

Engineering Physics
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Module-01
Lecture 02
Malus Law & Superposition of Waves

I am M.K. Srivastav of the Department of Physics, IIT Roorkee. This is the second lecture of the five lecture series on Polarization. In the last lecture, which was the first one, we considered Maxwell's equations leading to electromagnetic waves. We then considered the meaning of polarized light, polarization by reflection, refraction and scattering. We study Brewster's law and considered working of a Polaroid.

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In the last lecture we considered Maxwell's equations leading to electromagnetic wave. We then considered meaning of polarized light and polarization by reflection, refraction and scattering. We studied Brewster's law and considered working of a Polaroid.

In the present lecture we shall take up Malus law and then shall consider superposition of two electromagnetic waves under different conditions.

In the present lecture which is the second one of the series, we shall take up Malus law and then shall consider superposition of two electromagnetic waves under different conditions. The different condition means difference in amplitude, different phase differences and things like that. Law of Malus;

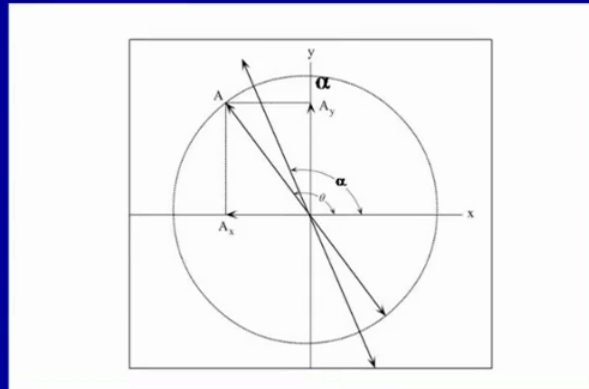
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III. Law of Malus

Consider an unpolarized beam of amplitude A and intensity I_0 (which is proportional to square of the amplitude) incident on a polarizer. Let the incident light have components $A_x = A \cos \theta$ and $A_y = A \sin \theta$ along the axes with θ varying randomly as the light is unpolarized.. Let the pass-axis of the polarizer make an angle α with the x-axis.

Considered an unpolarized beam of amplitude A and intensity I naught which is proportional to the square of the amplitude, as we know, incident on a polarizer, the incident light have components $A_x = A \cos \theta$ and $A_y = A \sin \theta$ along the axis the θ varying randomly as the light is unpolarized.

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Let the pass axis of the polarizer making an angle α with the x axis. The polarizer as we have seen is a device which allows the light wave with electric vibrations parallel to the pass direction to pass through without attenuation essentially and absorb the light wave electric vibrations perpendicular to the pass direction.

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Polarizer, as we have seen, is a device which allows the light wave with electric vibrations parallel to the pass-direction to pass through without attenuation and absorbs the light wave with electric vibrations perpendicular to the pass-direction.

The long arrow in the figure making an angle α with the x-axis indicates the pass-direction of the polarizer.

The long arrow in the figure making an angle α with the x axis indicates the pass direction of the polarizer.

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The intensity I of the transmitted beam through this polarizer can be obtained by taking the projection of the components A_x and A_y along the pass-direction. As the intensity is proportional to the amplitude, we have

$$\begin{aligned} I &\propto (A_x \cos \alpha + A_y \sin \alpha)^2 \\ &\propto A^2 (\cos \theta \cos \alpha + \sin \theta \sin \alpha)^2 \\ &\propto A^2 (\cos^2 \theta \cos^2 \alpha + \sin^2 \theta \sin^2 \alpha \\ &\quad + 2 \sin \theta \cos \theta \sin \alpha \cos \alpha) \end{aligned}$$

The intensity I of the transmitted beam through the polarizer can be obtained by taking the production of the components A_x and A_y along the pass direction and intensity is proportional to the amplitude we have, Intensity I is proportional to the square of $A_x \cos \alpha + A_y \sin \alpha$ it means the proportional to A square times the square of $\cos \theta \cos \alpha + \sin \theta \sin \alpha$.

It means the proportional to $A^2 \cos^2 \theta \cos^2 \alpha + \sin^2 \theta \sin^2 \alpha + 2 \sin \theta \cos \theta \sin \alpha \cos \alpha$. Now and the angle θ varies randomly in the unpolarized beam. The average values of term appearing there are $\cos^2 \theta$, the average value is half $\sin^2 \theta$ again.

