

**Engineering Physics 1**  
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**Module-04**  
**Lecture-03**  
**Diffraction Part 03**

In the previous lecture, we had discussed how the diffraction effects limits the resolving power of optical instruments and explain the relay criterion of just resolution. Using this criterion, I derived the expression for resolving power of rating. Now, in this lecture I am going to describe the resolving power of microscope and telescope. So, let us first discuss about the resolving power of microscope.

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**Resolving power of a Microscope:**

As we know the primary function of a microscope is not to magnify an object but to reveal those finer details in the object which are invisible to unaided eye. The extent to which finer details are revealed depends not on the magnifying power but on the resolving power of the microscope. Resolving power of a microscope is conveniently expressed in terms of the smallest linear separation between two points object which are just resolved, that is their diffraction patterns formed by the microscope.

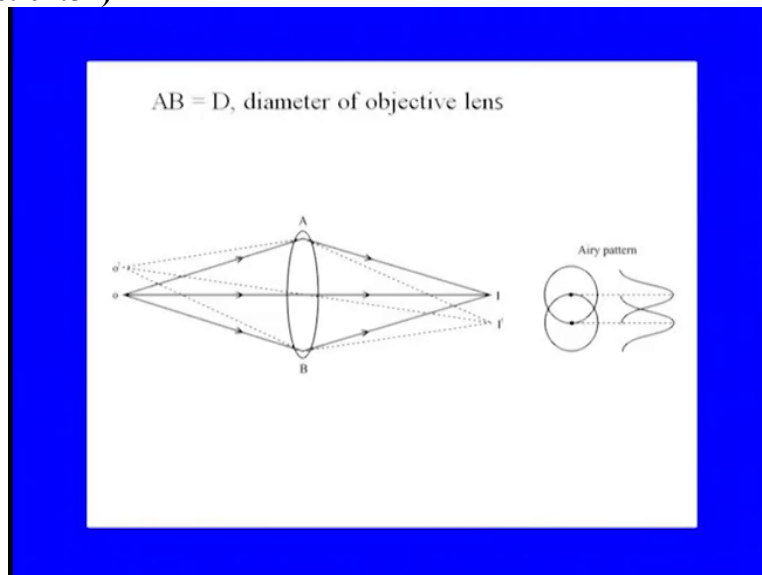
As we know primary function of a microscope is not to magnify an object but to reveal those final details in the object which are invisible to unaided eye. The extent to which final details are revealed depends not on the magnifying power, but on the resolving power of microscope. Resolving power of a microscope is conveniently expressed in terms of the smallest linear separation between two points object which are just resolved. That is the difference and patterns formed by the microscope.

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objects are just distinguishable. The smaller is this distance, the greater is said to be resolving power. Let us derive this least separation between two point objects which are self luminous having no phase relations.

Objects are just distinguishable, the smaller is this distance, the greater is said to be resolving power. Let us derive this least separation between two point objects which are self luminous having no phase relationship. The objective of the microscope is a simple convex lens as shown in this figure.

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And the waves of light assumed to be of the same wavelength  $\lambda$  starting from self luminous point objects  $O$  and  $O'$  are diffracted by the circular periphery of the lens. Therefore, a diffraction pattern corresponding to each object is formed in the focal plane of the objective. Each image is therefore a diffraction pattern consisting of a central bright disk bordered by alternate concentric circular dark and bright rings.

The centre of the bright disc in the diffraction pattern of O is at I while for O Prime it is at I prime. If these objects are to be resolved in their images then according to Rayleigh's criterion of resolution the distance I, I prime between the centre of 2 Central Maxima should be at least = the radius of either one. This simply means that the waves of light from O Prime after refraction by the lens should fall first dark ring passing through I.

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If these object are to be resolved in their images then, according to Rayleigh's criterion of resolution, the distance O' between the centre of two central maxima should be at least equal to the radius of either one. This simply means that the waves of light from II' after diffraction by the lens, should form first dark ring passing through I. This is possible only if, the path difference between the extreme rays O'BI and O'AI is equal to  $1.22\lambda_0$ , that is

$$O'BI - O'AI = 1.22\lambda_0$$

This is possible only if the path difference between the extreme rays O Prime BI and O prime AI is  $= 1.22 \lambda$ . That is O Prime BI - O prime AI is  $= 1.22 \lambda$ .

