

**Engineering Physics 1**  
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**Module-01**  
**Lecture-04**  
**Interference of Polarized Light**

This is the fourth lecture of these five lecture series on polarization. In the last lecture which was the third one we considered phenomena of double refraction, looked the dichroism and then studied the working of a Nicol prism. So, double refraction is a very typical phenomena of an isotropic crystals. We considered crystals of calcite and quartz.

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In the last lecture we considered phenomenon of double refraction .Looked at dichroism and studied the working of a Nicol prism.

In the present lecture we shall consider interference of polarized light. We shall analyze working of quarter- and half-wave plates of doubly refracting crystals. We shall consider different types of polarizations and their analysis.

What happens in these when a light beam is incident, there are two refracted rays one the so called ordinary; ordinary in the sense that, it obeys Snell's laws of refraction, the usual one. The speed is same in all directions, the phase front is a spherical that it ends the refracted beam called the extraordinary beam where the direction of propagation is different it moves with different speeds in different directions.

The result is that the phase fronts are ellipsoid of revolution about the optic axis. Optic axis is the axis of symmetry; the two beams ordinary and extraordinary, both of them are plane polarized with vibrations in perpendicular directions. In the case of ordinary, the direction of vibration is

perpendicular to the plane containing the propagation direction and the optic axis. In the case of extraordinary beam, these vibrations are in the same plane. Both of them these beams are plane polarized.

We also studied dichroism you see dichoric core crystals like tourmaline. They are also doubly refracting with the additional property that they have a very markedly unequal absorption of the ordinary and extraordinary beams, with the result that the merchant beam is plane polarized to a pretty good degree. Then, we considered Nicole prism. Here the ordinary beam is totally reflected things are arranged in such a way and the result is, the emergent beam is plane polarized and we have got a polarizing device.

In the present lecture which is the fourth one, we shall consider interference of polarized light. We shall analyze working of quarter and half wave plates of doubly refracting crystals. We shall consider different types of polarization and their analysis, Interference of polarized light.

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## **VI. Interference of polarized light**

We have already studied how a plane wave, incident on a doubly refracting crystal, splits up into two waves each characterized by a certain state of polarization. Their directions of polarization can be ascertained experimentally. The result of their superposition can also be analyzed experimentally.

We have already studied how a plane wave incident on a doubly refracting crystal, splits up into two waves within the crystal, each characterized by a resultant state of polarization. Their directions of polarization can be ascertained experimentally. The result of their superposition can also be analyzed experimentally.

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Now we will consider normal incidence of a plane-polarized beam on a calcite crystal which is cut such that optic axis is parallel to the surface of the crystal.

We will study the state of polarization of the beam emerging from the crystal.

Now, we will consider normal incidence of a plane polarized beam on a calcite crystal which is cut such that the optic axis is parallel to the surface of the crystal. The propagation here is perpendicular to the optic axis.

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For any other state of polarization of the incident beam, **both** the extra-ordinary and the ordinary components will be present.

For a negative crystal like calcite,  
 $n_e < n_o$  and the *e*-wave will travel faster than the *o*-wave.

We will study the state of polarization of the beam emerging from the crystal.

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Let us take the propagation direction as the  $x$ -axis.

We will assume the  $z$ -axis to be along the optic axis.

Now we know that, if the incident beam is  $y$ -polarized, the beam will propagate as an ordinary wave and the extra-ordinary wave will be absent. Similarly if the incident beam is  $z$ -polarized, the beam will propagate as an extra-ordinary wave and the ordinary wave will be absent.

Let us make the propagation direction as the  $x$ -axis, just choose it like that. We will assume the  $z$  axis to be along the optic axis which is the axis of symmetry in the crystal. Now, we know that if the incident beam is  $y$  polarized, the beam will propagate as an ordinary wave and the extraordinary wave will be absent. In this case the vibrations are in that explained. Similarly, if the beam is  $z$  polarized the beam will propagate as an extraordinary wave and the ordinary wave in this case will be absent.

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Let  $E_0$  be the amplitude of the electric vector associated with the incident beam and let it make an angle  $\phi$  with the  $z$ -axis (which is the optic axis).

Such a beam can be assumed to be a superposition of two linearly polarized beams (vibrating in phase at the entry point of the crystal), polarized along the  $y$ - and  $z$ -directions with amplitudes  $E_0 \sin \phi$  and  $E_0 \cos \phi$  respectively.

Let  $E$  be the amplitude of electric vector associated with the incident beam and let it make an angle  $\phi$  with  $z$  axis which is the optic axis, as we have chosen. Such a beam can be assumed to be the superposition of two linearly polarized beams vibrating in phase at the entry point of the



crystal and polarized along the y and z directions with amplitudes  $E \sin \Phi$  and  $E \cos \Phi$  respectively.

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The z-component, whose amplitude is  $E_0 \cos \phi$ , passes through as an extraordinary beam propagating with wave velocity  $c/n_e$ . Here the vibrations are parallel to the optic axis.

The y-component, whose amplitude is  $E_0 \sin \phi$ , passes through as an ordinary beam propagating with wave velocity  $c/n_o$ . Here the vibrations are perpendicular to the optic axis.

The z component whose amplitude is  $E \cos \Phi$  passes through as an extraordinary beam propagating with the wave velocity  $c/n_e$ , here the vibrations are parallel to the optic axis. The y component whose amplitude  $E \sin \Phi$ , passes through as an ordinary beam propagating with the wave velocity  $c/n_o$ . Here the vibrations are perpendicular to the optic axis, also perpendicular to the propagation direction.