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Now, as the angle θ varies randomly in the unpolarized beam, the average values of $\cos^2 \theta$, $\sin^2 \theta$ and $\sin \theta \cos \theta$ are needed to evaluate the above expression. These average values are

$$\langle \cos^2 \theta \rangle_{av} = 1/2$$

$$\langle \sin^2 \theta \rangle_{av} = 1/2$$

$$\langle \sin \theta \cos \theta \rangle_{av} = 0$$

The average value are half you see the averages are taken over a complete cycle then the average value of $\sin \theta \cos \theta$ again over a complete cycle is 0.

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This leads to

$$I \propto A^2 \left(\frac{1}{2} \cos^2 \alpha + \frac{1}{2} \sin^2 \alpha \right)$$

$$\propto \frac{1}{2} A^2$$

or

$$I = \frac{1}{2} I_0.$$

The intensity of the transmitted beam is thus half of the intensity of the incident unpolarized beam,

This leads to intensity is not proportional to $A^2 \cos^2 \alpha + \sin^2 \alpha$, which means it is proportional to half of A^2 . And we get the

final result that $I = \frac{1}{2} I_0$. The intensity of the transmitted wave means, remember, the incident beam is unpolarized; so, the intensity of the transmitted beam is thus half of the intensity of the incident unpolarized beam.

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and this result *does not* depend on the orientation of the pass-direction of the polarizer, and the transmitted beam is plane-polarized with electric vibrations along the pass-direction of the polarizer.

And another thing this result does not depend on the orientation of the pass direction of the polarizer. That is it is independent of the angle α . And the transmitted beam is plane polarized with electric vibrations along the pass direction of the polarizer.

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If the incident beam is plane-polarized with electric vibrations making an angle θ (now fixed) with the x -axis and, the pass direction of the polarizer makes an angle α with the x -axis as before, the intensity of the transmitted beam is again given by the same expression.

If the incident beam is plane polarized with electric vibrations making an angle θ as before but now this angle is fixed with the x axis and the pass direction of the polarizer makes an angle

alpha with the x axis as before the intensity of the transmitted beam is again given by the same expression.

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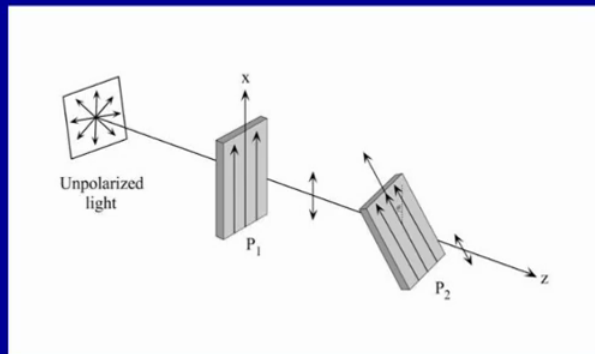
$$\begin{aligned} I &\propto (A_x \cos \alpha + A_y \sin \alpha)^2 \\ &\propto A^2 (\cos \theta \cos \alpha + \sin \theta \sin \alpha)^2 \\ &\propto A^2 \cos^2 (\theta - \alpha) \\ &= I_0 \cos^2 (\theta - \alpha) \end{aligned}$$

But, no averaging need be done now.

That is, the intensity proportional to the square of $A_x \cos \alpha + A_y \sin \alpha$ which means it is proportional to $A^2 \cos^2 (\theta - \alpha)$. So, it is proportional to $A^2 \cos^2 (\theta - \alpha)$. Therefore the intensity $I = I_0 \cos^2 (\theta - \alpha)$. You see, no averaging need be done now because the angle θ is fixed; the incident light beam is polarized.

The intensity of the transmitted beam now you find thus varies as the square of the cosine of the angle $\theta - \alpha$, the angle between the plane of vibration, of the incident beam and the pass direction of the polarizer, working here as an analyzer.

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This is called Law of Malus, Now if a plane polarized beam, polarized along the x axis by the polarizer P_1 , the incident on another polarizer P_2 here.