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where  $\lambda_0$  is the wavelength of light waves in a medium of refractive index  $\mu$ , which is supposed to be between the objects and the objective.

$$(O'B + BI) - (O'A + AI) = 1.22\lambda_0$$

If O is symmetrical w.r.t. A and B, then point I will also be symmetrical, that is  $BI = AI$ . Hence the above equation reduces to

$$O'B - O'A = 1.22\lambda_0 \quad (O'B - OB) + (OA - O'A) = 1.22\lambda_0$$

The objects O and O' are in practice, so close together that we can consider O'A essentially parallel to OA and O'B parallel to OB as shown in this figure. Therefore, drawing the line

, where  $\lambda$  is the wavelength of light waves in a medium of refractive Index  $\mu$ , which is supposed to be between the objects and the objective. Now this condition we can also write  $O'BI - O'AI = 1.22 \lambda$ . If O is symmetrical with respect to A and B then,

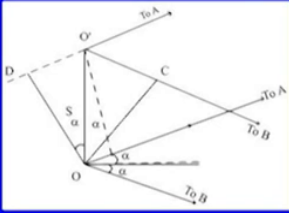
point I will also be symmetrical. That is BI will be = AI. Hence, the above equation reduces to O prime B - O Prime A = 1.22 Lambda. This equation we can also write O prime B - OB + OA - O Prime A = 1.22 Lambda.

The objects O and O prime are in practice so close together that we can consider O prime A is essentially parallel to OA and O Prime B parallel to OB as found in this figure;  
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$OC \perp OB$  and  $O'B$  (both)  
 and  $OD \perp OA$  and  $O'A$ ,  
 the above eqn gives  
 $O'C + O'D = 1.22\lambda_0$   
 $2S \sin \alpha = 1.22\lambda_0$   

$$S = \frac{1.22\lambda_0}{2 \sin \alpha}$$

where  $s$  is the linear distance  $OO'$  between the point objects and  $\alpha$  is the half the angle subtended at the axial point object  $O$  by the rim of the microscope objective and  $\lambda_0$  is the wavelength of light in vacuum. It would be observed that the entire difference in paths under consideration, that is  $(O'B - O'A)$ , is in the medium between object and the objective.



Therefore, drawing the line OC perpendicular to OB and O prime and OD perpendicular to both OA and O prime A, the above equation gives O Prime C + O Prime D = 1.22 Lambda. Therefore, we can we have  $2S \sin \alpha = 1.22 \text{ lambda}$ , so  $S$  will be =  $1.22 \text{ lambda}$  upon  $2 \sin \alpha$ , where  $S$  is the linear distance  $O, O$  prime between the point objects and  $\text{Alpha}$  is the half the angle subtended at the axial point object  $O$  by the rim of the microscope objective and  $\text{Lambda}$  is the wavelength of light in vacuum.

It would be observed that the entire difference and path under consideration that is O prime B - O prime A is in the medium between objects and the objective.  
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The above equation gives the linear distance between the just resolvable objects. The product of the index of refraction of medium in which the object is situated and the sine of half the angle of cone of rays admitted by the objective. i.e,  $\mu \sin \alpha$  is called the numerical aperture of the objective. Thus decreasing the wavelength and increasing the numerical aperture decrease  $s$ , i.e. increases the resolving power of the microscope. For air, the upper limit of numerical aperture of the microscope objective is about 0.95. Therefore with white light of effective wavelength  $5600 \times 10^{-8} \text{cm}$ , the least resolution distance in air is

$$S_{\text{air}} = \frac{1.22 \times 5600 \times 10^{-8}}{2 \times 0.95} = 3.6 \times 10^{-5} \text{cm}$$

The above equation gives the linear distance  $S$  between the just resolvable objects. The product of the index of refraction of medium in which the object is situated and the sine of half the angle of cone of rays admitted by the objective that is  $\mu \sin \alpha$  is called the numerical aperture of the objective, thus decreasing the wavelength and increasing the numerical aperture decrease  $S$  that is increases the resolving power of the microscope. For air, the upper limit of numerical aperture of the microscope objective is about 0.95.

