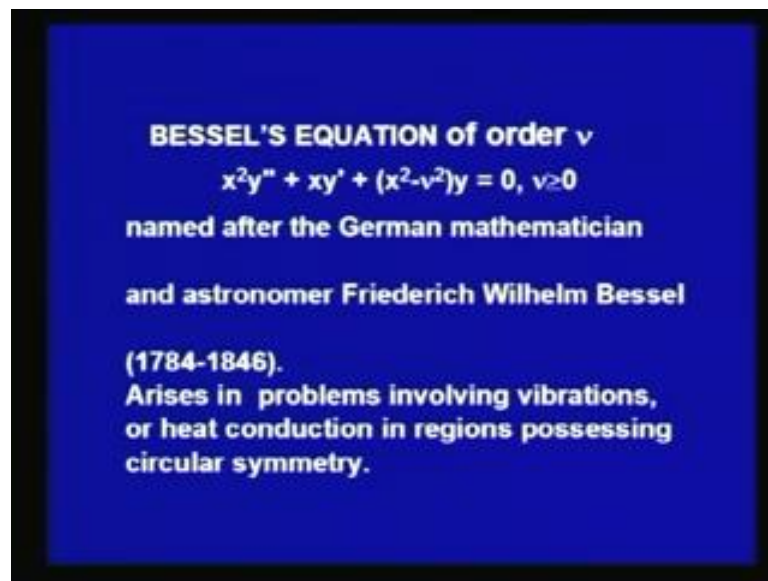


**Mathematics III**  
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**Indian Institute of Technology, Roorkee**

**Lecture - 6**  
**Bessel Functions and Their Properties**

Dear viewers, in my last lecture we had seen how to find the general solution of a homogeneous linear differential equation of second order with variable coefficients about a regular singular point. An important differential equation of this type is the Bessel's equation. So in my lecture today we shall be discussing the solutions of the Bessel's equation. The solutions of the Bessel's equations are called as Bessel functions, so we shall study Bessel functions and their properties.

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Let us first define the Bessel's equation, a Bessel's equation of order  $\nu$  is given by  $x^2y'' + xy' + (x^2 - \nu^2)y = 0$ , where  $\nu$  is greater than or equal to 0. This equation is, so called after the German mathematician and astronomer Friedrich Wilhelm Bessel from 1784 to 1846, now this differential equation arises in problems which involve vibrations are heat conduction in regions that possess circular symmetry.

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Since  $x=0$  is a regular singular point of the equation, by Frobenius method, substituting

$$y(x) = \sum_{m=0}^{\infty} c_m x^{m+r}, c_0 \neq 0.$$

in the Bessel's equation, we obtain

$$\sum_{m=0}^{\infty} [(m+r)^2 - \nu^2] c_m x^{m+r} + \sum_{m=0}^{\infty} c_m x^{m+r+2} = 0.$$

Hence the indicial equation is

$$r^2 - \nu^2 = 0, \text{ as } c_0 \neq 0.$$

Now, from the differential equation it is clear that  $x$  equal to 0 is a regular singular point of the differential equation and, so we can apply Frobenius method to this differential equation. Let us therefore, substitute  $y$  equal to  $\sum_{m=0}^{\infty} c_m x^{m+r}$ , where  $c_0$  is not equal to 0 in the Bessel's equation. After simplification we will obtain  $\sum_{m=0}^{\infty} (m+r)^2 c_m x^{m+r} - \nu^2 \sum_{m=0}^{\infty} c_m x^{m+r} + \sum_{m=0}^{\infty} c_m x^{m+r+2} = 0$ .

Now, this equation is an identity therefore, the coefficients of various powers of  $x$  can be equated to 0, when we equate it to 0 the coefficient of the least power of  $x$  in  $r$  that is we equate to 0 the coefficient of  $x$  to the power  $r$  to 0, we get the indicial equation, the indicial equation is therefore, given by  $r^2 - \nu^2 = 0$  as  $c_0$  is not equal to 0.

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Thus,  $r = \pm \nu$ . Equating the coeff. of  $x^{r+1}$  to zero, we have

$$[(r+1)^2 - \nu^2]c_1 = 0$$

$\Rightarrow c_1 = 0$  as  $r = \pm \nu$  except when  $r = -1/2$ .

In the case  $r = -1/2$ ,  $c_1$  is arbitrary but without any loss of generality, we can assume  $c_1 = 0$ .

And this gives us the roots of the indicial equation as  $r$  equal to plus minus  $\nu$ , now let us equate the coefficient of next higher power of  $x$  that is let us equate the coefficient of  $x$  to the power  $r$  plus 1 to 0, we shall have  $(r+1)^2 - \nu^2 = 0$ . Now, since  $r = \pm \nu$ , so  $r^2 = \nu^2$  and therefore, this equation  $(r+1)^2 - \nu^2 = 0$  implies  $c_1 = 0$  in all cases, except one  $r$  is equal to minus half.

In the case  $r$  is equal to minus half both sides become 0 and therefore,  $c_1$  can take any value, so  $c_1$  is an arbitrary constant. But, then what happens is that when we take  $c_1$  as an arbitrary constant and find the solution corresponding to  $r$  equal to minus half, then the terms which correspond to  $c_1$  are observed in the solution, which we find for the other value of  $r$  that is  $r$  equal to half. And, so no  $\nu$  part of the solution is obtained corresponding to  $c_1$  being arbitrary and therefore, we shall without any loss of generality we can assume that  $c_1$  is equal to 0 even in this case.

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**Recurrence formula**

$$c_{m+2} = -\frac{c_m}{(m+r+2-\nu)(m+r+2+\nu)}, m \geq 0$$

follows on equating the coeff. of  $x^{m+r+2}$  to zero .

$c_1 = 0$  implies  $c_m = 0$  for  $m = 3, 5, 7, \dots$

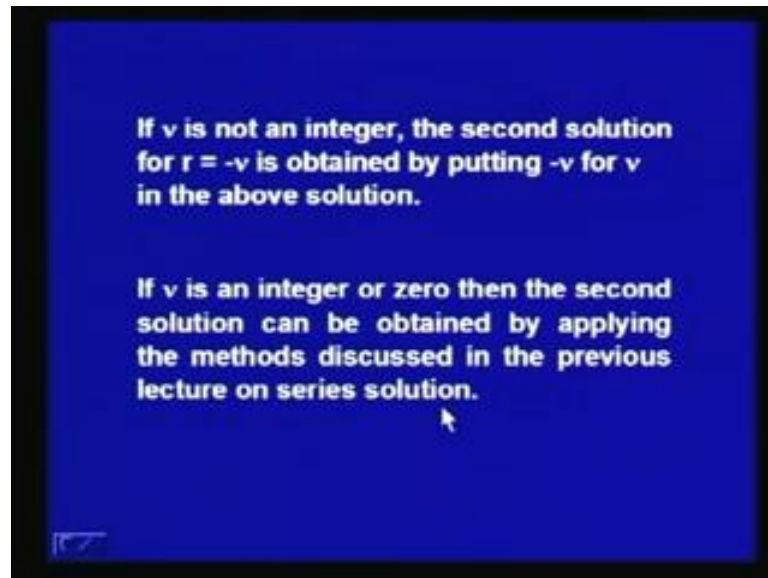
For  $r = \nu$ , we get one of the linearly independent solutions as

$$y_1(x) = c_0 x^\nu \left[ 1 - \frac{x^2}{2(2\nu+2)} + \frac{x^4}{2 \cdot 4(2\nu+2)(2\nu+4)} + \dots \right]$$

Next, let us equate to 0 the coefficient of  $x$  to the power  $m$  plus  $r$  plus 2 to 0 in the equation, we will have  $c_{m+2}$  equal to minus  $c_m$  over  $m$  plus  $r$  plus 2 minus  $\nu$  into  $m$  plus  $r$  plus 2 plus  $\nu$  for all  $m$  greater than or equal to 0 that is  $m$  equal to 0, 1, 2, 3 and so on. Now,  $c_1$  is equal to 0 implies that  $c_m$  is from this recurrence relation it follows that when  $c_1$  is 0,  $c_3$  is 0 and  $c_3$  is 0 implies  $c_5$  is 0. So,  $c_m = 0$  implies that  $c_m$  is equal to 0 for  $m$  equal to 3, 5, 7 and so on.

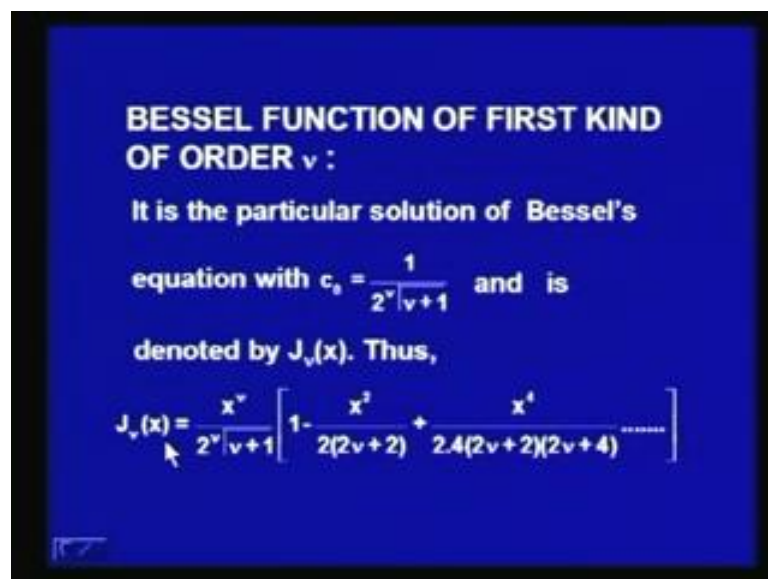
Thus for  $r$  equal to  $\nu$  we can get one of the linearly independent solutions as  $y_1(x)$  equal to  $c_0 x^\nu$  into  $1 - \frac{x^2}{2(2\nu+2)} + \frac{x^4}{2 \cdot 4(2\nu+2)(2\nu+4)} + \dots$  and so on. Where, the values of  $c_2$  and  $c_4$  and so on have been obtained from the recurrence relation  $c_{m+2}$  equal to minus  $c_m$  over  $m$  plus  $r$  plus 2 minus  $\nu$  into  $m$  plus  $r$  plus 2 plus  $\nu$ .