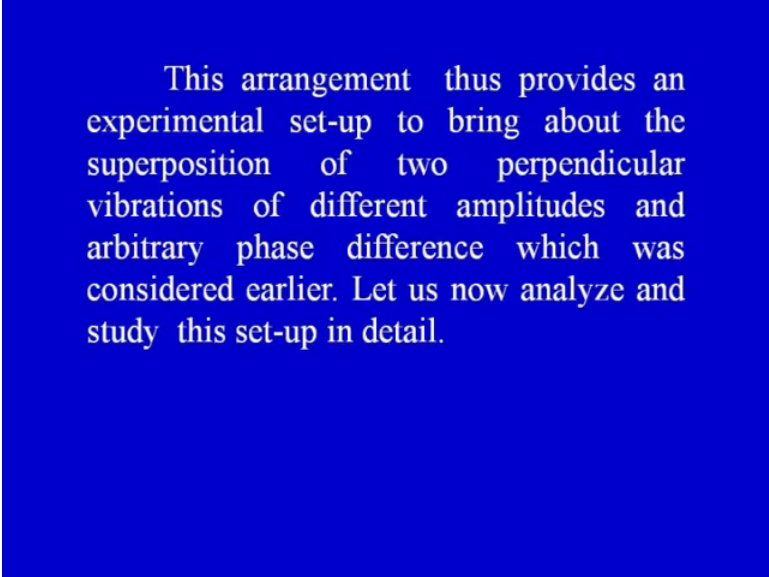
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Since  $n_e \neq n_o$ , the two beams will propagate with different velocities and, as such, when they come out of the crystal, they will not be in phase.

Consequently, the emergent beam, which will be a superposition of these two beams, will be, in general, elliptically polarized.

Since  $n_e$  is not  $= n_o$ , the two beams will propagate with different velocities but there would not be any separation in this case because they are incident normally and as such when they come out of the crystal naturally, they will not be in phase. Consequently, the emergent beam which will be a superposition of these two beams will be in general elliptically polarized. We have discussed this case in the earlier lecture as well. The arrangement thus provides an experimental setup and that is the main thing we are discussing the experimental setup.

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This arrangement thus provides an experimental set-up to bring about the superposition of two perpendicular vibrations of different amplitudes and arbitrary phase difference which was considered earlier. Let us now analyze and study this set-up in detail.

Now to bring about the superposition of two perpendicular vibrations of different amplitudes and arbitrary phase difference which was considered earlier, as I pointed out? You see, the different amplitudes can be obtained by varying the angle  $\Phi$  and the phase difference can be varied by varying the thickness of the crystal. Let us now analyze and study this setup in detail.

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Let the plane  $x = 0$  represent the surface of the crystal on which the beam is incident. The  $y$ - and  $z$ -components of the incident beam can be written in the form

$$E_y = E_0 \sin \phi \cos(kx - \omega t)$$

$$E_z = E_0 \cos \phi \cos(kx - \omega t),$$

where  $k (= \omega/c)$  represents the free space wave number.

Let the plane  $x = 0$  represent the surface of the crystal on which the beam is incident.  $x = 0$  plane is the incident. The  $y$  and  $z$  components of the incident beam now can be written in the form  $E_y = E_0 \sin \phi \cos(kx - \omega t)$ . That is the amplitude  $\cos kx - \omega t$   $E_z = E_0 \cos \phi \cos(kx - \omega t)$  and  $\cos$  of  $kx - \omega t$ . The propagation is along the  $x$  axis  $k = \omega/c$ . It represents the free space wave number.

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Thus, at  $x = 0$ , we have

$$E_y(x = 0) = E_0 \sin \phi \cos \omega t$$

$$E_z(x = 0) = E_0 \cos \phi \cos \omega t$$

At  $x = 0$  which was the entry point, we have  $E_y$  at  $x = 0 = E_0 \sin \phi \cos \omega t$  and  $E_z$  at  $x = 0 = E_0 \cos \phi \cos \omega t$ .

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Inside the crystal, the two components will be given by

$$E_y = E_0 \sin \phi \cos(n_o kx - \omega t), \text{ordinary wave}$$

$$E_z = E_0 \cos \phi \cos(n_e kx - \omega t), \text{extra-ordinary wave}$$

Inside the crystal, the two components will be given by  $E_y = E_0 \sin \phi \cos n_o kx - \omega t$  is for the ordinary,  $n_o$  is the refractive index for the ordinary. And  $E_z = E_0 \cos \phi \cos n_e kx - \omega t$ ,  $n_e$  the refractive index for the extraordinary. So, this is for the extraordinary wave.

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If the thickness of the crystal is  $d$ , then at the emerging surface, we will have

$$E_y = E_0 \sin \phi \cos(\omega t - \theta_o)$$

$$E_z = E_0 \cos \phi \cos(\omega t - \theta_e),$$

where  $\theta_o = n_o k d$  and  $\theta_e = n_e k d$ .

If the thickness of the crystal is  $d$  then, after traversing a distance  $= d$ , they were at the emerging surface, we will have,  $E_y = E_0 \sin \phi \cos \omega t - \theta_o$ .  $E_z = E_0 \cos \phi \cos \omega t - \theta_e$  where  $\theta_o = n_o k d$  and  $\theta_e = n_e k d$ .