Therefore, with white light of effective wavelength 5600 angstroms, the least resolution distance in air is of the order of  $3.6 \times 10^{-5} \text{cm}$ . But by filling the space between the object and the objective by air, the numerical aperture may be increased to 1.6.  
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But by filling the space between the object and the objective by oil, the numerical aperture may be increased to 1.6. In this case the least resolvable distance becomes

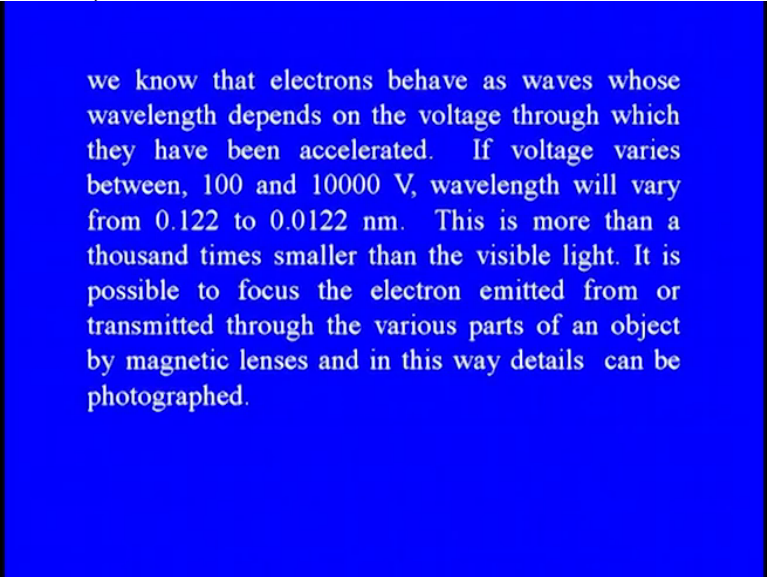
$$S_{\text{oil}} = \frac{1.22 \times 5600 \times 10^{-8}}{2 \times 1.6} = 2.14 \times 10^{-5} \text{cm}$$

When space between the object and objective is filled with oil then it is called oil immersion objectives. By using ultraviolet light in place of visible light resolving power can be further increased. One of the most remarkable steps in the improvement of microscopic resolution has been the development of electron microscope. From de Broglie concept,

In this case, the least resolvable distance  $S$  becomes  $= 2.14 \times 10^{-5} \text{ cm}$  when space between the object and objective is filled with oil then, it is called oil immersion objectives. So, by using oil immersion objective, we can increase the resolving power of microscope and by using ultraviolet light, in place of visible light, resolving power can be further increased.

One of the most remarkable steps in the improvement of microscopic resolution has been the development of electron microscope. From de Broglie concept, we know that electrons behave as wave, when wavelength where waves wavelength depends on the voltage through which they have been accelerated. A voltage where is between  $10^2$  to  $10^4$  volt wavelength will vary from  $0.122$  to  $0.0122$  nanometre.

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we know that electrons behave as waves whose wavelength depends on the voltage through which they have been accelerated. If voltage varies between,  $10^2$  and  $10^4$  V, wavelength will vary from  $0.122$  to  $0.0122$  nm. This is more than a thousand times smaller than the visible light. It is possible to focus the electron emitted from or transmitted through the various parts of an object by magnetic lenses and in this way details can be photographed.

This is more than a 1000 times smaller than the visible light. It is possible to focus the electron emitted from or transmitted through the various parts of an object, by magnetic lenses and in this way, details can be photographed. And in case of electron microscope since the wavelength of electron beam is much smaller than the visible light, so, we expect very high value of resolving power for electron microscope. Now let us discuss the resolving power of a telescope.