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Now, if  $\nu$  is not an integer the second solution for  $r$  equal to minus  $\nu$  is obtained by putting minus  $\nu$  for  $\nu$  in the above solution. However, if  $\nu$  is an integer or 0 then the second solution will have to be obtained by the methods, which we have discussed in the previous lecture on series solution that is if  $\nu$  is 0 then the two roots of the indicial equation will be both 0. So, we shall apply the case of the indicial equation having equal roots, if  $\nu$  is an integer we shall apply the case third, where we had discussed the case of indicial equation having roots which differ by an integer.

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Let us, now discuss Bessel function of first kind of order nu, Bessel function of first kind of order nu is the particular solution of Bessel's equation, where we take c naught equal to 1 over 2 to the power nu into gamma nu plus 1, it is denoted by J nu x. So, thus J nu x becomes x to the power nu over 2 to the power nu into gamma nu plus 1 1 minus x square over 2 into 2 nu plus 2 plus x to the power 4 2 into 4 2 nu plus 2 2 nu plus 4 and so on. This J nu x we have obtained by replacing c naught in y 1 x y 1 over 2 to the power nu gamma nu plus 1. Now, this choice of c naught equal to 1 over 2 to the power nu gamma nu plus 1 has be made in order to get complicated in order to get a simplified expression for J nu x.

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or

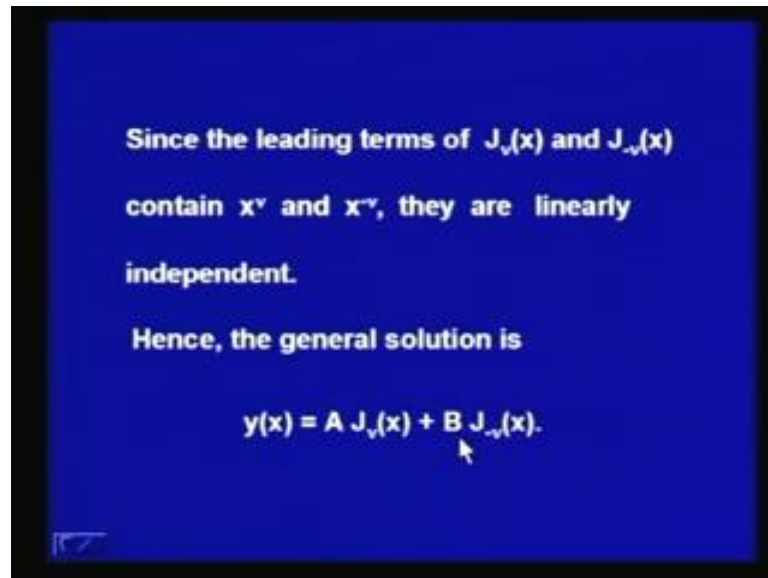
$$J_\nu(x) = \sum_{k=0}^{\infty} (-1)^k \frac{x^{\nu+2k}}{2^{\nu+2k} (k!) \Gamma(\nu+k+1)}$$

If  $\nu$  is not an integer, the second solution is

$$J_{-\nu}(x) = \sum_{k=0}^{\infty} (-1)^k \frac{x^{-\nu+2k}}{2^{-\nu+2k} (k!) \Gamma(-\nu+k+1)}$$

As we see now with this choice of c naught x, we can write J nu x in the form of summation as sigma k equal to 0 to infinity minus 1 to the power k x to the power nu plus 2 k over 2 to the power nu plus 2 k. Now, this can be combined and we may write as x over 2 raised to the power nu plus 2 k into 1 over then 1 over k factorial and 1 over gamma nu plus k plus 1. If nu is not an integer, then the second solution can be obtained by replacing nu by minus nu here. So, say J minus nu x will be equal to sigma k equal to 0 to infinity minus 1 to the power k x to the power minus nu plus 2 k over 2 to the power nu plus 2 k into k factorial gamma minus nu plus k plus 1.

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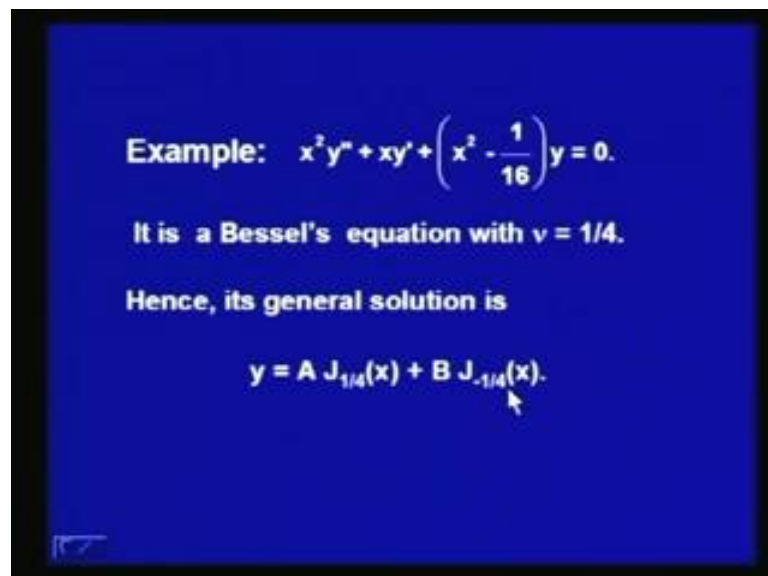
Since the leading terms of  $J_\nu(x)$  and  $J_{-\nu}(x)$  contain  $x^\nu$  and  $x^{-\nu}$ , they are linearly independent.

Hence, the general solution is

$$y(x) = A J_\nu(x) + B J_{-\nu}(x).$$

Now, in the expressions for  $J_\nu x$  and  $J_{\text{minus } \nu} x$ , we can see that the leading terms are containing  $x$  to the power  $\nu$  and  $x$  to the power  $\text{minus } \nu$ . Therefore, the two solutions  $J_\nu x$  and  $J_{\text{minus } \nu} x$  of the Bessel's equation are linearly independent, that is one is not a scalar multiple of the other. And hence, we can write the general solution of the Bessel's equation for the case where  $\nu$  is not an integer as  $y x$  equal to  $A$  times  $J_\nu x$  plus  $B$  times  $J_{\text{minus } \nu} x$ .

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Example:  $x^2 y'' + xy' + \left(x^2 - \frac{1}{16}\right)y = 0.$

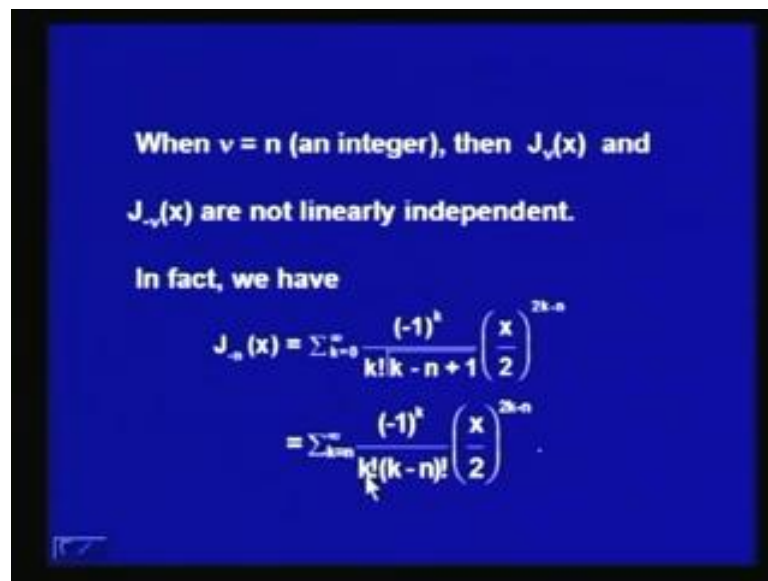
It is a Bessel's equation with  $\nu = 1/4.$

Hence, its general solution is

$$y = A J_{1/4}(x) + B J_{-1/4}(x).$$

Let us study an example,  $x^2 y'' + x y' + x^2 y = 0$ . It is a Bessel's equation, where we have taken  $\nu$  equal to  $1/4$  if you compare this equation with the standard form of the Bessel's equation, you can see that  $\nu$  is equal to  $1/4$  here. So,  $\nu$  is not an integer and therefore, the two solutions  $J_{\nu}(x)$  and  $J_{-\nu}(x)$  of this Bessel's equation will be linearly independent of each other and therefore, hence we may write its general solution as  $y = A J_{1/4}(x) + B J_{-1/4}(x)$ .

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When  $\nu$  is equal to  $n$  let us say an integer then  $J_{\nu}(x)$  and  $J_{-\nu}(x)$  are not linearly independent, we are going to show that, in fact  $J_{-n}(x)$  is equal to  $(-1)^n$  times  $J_n(x)$ . Let us find the value of  $J_{-n}(x)$ ,  $J_{-n}(x)$  can be written as  $\sum_{k=0}^{\infty} \frac{(-1)^k}{k! \Gamma(k-n+1)} \left(\frac{x}{2}\right)^{2k-n}$ , replacing  $\nu$  by  $n$  in the expression for  $J_{\nu}(x)$ .

