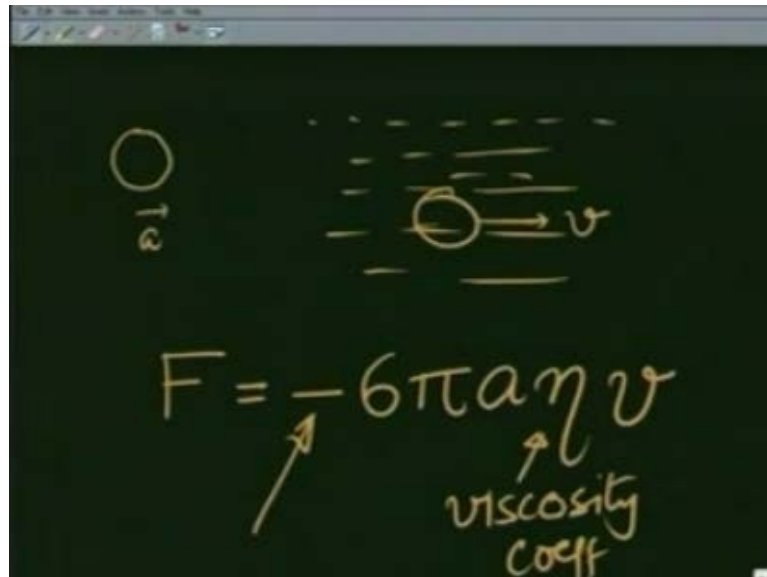


Engineering Mechanics
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Module - 05
Lecture - 04
Motion of Particles with Drag

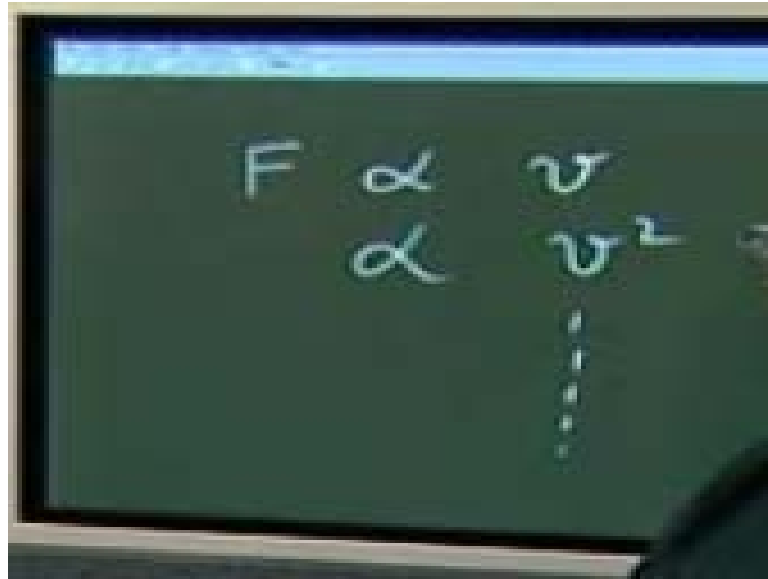
In the previous lecture, we saw how two solids when they come in contact apply frictional force on each other, and how that affects their motion. In this lecture we are going to look at another kind of frictional force, when a body moves through a fluid or gas. You are well familiar from Stokes's law that.

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When a spherical body of radius a , moves through a gas or fluid at speed v , it is moving with speed v then, there is a force on a due to viscosity which is F equals minus six pi a eta v. This minus sign denotes that, the force is opposite to its direction of motion, eta is the viscosity coefficient. This is 1 example of the drag force that anybody moving through a fluid or gas experiences. Of course, this is the lowest order force.

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We may have forces which are proportional to v , which are proportional to v square or higher powers. These forces oppose the motion, and in this lecture we are going to look at how to deal with these forces when they are included in the motion of a particle.

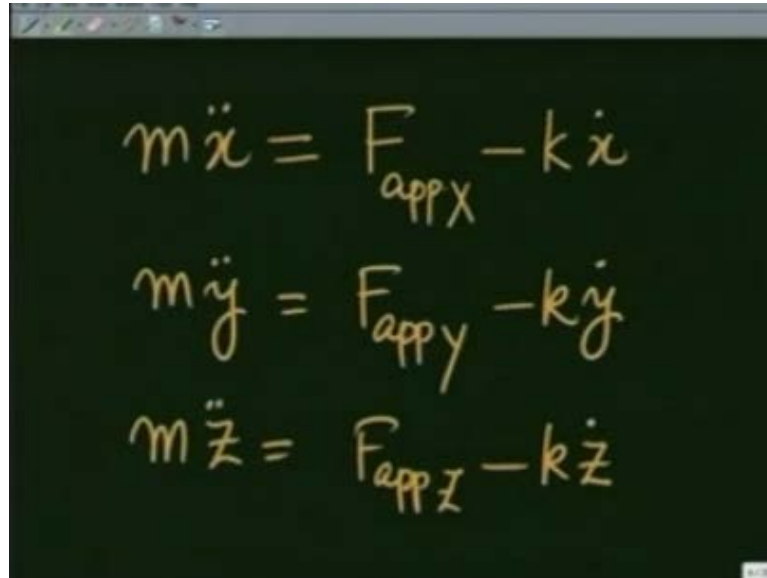
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$\vec{F}_{\text{drag}} = -k\vec{v}$
 $m\ddot{\vec{r}} = \vec{F}_{\text{applied}} - k\vec{v}$
Drag force

The simplest one as I said earlier is the force which is proportional to v , and I am going to write this as F equals minus kv . And I am going to call it drag force, k is the coefficient of drag, minus sign again denotes that, the force is opposite to the motion.

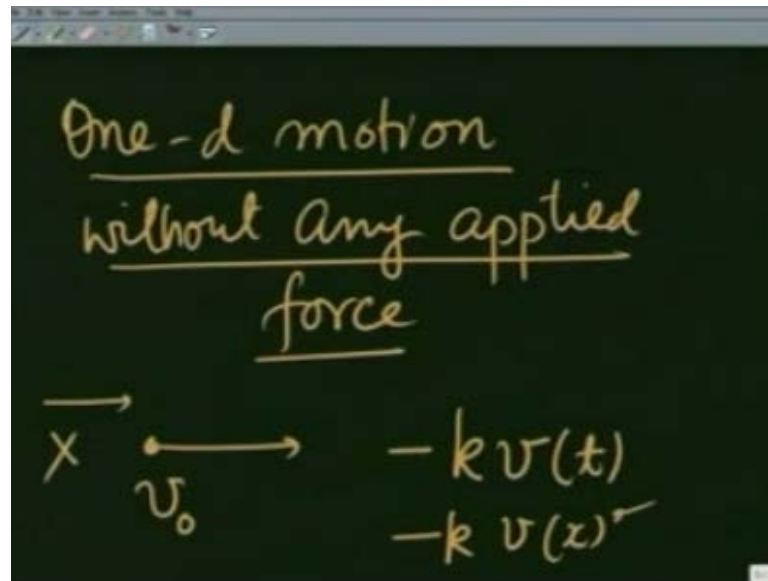
And therefore, the equation of motion is going to be $m\ddot{x}$, mass times acceleration equals F applied that, the force that I apply from outside, minus kx , which is the drag force.