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### Resolving power of a Telescope:

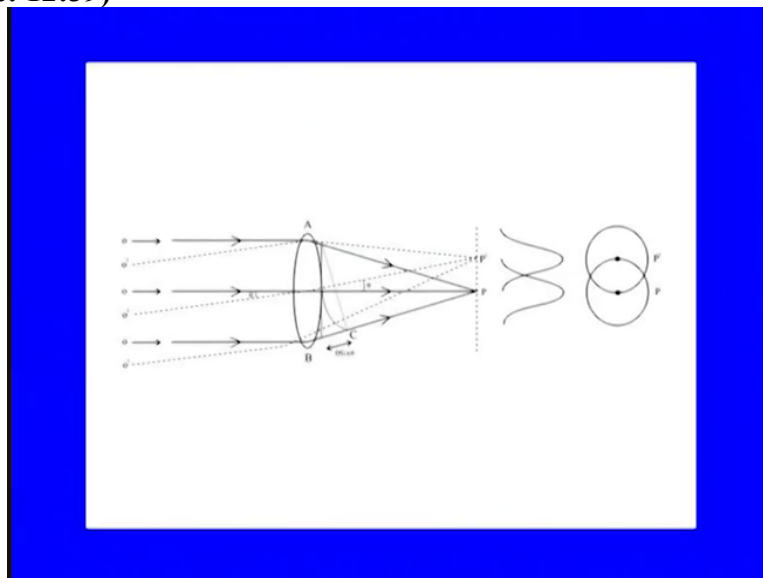
We know that the telescope is employed to view distant object and therefore the amount of detail which it reveals depends on the angle the two point objects subtend at the objective rather than on the linear separation between them. The resolving power of a telescope is therefore defined as the inverse of the least angle subtended at the objective by the distant point objects which can be just distinguished as separate in its focal plane.

Let  $O$  be a monochromatic point source of light of wavelength  $\lambda$

The telescope is employed to view distant object and therefore the amount of detail which it reveals depends on the angle the two point objects subtend at the objective rather than on the linear separation between them. The resolving power of a telescope is therefore defined as the inverse of the least angle subtended at the objective, by the distant point objects which can be just distinguished as separate in its focal plane.

Let  $O$  be a monochromatic point source of light of wavelength  $\lambda$  situated at a great distance from the telescope objective  $AV$  as shown in this figure.

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A beam of parallel rays from this object is incident normally on the objective. The ring supporting the lens as well as lens itself serves as a circular aperture and therefore from Fraunhofer diffraction pattern is produced in the focal plane of the lens. The refraction pattern of

O consists of a central bright spot surrounded by concentric dark and bright Rings, the common Centre being at the point P, where the point image would have occurred in the absence of diffraction phenomenon.

The image of a point object O prime situated close to A is also a similar diffraction pattern with its centre at the point P prime. These diffraction patterns overlap and two point objects can be just resolved when according to Rayleigh's criterion the separation PP Prime, between the centres of 2 Central Maxima is at least = the radius of either one. Under this condition, the angle subtended by 2 point objects and the objective is taken as the measure of resolving power of a telescope.

Actually, resolving power is inverse of this angle. Now, with P prime as centre and radius P prime A, we draw an arc cutting C prime B at the point C as shown in the figure. This arc AC therefore, represents the transmitted wave front which gives rise to the central maximum at P Prime, in the diffraction pattern of O prime. The second point object near to the first fastest resolution of O and O prime, the dark ring in the diffraction pattern of O should pass through the point P prime.

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Therefore AC should also be a diffraction wavefront for O, so as to form in the diffraction pattern the first dark ring passing through P'. Hence, the path difference (BP' - AP') between the extreme diffracted rays reaching P' according to theory of Fraunhofer diffraction at a circular aperture, should be equal to  $1.22\lambda$ . Thus we have

$$BP' - AP' = BC = 1.22\lambda$$

$$\text{Hence } \angle BAC = \theta = \frac{BC}{AB} = \frac{1.22\lambda}{D}$$

Therefore, AC should also be a diffraction wavefront for O, so as to form in the diffraction pattern, the first dark Line passing through P prime. Hence the path difference BP Prime - AP prime between the extreme different rays reaching p prime according to theory of Fraunhofer diffraction at a circular aperture should be =  $1.22\lambda$ . Thus, we have BP prime - AP prime =



BC, which is  $= 1.22 \lambda$ . Hence angle BAC =  $\theta$  will be  $= \frac{1.22 \lambda}{D}$  upon capital D.

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Where D is the diameter of the objective. Since rays QP' and QP are normal to the respective wave fronts which form central maxima at P' and P, hence  $\angle PQP' = \angle BAC$ .