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Now, if a plane-polarized beam (polarized along the x -axis by the polarizer P_1) is incident on another polarizer (P_2 here) and if this polarizer is rotated about the beam direction, then the intensity of the emergent wave will vary according to the above law.

And if this polarizer is rotated about the beam direction then, the intensity of the emergent wave will vary according to the above law.

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For example, if the polarizer P_2 is rotated in the clockwise direction, then the intensity will increase till the pass-axis is parallel to the x-axis; a further rotation will result in a decrease in intensity till the pass-axis is perpendicular to the x-axis, where the intensity will be almost zero.

For example, the polarizer P_2 is rotated in the clockwise direction, then, the intensity will increase, till the pass axis is parallel to the x axis. Further rotation will result in a decrease in intensity till the pass axis is perpendicular to the x axis where the intensity will be almost 0. If we further rotate it will pass through maximum and again a minimum before it reaches its original position.

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If we further rotate it, it will pass through a maximum and again a minimum before it reaches its original position.

The emergent light after passing through the polarizer P_2 is now polarized with direction of electric vibrations parallel to the pass direction of P_2 (working as an analyzer).

The emergent light after passing through the polarizer P_2 is now polarized with the direction of electric vibrations parallel to the pass direction of P_2 working here as an analyzer. Let us now consider superposition of two waves.

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IV. Superposition of two waves

Let us consider the propagation of two plane polarized electromagnetic waves (both propagating along the z -axis) with their electric vectors oscillating along the same direction, *i.e.* x -axis.

The electric fields associated with the waves can be written in the form

$$\vec{E}_1 = a_1 \cos(kz - \omega t + \theta_1) \hat{x}$$

$$\vec{E}_2 = a_2 \cos(kz - \omega t + \theta_2) \hat{x}$$

Consider the propagation of two plane polarized electromagnetic waves both propagating along the z axis that is the same direction with their electric vectors oscillating along the same direction that is the x axis. For the electric fields associated with the waves can be written in the form, $E_1 = a_1 \cos(kz - \omega t + \theta_1)$ and E_2 which is $a_2 \cos(kz - \omega t + \theta_2)$.

Both of these fields are along the x axis. a_1 and a_2 representing the amplitudes of the waves and \hat{x} represents the unit vector along the x axis, θ_1 and θ_2 are arbitrary constants. The resultant of these two waves would be given by

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where a_1 and a_2 represent the amplitudes of the waves; \hat{x} represents the unit vector along the x -axis, and θ_1 and θ_2 are arbitrary phase constants.

The resultant of these two waves would be given by $\vec{E} = \vec{E}_1 + \vec{E}_2$

which can always be written in the form

$$\vec{E} = a \cos(kz - \omega t + \theta) \hat{x},$$

Just the vector addition, so $E = E_1 + E_2$. Now, this can always be written in a form like this: $E = a \cos(kz - \omega t + \theta)$ along the x axis naturally,

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where

$$a = \left[a_1^2 + a_2^2 + 2a_1a_2 \cos(\theta_1 - \theta_2) \right]^{1/2}$$

represents amplitude of the resultant wave.

We observe that the resultant is also a plane-polarized wave with its electric vector oscillating along the same direction.

Where a is given by the square root of $a_1^2 + a_2^2 + 2a_1a_2 \cos(\theta_1 - \theta_2)$, a represents the amplitude of the resultant wave; resultant of the superposition of the original two waves. We observe that the resultant is also a plane polarized wave with its electric vector oscillating along the same direction, along the x axis.

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Let us now consider superposition of two plane-polarized waves (both propagating along the z-axis) but with their electric vectors oscillating along two mutually perpendicular directions. Thus, we have

$$\vec{E}_1 = a_1 \cos(kz - \omega t) \hat{x}$$

$$\vec{E}_2 = a_2 \cos(kz - \omega t + \theta) \hat{y}.$$

Angle θ is the phase difference between them.

Let us now consider superposition of two plane polarized waves both propagating along the z axis but with their electric vectors oscillating along two mutually perpendicular directions.

Earlier they were along the same direction; now, they are mutually perpendicular directions, say along the x axis and along the y axis. Remember, both the waves are propagating along the z axis.

So, we can write $E_1 = a_1 \cos(kz - \omega t)$ this is along the x axis. $E_2 = a_2 \cos(kz - \omega t + \theta)$ along the y direction. The angle θ is the phase difference between these two waves.