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By appropriately choosing the instant  $t = 0$ , these components may be written as

$$E_y = E_0 \sin \phi \cos(\omega t - \theta)$$

$$E_z = E_0 \cos \phi \cos \omega t,$$

By appropriately choosing the instant  $t = 0$ , these components may also be written as in a simpler way  $E_y = E_0 \sin \phi \cos \omega t - \theta$  and  $E_z = E_0 \cos \phi \cos \omega t$ .

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where

$$\begin{aligned} \theta = \theta_o - \theta_e &= kd(n_o - n_e) = \frac{\omega}{c}(n_o - n_e)d \\ &= \frac{2\pi}{\lambda_0}(n_o - n_e)d \end{aligned}$$

represents the phase difference between the ordinary and the extra-ordinary beams at emergence.  $\lambda_0$  is the free space wavelength of the light.

Where  $\theta$  is the difference between  $\theta_o$  and  $\theta_e$  which is  $= kd$  times  $n_o - n_e$  which  $= \frac{\omega}{c}(n_o - n_e)d$  which  $= \frac{2\pi}{\lambda_0}(n_o - n_e)d$ , is represented the phase difference at the emergence, between the ordinary and the extraordinary beams.  $\lambda_0$  is the free space wavelength of the light used.

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The resultant of these two components can be obtained by eliminating  $t$  between them,

$$\frac{E_y}{E_0 \sin \phi} = \frac{E_z}{E_0 \cos \phi} \cos \theta + \sqrt{1 - \left( \frac{E_z}{E_0 \cos \phi} \right)^2} \sin \theta$$

or

$$\frac{E_y}{E_0 \sin \phi} - \frac{E_z}{E_0 \cos \phi} \cos \theta = \sqrt{1 - \left( \frac{E_z}{E_0 \cos \phi} \right)^2} \sin \theta$$

The resultant of these two components can be obtained by eliminating  $t$  between them and that gives  $E_y$  upon  $E_0 \sin \phi$ ,  $= E_z$  upon  $E_0 \cos \phi$  times  $\cos \theta$  + the square root of  $1 - E_z$  upon  $E_0 \cos \phi$  squared times  $\sin \theta$ . We can rearrange that and that gives  $E_y$  upon  $E_0 \sin \phi$  -  $E_z$  upon  $E_0 \cos \phi$  times  $\cos \theta$  which is = the square root of  $1 - E_z$  upon  $E_0 \cos \phi$  squared times  $\sin \theta$

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or

$$\frac{E_y^2}{(E_0 \sin \phi)^2} + \frac{E_z^2}{(E_0 \cos \phi)^2} - 2 \frac{E_y E_z}{E_0^2 \sin \phi \cos \phi} \cos \theta = \sin^2 \theta$$

This is the equation of an ellipse which represents the general state of polarization resulting from the superposition of two perpendicularly polarized plane-polarized beams traveling in the same direction with unequal amplitudes  $E_0 \sin \phi$  and  $E_0 \cos \phi$ , and arbitrary phase difference  $\theta$ .

or  $E_y^2$  upon  $E_0^2 \sin^2 \phi$  +  $E_z^2$  upon  $E_0^2 \cos^2 \phi$  - twice of  $E_y E_z$  upon  $E_0^2 \sin \phi \cos \phi$  times  $\cos \theta$  =  $\sin^2 \theta$ . This is the equation of an ellipse which represents the general state of polarization resulting from the superposition of two



perpendicularly polarized, plane polarized beams, traveling in the same direction, with an equal amplitudes  $E \sin \Phi$  and  $E \cos \Phi$  and an arbitrary phase difference  $\theta$ .

Various cases corresponding to different values of  $\theta$  were in a general way discussed in our earlier lecture. We shall now, in this lecture consider in detail, especial situations where  $\theta = 2\pi$  or  $\pi$  or  $\pi/2$  and  $\Phi = 45^\circ$  and how to obtain them how to realize them.

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We shall now consider special situations when  $\theta = 2\pi, \pi$  or  $\pi/2$  and  $\phi = 45^\circ$  and how to obtain them.

In the lab performing the experiment that is the main idea of today's lecture, let us consider these full half and quarter wave plates.

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### **VI.1 Full- half- and quarter-wave plates**

If the thickness of the crystal is such that  $\theta = 2\pi, 4\pi, 6\pi, \dots$ , the emergent beam will have the same state of polarization as the incident beam. Such a crystal plate is called a full-wave plate. Such a plate is, however, not very useful.

Remember, we are talking about a crystal plate where the optic axis is in the plane of the crystal. And the propagation is perpendicular to the optic axis and the incidence is normal under those conditions. If the thickness of the crystal is such, the  $\theta = 2\pi, 4\pi, 6\pi$ , I mean one complete vibration, 2 complete vibration, 3 complete vibrations, the emergent beam will have the same state of polarization as the incident beam. Such a crystal plate is called a full wave plate. Such a plate however is not very useful.

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If the thickness  $d$  of the crystal is such that  $\theta = \pi/2$ , the crystal plate is called a quarter-wave plate (usually abbreviated as QWP). A phase difference of  $\pi/2$  implies a path difference of a quarter of a wavelength between the ordinary and the extraordinary waves.

If the thickness  $d$  of the crystal is such that  $\theta = \pi/2$ , a quarter vibration, the crystal plate is called a quarter wave plate. Usually abbreviated as QWP; Phase difference of  $\pi/2$  implies a path difference of a quarter of wavelength that is  $\lambda/4$ , between the ordinary and the extraordinary waves.