Here  $\angle PQP'$  is the angle subtended at the objective by the two point objects. Hence the minimum angular separation  $\theta$ , which two point object should have so that they can be just resolved is given by the relation

$$\theta = \frac{1.22 \lambda}{D}$$

Where capital D is the diameter of the objective; Since rays QP prime and QP are normal to the respective wave fronts which form central maxima at P prime and P hence, angle PQP prime will be = Angle BAC. If angle PQP prime is the angle subtended at the objective by the two point objects. Hence the minimum angular separation which two point object should have so that they can be just resolved is given by the relation  $\theta = \frac{1.22 \lambda}{D}$ . This angle is therefore the measure of angular resolving power of the telescope objective.

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This angle is, therefore the measure of angular resolving power of the telescope objective. The smaller the value of  $\theta$ , the greater is said to be resolving power. Therefore, telescope with larger objective have higher resolving power.

We may now calculate the value of limiting angle  $\theta$  which the two distant stars should subtend on the objective of one inch diameter so as to be just resolved by it.

The smaller the value of  $\theta$ , the greater is said to be resolving power. Therefore, telescope with larger objective and higher resolving power. We may now calculate the value of limiting angle

theta with the two distant stars should subtend on the objective of 1 inch diameter so as to be just resolve by it. Now putting the value of Lambda and capital D.

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$$\begin{aligned}\theta &= \frac{1.22 \times 5500 \times 10^{-8} \text{ cm}}{2.54} \text{ radian} && \text{Take } 1'' = 2.54 \text{ cm} \\ &= \frac{180}{\pi} \times \frac{1.22 \times 5500 \times 10^{-8}}{2.54} \text{ degree of arc} \\ &= 5.52 \text{ seconds of arc}\end{aligned}$$

Thus, the limiting angle  $\theta$  for 1" telescope is of the order 5 sec. Therefore, a telescope with a 200" objective will resolve them even if their separation be only  $\frac{1}{40}$ th of a second of arc. So it is obvious that the larger the diameter of the objective, the better will be its resolving power.

We get theta = 5.52 seconds of arc. Thus, the limiting angle theta for 1 inch telescope is of the order of 5 second. Therefore a telescope with a 200 inch objective will resolve them even if their separation be only one by 40th of a second of arc. So, it is obvious that the larger the diameter of the objective, the better will be its resolving power. Now, let us take some numerical problems.

Problem number 1:

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1. Find the separation of two points on the moon's surface that can be resolved by the 508 cm telescope, assuming that this distance is determined by diffraction effects. The distance from the earth to the moon is  $3.84 \times 10^5$  km.

Solution:

Given  $D = 508 \text{ cm} = 5.08 \text{ m}$

Distance from earth to moon =  $3.84 \times 10^5 \text{ km}$   
 $= 3.85 \times 10^8 \text{ m}$

If  $r$  is the minimum resolvable distance between the points, then corresponding angle formed by the points at telescope objective will be,

Find the Separation of two points on the moon surface that can be resolved by the 508 cm telescope, assuming that this distance is determined by diffraction effects. The distance from the earth to the moon is about 3.84 into 10 to the power 5 km. Now, in this problem, the diameter of

Now, if  $r$  is the minimum resolvable distance between the points then corresponding angle formed by the points at telescope objective will be  
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So this is the minimum resolvable distance between the two points.

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This problem is similar to the previous one.

Suppose a spy satellite circles the earth at an altitude of 150 km and is fitted with a camera having a lens diameter of 35cm. If diffraction limit the resolution, then estimate the least separation of two small object on the earth here we can take  $\lambda = 5500$  angstrom. In this problem, the distance between telescope objective and object is given; diameter of objective is also given.

So, like previous problem this in this problem also we can find the separation between the two objects. So, thus we have finished the discussion related to resolving power of microscope and satellite and we have taken, we have discussed some numerical example also. Now, I will discuss some problem related to refraction of single slit, double slit and grating.

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#### **Problems on Diffraction of light**

1. Suppose parallel light of wavelength  $6463 \text{ \AA}$  is incident normally on a slit of width  $0.3850 \text{ mm}$ . A lens with a focal length of  $50.0 \text{ cm}$  is located just behind the slit bringing the diffraction pattern to focus on a screen. What is the distance from the centre of the principal maximum to the first minimum and the fifth minimum?