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In order to find their resultant, we consider time variation of the resultant electric field at an arbitrary plane perpendicular to the propagation direction (here, z-axis) which we may, without any loss of generality, assume to be $z = 0$.

If E_x and E_y represent the x- and y-components of the resultant field $\vec{E} = \vec{E}_1 + \vec{E}_2$, then

$$E_x = a_1 \cos \omega t$$

$$E_y = a_2 \cos(\omega t - \theta)$$

Now, in order to find the resultant, we consider the time variation of the resultant electric field at an arbitrary plane, just for simplicity, at an arbitrary plane perpendicular to the propagation direction which is the z axis. This arbitrary plane without any loss of generality we can assume it to be $z = 0$ just for simplicity. E_x and E_y represent the x and y components of the resultant field $E_1 + E_2$, then E_x will be $= a_1 \cos \omega t$ and E_y is $= a_2 \cos(\omega t - \theta)$.

Now eliminate t between these two equations. And for that we rewrite the above expression for E_y as;

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We eliminate t between these two equations. For that, we rewrite the above expression for E_y as

$$\begin{aligned} E_y &= a_2 \cos \omega t \cos \theta + a_2 \sin \omega t \sin \theta \\ &= a_2 \cos \omega t \cos \theta + a_2 \sqrt{1 - \cos^2 \omega t} \sin \theta \\ &= \frac{a_2}{a_1} E_x \cos \theta + a_2 \sqrt{1 - \left(\frac{E_x}{a_1}\right)^2} \sin \theta \\ \text{or } \left(E_y - \frac{a_2}{a_1} E_x \cos \theta\right)^2 &= a_2^2 \sin^2 \theta \left[1 - \left(\frac{E_x}{a_1}\right)^2\right] \end{aligned}$$

$E_y = a_2 \cos \Omega t \cos \theta + a_2 \sin \Omega t \sin \theta$. We can write it as first term remaining unchanged, $a_2 \cos \Omega t \cos \theta + a_2$ times the square root of $1 - \cos^2 \Omega t$ times $\sin \theta$. This can be written as $= a_2$ upon a_1 times, $E_x \cos \theta + a_2$ times the square root of $1 - E_x$ upon a_1 square times $\sin \theta$, or this means $E_y - a_2$ upon a_1 $E_x \cos \theta$ whole squared is $= a_2^2 \sin^2 \theta$ times $1 - E_x$ upon a_1 square.

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Rearranging this, we get

$$\frac{E_x^2}{a_1^2} + \frac{E_y^2}{a_2^2} - \frac{2E_x E_y}{a_1 a_2} \cos \theta = \sin^2 \theta.$$

This is a general equation of an ellipse and represents, in general, the result of superposition of two perpendicular vibrations of different amplitudes and a phase difference of θ .

We can rearrange this and when we do this, we get, E_x is squared upon a_1 square + E_y square + a_2 square - twice $E_x E_y$ upon $a_1 a_2$ times $\cos \theta$. And all of this is $= \sin^2 \theta$. This is the general equation of an ellipse. The general equation of an ellipse, is an ellipse is at the center of the origin which is the major minor axis making some angle with the coordinate axis.

So, this represents in general the result of superposition of two perpendicular vibrations of different amplitudes and the phase difference of theta. For the phase difference theta = n Pi where n is 0, 1, 2. This means an integral multiple of Pi 1, Pi 2, Pi 3, Pi 4.

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For the phase difference $\theta = n\pi$, ($n = 0, 1, 2, \dots$),
above equation becomes

$$\text{i.e.,} \quad \left[\frac{E_x}{a_1} - (-1)^n \frac{E_y}{a_2} \right]^2 = 0$$

$$\frac{E_y}{E_x} = (-1)^n \frac{a_2}{a_1}.$$

The above equation becomes E_x upon $a_1 - 1$ raised to the power n, E_y upon a_2 and the whole square = 0 which really means E_y upon E_x . The ratio of the Y component to the X component is = -1 raised to the power n, a_2 upon a_1 . In the E_x, E_y plane the above expression represents a straight line. The angle Phi that this line makes with the x axis depends naturally on the ratio a_2 upon a_1 .