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$$\begin{aligned}m\ddot{x} &= F_{\text{app}x} - kx \\m\ddot{y} &= F_{\text{app}y} - ky \\m\ddot{z} &= F_{\text{app}z} - kz\end{aligned}$$

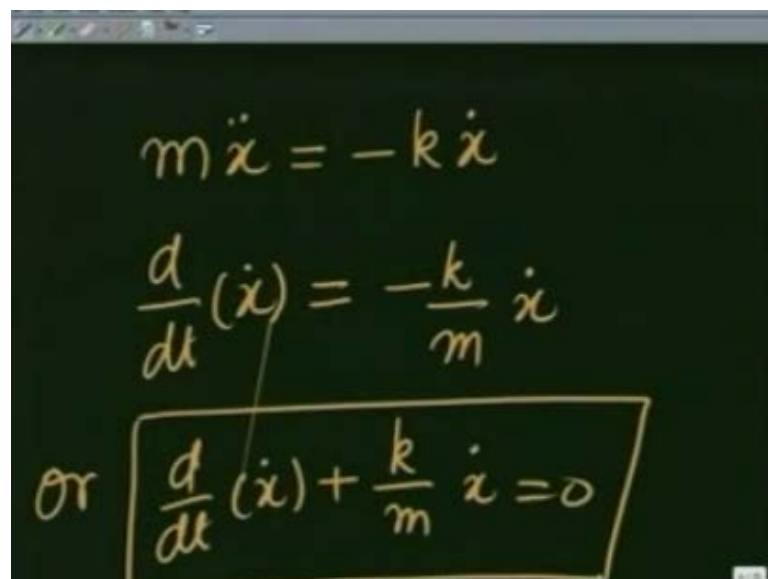
If I write these in component form I get $m\ddot{x}$ equals F applied x component, minus kx , $m\ddot{y}$ equals F applied y component, minus ky and similarly, $m\ddot{z}$ equals F applied z, minus kz . This is the simplest of drag forces and in the next few examples, we are going to see how to incorporate this in the motion.

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The first example I am going to take is going to be 1 dimensional motion without any applied force. There is no applied force, but drag is there. Therefore, initially I have to throw the particle with some initial speed v_0 , and as it moves along it experiences a force which is minus kv , v is changing with time or I can also write v is changing with the distance. Let this direction be x direction that is I have written x here.

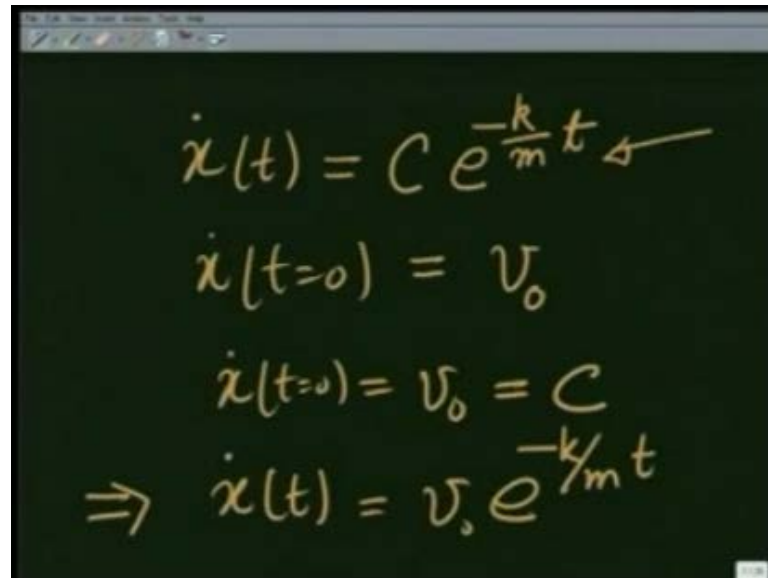
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The equation of motion for this is going to be $m\ddot{x} = -k\dot{x}$. In more familiar terms, it is going to be $\frac{d}{dt}(\dot{x}) + \frac{k}{m}\dot{x} = 0$ which is nothing but $\ddot{x} + \frac{k}{m}\dot{x} = 0$.

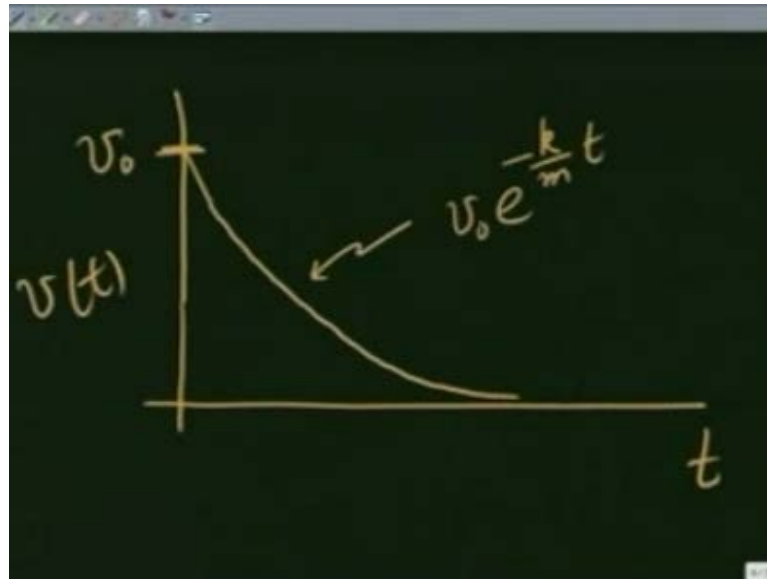
is equal to minus k over m times \dot{x} or $\frac{d}{dt}$ of \dot{x} , plus k over m times \dot{x} equal to 0 . This is my equation of motion that, I want to solve for this particle. You can check by direct substitution that the solution is going to be.