Solution: Given slit width  $b = 0.3850 \text{ mm}$

Wavelength  $\lambda = 6463 \text{ \AA}$

If the lens is close to slit, then the separation between slit and screen will be  $\sim$  focal length of lens, i.e.  $50.0 \text{ cm}$ .

Problem number 1: Suppose parallel light of wavelength  $6463$  angstrom is incident normally on a slit of width  $0.3850$  millimetres with a focal length of  $50 \text{ cm}$  is located just behind the slit bringing the diffraction pattern to the focus on a screen. What is the distance from the centre of the principal maximum to the first minimum and 5th minimum? In this problem, slit width  $b$  is = zero point 3850 millimetre and the wavelength of light  $\lambda = 6463$  angstrom.

If the lens is close to slit then, the separation between the slits and the screen will be = focal length of the lens that is  $50 \text{ cm}$ .

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For single slit the condition for minima is  $b \sin \theta = n\lambda$

So for first minimum  $b \sin \theta = \lambda$

Therefore,  $\sin \theta = \lambda/b$

and if  $\theta$  is small,  $\sin \theta \approx \tan \theta$

If  $r_1$  is the distance between centre maximum and first minimum, then  $r_1 = f \times \lambda/b$

Putting the values of  $f$ ,  $\lambda$  and  $b$ , we get

$r_1 = 0.8523 \text{ mm}$ .

We have for fifth minimum  $b \sin \theta = 5 \lambda$ .

By using this equation, the distance between Principal maximum and fifth minimum  $r_5$  comes out to be  $4.2615 \text{ mm}$ .

For single slit the condition for minima is  $b \sin \theta = n\lambda$ . So, for first minimum  $b \sin \theta = \lambda$ . Therefore  $\sin \theta = \lambda/b$ . Now, if  $\theta$  is small then,  $\sin \theta \approx \tan \theta$ . If  $r_1$  is the distance between Central maximum and first minimum then  $r_1$  will be  $= f \times \lambda/b$ . Putting the value of focal length  $f$ ,  $\lambda$  and  $b$  in this equation, we get,  $r_1 = 0.8523 \text{ mm}$ .

For 5th minimum  $b \sin \theta = 5\lambda$  and by using this equation the distance between principal maximum and 5th minimum,  $r_5$  comes out to be  $4.2615 \text{ mm}$ .

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2. Suppose plane waves of light of wavelength  $4340 \text{ \AA}$ , fall on a single slit, then pass through a lens with focal length of  $85.0 \text{ cm}$ . If the central band of the diffraction pattern on the screen has a width of  $2.450 \text{ mm}$ , find the width of the slit.

Solution: In this problem the width of the central maximum is given. This width will be twice of the distance between the principal maximum position and the first minimum position, i.e.  $2r_1$ . So here also we have to proceed like Q. 1.

Slit width  $b = 2\lambda f / 2r_1$ , substituting the values of  $\lambda$ ,  $f$  and  $2r_1$ , we get  $b = 0.3011 \text{ mm}$ .

Problem number 2: Suppose plane waves of light of wavelength  $4340 \text{ \AA}$  fall on a single slit then pass through a lens with focal length of  $85 \text{ cm}$  if the central band of the diffraction pattern on the screen has a width of  $2.450 \text{ mm}$ , find the width of the slit. Now in this problem,



the width of the central maximum is given. This width will be twice of the distance between the principal maximum position and the first minimum position that is 2 times  $r_1$ .

So, here also we have to proceed like question number 1. So, slit width  $b$  is  $= 2 \lambda f$  divided by  $2 r_1$  where  $f$  is the focal length of the lens substituting the value of  $\lambda f$  and  $2 r_1$  we get  $b = 0.3011$  millimetre. That is the width of the slit is 0.3011 millimeter.

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3. Suppose parallel beam of light is incident on two slit arrangement. The width of each slit is 0.140 mm and the distance between the centers of two slits is 0.840 mm. Find the missing orders and approximate intensity of interference fringes of order 0 to 6.