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If the thickness of the crystal is such that  $\theta = \pi$ , the crystal plate is called a half-wave plate (usually abbreviated as HWP). This plate corresponds to a path difference of half a wavelength between the ordinary and the extra-ordinary waves.

Similarly, if the thickness of the crystal is such the  $\theta = \pi/2$ , the crystal plate is called a half wave plate, usually abbreviated as HWP. This plate corresponds to a path difference of half a wavelength that is  $\lambda/2$ , between the ordinary and the extraordinary waves. Let us consider in detail the working of a quarter wave plate.

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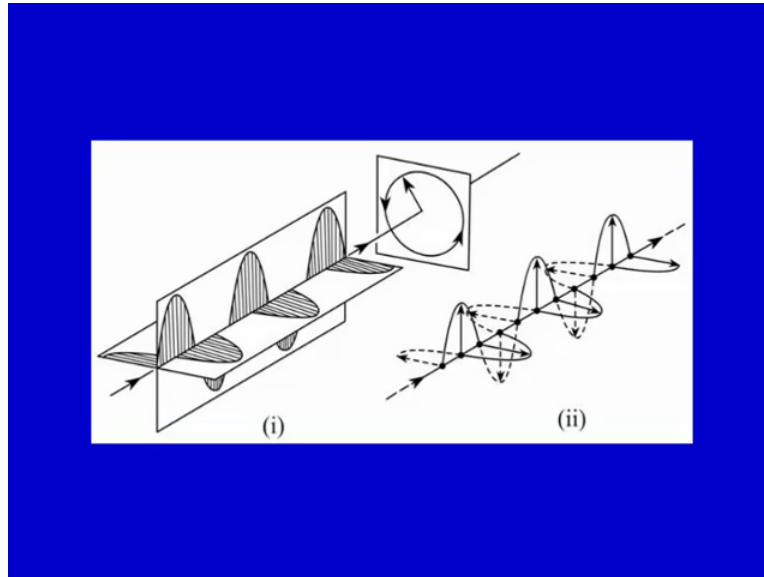
## VI.2 Quarter-wave plate

Let us consider the case when the phase difference  $\theta = \pi/2$  and the angle  $\phi$  which the electric vector (of amplitude  $E_0$ ) of the incident wave makes with the z-axis (which is the optic axis) is equal to  $\pi/4$ . Now the y- and z-components of the incident wave have equal amplitudes ( $E_0 \cos(\pi/4) = E_0/\sqrt{2}$ ) and the crystal introduces a phase difference of  $\pi/4$  during the passage.

Let us consider the case when  $\theta = \pi/2$  and the angle  $\phi$  with the electric vector of amplitude  $E_0$  of the incident wave, makes with the z axis, which is the optic axis. Remember, the optic axis is the symmetry axis and the propagation here is perpendicular to the optic axis. So, the angle  $\phi = \pi/4$ . Now the y and z components of the incident wave have equal amplitudes in hand course of  $\pi/4$  or  $E_0 \sin(\pi/4)$ .

Both of them are  $= E$  naught upon a square root of 2 and the crystal introduces a phase difference of  $\pi/4$  during the passage.

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The figure shows the progressive phase difference between the ordinary and the extraordinary beams.

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The figure shows the progressive phase difference between the ordinary and the extra-ordinary beams as they travel in the crystal. At emergence this phase difference becomes  $\pi/2$  as shown in the figure.

Thus, for the emergent beam, we have

$$E_y = \frac{E_0}{\sqrt{2}} \sin \omega t, \quad E_z = \frac{E_0}{\sqrt{2}} \cos \omega t,$$

which represents a circularly polarized wave because  $E_y^2 + E_z^2 = \frac{E_0^2}{2}$

As they travel in the crystal in the beginning they are in the same phase at the entry point. At emergence at the other end of the crystal, after traveling a distance  $d$ , this phase difference in this case becomes  $\pi/4$  as shown in the figure. Just for the emergent beam, we have,  $E_y = E$

naught upon a square root of 2 times sine omega t and  $E_z = E_0 \cos \omega t$  upon the square root of 2 which represents circularly polarized wave. Because  $E_y^2 + E_z^2 = E_0^2$  which is the equation of a circle.

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This is the equation of a circle.

In order to determine the direction of rotation of the electric vector, let us consider the electric fields of the components at time  $t$  equal to zero,

$$E_y = 0, \quad E_z = \frac{E_0}{\sqrt{2}},$$

and a little later at  $t = \Delta t$ ,

$$E_y = \frac{E_0}{\sqrt{2}} \omega \Delta t, \quad E_z = \frac{E_0}{\sqrt{2}} \left( 1 - \frac{(\omega \Delta t)^2}{2} \right),$$

In order to determine the direction of a rotation, I mean the sense of rotation of the electric vector. Let us consider the electric fields of the components initially at time  $t = 0$   $E_y = 0$  and  $E_z = E_0 / \sqrt{2}$ . And a little time later at  $t = \Delta t$ ,  $E_y$  becomes  $E_0 \omega \Delta t / \sqrt{2}$  upon the square root of 2 times omega times delta t and  $E_z = E_0 / \sqrt{2} (1 - \omega^2 \Delta t^2 / 2)$ .