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$$\begin{aligned}\dot{x}(t) &= C e^{-\frac{k}{m}t} \\ \dot{x}(t=0) &= v_0 \\ \dot{x}(t=0) = v_0 &= C \\ \Rightarrow \dot{x}(t) &= v_0 e^{-\frac{k}{m}t}\end{aligned}$$

\dot{x} as a function of time equal to some constant C e raise to minus k over m t , what is C ? Let us write \dot{x} at t is equal to 0 , that is initial time, through the particle with speed v_0 . And therefore, I must have \dot{x} at t equal to 0 , is equal to v_0 , and when I substitute t equal to 0 here, I get this, this is equal to C . And therefore, my general solution is going to be \dot{x} , that is the velocity of the particle equal to initial speed, initial velocity minus k over m t . Let us plot it and see, how does it look?

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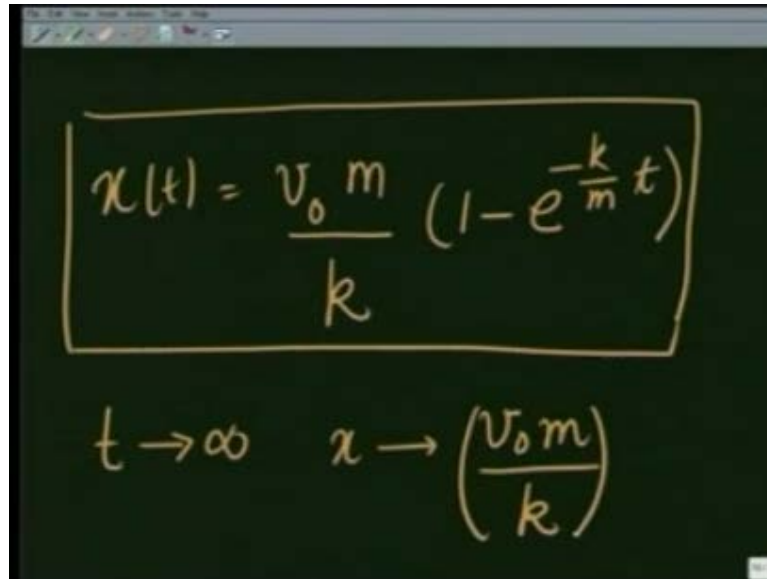
So, initially the particle started moving with some speed v_0 , I am plotting v versus t , and as time passed, velocity decreases exponentially. This curve is $v_0 e^{-k/m t}$. How about the distance travelled by the particle? That is also easy to calculate once I know the velocity.

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$$\begin{aligned} \dot{x}(t) &= v_0 e^{-\frac{k}{m}t} \\ x(t) &= \int_0^t \dot{x}(t') dt' \\ &= v_0 \int_0^t e^{-\frac{k}{m}t'} dt' \end{aligned}$$

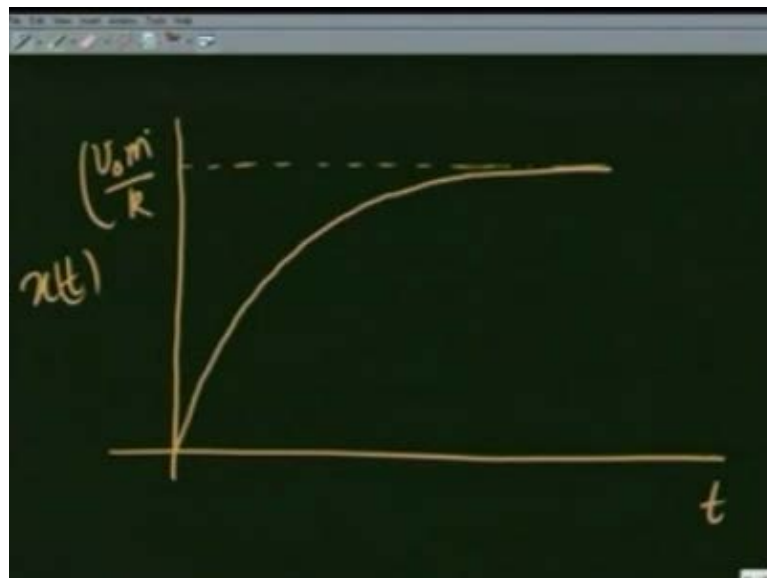
\dot{x} is given to be $v_0 e^{-k/m t}$. So, the distance travelled in time t , is going to be $\int_0^t \dot{x}(t') dt'$ which is equal to $v_0 \int_0^t e^{-k/m t'} dt'$.

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$$x(t) = \frac{v_0 m}{k} \left(1 - e^{-\frac{k}{m} t}\right)$$
$$t \rightarrow \infty \quad x \rightarrow \left(\frac{v_0 m}{k}\right)$$

After integration I get $x(t)$ is equal to $v_0 m$ over k $1 - e^{-\frac{k}{m} t}$. So, I have also calculated, how much is the distance travelled in time t . As t goes to infinity, x goes to $v_0 m$ over k . So, that is the distance that the particle will finally, travel and stop after this. Let us see how does the plot of this look.

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If I plot $x(t)$ as a function of time. Initially obviously, distance travelled is 0, and slowly it moves up and goes like this, distance is being $v_0 m$ over k .

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$$x(t) = \frac{v_0 m}{k} \left(1 - e^{-\frac{k}{m} t}\right)$$

Let $k \rightarrow 0$ $x(t) = v_0 t$

~~$\left(\frac{0}{0}\right)$~~ $x = \frac{v_0 m}{k} \left(1 - 1 + \frac{k}{m} t\right) = v_0 t$

So, let us look at this expression once more. Just $v_0 m$ over k one, minus e raise to minus k over m t . And ask if this is a correct answer, 1 way to check this is, let k go to 0 , and I should get my familiar x equals $v_0 t$, but if I substitute k directly here, I get an answer like 0 over 0 which is not correct. I have to be more careful and take the limit, k going to 0 keeping terms up to order k in this expression, and see if I get the correct answer. If I do that, I find x is equal to $v_0 m$ over k 1 minus, 1 plus k over m t , this cancels and I get $v_0 t$, which is correct. So, this is 1 way, when you get an expression in these, when you apply these forces, you can check your answer.

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(2)

$\bullet \rightarrow F$, $-kv$
 $-kx$

$m\ddot{x} = F - kx$

$v(t=0) = 0 \checkmark$

Next let us consider a particle still moving in 1 dimension, but applied on it is a constant force. So, I have a particle, I apply a constant force F again an x direction, and let there be a drag force minus kv or minus $k \dot{x}$. And I want to solve how, I want see how this particle moves. The equation of motion is going to be because $m \ddot{x}$ is equal to F minus $k \dot{x}$.

In this case notice that I do not really have to start by shooting a particle with some initial speed. I can start with v at $t = 0$ still 0 because I am applying a force, the particle will start moving and this is what I want to see how the particle moves. So, this is going to be 1 of my initial conditions. The other 1 is going to be, let us also say that x at $t = 0$ it starts from the origin.