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In the $E_x - E_y$ plane, the above expression represents a straight line; the angle ϕ that this straight line makes with the E_x -axis depends on the ratio a_2/a_1 .

In fact

$$\phi = \tan^{-1} \left(\pm \frac{a_2}{a_1} \right).$$

The condition $\theta = n\pi$ implies that the two vibrations are either in phase ($n = 0, 2, 4, \dots$) or completely out of phase ($n = 1, 3, 5, \dots$).

In fact, angle Phi is tangent inverse + - a2 upon a1. If $a_1 = a_2$, this angle will be 45 degrees. The condition $\theta = n\pi$ implies that the two vibrations are, either in phase, if n is even 0, 2, 4 which means $\theta = 0$ or the phase difference of $2\pi, 4\pi, 6\pi$ or completely out of phase which means $n = 1, 3, 5$; the phase difference $\theta = \pi, 3\pi, 5\pi$, like this. This means the phase difference of half a vibration, 3 half vibrations or 5 half vibrations;

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Thus, the superposition of two linearly-polarized light waves with their electric fields at right angles to each other and oscillating in phase or completely out of phase, is again a linearly-polarized wave with its electric vector, in general oscillating in a direction which is different from the fields of either of the two.

Thus the super position of two linearly polarized light waves with their electric fields at right angles to each other. And oscillating in phase are completely out of phase, ok. These just these two conditions: digital vibrations are perpendicular to each other; but they are in phase, or completely out of phase. Then, the resultant is again the linearly polarized wave with a selective

vector in general oscillating in a direction which is different from the fields wider of the two making an angle of Φ .

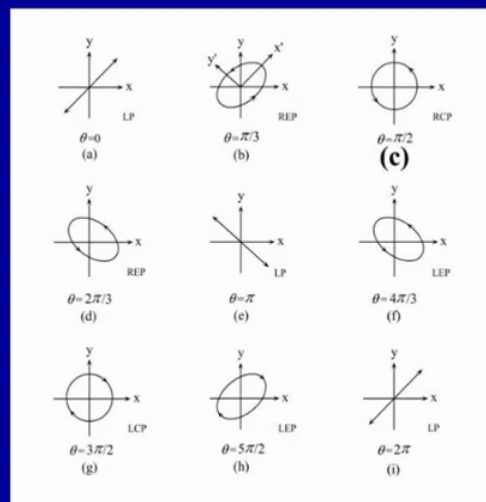
We would say the first vibrations and this angle Φ depends on the ratio of the two amplitudes a_2 to a_1 .

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For $\theta \neq n\pi$ ($n = 0, 1, 2, \dots$), the resultant electric vector *does not*, in general, oscillate along a straight line.

For $\theta \neq n\pi$, it is not if the phase difference is not a integer multiple of π is not $= \pi$ or 2π or 3π , the resultant electric vector in this case does not oscillate along a straight line.

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This figure shows the resultant for different values of θ ; $\theta = 0, \pi/3, \pi/2, 2\pi/3, \pi, 4\pi/3, 3\pi/2, 5\pi/3, 2\pi$ is a complete range of θ varying from 0 to 2π .

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This figure shows states of polarization of the resultant for various values of θ corresponding to equal amplitudes, $a_1 = a_2$.

Corresponding to the simpler case when the amplitudes are equal $a_1 = a_2$. We shall consider these cases in a little more detail.

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Sections (a), (e) and (i) of this figure correspond to $\theta = 0, \pi$ and 2π respectively.

The sections a, e and i of this figure correspond to $\theta = 0, \pi$ and 2π respectively, a, corresponds to $\theta = 0$. That is the two vibrations are in the same phase. Remember here the amplitudes are equal $a_1 = a_2$. The E corresponds to $\theta = \pi$ they are completely out of

phase. The phase difference of half a vibration and π is again when the phase difference is 2π which is essentially means phase difference is 0.

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These are the situations when the vibrations are in phase or completely out of phase. We have already considered these cases.

These are the situations as I pointed out earlier, when the vibrations are in phase or completely out of phase. We have already considered these cases. Let us consider now, a simple case corresponding to $\theta = \pi/2$. That is very interesting. The difference is of a quarter vibrations.