Now, here  $\Gamma(k-n+1)$  can be replaced by  $(k-n)!$  because,  $k$  is taking values  $0, 1, 2, 3$  and so on it is second integral values  $n$  also  $\nu$  as is an integer. So,  $k-n+1$  is an integer and, so  $\Gamma(k-n+1)$  will be equal to  $(k-n)!$ , but then  $k$  will start with  $n$  onwards  $k$  will go from  $n$  to infinity.



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$$\begin{aligned} J_{-n}(x) &= \sum_{m=0}^{\infty} \frac{(-1)^{n+m}}{(n+m)!m!} \left(\frac{x}{2}\right)^{2m+n} \\ &= (-1)^n \sum_{m=0}^{\infty} \frac{(-1)^m}{m!(n+m)!} \left(\frac{x}{2}\right)^{2m+n} \\ &= (-1)^n J_n(x). \end{aligned}$$

So, replacing  $k$  by  $k$  minus  $n$  by  $m$  we will have the value of  $J$  minus  $n$   $x$  as sigma  $m$  equal to  $0$  to infinity minus  $1$  to the power  $n$  plus  $m$  over  $n$  plus  $m$  factorial into  $m$  factorial  $x$  by  $2$  rest to the power  $2$   $m$  plus  $n$ , which is equal to minus  $1$  to the power  $n$  into sigma  $m$  equal to  $0$  to infinity minus  $1$  to the power  $m$  over  $m$  factorial  $n$  plus  $m$  factorial  $x$  by  $2$  raised to the  $2$   $m$  plus  $n$ .

And which as we know it is minus  $1$  to the power  $n$  into  $J$   $n$   $x$  this is the expression for  $j$   $n$   $x$ . So, when  $n$  is an integer equal to  $m$ , then  $J$  minus  $n$   $x$  and  $J$   $n$   $x$  are not linearly independent one is a scalar multiple of the other it is  $J$  minus  $n$   $x$  is equal to minus  $1$  to the power  $n$  into  $J$   $n$   $x$ . So, when  $n$  is an integer the second solution has to be obtained by the cases that we have discussed in the last lecture.

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**Bessel function of order 0:**

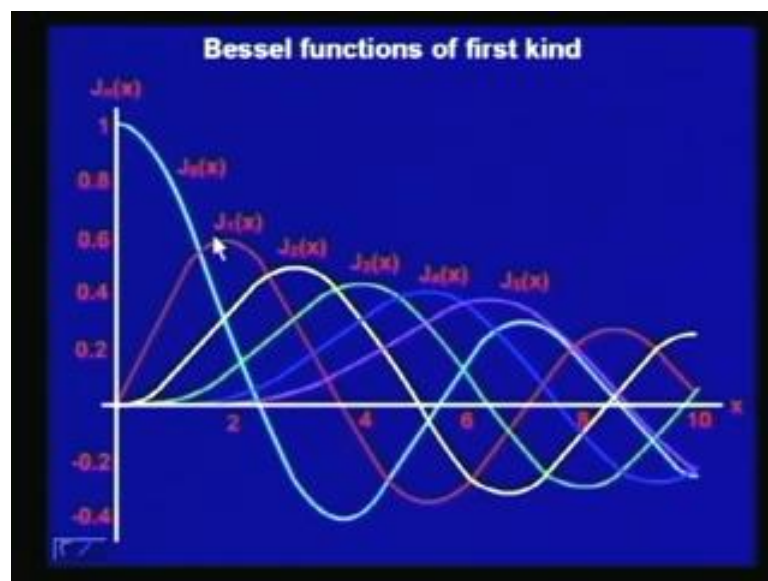
$$J_0(x) = 1 - \frac{x^2}{2^2} + \frac{x^4}{2^2 \cdot 4^2} - \frac{x^6}{2^2 \cdot 4^2 \cdot 6^2} + \dots$$

**Bessel function of order 1:**

$$J_1(x) = \frac{x}{2} - \frac{x^3}{2^2 \cdot 4} + \frac{x^5}{2^2 \cdot 4^2 \cdot 6} - \dots$$

Let us, take an nu equal to 0 we will have the Bessel function of order 0, Bessel function of order 0 is given by 1 minus x square by 2 square plus x to the power 4 by 2 square into 4 square minus x to the power 6 over 2 square 4 square 6 square and so on. This infinite series expression for a J naught x we have obtained from the corresponding expression for J nu x by replacing nu y 0. Next let us take nu equal to 1 we have the expression for J 1 x as x by 2 minus x cube over 2 square into 4, then x to the power 5 over 2 square 4 square into 6 and so on.

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In this picture, we see the graphs of Bessel functions of first kind, this the graph of  $J_0 x$ , this is the graph of  $J_1 x$ , then this the graph of  $J_2 x$ ,  $J_3 x$ ,  $J_4 x$ ,  $J_5 x$  and so on. We can see the close resemblance of the graphs of  $J_0 x$  and  $J_1 x$  with the functions  $\cos x$  and  $\sin x$  and that is quite interesting, they closely resemble the graphs of cosine  $x$  and sin  $x$  functions.

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**Bessel functions of order  $\frac{1}{2}$ :**

$$J_{1/2}(x) = \frac{\sqrt{x}}{\sqrt{2} \Gamma(3/2)} \left( 1 - \frac{x^2}{2 \cdot 3} + \frac{x^4}{2 \cdot 4 \cdot 3 \cdot 5} - \dots \right)$$

$$= \frac{1}{\sqrt{x} \sqrt{2} \frac{1}{2} \Gamma(1/2)} \left( x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots \right)$$

$$= \sqrt{2/\pi x} \sin x$$

Now, let us study the Bessel functions of order half, we can see that the infinite series expansion of  $J_{1/2} x$  can be put in the closed form that is in terms of the  $\sin x$  function  $J_{1/2} x$  is equal to  $\frac{\sqrt{x}}{\sqrt{2} \Gamma(3/2)} \left( 1 - \frac{x^2}{2 \cdot 3} + \frac{x^4}{2 \cdot 4 \cdot 3 \cdot 5} - \dots \right)$  which is equal to  $\frac{1}{\sqrt{x} \sqrt{2} \frac{1}{2} \Gamma(1/2)} \left( x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots \right)$  and this becomes after multiplying by  $\sqrt{x}$  in the numerator and denominator, this becomes  $x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots$  and so on. And  $x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots$  is the  $\sin x$  function. So, we can write  $J_{1/2} x$  as equal to  $\frac{\sqrt{2}}{\sqrt{\pi x}} \sin x$ .

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and

$$J_{-1/2}(x) = \frac{x^{-1/2}}{2^{-1/2} \Gamma(1/2)} \left[ 1 - \frac{x^2}{2 \cdot 1} + \frac{x^4}{2 \cdot 4 \cdot 1 \cdot 3} \dots \right]$$
$$= \sqrt{\frac{2}{\pi x}} \cos x.$$

And  $J_{-1/2} x$  is equal to  $x$  to the power minus half over  $2$  to the power minus half gamma half into  $1 - \frac{x^2}{2 \cdot 1} + \frac{x^4}{2 \cdot 4 \cdot 1 \cdot 3}$ , which is equal to square root  $2$  over  $\pi x$  into  $\cos x$  the expression inside the bracket as you can see is the expression of  $\cos x$  it is  $1 - \frac{x^2}{2 \text{ factorial}} + \frac{x^4}{4 \text{ factorial}}$  and so on. So,  $J_{-1/2} x$  can be expressed in the closed form as the square root  $2$  over  $\pi x$  into cosine  $x$ .

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**RECURRENCE FORMULAE:**

Bessel's function of the first kind is one of the most important special functions and satisfies a large no. of relationships.

**Derivatives and integrals of Bessel functions:**

(1)  $\frac{d}{dx} (x^{-\nu} J_{\nu}(x)) = -x^{-\nu} J_{\nu+1}(x)$

The Bessel functions of first kind possess some very important relationships, we are going to study the relationships, which the Bessel's functions of first kind satisfied. So, let us first see that  $\frac{d}{dx} (x^\nu J_\nu(x))$  is equal to  $x^\nu J_{\nu-1}(x)$ .

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We have

$$x^\nu J_\nu(x) = \sum_{m=0}^{\infty} \frac{(-1)^m x^{2m+\nu}}{m!(m+\nu+1)2^{2m+\nu}}$$

Hence,

$$\begin{aligned} \frac{d}{dx} (x^\nu J_\nu(x)) &= \sum_{m=0}^{\infty} \frac{(-1)^m (2m+2\nu)x^{2m+2\nu-1}}{m!(m+\nu)(m+\nu+1)2^{2m+\nu}} \\ &= x^\nu \sum_{m=0}^{\infty} \frac{(-1)^m}{m!(m+\nu)} \left(\frac{x}{2}\right)^{2m+\nu-1} \\ &= x^\nu J_{\nu-1}(x). \end{aligned}$$

Let us multiply the expression for  $J_\nu(x)$  by  $x^\nu$ , we shall have  $x^\nu J_\nu(x)$  is equal to  $\sum_{m=0}^{\infty} \frac{(-1)^m x^{2m+\nu}}{m!(m+\nu+1)2^{2m+\nu}}$ . Now, let us differentiate this with respect to  $x$ , so  $\frac{d}{dx} (x^\nu J_\nu(x))$  will give us  $\sum_{m=0}^{\infty} \frac{(-1)^m (2m+2\nu)x^{2m+2\nu-1}}{m!(m+\nu)(m+\nu+1)2^{2m+\nu}}$ .

Because of the property of gamma function,  $\Gamma(m+\nu+1)$  can be written as  $(m+\nu) \Gamma(m+\nu)$ . Now, we may cancel  $(m+\nu)$  in the numerator and denominator and then we shall have the right hand side equal to  $x^\nu \sum_{m=0}^{\infty} \frac{(-1)^m}{m!(m+\nu)} \left(\frac{x}{2}\right)^{2m+\nu-1}$ , which is equal to  $x^\nu J_{\nu-1}(x)$ .