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$$\ddot{x} = \frac{F}{m} - \frac{k}{m} \dot{x}$$

$$\frac{d}{dt}(\dot{x}) + \frac{k}{m} \dot{x} = \frac{F}{m}$$

$$F=0$$

If I rewrite the equation will come out in the form of \ddot{x} is equal to F over m , minus k over $m \dot{x}$, or $\frac{d}{dt}$ of \dot{x} , plus k over m of \dot{x} is equal to F over m . Notice, I have already solved this problem for F equals 0, that was 1 dimensional motion without any force. In such situations, where I have 1 part like this and another part like this, I can write my solution as the previously obtained solution with F equals 0, which was nothing but $C e^{-kt/m}$.

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The image shows a chalkboard with handwritten mathematical expressions. At the top, the general solution is written as $\dot{x}(t) = C e^{-k/m t} + \text{particular solution}$. A bracket under the exponential term has an arrow pointing down to the differential equation $\frac{d}{dt}(\dot{x}) + \frac{k}{m} \dot{x} = \frac{F}{m}$. The term $\frac{F}{m}$ on the right side of the equation is circled. Below the equation, there is a double line under the letter 'x' with an arrow pointing to the expression $\dot{x} = F/k$.

This was \dot{x} of t plus a particular solution, that corresponds to the term on the right hand side. Remember I am solving the equation $\frac{d}{dt} \dot{x}$, plus $\frac{k}{m} \dot{x}$, is equal to $\frac{F}{m}$. This is a solution corresponding to this part equal to 0 the particular solution would refer to this term. And you can see, the particular solution is going to be \dot{x} is equal to F/k . I substitute this here and I get on the right hand side F/m .

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The image shows a chalkboard with handwritten mathematical expressions. The general solution is written as $\dot{x}(t) = C e^{-k/m t} + \frac{F}{k}$. Below this, the initial condition $\dot{x}(t=0) = 0$ is written. An arrow points from the 'C' in the general solution to the equation $\Rightarrow C = -F/k$. Finally, the specific solution is written as $\dot{x}(t) = \frac{F}{k} (1 - e^{-k/m t})$.

And therefore, my general solution for \dot{x} is going to be \dot{x} of t is equal to $C e$ raise to minus $k/m t$, plus F/k . This constant C is again determined by initial

condition, which is nothing but initially a particle was at rest. And therefore, I get C equals minus F over k, and the general solution therefore, comes out to be $x \dot{t}$, other solution with this initial condition comes out to be $x \dot{t}$, equals F over k $1 - e^{-k/m t}$. Let us now look at the solution carefully.

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The image shows a chalkboard with the following handwritten content:

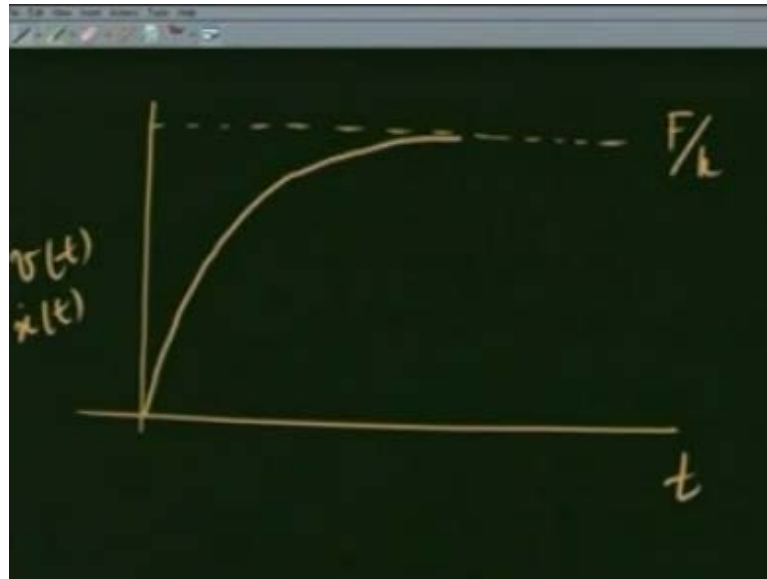
$$\dot{x}(t) = \frac{F}{k} (1 - e^{-k/m t})$$

$t \rightarrow \infty \quad \dot{x} \rightarrow \frac{F}{k}$
terminal speed

$$F = k v_{\text{terminal}}$$

So, I have $x \dot{t}$ equals F over k $1 - e^{-k/m t}$. We notice as t goes to infinity, that is after a long time I have applied the force $x \dot{t}$ goes to F over k . This is nothing but the terminal speed, at this speed the applied force F becomes equal to k times v_{terminal} . The two forces balance and therefore, the particle moves with a constant speed of v_{terminal} .

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If I plot this speed I am going to get vt or $x \dot{t}$ same thing versus t , the final speed is F over k and the speed slowly builds up. How about the distance travelled? Well for that, I again integrate it.

(Refer Slide Time: 16:14)

$$\begin{aligned}\dot{x}(t) &= \frac{dx}{dt} = \frac{F}{k} (1 - e^{-k/m t}) \\ x(t) &= \int_0^t \frac{F}{k} (1 - e^{-k/m t'}) dt' \\ &= \frac{F}{k} \left[t - \frac{m}{k} \{1 - e^{-\frac{k}{m} t}\} \right]\end{aligned}$$

So, I have $x \dot{t}$, equals $dx dt$ is equal to F over k $1 - e$ raise to minus k over $m t$. So, x at t is again going to be 0 to t , F over k $1 - e$ raise to minus k over $m t$ prime, dt prime integrating, I get F over k , t minus m over k , $1 - e$ raise to minus k over $m t$. So, this is going to be the distance travelled by the particle. Now, that we have gotten

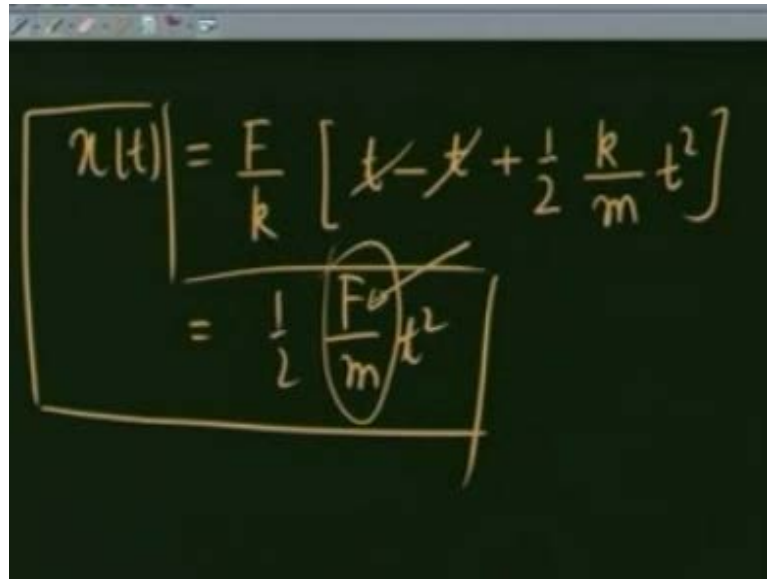
the expression for the distance travelled, let us play around with it a bit, and see if it goes to correct limits.

