And that is used in the study of electric circuits, then dynamical systems and then, many other problems, so we will be discussing Laplace transformation in our next lecture.

Thank you.