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Hence,

$$\int x^\nu J_{\nu-1}(x) dx = x^\nu J_\nu(x) + c$$

Setting  $\nu = 0$  in (1), we get  $J'_0(x) = J_{-1}(x) = -J_1(x)$ , as  $J_{-n}(x) = (-1)^n J_n(x)$  if  $n$  is an integer.

Similarly,

$$(2) \quad \frac{d}{dx} (x^{-\nu} J_\nu(x)) = -x^{-\nu} J_{\nu+1}(x)$$

or

$$\int x^{-\nu} J_{\nu+1}(x) dx = -x^{-\nu} J_\nu(x) + c.$$

Now, integrating both sides of this relationship we obtain integral of  $x$  to the power  $\nu$  into  $J_{\nu-1}$   $x$   $dx$  equal to  $x$  to the power  $\nu$  into  $J_\nu$   $x$  plus  $c$ . If you set  $\nu$  equal to  $0$  in that relationship  $d$  over  $dx$  of  $x$  to the power  $\nu$  into  $J_\nu$   $x$  equal to  $x$  to the power  $\nu$  into  $J_{\nu-1}$   $x$ , we obtain that the derivative of  $J_\nu$   $x$  that is  $J_\nu$  dash  $x$  is equal to  $J_{\nu-1}$   $x$ . And we know that  $J_{-n}$   $x$  is equal to  $(-1)^n$  to the power  $n$  into  $J_n$   $x$  where  $n$  is an integer.

So,  $J_{-1}$   $x$  will be equal to  $(-1)^1 J_1$   $x$  and therefore,  $J_\nu$  dash  $x$  is equal to  $(-1)^{\nu-1} J_{\nu-1}$   $x$ . Now, the next result that we have is  $d$  over  $dx$  of  $x$  to the power  $\nu$  minus  $\nu$   $J_\nu$   $x$  is equal to  $(-1)^{\nu-1} J_{\nu+1}$   $x$ , so here the order of the Bessel function increases by  $1$   $J_\nu$   $x$  becomes  $J_{\nu+1}$   $x$ , in the previous relationship  $J_\nu$   $x$  becomes  $J_{\nu-1}$   $x$ . So, here when we will integrate we will get integral of  $x$  to the power  $\nu$  minus  $\nu$   $J_{\nu+1}$   $x$   $dx$  equal to  $(-1)^{\nu-1} x$  to the power  $\nu$  minus  $\nu$   $J_\nu$   $x$  plus  $c$ .

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From (1) and (2), we obtain

$$J_v'(x) = J_{v-1}(x) - \frac{v}{x} J_v(x)$$

and

$$J_v'(x) = \frac{v}{x} J_v(x) - J_{v+1}(x)$$

on adding and subtracting these equations we get

$$2J_v'(x) = J_{v-1}(x) - J_{v+1}(x)$$

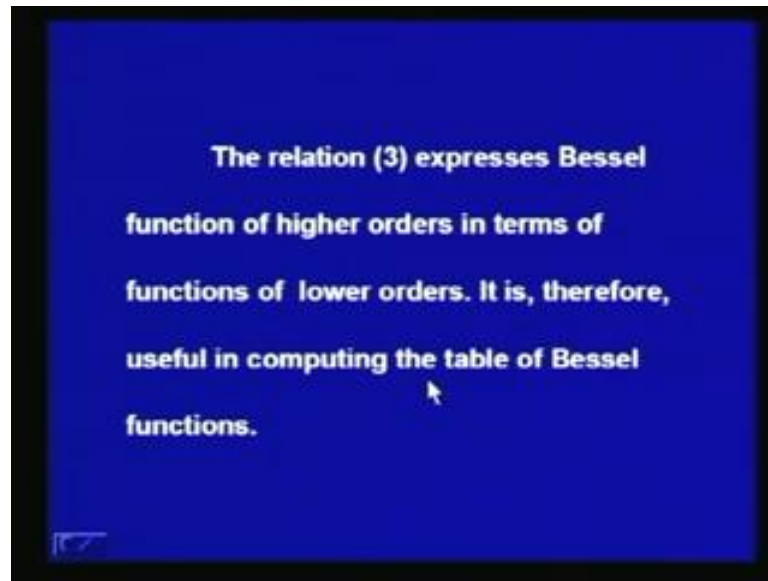
and

$$(3) \quad \frac{2v}{x} J_v(x) = J_{v-1}(x) + J_{v+1}(x)$$

From these two relationships we will obtain  $J_v'(x)$  the value of  $J_v'(x)$  as  $x$  is equal to  $J_{v-1}(x) - \frac{v}{x} J_v(x)$  and  $J_v'(x) = \frac{v}{x} J_v(x) - J_{v+1}(x)$ . Now, adding these two equations and subtracting we shall have the following, if you add these two equations you get  $2J_v'(x) = J_{v-1}(x) - J_{v+1}(x)$  and if you subtract this equation, this equation from this one you get the left hand side as 0.

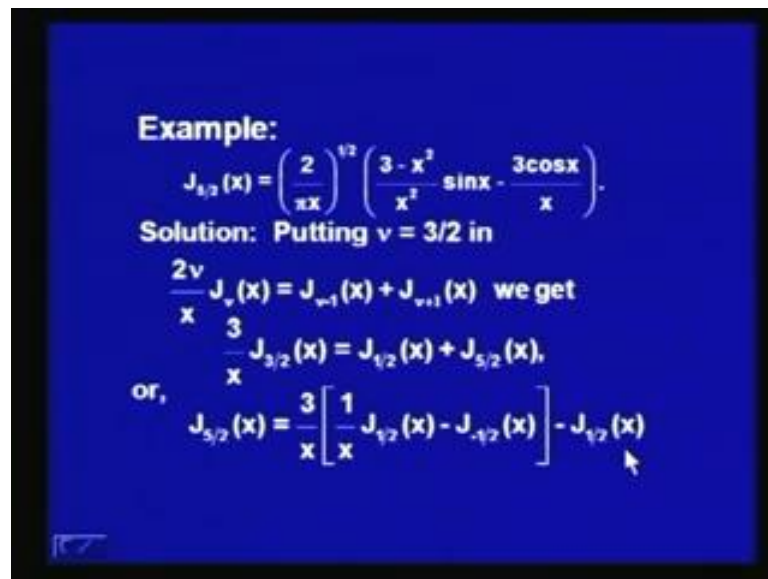
And therefore, we will be getting  $2 \frac{v}{x} J_v(x) = J_{v-1}(x) + J_{v+1}(x)$ , which would give us  $2v J_v(x) = x J_{v-1}(x) + x J_{v+1}(x)$ . In three expresses the Bessel function of order  $v+1$  in terms of the Bessel's functions of order  $v$  and  $v-1$  and, so if we know the Bessel functions of lower orders that is  $J_{v-1}(x)$  and  $J_v(x)$ , then we can get the Bessel function of  $J_{v+1}(x)$  from this equation and therefore, we call this formula as the recurrence formula.

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This formula is therefore, useful in computing the table of Bessel functions.

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Let us take an example based on this recurrence relation, let us show that  $J_{5/2}(x) = \frac{2}{\pi x} \left( \frac{3-x^2}{x^2} \sin x - \frac{3\cos x}{x} \right)$ . So, let us put  $\nu$  equal to  $3/2$  in the recurrence relation, we will have the recurrence relation as we know is  $2\nu/x J_\nu(x) = J_{\nu-1}(x) + J_{\nu+1}(x)$  we have chosen  $\nu$  equal to  $3/2$ .



Because, the  $J_{\nu+1}(x)$  will then become  $J_{5/2}(x)$  and it will be expressible in terms of the Bessel functions of lower order, that is  $J_{3/2}$  and Bessel function of order half. So, making use of this recurrence relation and taking  $\nu$  equal to  $3/2$ , we can see that we get  $J_{3/2}(x) + J_{5/2}(x) = 3/x J_{3/2}(x) - J_{1/2}(x)$  or we may write  $J_{5/2}(x) = 3/x J_{3/2}(x) - J_{1/2}(x)$ , the value of  $J_{3/2}(x)$  will turn out to be from this recurrence relation taking  $\nu$  equal to half as  $J_{1/2}(x) - J_{-1/2}(x)$ . So, the value of  $J_{5/2}(x)$  becomes  $3/x J_{1/2}(x) - J_{-1/2}(x)$ .

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$$= \left( \frac{3-x^2}{x^2} \right) J_{1/2}(x) - \frac{3}{x} J_{-1/2}(x)$$

$$= \left( \frac{2}{\pi x} \right)^{1/2} \left( \frac{3-x^2}{x^2} \sin x - \frac{3}{x} \cos x \right).$$

And when we simplify this we get the right hand side as  $3/x^2 J_{1/2}(x) - 3/x J_{-1/2}(x)$ . Now, we know the closed forms of  $J_{1/2}(x)$  and  $J_{-1/2}(x)$ , so let us make use of them we will get the right hand side as  $2/\sqrt{\pi x} \left( (3-x^2)/x^2 \sin x - 3/x \cos x \right)$ .