Solution: Given  $b = 0.140$  mm and  $(a+b) = 0.840$  mm

$$n/m = (a+b)/b = 6$$

$$\text{i.e. } n = 6m, m = 1, 2, 3, \dots$$

Therefore, missing orders are: 6, 12, 18, 24, --

We know that the intensity  $I$  of double slit diffraction pattern is

$$4I_0(\sin^2\beta/\beta^2)\cos^2\gamma, \text{ where } \beta = (\pi\lambda)b\sin\theta$$

$$\text{and } \gamma = (\pi\lambda)(a+b)\sin\theta.$$

Problem number 3: Suppose, parallel beam of light is incident on two slit arrangement. The width of a slit is 0.140 millimetre and the distance between the centres of two slit is 0.840 millimetre. Find the missing Orders and their approximate intensity of interference fringes of order 0 to 6. In this problem, slit width  $b$  is  $= 0.140$  millimetre and  $a + b = 0.840$  millimetre, so we can find out  $n$  by  $m$  which is  $a + b$  by  $b = 6$ .

That is  $n$  is  $= 6$  time  $m$  where  $m$  is  $= 1, 2, 3$  and so on. Therefore missing order are: 6, 12, 18, 24 and so on. We know that the intensity  $I$  of the double slit diffraction pattern is  $4I_0 \sin^2\beta/\beta^2 \cos^2\gamma$ , where  $\beta$  is  $= \pi \lambda b \sin\theta$  and  $\gamma$  is  $= \pi \lambda (a + b) \sin\theta$ . We have already discussed about these expressions while discussing the diffraction pattern of double slit.

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We know that for maxima  $(a+b)\sin\theta = m\lambda$ , where  $m = 0, 1, 2, 3, \dots$ .

From this equation, we can find the values of  $\sin\theta$  for zero, first, second, ---- order interference maxima. For these values of  $\sin\theta$ ,  $\cos^2\gamma = 1$ . So the intensity in these direction is decided by the values of  $4I_0(\sin^2\beta/\beta^2)$ . So after determining the values of  $\beta$  in the desired direction we can calculate the values of  $4I_0(\sin^2\beta/\beta^2)$  to get the intensity of interference maxima for  $m = 0, 1, 2, 3, 4, 5$  and  $6$ .

For  $m = 0$ ,  $I = 4I_0$ , for  $m = 1$ ,  $I_1 = 0.91I$ ,  $I_2 = 0.68I$ ,  $I_3 = 0.40I$ ,  $I_4 = 0.17I$ ,  $I_5 = 0.036I$  and  $I_6 = 0$ .

And we know that for maxima  $a + b \sin \theta = m \lambda$  where  $m$  is  $= 0, 1, 2, 3$  and so on, from this equation, we can find the values of  $\sin \theta$  for zero, first, second, third, so on, order in reference maxima. For these values of  $\sin \theta$ ,  $\cos^2 \gamma$  will be  $1$ , so, intensity in this direction is decided by the values of  $4I \sin^2 \beta / \beta^2$ . So, after determining the values of  $\beta$  in the desired direction, we can calculate the values of  $4I \sin^2 \beta / \beta^2$ , to get the intensity of interference maxima for  $m$  is  $= 0, 1, 2, 3, 4, 5$  and  $6$ .

For  $m$  is  $= 0$   $I$  will be  $= 4I$  naught, for  $m$  is  $= 1$ ,  $I_1$  will be  $= 0.91$  times  $I$ , that is intensity will be about  $90\%$  off Central Maxima. Similarly,  $I_2$  will be  $= 0.68I$ ,  $I_3$  will be  $= 0.40I$ ,  $I_4$  will be  $= 0.17I$ ,  $I_5$  is  $= 0.036 I$  and  $I_6$  will be  $0$ . Thus, we see that the intensity of interference maxima is gradually decreasing as order increases.

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4. Light of two wavelengths 5600 Å and 5650 Å, fall normally on a plane transmission grating having 2500 lines per centimeter. The emerging parallel light is focused on a flat screen by a lens of 120 cm focal length. Find the distance on the screen in centimeters between the two spectrum lines in the first order and in the second order.