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The image shows a chalkboard with handwritten mathematical equations. The top equation is $x(t) = \frac{F}{k} \left[t - \frac{m}{k} \left\{ 1 - e^{-\frac{k}{m}t} \right\} \right]$. Below it, the limit $k \rightarrow 0$ is indicated. The next equation shows the expansion of the exponential term: $1 - e^{-\frac{k}{m}t} \approx 1 - \left(1 - \frac{k}{m}t + \frac{1}{2} \frac{k^2 t^2}{m^2} \right)$. An arrow points from this expansion back to the curly braces in the first equation. The final equation is $x(t) = \frac{F}{k} \left[t - \frac{m}{k} \left(\frac{k}{m}t - \frac{1}{2} \frac{k^2 t^2}{m^2} \right) \right]$.

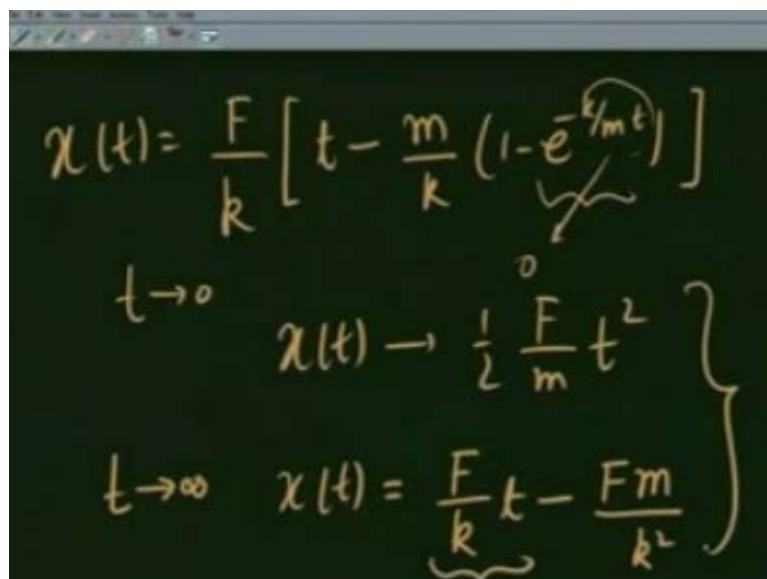
So, I have $x(t)$ equals F over k , t minus m over k , 1 minus e raise to minus k over m t . So, first question I ask is, what happens if k goes to 0 . Remember I said earlier you have to be careful in taking limits. So, I am going to take expand this term keep terms up to order k and see what happens. In the limit of k going to 0 , I can write 1 minus e raise to 0 , minus k over m t , if k is really small for any t , I can write this approximately as 1 minus, 1 minus k over m , t plus half, k square t square over m square. Substituting this in here, I get $x(t)$ is equal to F over k , t minus m over k , k over m , t minus 1 half, k square t square over m square.

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$$x(t) = \frac{F}{k} \left[t - \frac{F}{mk} + \frac{1}{2} \frac{k}{m} t^2 \right]$$
$$= \frac{1}{2} \frac{F}{m} t^2$$

And expanding it further I get x equals F over k , t minus t plus 1 half, k over m , t square, which is nothing but you are familiar 1 half F over m , t square which is the correct answer in the limit of k going to 0 . That is I am taking particle and applying a constant force F , its acceleration is F over m , and there it moves distance 1 half F over m , t square in time t . Let us try to plot this and see how it goes.

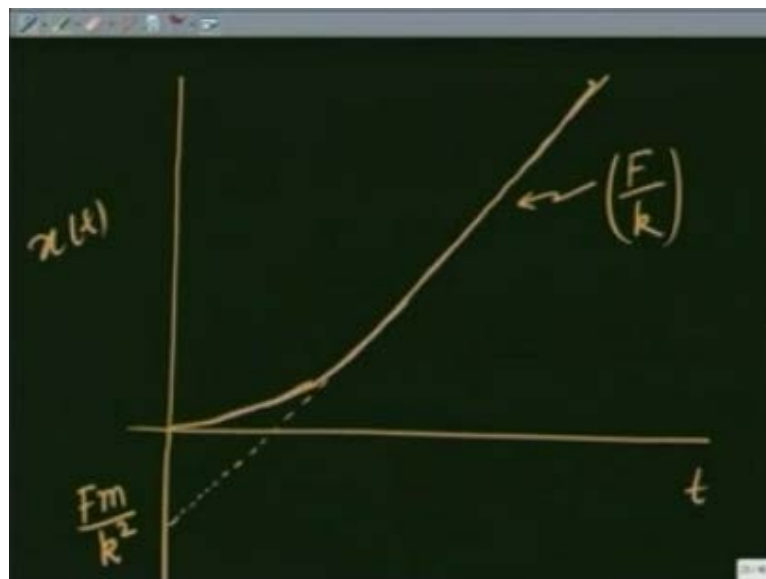
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$$x(t) = \frac{F}{k} \left[t - \frac{m}{k} (1 - e^{-k/mt}) \right]$$
$$\left. \begin{array}{l} t \rightarrow 0 \quad x(t) \rightarrow \frac{1}{2} \frac{F}{m} t^2 \\ t \rightarrow \infty \quad x(t) = \frac{F}{k} t - \frac{Fm}{k^2} \end{array} \right\}$$

So, $x(t)$ was $\frac{F}{k} t - \frac{m}{k} \left(1 - e^{-\frac{k}{m} t} \right)$, again in the limit of very small t . Basically making this term very small, we just saw that $x(t)$ goes as $\frac{1}{2} \frac{F}{m} t^2$.