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**Example:**

$$\frac{d}{dx} (x J_n(x) J_{n+1}(x)) = x (J_n^2(x) - J_{n+1}^2(x)).$$

**Solution:**

$$\begin{aligned} \text{L.H.S.} &= \frac{d}{dx} \{ (x^{-n} J_n(x)) (x^{n+1} J_{n+1}(x)) \} \\ &= (x^{-n} J_n(x)) \frac{d}{dx} (x^{n+1} J_{n+1}(x)) + \frac{d}{dx} (x^{-n} J_n(x)) (x^{n+1} J_{n+1}(x)) \\ &= (x^{-n} J_n(x)) (x^{n+1} J_{n+1}(x)) - (x^{-n} J_{n+1}(x)) (x^{n+1} J_n(x)) \\ &= x (J_n^2(x) - J_{n+1}^2(x)). \end{aligned}$$

Let us take an example based on the other relationships, which involve the derivatives of Bessel functions. Let us take the example of  $\frac{d}{dx} (x J_n(x) J_{n+1}(x))$ , we are going to show that it is equal to  $x (J_n^2(x) - J_{n+1}^2(x))$ . The left hand side of this example can be expressed as  $\frac{d}{dx} (x^{-n} J_n(x) x^{n+1} J_{n+1}(x))$ . This we have done in order to be able to find the derivatives that will arise when we differentiate the product of these two functions of  $x$ .

So,  $x^{-n} J_n(x)$  that is the first function of  $x$  multiplied by the derivative of the second function  $\frac{d}{dx} (x^{n+1} J_{n+1}(x))$  and then plus derivative of the first function  $\frac{d}{dx} (x^{-n} J_n(x)) x^{n+1} J_{n+1}(x)$ . Now, let us make use of the properties for derivatives of Bessel functions  $x^{-n} J_n(x)$  is same and then the derivative of  $x^{n+1} J_{n+1}(x)$ , this becomes  $x^{n+1} J_n(x)$ .

Because, we had the one result as  $\frac{d}{dx} (x^{\nu} J_{\nu}(x)) = x^{\nu} J_{\nu-1}(x)$  equal to  $x^{\nu} J_{\nu}(x) - x^{\nu+1} J_{\nu+1}(x)$  the derivative of that we had seen was  $x^{\nu} J_{\nu-1}(x)$ . So, making use of that the derivative of  $x^{n+1} J_{n+1}(x)$  will be  $x^{n+1} J_n(x)$ . And then  $x^{-n} \frac{d}{dx} (x^{-n} J_n(x))$  will be  $-x^{-n} J_{n+1}(x)$  using the other property of the Bessel function, which involve the

derivatives and multiplied by  $x$  to the power  $n + 1$  into  $J_{n+1} x$ . Now, let us simplify this it will give us  $x$  into  $J_n$  square  $x$  minus  $J_{n+1}$  square  $x$  that is the right hand side of the given problem.

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**Example:** Evaluate  $\int x^{-1} J_4(x) dx$ .

**Solution:**

$$\begin{aligned} \int x^{-1} J_4(x) dx &= \int x^2 \{x^{-3} J_4(x)\} dx \\ &= -x^2 \{x^{-3} J_3(x)\} + \int 2x \{x^{-3} J_3(x)\} dx \\ &\quad \text{(using } \int x^{-3} J_4(x) dx = -x^{-3} J_3(x) + c) \\ &= -x^{-1} J_3(x) + 2 \int x^{-2} J_3(x) dx \\ &= -x^{-1} J_3(x) + 2(-x^{-2} J_2(x)) + c \\ &= -x^{-1} J_3(x) - 2x^{-2} J_2(x) + c. \end{aligned}$$

Now, let us evaluate the integral of  $x$  to the power minus 1  $J_4 x dx$ , so integral of  $x$  to the power minus 1  $J_4 x$  can be then written as integral of  $x$  square into  $x$  to the power minus 3  $J_4 x dx$  this arrangement also we have done in order to be able to find the integral of  $x$  square into  $x$  to the power minus 3  $J_4 x dx$ . So, we have first function that is into integral of  $x$  to the power minus 3  $J_4 x$  we know, we have seen earlier it is equal to minus  $x$  to the power minus 3  $J_3 x$  and then we have plus integral derivative of  $x$  square is  $2x$  then  $x$  to the power minus 3  $J_3 x dx$ .

So, we are doing finding the integral  $y$  integration by parts and this becomes minus  $x$  to the power minus 1  $J_3 x$  plus 2 times integral of  $x$  to the power minus 2  $J_3 x dx$  integral of  $x$  to the power minus 2  $J_3 x$  we know is equal to minus  $x$  to the power minus 2  $J_2 x$ . So, we have the right hand side as minus  $x$  to the power minus 1  $J_3 x$  plus 2 times minus  $x$  to the power minus 2  $J_2 x$  plus  $c$  or we may write it as minus  $x$  to the power minus 1  $J_3 x$  minus  $2x$  to the power minus 2  $J_2 x$  plus  $c$ , which is the integral of  $x$  to the power minus 1  $J_4 x dx$ .

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**SECOND SOLUTION OF BESSEL'S EQUATION :**

When  $\nu$  is an integer  $n$ , the second independent solution can be obtained by using the methods discussed earlier for finding the series solution, we illustrate it for the case  $n = 0$ , i.e. Bessel's equation of order zero.

(4)  $xy'' + y' + xy = 0$

Now, let us study the second solution of Bessel's equation, we had seen that when  $\nu$  is an integer, the second solution of the Bessel's equation has to be obtained by the methods, which we had discussed in the previous lecture. So, we will be illustrating how to find the second solution of the Bessel's equation for the case  $\nu$  equal to 0 for the other integral values of  $\nu$  the second solution of the Bessel's equation can be obtained similarly. Now, when  $\nu$  is an integer equal to 0 the Bessel's equation of order 0 will be  $xy'' + y' + xy = 0$ , you can see from the Bessel's equation, if you put  $\nu$  equal to 0 there, then it reduces to  $xy'' + y' + xy = 0$ .

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Substituting  $y(x) = \sum_{m=0}^{\infty} c_m x^{m+r}$

in (4), we find the roots of the indicial equation as  $r = 0, 0$ . The first solution is  $y_1(x) = J_0(x)$  while the second solution is of the form

$$y_2(x) = J_0(x) \ln x + \sum_{k=1}^{\infty} d_k x^k.$$

Substituting  $y_2$  and its derivatives in (4) we obtain

$$2J_0' + \sum_{k=1}^{\infty} k(k-1)d_k x^{k-1} + \sum_{k=1}^{\infty} k d_k x^{k-1} + \sum_{k=1}^{\infty} d_k x^{k+1} = 0.$$

So, let us put again  $x$  equal to 0 is a regular singular point here, so let us put  $y = x^r$  equal to  $\sum_{m=0}^{\infty} c_m x^{m+r}$  where  $c_0$  is not equal to 0 in this equation, we find that the roots of the indicial equation are 0 comma 0 that is it is the case of indicial equation having equal roots. Now, one solution of the indicial equation, one solution of the Bessel equation can be then obtain it will be  $y_1(x) = J_0(x)$ . We had seen earlier that corresponding to the indicial equation having roots  $\nu$  and  $\nu - 1$  linearly independent solution was of the Bessel's equation were  $J_\nu(x)$ .

So, in that you put  $\nu$  equal to 0 you get one solution of the Bessel's equation of order 0 here, as  $y_1(x) = J_0(x)$ , while the second solution as we know from the previous lecture on series solution can be expressed as  $y_2(x) = J_0(x) \ln x + \sum_{k=1}^{\infty} d_k x^k$ . So, substituting this second solution into the and its derivatives into the Bessel's equation of order 0, we obtain  $2J_0' + \sum_{k=1}^{\infty} k(k-1)d_k x^{k-1} + \sum_{k=1}^{\infty} k d_k x^{k-1} + \sum_{k=1}^{\infty} d_k x^{k+1} = 0$ .

Now, you can see in this equation the terms containing logarithm do not occur, this is because of the fact that  $J_0(x)$  is a solution of the Bessel's equations. So, all the terms which involve logarithm  $\ln x$  function they banish and therefore, when you substitute  $y_2(x)$  and its derivatives in the equation 4 it reduces to this form.

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$$\Rightarrow \sum_{k=1}^{\infty} \frac{(-1)^k x^{2k-1}}{2^{2k-2} k!(k-1)!} + \sum_{k=1}^{\infty} k^2 d_k x^{k-1} + \sum_{k=1}^{\infty} d_k x^{k+1} = 0.$$

Equating to zero, the coefficient of  $x^{2n}$   
 $(2n+1)^2 d_{2n+1} + d_{2n-1} = 0.$

The coefficient of  $x^0$  equated to zero  
 $\Rightarrow d_1 = 0.$   
Hence,  $d_3 = 0 = d_5 = d_7 = \dots$

Now, let us find the derivative of  $J_n \text{ naught } x$  we know the infinite series expansion of  $J_n \text{ naught } x$ . So, let us differentiate that infinite series with respect to  $x$  and hence find  $J_n \text{ naught dash } x$  and put that value here, we get the previous equation like this. So, it is an identity and therefore, the coefficients of various powers of  $x$  can be equated to 0 in order to find the values of the unknown coefficients  $d_k$ .

Now, you can see that in this first term only odd powers of  $x$  are occurring, the power of  $x$  is  $2k - 1$  where  $k$  is running from 1 to infinity. So, only odd powers of  $x$  are occurring here while here power of  $x$  begins with 0 where  $k$  equal to 1, so power of  $x$  begins with 0 and here the powers of  $x$  begin with 2. So, when we equate to 0, so what we will do, we will first equate to 0 the coefficients of even powers of  $x$ .

Now, when you put the coefficient of  $x$  to the power  $n$  to 0, the coefficient of  $x$  to the power  $n$  will not be found here because, it involves only odd powers of  $x$  terms. So, coefficient of  $x$  to the power  $n$  if we want to equate to 0, we will have to find that from this second and third terms. So, coefficient of  $x$  to the power  $2n$  here will be  $(2n+1)^2 d_{2n+1}$  where the coefficient of  $x$  to the power  $2n$  here will be  $d_{2n-1}$ , so we get the equation as  $(2n+1)^2 d_{2n+1} + d_{2n-1} = 0$ .