Solution:

From the given data we can find the value of grating element (a+b).

$$(a+b) = 1/2500 \text{ cm} = 4.0 \times 10^{-4} \text{ cm}$$

We know that angular dispersive power

$$d\theta/d\lambda = n/(a+b)\cos\theta$$

Question number 4: light of two wavelengths 5600 angstrom and 5650 angstrom fall normally on a plane transmission grating having 2500 lines per centimetre. The emerging parallel light is focused on a flat screen by a lens of 120 CM focal length. Find the distance on the screen in cm between the two spectrum lines in the first order and in the second order. In this problem, first we have to find out the grating element a + b which will be = 1 upon 2500 CM which will be = 4.0 into 10 to the power - 4 cm. We know that angular dispersive power d theta upon d lambda is = n divided by a + b cos theta.

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If f is the focal length of the lens, then linear dispersive power

$$dl/d\lambda = f n / [(a+b)\cos\theta]$$

$$\text{linear separation } dl = (f n d\lambda) / [(a+b)\cos\theta]$$

Given  $d\lambda = 50\text{Å}$ .

For first order diffraction

$$\sin\theta_1 = [5600 \times 10^{-8}] / [4.0 \times 10^{-4}] = 0.14$$

$$\cos\theta_1 = 0.99$$

$$dl \approx 0.15 \text{ cm}$$

For second order:  $\sin\theta_2 = 0.28$ , so  $\cos\theta_2 =$

In second order this separation will be  $\approx 0.30 \text{ cm}$

If f is Focal length of the lens then linear dispersive power  $dl$  by  $d\lambda$  will be =  $f n$  divided by  $a + b \cos\theta$ . From this equation, we can calculate the linear separation  $dl = f$  into  $n d\lambda$  divided by  $a + b \cos\theta$ . In the problem,  $d\lambda$  is = 50 angstrom. So, for first order

diffraction, we can calculate what is the value of  $\sin \theta_1$  and after knowing the value of  $\sin \theta_1$ , we can calculate what is the  $\cos \theta_1$ .

And then, we can find out the value of  $d$ . So, in the first order value of  $d$  is approximately 0.15 cm. Similarly, for the second order, we can calculate what is the value of  $\sin \theta_2$ , then, we can find out the corresponding value of  $\cos \theta_2$ . And after substituting the value in the above equation we can find the value of  $d$  here. So,  $d$  comes out to be 0.30 cm.

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5. A diffraction grating used at normal incidence gives a green line  $\lambda_1 = 5400 \text{ \AA}$  in a certain order superimposed on the violet line  $\lambda_2 = 4500 \text{ \AA}$  of next higher order. If the angle of diffraction is  $30^\circ$ , how many ruling are there per centimeter in the Grating.

Solution:  $(a+b)\sin\theta = n\lambda_1 = (n+1)\lambda_2$

using the values of  $\lambda_1$  and  $\lambda_2$

we get  $n = 3$ . Since angle of diffraction is given, we can find  $(a+b) = 3.24 \times 10^{-4} \text{ cm}$ .

Therefore, no. of lines per cm =  $1/(a+b)$   
 $= 3086 \text{ lines cm}^{-1}$

Problem Number 5: A diffraction grating used at normal incidence gives a green line  $\lambda_1 = 5400 \text{ \AA}$  in a certain order superimposed on the violet line  $\lambda_2 = 4500 \text{ \AA}$  of next higher order. If the angle of diffraction is  $30^\circ$ , how many ruling are there per centimetre in the grating? So, according to the problem here,  $a + b \sin \theta$  will be  $= n\lambda_1$  which will be  $= (n+1)\lambda_2$ .

Now putting the value of  $\lambda_1$  and  $\lambda_2$  we get,  $n = 3$  here, since the angle of diffraction is given, we can find  $a + b = 3.24 \times 10^{-4} \text{ cm}$ . With the value of  $a + b$ , we can find the number of lines per centimetre is comes out to be 3086 lines per centimetre. So, here we have solved some problem related to single slit, double slit and grating.

Some more problems can be found in any standard textbook of optics. Thus, we have finished the discussion related to refraction of light and here we have mainly discussed the Fraunhofer diffraction pattern of single plate, double plate and grating. I also discussed how by using grating

we can find out the wavelength of light in the laboratory. And I also discussed how the diffraction effect, limit in the resolving power of certain institute, optical instruments.

And explained the Rayleigh criteria for just a resolution and by using this criteria I derive the expression for resolving power of grating microscope and telescope. Thank you.