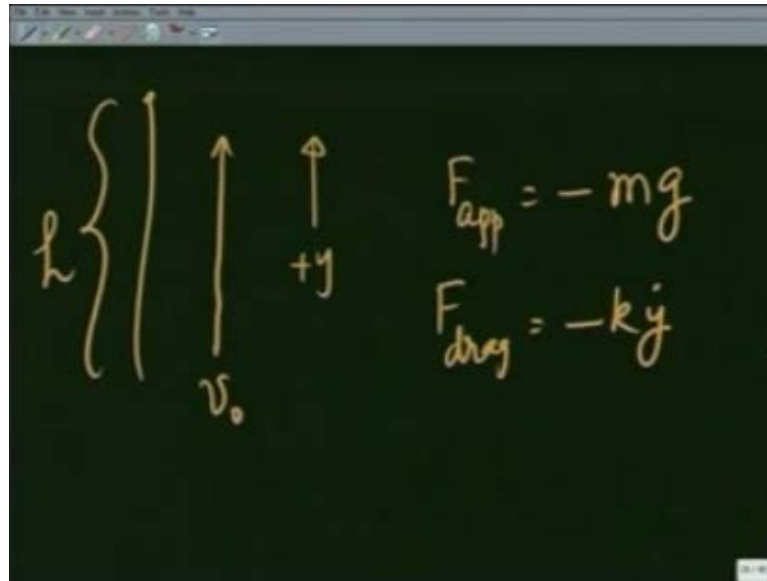
How about large t ? In the large t limit, this term would tend to 0 and therefore, I would have $x(t)$ as $\frac{F}{k} t - \frac{Fm}{k^2}$. This is nothing but movement with the constant speed, $\frac{F}{k}$ you call from earlier is nothing but the terminal speed. So, the particle has now, after long time started moving with terminal speed $\frac{F}{k}$. So, therefore, if I combine the two, and plot $x(t)$ is going to look like.

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At $x(t)$ initially it is going as t^2 and finally, it goes as a straight line. a straight line with slope $\frac{F}{k}$. which is nothing but the terminal speed and which cuts the y axis somewhere here, which is nothing but $\frac{Fm}{k^2}$. So, this is how the particle is going to cover the distance with time. This is a problem where we applied a constant force. So, the most familiar problem in everyday life of this is a ball being thrown up with say some initial speed v_0 .

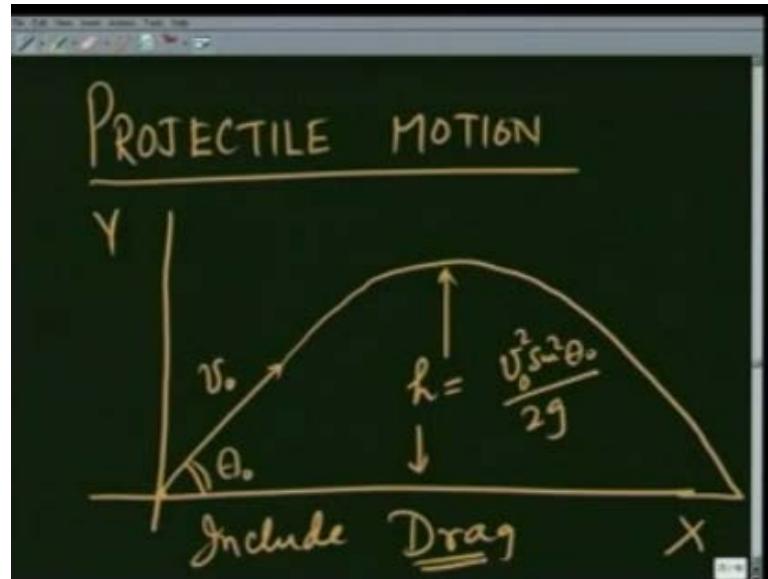
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The problem that I just solved, I took initial speed to be 0, but suppose I throw a ball up then, the force and take this direction to be plus y then, the force applied is going to be minus mg. Drag force is going to be minus k y dot. and the problem becomes similar to what I solved.

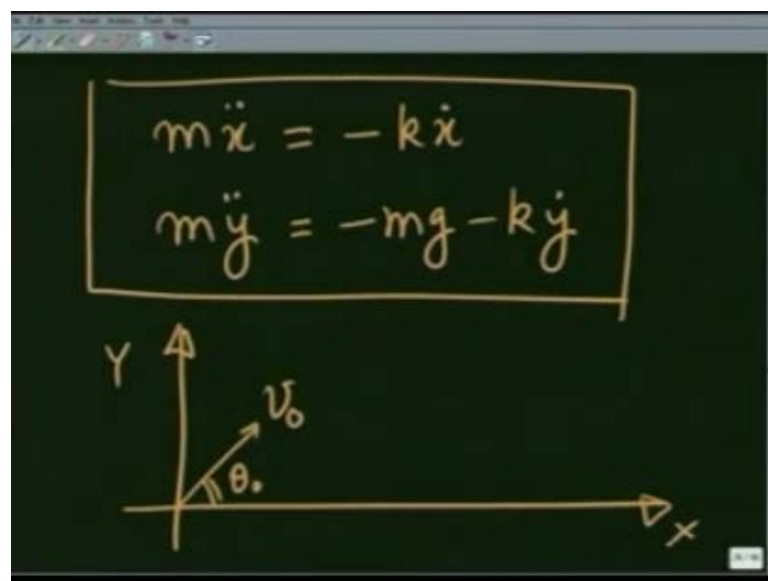
I leave this as an exercise to find how long does it take for the particle to reach up. how high does it go. Because I am going to solve a slightly more complicated problem which is related to this, and that is going to be the motion in two dimension, the projectile motion in gravitational field. So, with I am done 1 dimensional motion now, let us go to the 2 dimensional motion with drag force included.

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You are familiar with projectile motion, which is in the vertical plane. In this motion, I take a particle and project it, add some angle θ_0 from the horizontal, with some initial speed v_0 , and without drag it follows a perfect parabola. And this problem you solved in your 12th grade height comes out to be $\frac{v_0^2 \sin^2 \theta_0}{2g}$. You can calculate the time, the range, and everything. What we are going to do now is make the problem slightly more complicated by including drag force in it, in a simplest form. And therefore, my equation of motion for the projectile is going to be.

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$m\ddot{x}$ is equal to $-kx$ because there is no horizontal applied force. And in the y direction, I am going to have $m\ddot{y}$ is equal to $-mg - k\dot{y}$. These are my two equations of motion. When I am solving for a particle which is being thrown at an angle θ_0 with initial speed v_0 , this I am taking to be positive y direction, this to be x direction. These two I have solved separately when I was solving the 1 d problem so, that I know the solution.