Now, next we equate to 0 the coefficient of  $x$  to the power 0 which is available only in this term. So, it gives us  $d_1 = 0$  or you can say  $d_1 = 0$ . Now,

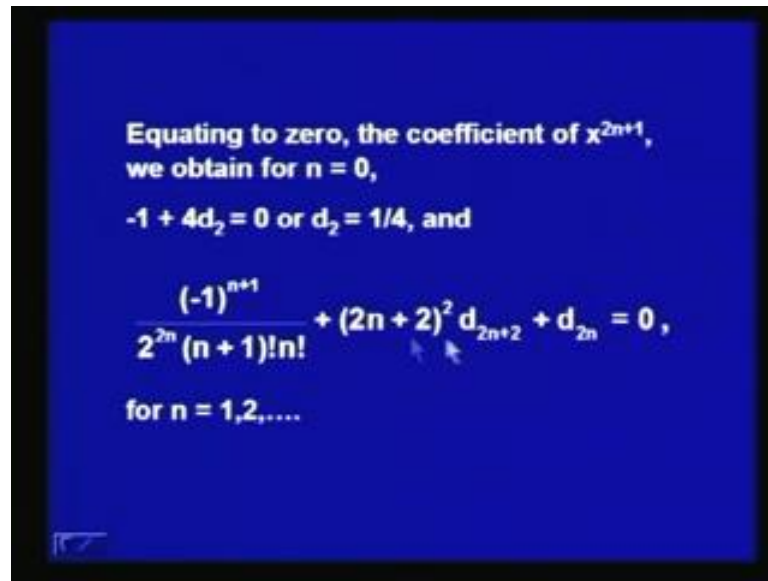
while  $d_1$  is equal to 0 from this equation, we find that  $d_3$  is 0,  $d_5$  is 0,  $d_7$  is 0 and so on. So, when  $d_1$  is 0,  $d_3$ ,  $d_5$ ,  $d_7$  all are 0's it follows from this equation, next we equate to 0 the coefficients of odd powers of  $x$ , we have already equated to 0 the coefficients of even powers of  $x$  and also the power of  $x$  as 0.

So, let us now equate to 0 the coefficients of  $x$  to the power  $2n + 1$  to 0, but we noticed that when  $n$  will be 0 that is if we 1 to equate to 0 the coefficient of  $x$ , then the coefficients of  $x$  are available in this term, this term, but it is not here. So, let us equate to 0 the coefficients of  $x^0$  separately and the coefficients of  $x$  to the power  $2n + 1$  where  $n$  will take values 1, 2, 3 and so on from as those terms will be available in all these three terms.

So, let us first equate to 0 the coefficient of  $x$  to 0, the coefficient of  $x$  if we want to equate to 0 put  $k$  equal to 1 here, we will get minus 1 over 2 to the power 0, then we have 1 factorial then we have 0 factorial. So, we will have minus 1 over 1 that is the coefficient of  $x$  here and the coefficient of  $x$  here will be  $k$  equal to 2 we put  $k$  equal to 2. So, we get  $4d_2$ , so  $1 + 4d_2$  equal to 0 that will be the coefficient of  $x$ , when we put equal to 0 and the coefficient of  $x$  to the power  $2n + 1$  when we want to find, we will have to put here  $k$  equal to  $n + 1$ , so that  $2^k - 1$  becomes  $2^{n+2} - 1$  that is  $2^{n+1}$ .

So, when you take  $k$  equal to  $n + 1$ , then we will have minus 1 to the power  $n + 1$  over 2 to the power  $2n + 2 - 2$ . So,  $2$  to the power  $2n$  over  $n + 1$  factorial and then  $n$  factorial and then here  $x$  to the power  $2n + 1$  the coefficient of that will be you put  $k$  equal to  $2n + 2$  factorial  $2n + 2$  whole square into  $d_{2n+2}$  and here it will be, if you want the coefficient of  $x$  to the power  $2n + 1$  you have to take  $k$  equal to  $2n$ . So, we will get the coefficient of  $x$  to the  $2n + 1$  as  $d_{2n}$ , so we will put them equal to 0.

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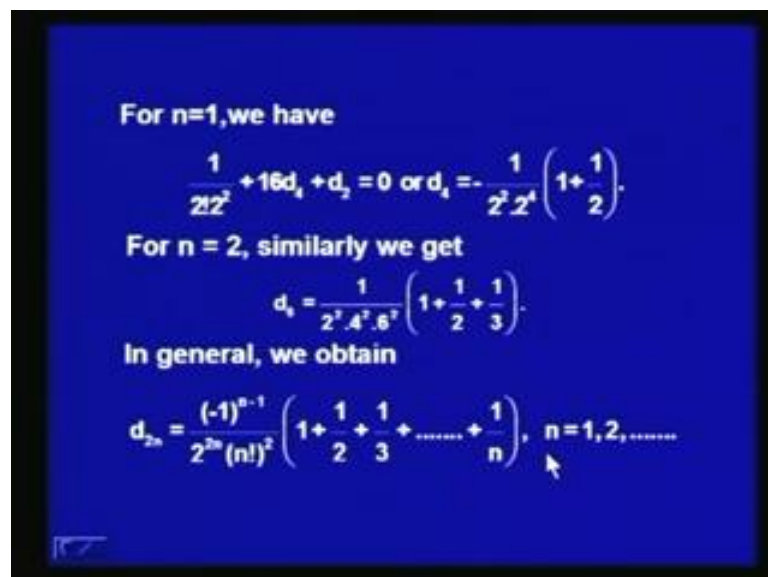
Equating to zero, the coefficient of  $x^{2n+1}$ ,  
we obtain for  $n = 0$ ,  
 $-1 + 4d_2 = 0$  or  $d_2 = 1/4$ , and

$$\frac{(-1)^{n+1}}{2^{2n} (n+1)! n!} + (2n+2)^2 d_{2n+2} + d_{2n} = 0,$$

for  $n = 1, 2, \dots$

So, when we put the coefficient of  $x$  to the power  $2n + 1$  equal to 0 for  $n$  equal to 0 as we have seen just now, we will get minus 1 plus 4  $d_2$  equal to 0, which will give us the value of  $d_2$  as 1 by 4 and the other coefficients of  $x$ , that is coefficients of  $x$  to the power  $2n + 1$  where  $n$  takes values going to on 1, 2, 3 and so on, we have just now seen that this is the equation that we will get.

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For  $n=1$ , we have

$$\frac{1}{2 \cdot 2^2} + 16d_4 + d_2 = 0 \text{ or } d_4 = -\frac{1}{2^2 \cdot 2^4} \left(1 + \frac{1}{2}\right).$$

For  $n = 2$ , similarly we get

$$d_6 = \frac{1}{2^3 \cdot 4^2 \cdot 6^2} \left(1 + \frac{1}{2} + \frac{1}{3}\right).$$

In general, we obtain

$$d_{2n} = \frac{(-1)^{n-1}}{2^{2n} (n!)^2} \left(1 + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{n}\right), \quad n=1, 2, \dots$$

So, from this equation for  $n$  equal to 1 we shall have 1 over 2 factorial into 2 square plus 16  $d_4$  plus  $d_2$  equal to 0 or we will have  $d_4$  equal to minus 1 over 2 square into 2 to the



power 4 into 1 plus half making use of the value of d 4 that is 1 by 4. And similarly for n equal to 2 we shall have d 6, d 6 will be equal to 1 over 2 square 4 square 6 square into 1 plus 1 by 2 plus 1 by 3.

In general we will have the value of d 2 n, d 2 n will be equal to minus 1 to the power n minus 1 2 to the power 2 n into n factorial square 1 plus 1 by 2 plus 1 by 3 and so on plus 1 by n, where n takes values 1, 2, 3, and so on. Thus we have obtain the values of all the unknown coefficients degage that occur in our second solution.

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Therefore,

$$y_2(x) = J_0(x) \ln x + \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{2^{2n} \cdot (n!)^2} \left( 1 + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{n} \right) x^{2n}$$

is the second linearly independent solution of (4). Hence  $J_0(x)$  and  $y_2(x)$  form a fundamental system of (4). Another linearly independent solution taken is

$$y_2^*(x) = \frac{2}{x} [y_2(x) + (\gamma - \ln 2) J_0(x)]$$

And let us put the values of these unknowns in the second solution we find that  $y_2 x$  now becomes  $J_0(x) \ln x + \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{2^{2n} \cdot (n!)^2} \left( 1 + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{n} \right) x^{2n}$ . Now, these two solutions are linearly independent of each other because,  $y_2 x$  involves the logarithmic function. So, they form a fundamental system of the Bessel's equation of order 0.

Now, another linearly independent solution that is usually taken is given by  $y_2^* x = \frac{2}{x} [y_2(x) + (\gamma - \ln 2) J_0(x)]$ . Now, you can see that this solution is independent of the first solution that is  $J_0(x)$ , so it could also be taken as the second linearly independent solution of the Bessel's equation. This is a linear combination of the functions  $J_0(x)$  and  $y_2 x$ , which we have seen or solutions of the Bessel's equation already.

So, it is also a solution of the Bessel's equation and clearly it is independent of  $J_0(x)$ , so it may also be taken as the second linearly independent solution of Bessel's equation.