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Let us first consider the simple case corresponding to $\theta = \pi/2$. Thus

$$E_x = a_1 \cos \omega t,$$

$$E_y = a_1 \sin \omega t,$$

$$\frac{E_y}{E_x} = \tan \omega t.$$

E_x , $a_1 \cos \Omega t$; E_y $a_1 \sin \Omega t$. Remember $a_1 = a_2$. E_x is $a \cos \Omega t$ and E_y is $a \sin \Omega t$. The phase difference of $\pi/2$, E_y upon E_x , comes out to be $\tan \Omega t$. If

we plot the time variation of the resultant of these two, we would find that the tip of the electric vector rotates from the circumference of a circle of radius a_1 , $E_x^2 + E_y^2 = a_1^2$ in the anti clockwise direction.

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If we plot the time variation of the resultant of these, we would find that the tip of the electric vector rotates on the circumference of a circle of radius a_1

$$E_x^2 + E_y^2 = a_1^2.$$

in the anticlockwise direction.

The propagation of light in the pole is in the positive direction of that axis. This had been taken here to be coming out of the screen.

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The propagation of light is in the $+z$ -direction which has been taken here to be coming out of the screen. Such a wave is known as a right circularly polarized wave (usually denoted as a RCP wave) and is shown in section (c) of the figure.

Such a wave is known as right circularly polarized wave usually denoted as RCP wave, right circularly polarized and is shown in Section c of the figure which is a section c. This shows a circular vibration. Circular because $a_1 = a_2$ and it is anti clockwise direction. Remember the

beam propagation is coming out of the screen. Directive vector is rotating in the xy plane in the anti clockwise direction. And $\theta = 3\pi/2$.

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For $\theta = 3\pi/2$, we have

$$E_x = a_1 \cos \omega t$$

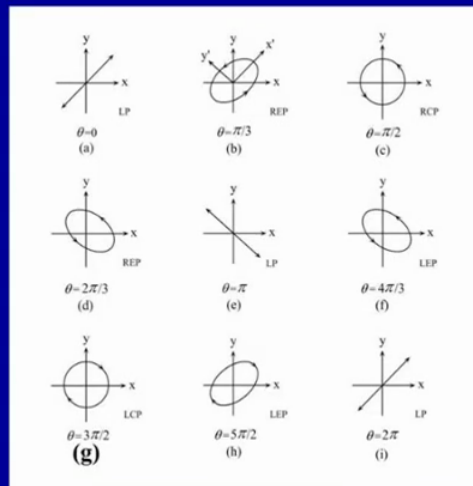
$$E_y = -a_1 \sin \omega t$$

which also represents a circularly polarized wave.

However, in this case, the electric vector rotates in the clockwise direction. Such a wave is known as left circularly polarized wave (usually denoted as LCP wave) and is shown in section (g) of the figure.

Now we have $E_x = a_1 \cos \Omega t$ as before. But now $E_y = -a_1 \sin \Omega t$. It again represents a circularly polarized wave as before. However, in this case, the electric vector rotates in the clockwise direction. Such a wave is known as left circularly polarized wave usually denoted as LCP wave and is shown in Section g of the figure.

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Again a circular vibration, the only difference with the earlier case c is that here the electric vector rotates in the clockwise direction. Remember, the beam is propagation Direction is coming out of the screen. This is called a left circularly polarized.

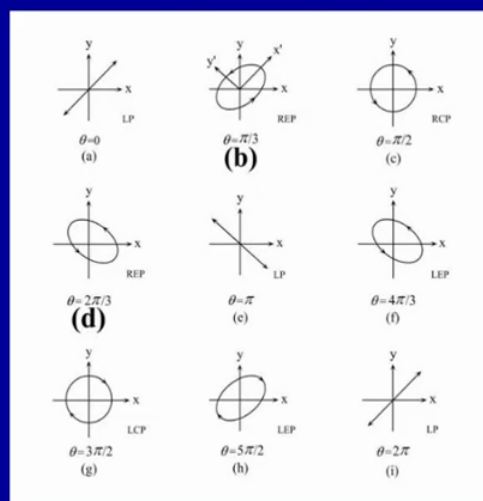
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For $\theta \neq m\pi/2$ ($m = 0, 1, 2, 3, \dots$), the tip of the electric vector rotates on the circumference of an ellipse.