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The image shows three equations written on a blackboard, grouped by a large right-facing curly brace:

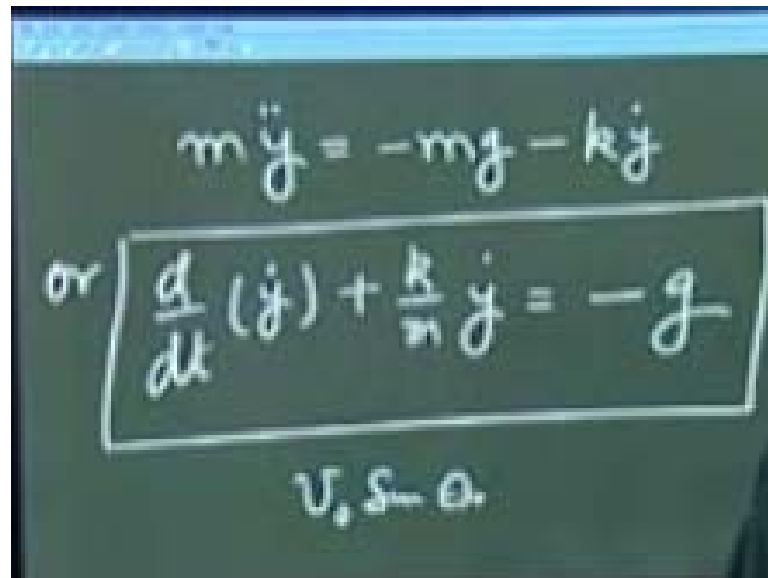
$$m\ddot{x} = -kx$$

$$\dot{x} = v_0 \cos \theta_0 e^{-\frac{k}{m}t}$$

$$x(t) = \frac{m v_0 \cos \theta_0}{k} (1 - e^{-\frac{k}{m}t})$$

$m\ddot{x} = -kx$ gives me \dot{x} , is equal to $v_0 \cos \theta_0 e^{-\frac{k}{m}t}$. And also x as a function of time comes out to be $\frac{m v_0 \cos \theta_0}{k} (1 - e^{-\frac{k}{m}t})$. This is the motion in x direction, and this is related, this is going on concurrently with motion in y direction, for which I have.

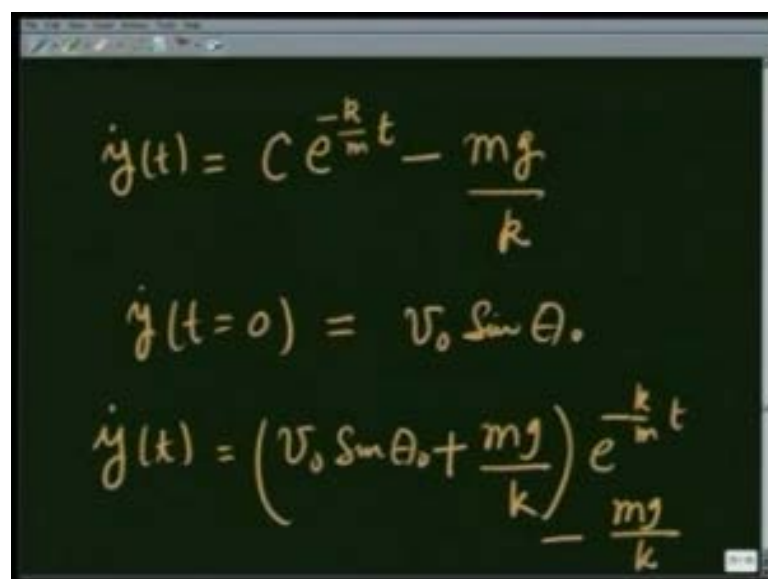
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A chalkboard with a dark green background and a blue header. The text is written in white chalk. The top equation is $m\ddot{y} = -mg - k\dot{y}$. Below it, the word "or" is written to the left of a rectangular box containing the equation $\frac{d}{dt}(\dot{y}) + \frac{k}{m}\dot{y} = -g$. Below the box, the text $v_0 \sin \theta$ is written.

My double dot, is equal to minus mg, minus k y dot or d over dt of y dot, plus k over m, y dot is equal to minus g. In our 1 dimension example, we have already solved this. The slightly different parameters instead of this minus g, I had an F there, the initial velocity was 0, here initial velocity is $v_0 \sin \theta$. So, let us solve this for that. As I said earlier in the case of 1 dimension, the general solution is going to be summation of the particular solution, and the solution of the homogenous part.

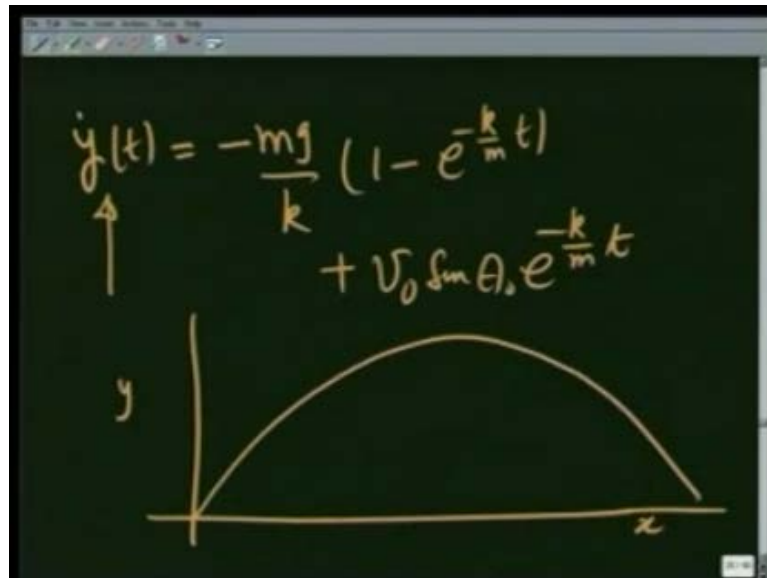
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A chalkboard with a dark green background and a grey header. The text is written in white chalk. The first equation is $\dot{y}(t) = C e^{-\frac{k}{m}t} - \frac{mg}{k}$. The second equation is $\dot{y}(t=0) = v_0 \sin \theta$. The third equation is $\dot{y}(t) = \left(v_0 \sin \theta + \frac{mg}{k} \right) e^{-\frac{k}{m}t} - \frac{mg}{k}$.

And therefore, \dot{y} as a function of t is going to look like $C e^{-\frac{k}{m}t}$, minus $\frac{mg}{k}$. And C is determined by the initial condition which is that \dot{y} at $t=0$ is equal to $v_0 \sin \theta_0$. And therefore, $\dot{y}(t)$ is equal to $v_0 \sin \theta_0 e^{-\frac{k}{m}t} - \frac{mg}{k}$.

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Or I can also write this as $\dot{y}(t)$ is equal to $-\frac{mg}{k} (1 - e^{-\frac{k}{m}t}) + v_0 \sin \theta_0 e^{-\frac{k}{m}t}$. This is how the part the velocity a particle in y direction varies with time. If I want to find the trajectory of the particle that is, I want to find, how does it go in the space then, I have to solve for x and y and plot them y versus x , one versus the other. However, before we find the trajectory, let us play around with the expression for the velocities that we have already obtained.