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where  $\gamma$  ( $\gamma = 0.577215\dots$ ) is the Euler constant. This solution is called the Bessel's function of second kind of order zero and is denoted by  $Y_0(x)$ . It is also named as Neumann's function of second kind. Hence

$$Y_0(x) = \frac{2}{\pi} \left[ \left( \ln \frac{x}{2} + \gamma \right) J_0(x) + \sum_{k=1}^{\infty} \frac{(-1)^{k-1} h_k x^{2k}}{2^{2k} (k!)^2} \right],$$

where

$$h_k = 1 + \frac{1}{2} + \dots + \frac{1}{k}, \quad k = 1, 2, \dots$$

Here, gamma is equal to 0.577215 and so on it is called the Euler's constant, this solution is called the Bessel's function of second kind of order 0 and is denoted by  $Y_0(x)$ . We also call it Neumann's function of second kind, so with this as the second solution  $Y_0(x)$  will then assume the form  $\frac{2}{\pi} \left[ \ln \frac{x}{2} + \gamma \right] J_0(x) + \sum_{k=1}^{\infty} \frac{(-1)^{k-1} h_k x^{2k}}{2^{2k} (k!)^2}$  and as we have seen our  $y = x$  becomes  $Y_0(x)$ . So, here this  $h_k$  denotes  $1 + \frac{1}{2} + \dots + \frac{1}{k}$  where  $k$  takes values 1, 2, 3, and so on.

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It follows that  $Y_0(x) \rightarrow -\infty$  as  $x \rightarrow 0$ . For small values of  $x$ ,  $Y_0(x)$  behaves like  $\ln(x)$ .

The solution for other integral values of  $n$  can be obtained similarly.

To ensure uniformity for the two cases, viz. where  $\nu$  is not an integer and  $\nu$  is an integer, the second solution is defined by the formula

(5)  $Y_\nu(x) = \frac{1}{\sin \nu\pi} [J_\nu(x)\cos \nu\pi - J_{-\nu}(x)]$ ,  $\nu \neq \text{integer}$

Now, from this expression  $Y_\nu(x)$  you can see that  $Y_\nu(x)$  goes to minus infinity as  $x$  goes to 0, we know that  $\ln x$  goes to minus infinity as  $x$  goes to 0. So,  $Y_\nu(x)$  goes to minus infinity as  $x$  goes to 0 and for small values of  $x$   $Y_\nu(x)$  we have like the logarithmic function  $\ln x$ . So, this how we find the second solution of the Bessel's equation of order 0, the second solution of the Bessel's equation for other integral values of  $n$  can be obtained in a similar manner.

Now, in practise when we make use of the Bessel's equation the two solutions of the Bessel's equation are given by this for uniformity in the two cases, that is when  $\nu$  is not an integer and when  $\nu$  is an integer. The second solution is given by the following formula  $Y_\nu(x) = \frac{1}{\sin \nu\pi} [J_\nu(x)\cos \nu\pi - J_{-\nu}(x)]$ , when  $\nu$  is not an integer, you see we have seen earlier that when  $\nu$  is not an integer the two solutions of the Bessel's equation  $J_\nu(x)$  and  $J_{-\nu}(x)$  are linearly independent.

So, when  $\nu$  is not an integer this  $Y_\nu(x)$  being a linear combination of  $J_\nu(x)$  and  $J_{-\nu}(x)$  is also a solution of the Bessel's equation and therefore, may be taken as the second solution of the Bessel's equation.

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and  $Y_\nu(x) = \lim_{k \rightarrow \nu} Y_k(x), \nu = \text{integer}.$

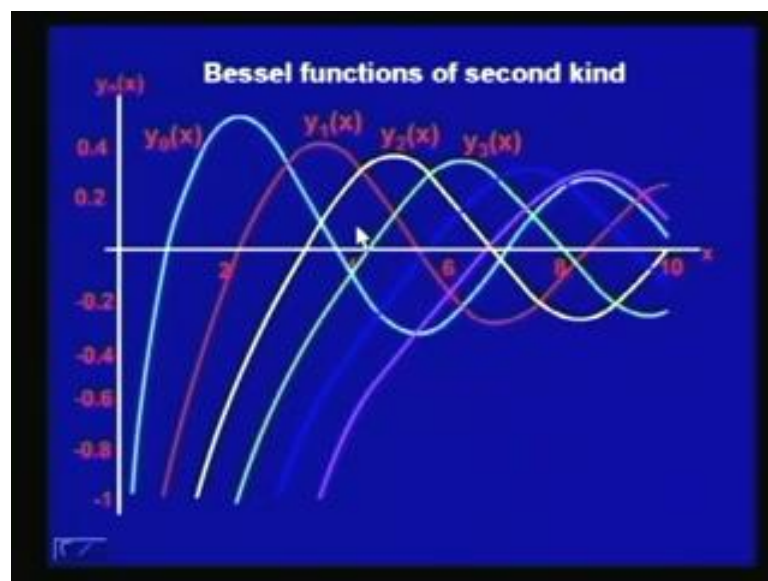
For non – integral order,  $Y_\nu(x)$  is a solution of Bessel's equation being a linear combination of  $J_\nu(x)$  and  $J_{-\nu}(x)$ . When  $\nu$  is an integer, by L' Hospital's rule  $\lim_{k \rightarrow \nu} Y_k(x)$  exists.

Hence, the general solution for all  $\nu$ , is

$$Y(x) = AJ_\nu(x) + BY_\nu(x).$$

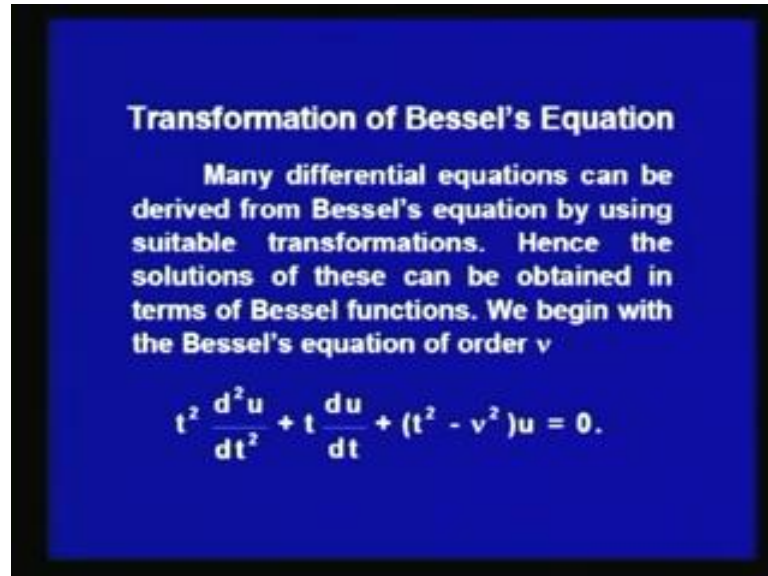
When  $\nu$  is an integer then the second solution of the Bessel's equation  $Y_\nu(x)$  is the limit of  $Y_k(x)$  as  $k$  goes to  $\nu$ , this  $Y_k(x)$  is given by the expression for  $Y_\nu(x)$  where  $\nu$  is not an integer. Now, for non-integral order  $Y_\nu(x)$  is a solution of Bessel's equation being a linear combination of  $J_\nu(x)$  and  $J_{-\nu}(x)$ , when  $\nu$  is an integer we find that by L Hospital's rule, this limit exists the limit of  $Y_k(x)$  as  $k$  goes to  $\nu$  it exist. And therefore, the general solution for all  $\nu$  of the Bessel's equation is given by  $Y(x)$  equal to  $A$  into  $J_\nu(x)$  plus  $B$  into  $Y_\nu(x)$ .

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This picture shows us the graphs of Bessel functions of second kind, this is the graph of  $y_0(x)$ , this is the graph of  $y_1(x)$ , then we have the graph of  $y_2(x)$ ,  $y_3(x)$  and so on.

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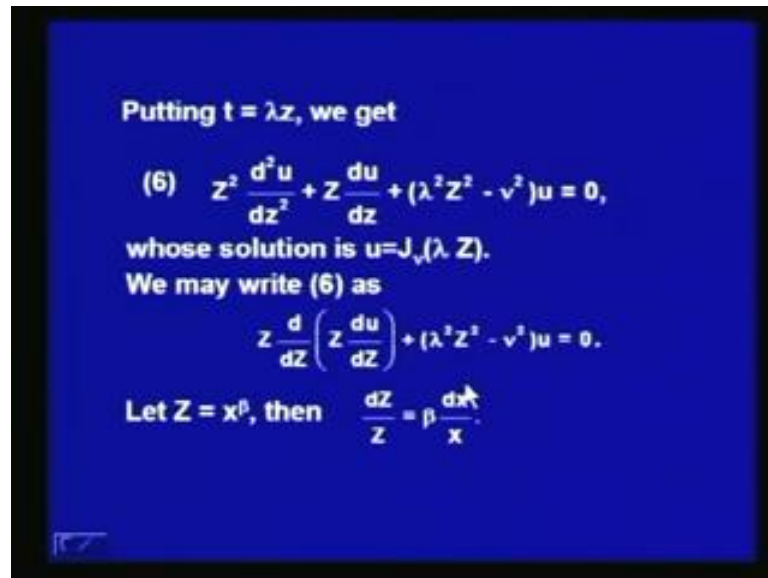
**Transformation of Bessel's Equation**

Many differential equations can be derived from Bessel's equation by using suitable transformations. Hence the solutions of these can be obtained in terms of Bessel functions. We begin with the Bessel's equation of order  $\nu$

$$t^2 \frac{d^2 u}{dt^2} + t \frac{du}{dt} + (t^2 - \nu^2)u = 0.$$

Now, many differential equations can be derived from the Bessel's equations by making use of some transformations. And the solutions of such equations can therefore, be obtained in terms of Bessel functions, we shall begin with the Bessel equation of order  $\nu$  and arrive at a fairly general linear differential equation of second order whose solution can be expressed in terms of Bessel functions by comparing that differential equation, we will be able to find the solutions of many differential equations that have solution in terms of Bessel functions. So, let us consider the Bessel's equation of order  $\nu$   $t^2 \frac{d^2 u}{dt^2} + t \frac{du}{dt} + (t^2 - \nu^2)u = 0$ .