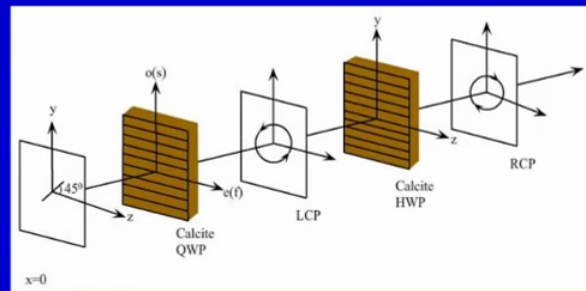
The y component has slightly increased and the z component has slightly decreased. That is why the above values show that as the time increases, the electric vector rotates in the anti clockwise direction.

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The above values show that as time increases, the electric vector rotates in the anticlockwise direction and hence the beam is left circularly polarized (LCP) as shown in the figure.

And hence the beam is left circularly polarized LCP as shown in the figure.

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Let us not consider the half-wave plate (HWP) part of the figure at the moment.

The passage through the quarter wave plate makes the initially plane polarized light. With equal amplitudes, the quarter wave plate introduces a phase difference of  $\pi/2$ , of quarter vibration the result is we get a left circularly polarized light.

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In order to introduce a phase difference of  $\pi/2$ , the thickness of the crystal plate should have a value given by the following relation

$$d = \frac{c}{\omega(n_o - n_e)} \frac{\pi}{2} = \frac{1}{4} \frac{\lambda_0}{(n_o - n_e)}$$

For calcite ,

$$n_o = 1.65836, \quad n_e = 1.48641$$

corresponding to light of  $\lambda_0 = 589.3 \text{ nm}$ .

Substituting these values gives

$$d = 0.000857 \text{ mm}.$$

In order to introduce a phase difference of  $\pi$  by 2, the thickness of the crystal plate should have a value given by the following relation  $d = c \text{ upon } \Omega \text{ into } n_o - n_e \text{ times } \pi \text{ by } 2 \text{ which} = \lambda_0 \text{ naught by } 4 \text{ divided by } n_o - n_e$ . For calcite,  $n_o$  was given by 1.65836 and  $n_e = 1.48641$ . This is corresponding to a light of wavelength 589.3 nanometers. And substituting these values gives  $d = 0.000857$  millimeters.

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Thus a calcite quarter-wave plate (at  $\lambda_0 = 589.3 \text{ nm}$ ) will have a thickness of 0.000857 mm and will have its optic axis parallel to the surface. Such a QWP will introduce a phase difference of  $\pi/2$  between the ordinary and the extra-ordinary components.

Thus a calcite quarter wave plate at  $\lambda = 589.3$  nanometers, will have a thickness of 0.000857 millimeters and will have its optic axis parallel to the surface. Remember, all the time the propagation is perpendicular to the optic axis and there is no separation between the two

beams. Such a quarter wave plate will introduce a phase difference of  $\pi$  between the ordinary and extraordinary components.

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Note that the y-polarized o-wave in calcite has a smaller wave velocity ( $= c/n_o$ ) and hence it is referred to as a slow wave and the extra-ordinary wave is the fast wave in calcite.

Note that y polarized ordinary wave in calcite has a smaller wave velocity  $c/n_o$  and hence it is referred to a slow wave and the extraordinary wave is the fast one in calcite.

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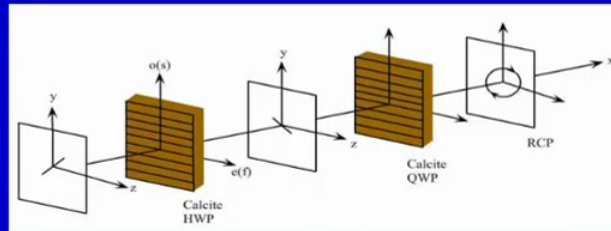
### VI.3 Half-wave plate

Let us now consider the case when a plane-polarized beam is incident on a HWP so that the phase difference  $\theta$  caused by the plate is equal to  $\pi$ . Also let us take the inclination  $\phi$  of the electric vibration of the incident plane-polarized light with the optic axis (z-axis here) equal to  $\pi/4$ .

Let us consider now a half wave plate. It is the case when a plane polarized beam with incident on HWP, half wave plate so that the phase difference  $\theta$  caused by the plate =  $\pi$  half a vibration, path difference of  $\lambda/2$ . Again let us take the inclination  $\phi$  of the electric

vibrations of the incident plane polarized beam with the optic axis which we have taken it z axis here again =  $\pi$  by 4.

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Now the y- and z-components of the incident wave have equal amplitudes and the crystal introduces a phase difference of  $\pi$ . Ignore the QWP part of the figure for the time being.

So, the ordinary and extraordinary again have equal amplitudes. The crystal introduces the phase difference of  $\pi$ . Let us not consider the quarter wave plate part at the moment. Just for the emergent light beam in this case we have,

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Thus, for the emergent beam we have

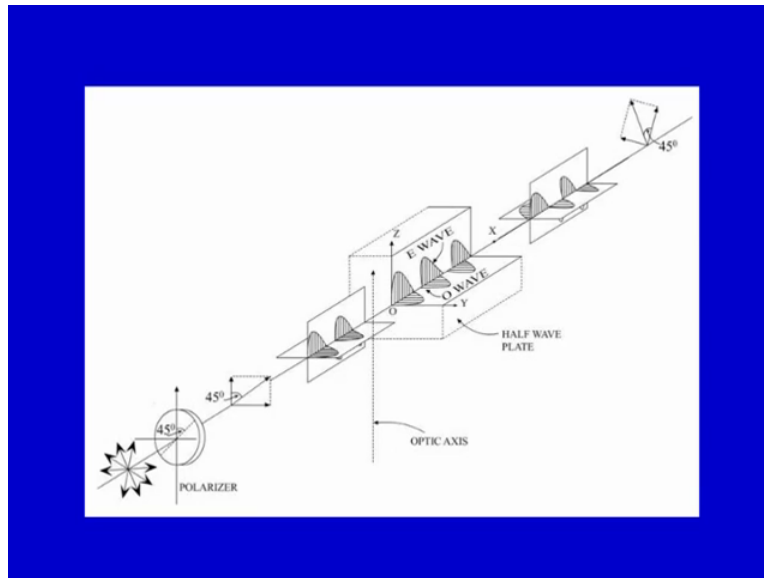
$$E_y = -\frac{E_0}{\sqrt{2}} \cos \omega t, \quad E_z = \frac{E_0}{\sqrt{2}} \cos \omega t,$$