Sections (b) and (d) of the figure for $\theta = \pi/3$ and $\theta = 2\pi/3$ correspond to right elliptically polarized (REP) light

Now for $\theta \neq m\pi/2$; so, this means, this is a case where it is not = an integer multiple of $\pi/2$. In this case, the tip of the electric vector rotates on the circumference of an ellipse. Sections b and d of the figure for $\theta = \pi/3$ and $\theta = 2\pi/3$ correspond to right polarized elliptical light. This is the REP light elliptically polarized.

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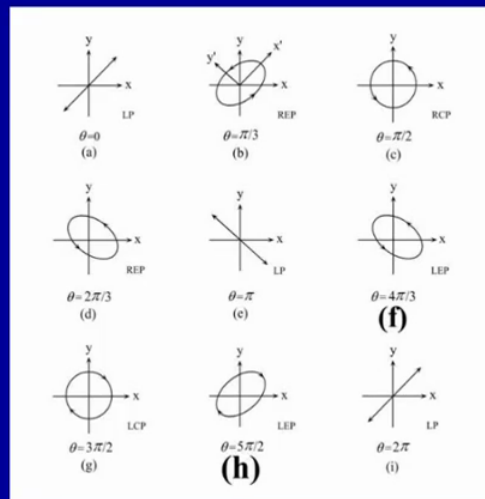
Remember again the $a_1 = a_2$. But the only difference here is the phase difference is not just $\pi/2$ or $3\pi/2$, it is $\pi/3$ or in the other. In the in the section d, it is $2\pi/3$. And that is why the resultant is, is a general ellipse whose major minor axis makes some angle with the coordinate axis. So, this is the situation when the phase difference is $\pi/3$ and $2\pi/3$.

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and sections (f) and (h) for
 $\theta = 4\pi/3$ and $\theta = 5\pi/3$
 correspond to left elliptically
 polarized (LEP) light.

Similarly, now, the sections f and h for $\theta = 4\pi/3$ and $\theta = 5\pi/3$, they correspond to left elliptically polarized LEP light.

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These are all in here section F again we have elliptical the only see the sense of rotation is different compared to the earlier case. And again remember the amplitudes are equal. So it is the

wide change in the resultant is found depending on the phase difference between the two vibrations. The two vibrations are perpendicular to each other. Their amplitudes are same, equal. $a_1 = a_2$. The only difference is the phase difference.

If the phase difference is 0 or π , 2π , 3π or 4π completely in phase or out of phase, the resultant is again a plane polarized wave linear vibrations making an angle of 45 degrees or 135 degrees in this case, $a_1 = a_2$. If the phase difference is $\pi/2$ or $3\pi/2$, this means half a vibration or three half vibrations the result is a circular vibration this is called a circularly polarized light; right circularly or left circularly depending on where the tip of the electric vector rotates in anti-clockwise direction or clockwise direction and the general case when the phase difference is not that simple.

If it is $\pi/3$, in this case, I showed in the figure or $2\pi/3$ or $5\pi/3$, the result is elliptically polarized light, right elliptically or left elliptically depending on whether the electric vector rotates as can be seen from the figure.

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As can be seen from the figure, this ellipse degenerates into a straight line or a circle when θ becomes an even or an odd multiple of $\pi/2$.

In general, when $a_1 \neq a_2$, we get an elliptically polarized wave which degenerates into a straight line for $\theta = 0, \pi, 2\pi, 3\pi, \dots$, etc.

Note again that all these different states of polarization are characteristics of transverse waves only.

Just let us in general degenerate into a straight line or a circle, as we have seen, whenever θ becomes an even or an odd multiple of $\pi/2$ as we have seen. But in general, when $a_1 \neq a_2$, we get an elliptically polarized light waves which again degenerates into a straight line for $\theta = 0, \pi, 2\pi, 3\pi$, etcetera. Remember, $\theta = 0, 2\pi, 4\pi$, means they are in the same phase.

And $\theta = \pi, 3\pi, 5\pi$ means they are completely out of phase.

Again note one thing all these different states of polarization which we have considered here depending on the value of θ are characteristics of transverse waves only and electromagnetic waves are transverse waves. So, this is the end of the second lecture. I hope you enjoyed this lecture and also the earlier one. Thank you.