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Putting  $t = \lambda z$ , we get

$$(6) \quad z^2 \frac{d^2 u}{dz^2} + z \frac{du}{dz} + (\lambda^2 z^2 - \nu^2)u = 0,$$

whose solution is  $u = J_\nu(\lambda Z)$ .

We may write (6) as

$$z \frac{d}{dz} \left( z \frac{du}{dz} \right) + (\lambda^2 z^2 - \nu^2)u = 0.$$

Let  $Z = x^\beta$ , then  $\frac{dZ}{Z} = \beta \frac{dx}{x}$ .

The first substitution that we make in this Bessel's equation of order  $\nu$  is  $t$  equal to  $\lambda z$  that is we change from the independent variable  $t$  to the independent variable  $z$   $\lambda$  is a constant here. So, when we put  $t$  equal to  $\lambda z$  in the Bessel's equation of order  $\nu$ , we get the equation  $Z^2 \frac{d^2 u}{dZ^2} + Z \frac{du}{dZ} + (\lambda^2 Z^2 - \nu^2)u = 0$ , now we know that a solution of the Bessel's equation of order  $\nu$  is  $J_\nu t$ .

So, and we have put  $t$  equal to  $\lambda Z$ , so a solution of this equation 6, this transformed equation is therefore,  $J_\nu \lambda Z$  we can we may also note that the first two terms of this equation 6 can be expressed as  $Z \frac{d}{dZ} \left( Z \frac{du}{dZ} \right)$ . Because, when you differentiate  $Z \frac{du}{dZ}$  with respect to  $Z$  what you get is  $\frac{du}{dZ} + Z \frac{d^2 u}{dZ^2}$  and when you multiply  $Z$  to that you get  $Z^2 \frac{d^2 u}{dZ^2} + Z \frac{du}{dZ}$ .

So, equation 6 can be written in an alternate form as  $Z \frac{d}{dZ} \left( Z \frac{du}{dZ} \right) + (\lambda^2 Z^2 - \nu^2)u = 0$ . Let us now make another substitution to this differential equation of second order, let us put  $Z$  equal to  $x$  to the power  $\beta$ , where  $\beta$  is a constant. When we take the logarithm on both sides we get  $\log Z$  equal to  $\beta \log x$  and when we differentiate on both sides what we get is  $\frac{dZ}{Z}$  equal to  $\beta \frac{dx}{x}$ . So, making this substitution  $Z$  equal to  $x$  to the power  $\beta$  and hence  $\frac{dZ}{Z}$  equal to  $\beta \frac{dx}{x}$  in this equation.

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Hence  $\frac{x}{\beta} \frac{d}{dx} \left( \frac{x}{\beta} \frac{du}{dx} \right) + (\lambda^2 x^{2\beta} - \nu^2) u = 0.$

Next, taking  $u = x^{-\alpha} y$  one has

(7)  $x \frac{d}{dx} \left\{ x \left( x^{-\alpha} \frac{dy}{dx} - \alpha x^{-\alpha-1} y \right) \right\} + \beta^2 (\lambda^2 x^{2\beta} - \nu^2) x^{-\alpha} y = 0$

since  $\frac{du}{dx} = -\alpha x^{-\alpha-1} y + x^{-\alpha} \frac{dy}{dx}.$

On simplification, (7) yields us a fairly general differential equation

We will get the transformed equation as  $x$  over  $\beta$  into  $d$  over  $d x$  of  $x$  over  $\beta$   $d u$  by  $d x$  plus  $\lambda$  square  $x$  to the power  $2 \beta$  minus  $\nu$  square into equal to  $0$ . Let us, now put  $u$  equal to  $x$  to the power minus  $\alpha$  into  $y$  in this differential equation, so now, at the these two transformations that we have done to the Bessel's equation, there we change the independent variables first from  $t$  to  $z$  and then from  $z$  to  $x$ , now we are changing the dependent variable from  $u$  to  $y$ .

So, let us put  $u$  equal to  $x$  to the power minus  $\alpha$  into  $y$ , in this differential equation what we shall have  $x \frac{d}{dx}$  in this can be combined and you get  $\beta$  square. So,  $\beta$  square be are multiplying to this whole equation, this  $\beta$  square is coming up here and so we get by  $x$  into  $d$  over  $d x$  then  $x$  over  $\beta$   $d u$  by  $d x$  changes into  $x$  times  $d u$  by  $d x$ . So,  $d$  over  $d x$  of  $u$  will be  $x$  to the power minus  $\alpha$   $d y$  by  $d x$  minus  $\alpha$  times  $x$  to the power minus  $\alpha$  minus  $1$  into  $y$  plus  $\beta$  square into  $\lambda$  square  $x$  to the power  $2 \beta$  minus  $\nu$  square into  $x$  to the power minus  $\alpha$  into  $y$  equal to  $0$ . There we have written the derivative of  $u$  like this, so the value of  $d$  over  $d x$  we have put here, now on simplification this equation 7 yields us a fairly general differential equation of second order.

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(8)  $\frac{d^2 y}{dx^2} + \frac{(1-2\alpha)}{x} \frac{dy}{dx} + \left[ (\lambda \beta x^{\beta-1})^2 + \frac{\alpha^2 - \nu^2 \beta^2}{x^2} \right] y = 0.$

A solution of this equation is  $x^{-\alpha} y = J_\nu(x^\beta \lambda)$   
or  $y = x^\alpha J_\nu(\lambda x^\beta).$

By comparing a given differential equation with (8), we can find its solution in the form  $y = x^\alpha J_\nu(\lambda x^\beta).$

It may be written as  $\frac{d^2 y}{dx^2} + \frac{1-2\alpha}{x} \frac{dy}{dx} + \left[ (\lambda \beta x^{\beta-1})^2 + \frac{\alpha^2 - \nu^2 \beta^2}{x^2} \right] y = 0$  or it may also be written as if you multiply by  $x^2$  in the whole of this equation you get  $x^2 \frac{d^2 y}{dx^2} + (1-2\alpha)x \frac{dy}{dx} + (\lambda^2 \beta^2 x^{2\beta} + \alpha^2 - \nu^2 \beta^2) y = 0$ .

So, it may also in that form, now a solution of this equation will there for be we had earlier  $u$  equal to  $J_\nu \lambda x^\beta$ . But, then we adds up change the independent variable from  $u$  to  $y$  by making this substitution  $u$  equal  $x^\beta$  into  $y$ . So, a solution of this equation is therefore,  $x^\alpha J_\nu \lambda x^\beta$  by multiplying by  $x^\alpha$  both sides, we get a solution of this equation as  $y$  equal to  $x^\alpha J_\nu \lambda x^\beta$ .

So, whenever we are given a differential equation whose solution is to be expressed in the form of the Bessel's functions, we can compare that differential equation with this fairly general differential equation of second order and find the values of the constants  $\alpha$ ,  $\beta$ ,  $\lambda$  and  $\nu$ . And we can then find the solution of the Bessel's the given differential equation in terms of the Bessel functions, so now, let us take an example based this article.



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**Example:**

$$(9) \quad y'' + \left(1 + \frac{1 - 4\nu^2}{4x^2}\right)y = 0.$$

Comparing the given ODE with (8), we find that  $\alpha = 1/2$ ,  $\beta = 1$  and  $\lambda = 1$ .

Therefore, a solution of (9) is

$$y_1(x) = x^{1/2}J_\nu(x).$$

Let us consider the differential equation  $y'' + \left(1 + \frac{1 - 4\nu^2}{4x^2}\right)y = 0$ , when we compare this differential equation with the transformed Bessel's equation that is the equation number 8, we find that  $\alpha$  is equal to half,  $\beta$  is equal to 1 and  $\lambda$  is equal to 1. And therefore, a solution of this equation 9, may be expressed as  $y_1(x) = x^{1/2}J_\nu(x)$ , where  $\nu$  can take integral or non-integral values here.

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If  $\nu$  is not an integer, then another independent solution of (9) is

$$y_2(x) = x^{1/2}J_{-\nu}(x).$$

Hence  $y = x^{1/2}[AJ_\nu(x) + BJ_{-\nu}(x)]$  is the general solution of (9).

When  $\nu$  is an integer, the general solution is given by  $y = x^{1/2}[AJ_\nu(x) + BY_\nu(x)]$ .

So, if the differential equations such that  $\nu$  is not an integer, then another independent solution of the equation 9 can be written as  $y = x^{\frac{1}{2}} J_{-\nu}(x)$ . And hence we may write the general solution of equation 9 as  $y = x^{\frac{1}{2}} [A J_{\nu}(x) + B J_{-\nu}(x)]$ , now if  $\nu$  is an integer then the second solution of the Bessel's equation will have to be obtained as we have discussed earlier for  $\nu = 0$  and for other integral values of  $\nu$  we can similarly find the solution.

So, the second solution where the equation 9 will have to be obtained using those methods and so then  $y$  will be equal to  $x^{\frac{1}{2}} [A J_{\nu}(x) + B Y_{\nu}(x)]$ , we have represented the second solution of the Bessel's equation by  $Y_{\nu}(x)$ . This lecture today we have discussed the solutions of the Bessel's equation, which are called as Bessel function of first kind that is  $J_{\nu}(x)$  and the Bessel function of the second kind that is  $Y_{\nu}(x)$ .

And also seen how we can find the solution in terms of Bessel function of a fairly general differential equation, which we obtained by making some transformations to the Bessel's equation of order  $\nu$ . In our next lecture, we shall be discussing the orthogonality of Bessel functions, then how to expand a function, which is continuous over a say an interval  $0 < x < a$  and has oscillations in that interval in terms of Bessel functions, we shall discuss the generating function. And we shall also discuss the representation of Bessel function of first kind in terms of an integral, which we call as the Bessel integral. So, in that the Bessel function of order  $\nu$  will consider where  $\nu$  will be an integer, so all these we will be discussing in our next lecture on Bessel's functions and their properties.