which represents a plane-polarized wave with the direction of polarization making an angle of  $135^\circ$  with the z-axis as shown in the figure.

$E_y = -E_0$  upon a square root of  $t \cos \Omega t$  and  $E_z = E_0$  upon the square root of  $2 \cos \Omega t$ . Note the - sign which is appearing in the expression for  $E_y$ . This represents a plane

polarized wave with the direction of polarization making an angle of  $135^\circ$  with the  $z$  axis. Remember, in the beginning, at the entry point the angle was  $45^\circ$ . Now it is  $135^\circ$ .

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In the beginning okay, remember they are in the same phase.

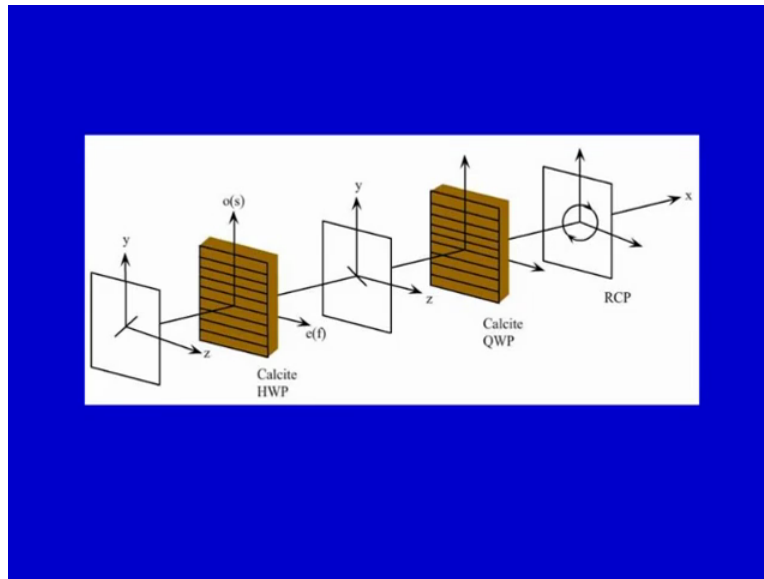
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This figure again shows the build up of the phase difference between the ordinary and the extra-ordinary. This phase difference ultimately becomes equal to  $135^\circ$ .

If we now pass this beam through a calcite QWP, the emergent beam will be right circularly polarized (RCP) as shown in the figure.

This figure again shows the buildup of the phase difference between the ordinary and extraordinary. The phase difference ultimately becomes  $= 135^\circ$ . If we now pass this beam through a calcite quarter wave plate, the quarter wave plate introduces a phase difference of  $\pi/2$ . The emergent beam will be right circularly plane polarized, right circularly polarized as shown in the figure.

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The incident light with a plane polarized with angle of 45 degrees passes through a half wave plate. The phase difference becomes 135 degrees passes through quarter wave plate which introduces another phase difference of  $\pi/2$ . And the result is the light wave right circularly polarized.

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This figure shows the passage of a plane-polarized light first through a half-wave plate which introduces a phase difference of  $\pi$  and then through a quarter-wave plate which introduces an additional phase difference of  $\pi/2$ , making a total of  $3\pi/2$ .

So, this is what this figure shows passage of a plane polarized light first through a half wave plate which as I said introduces a phase difference of  $\pi$ . And then, through a quarter wave plate which introduces an additional phase difference of  $\pi/2$  making a total of  $3\pi/2$ .

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Now

$$E_y = -\frac{E_0}{\sqrt{2}} \sin \omega t, E_z = \frac{E_0}{\sqrt{2}} \cos \omega t.$$

$$\text{Eliminating } t, \text{ we get } E_y^2 + E_z^2 = \frac{E_0^2}{2}.$$

Now finally at this stage  $E_y$  is given by  $-\frac{E_0}{\sqrt{2}} \sin \omega t$  and  $E_z = \frac{E_0}{\sqrt{2}} \cos \omega t$ , eliminating  $t$ . We again get the equation of a circle  $E_y^2 + E_z^2 = \frac{E_0^2}{2}$  again as we did it earlier. In order to determine the sense of rotation of the electric vector  $E$ ,

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In order to determine the sense of rotation of the vector  $\vec{E}$ , consider the values of  $E_y$  and  $E_z$  at time  $t = 0$ ,

$$E_y = 0 \quad E_z = \frac{E_0}{\sqrt{2}}$$

and a little later at  $t = \Delta t$ ,

$$E_y = -\frac{E_0}{\sqrt{2}} \omega \Delta t, E_z = \frac{E_0}{\sqrt{2}} \left( 1 - \frac{(\omega \Delta t)^2}{2} \right)$$

Consider the values of  $E_y$  and  $E_z$  at time  $t = 0$ . So, at time  $t = 0$   $E_y = 0$  and  $E_z = \frac{E_0}{\sqrt{2}}$ . A little time later, at  $t = \Delta t$ ,  $E_y = -\frac{E_0}{\sqrt{2}} \omega \Delta t$  and  $E_z = \frac{E_0}{\sqrt{2}} \left( 1 - \frac{(\omega \Delta t)^2}{2} \right)$ .

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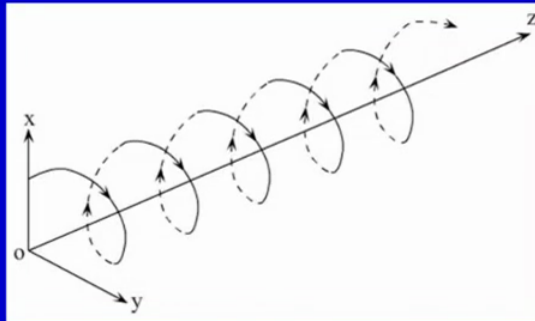


This shows that the electric vector rotates in the clockwise direction.

Note the negative sign appearing here in the expression for  $E_y$ .

The electric vector  $E$  has again increased, but in the other direction and  $E_z$  component has slightly decreased. And this shows the electric vector now in this case, rotates in the clockwise direction.

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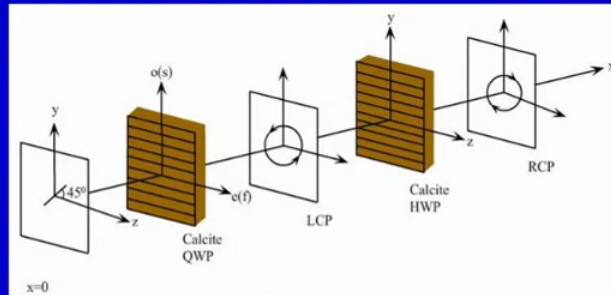


This figure shows the electric vector of a progressing (along the z-direction) right circularly polarized beam.

This figure shows the electric vector of a progressing right circularly polarized beam.

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On the other hand, if a left circularly polarized (LCP) beam is incident normally on a calcite HWP, the emergent beam will be right circularly polarized



Now, on the other hand, if your left circularly polarized beam is incident normally on a half wave plate, the emergent beam will be right circularly polarized; initial plane polarized light passes through a quarter wave plate becomes a left circularly polarized light passes through a half wave plate becomes right circularly polarized light.

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Only the intermediate state of polarization between the two plates is different here. First the quarter-wave plate introduces a phase difference of  $\pi/2$  and then the half-wave plate introduces an additional phase difference of  $\pi$  making a total of  $3\pi/2$  as before.

The thickness of the crystal plate for a HWP is given by

$$d = \frac{\lambda_0}{2(n_o - n_e)}$$

Only the intermediate state of polarization between the two plates is different here. First as I said, the quarter wave plate introduces the phase difference of  $\pi/2$ , makes the light beam left circularly polarized. And then, the half wave plate introduces an additional phase difference of  $\pi$  making a total of  $3\pi/2$  as before. And making a light beam right circularly polarized, the

thickness of the crystal plate for half a plate is given by, in a very natural way,  $d$  is given by  $\lambda n_o n_e / 2$ .

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## VII. Analysis of polarized light

A plane wave can be characterized by different states of polarization, which may be any one of the following:

- (a) Linearly or plane-polarized.
- (b) Circularly polarized.
- (c) Elliptically polarized.
- (d) Unpolarized.

Let us now consider the analysis of polarized light. The plane waves can be cut by different states of polarization as we have seen during the course of this lecture. It may be any one of the following. It could be a linearly polarized light, b. it could be a circularly polarized light, c. it could be elliptically polarized light and d. it will be just unpolarized.

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- (e) Mixture of plane-polarized and unpolarized.
- (f) Mixture of circularly polarized and unpolarized.
- (g) Mixture of elliptically polarized and unpolarized

Or in addition to all these 4 cases, e. it could be a mixture of plane polarized light and unpolarized light, it could be a mixture of circularly polarized light and unpolarized light or it

could be a mixture of elliptically polarized light and unpolarized light. So, we have 7 different possibilities. One thing should be noted.

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To the naked eye, all the states of polarization will appear to be the same. We shall now study the procedure for determining the state of polarization of a light beam.

To the naked eye, all these states of polarization will appear to be the same, absolutely no difference. We shall now study the procedure, what are the steps are to be taken, for determining the state of polarization of a light beam.

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If we introduce a polarizer (working here as an analyzer) in the path of the beam and rotate it about the direction of propagation, then either of the following three possibilities can occur.

(i) If there is complete extinction at two positions ( $180^\circ$  apart) of the analyzer, then the beam is plane-polarized.

If we introduce a polarizer, working here as an analyzer, in the path of the beam and we rotate it about the direction of propagation, then either of the following three possibilities can occur. You

see, the incident light beam with the air which is to be analyzed. We pass it through a polarizer and we rotate the polarizer about the direction of the beam and watch what happens.

First case if there is complete extinction at 2 positions, which are 180 degrees apart of the analyzer then, the beam is plane polarized. This means we rotate it we get 0 intensity at some position; rotate it 90 degrees, becomes maximum for the rotation of 90 degrees, again the intensity becomes 0. So, in a complete rotation, there are two positions 180 degrees apart, where there is a complete extinction. Remember, Malus law. And that shows that incident beam is plane polarized.

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(ii) If there is no variation of intensity, then the beam is either unpolarized or circularly polarized or a mixture of unpolarized and circularly polarized light.

We now put a quarter-wave plate in the path of the beam followed by the analyzer which is rotated as before about the direction of propagation.

Number two, if there is no variation of intensity then the beam is either unpolarized again we have seen this thing earlier. If the beam is unpolarized and you pass it through a polarizer, let it be call it as an analyzer here; and rotate it, there is no variation in the intensity. So, the beam could be unpolarized but we also know that the same thing will happen if the beam is circularly polarized or is a mixture of unpolarized and circularly polarized light.

To analyze the differentiate between these two possibilities that the light beam is unpolarized or is a mixture of unpolarized and circularly polarized. We now put a quarter wave plate in the path of the beam. So, the incoming beam we allow it to pass through a quarter wave plate and again by an analyzer, which is rotated as before about the direction of propagation.

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Three situations are now possible:

A. If there is no variation of intensity then the incident beam is unpolarized.

B. If there is complete extinction at two positions ( $180^\circ$  apart), then the original beam is circularly polarized. This is due to the fact that a quarter wave plate will transform a circularly polarized light into a plane-polarized light.

When we do this, three situations are possible: A. Now, there is no variation of intensity. The result is this means the incident beam is unpolarized. The reason is a quarter wave plate does essentially nothing to an unpolarized beam. It was showing no variation when analyzer is rotated even after passes through a quarter wave plate, the same situation will prevail. There will be no variation in intensity so that so that incident beam is unpolarized.

B. If there is a complete extension at two positions  $180$  degrees apart, this means the original beam is circularly polarized this is due to the fact that a quarter wave plate will transform circularly polarized light into a plane polarized light. You see, the incident incident beam is circularly polarized means, a combination of two perpendicular plane polarized light with a phase difference of  $\pi$  by  $2$ . The quarter wave plate introduces an additional phase difference of  $\pi$  by  $2$  making a total of  $\pi$   $0$ .

In both the cases, the emergent light becomes a plane polarized and there will be complete extension right to the  $2$  positions following Malu's law, okay. So, in case B clearly shows that the beam is circularly polarized.

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C. If there is a variation of intensity but without complete extinction, then the beam is a mixture of unpolarized and circularly polarized light.

C. A third possibility C, C is if there is a variation of intensity but without complete extinction naturally in this case at the incident beam is a mixture of unpolarized and circularly polarized. Passage to the quarter of a plate does nothing to the unpolarized part of the beam, part of the mixture, circularly polarized light becomes plane polarized. So, the result after passage through the quarter wave plate is something like a mixture of unpolarized and plane polarized. When we pass it to an analyzer and rotate it, the variation in intensity but complete extinction is not there.

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(iii) Now we come to the third possibility:

If there is a variation of intensity but without complete extinction, then the beam is either elliptically polarized or a mixture of plane-polarized and unpolarized or a mixture of elliptically polarized and unpolarized.

Now, we come to the third possibility. If there is a variation of intensity, but without complete extinction, the incident beam could be either elliptically polarized, given the same variation or a

mixture of plane polarized and unpolarized, but it could be a mixture of elliptically polarized and unpolarized. So, here again, its a quarter wave plate is a very useful set, very useful piece.

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We now put a quarter-wave plate in the path of the beam with its optic axis parallel to the pass-axis of the analyzer at the position of maximum intensity . This is followed by the analyzer which is a rotated as before about the direction of propagation.

We again put a quarter wave plates in the path of the beam with its optic axis parallel to the pass axis of the analyzer, at the position of maximum intensity. This is a change from the earlier case we put the quarter wave plate, all right deep in the path of the beam. There is optic axis parallel to the pass success of the analyzer, led the position of maximum intensity.

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Three situations may now occur:

A. The elliptically polarized light will transform to a plane-polarized light. Thus, if one obtains two positions ( $180^\circ$  apart) of the analyzer where complete extinction occurs, then the original beam is elliptically polarized.

This is followed by the analyzer which is rotated again as before about the direction of propagation. Three situations may occur: A. The elliptically polarized light will transform to a

plane polarized light. Remember, elliptically polarized light, there is also a path difference of a quarter vibration. Quarter wave plate will introduce an additional quarter vibration. The result is the emergent light will become plane polarized.

And one will get two positions and the 180 degree apart of the analyzer where the complete extinction occurs, the Malus Law. Then, the original beam is elliptically polarized.

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B. If complete extinction does not occur and the position of maximum intensity occurs at the same orientation as before, the original beam is a mixture of unpolarized and plane-polarized light.

C. Finally, if the position of maximum intensity occurs at a different orientation of the analyzer, the original beam is a mixture of elliptically polarized and unpolarized light.

B. If complete extinction does not occur and the position of maximum intensity occurs at the same orientation as before, before putting the quarter wave plate, same position as before, the original beam is a mixture of unpolarized and plane polarized light. The case C is finally the position of maximum intensity. Again there is a variation but, the position of maximum intensity occurs at a different orientation of the analyzer. In this case, the original beam is a mixture of elliptically polarized and unpolarized light.

So, these are the various steps which are needed to analyze the polarization, various polarization of the light. The 7 cases which we have considered, the unpolarized plane polarized, circularly polarized, elliptically polarized or the mixture when you have unpolarized along with circularly polarized or unpolarized along with plane polarized or unpolarized along with elliptically polarized case. With this, we have come to the end of the present lecture.