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$$\dot{x}(t) = v_0 \cos \theta_0 e^{-k/m t}$$
$$\dot{y}(t) = -\frac{mg}{k} (1 - e^{-k/m t}) + v_0 \sin \theta_0 e^{-k/m t}$$
$$T, \quad \underline{\underline{\dot{y}(t) = 0}}$$

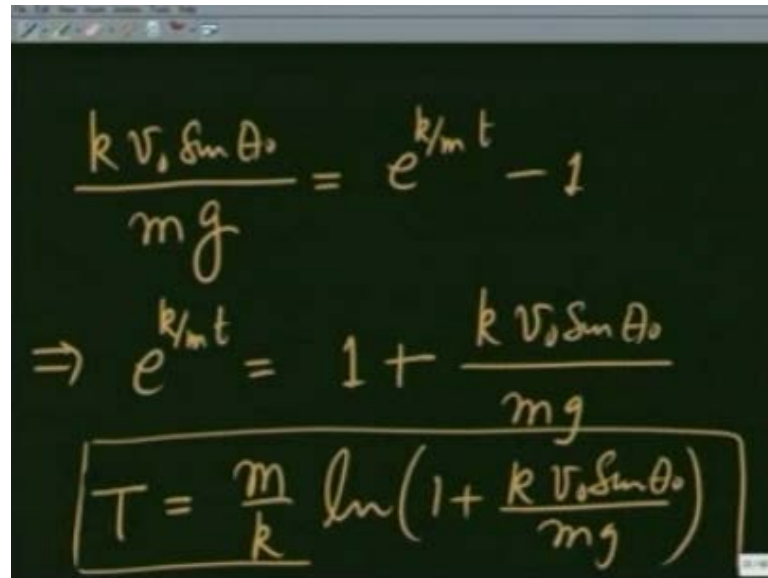
x dot t is $v_0 \cos$ of θ_0 , e raise to minus k over m , t , and we have just found that y dot t is equal to minus mg over k , 1 minus e raise to minus k over m , t , plus $v_0 \sin$ θ_0 e raise to minus k over m , t . At this point I may ask, how long does it take time T before y dot becomes 0 . That is after I throw the particle, after I project the particle when does it come to stop or when is when does it reach its maximum height? So, for that I substitute this is equal to 0 , and corresponding time would get me the answer. Let us do that, if I do that I get.

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$$\frac{mg}{k} (1 - e^{-k/m t}) = v_0 \sin \theta_0 e^{-k/m t}$$
$$\Rightarrow \frac{k v_0 \sin \theta_0}{mg} = \frac{1 - e^{-k/m t}}{e^{-k/m t}}$$

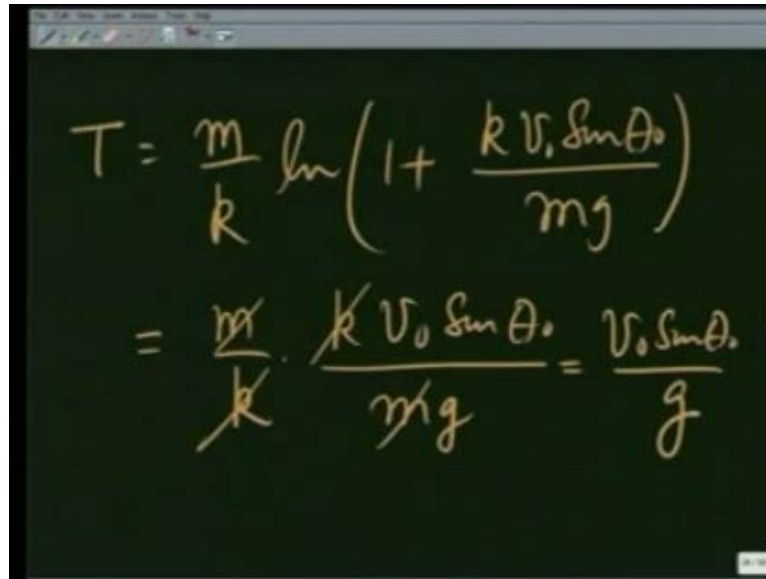
$\frac{mg}{k}$, $1 - e^{-\frac{k}{m}t}$, is equal to $v_0 \sin \theta_0 e^{-\frac{k}{m}t}$. And this gives me $v_0 \sin \theta_0$, $\frac{mg}{k}$ is equal to, $1 - e^{-\frac{k}{m}t}$.

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$$\frac{k v_0 \sin \theta_0}{mg} = e^{\frac{k}{m}t} - 1$$
$$\Rightarrow e^{\frac{k}{m}t} = 1 + \frac{k v_0 \sin \theta_0}{mg}$$
$$T = \frac{m}{k} \ln \left(1 + \frac{k v_0 \sin \theta_0}{mg} \right)$$

Or $\frac{k v_0 \sin \theta_0}{mg}$ is equal to $e^{\frac{k}{m}t} - 1$, which implies $e^{\frac{k}{m}t}$ is equal to $1 + \frac{k v_0 \sin \theta_0}{mg}$ or time T , when its velocity becomes 0 or reaches highest point is equal to $\frac{m}{k} \ln \left(1 + \frac{k v_0 \sin \theta_0}{mg} \right)$. Let me again ask and check my answer. Is this correct in the limit of k going to 0? Sure it is.

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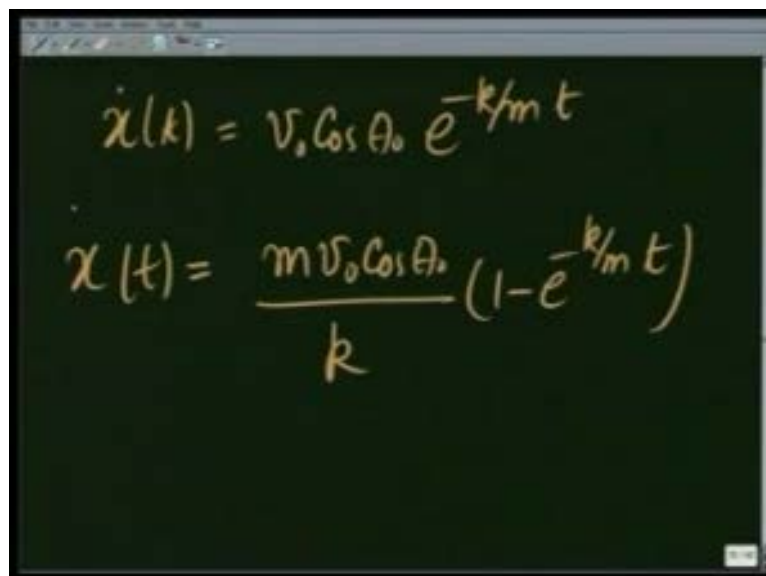


A photograph of a chalkboard showing the derivation of terminal velocity. The first line is $T = \frac{m}{k} \ln\left(1 + \frac{k v_0 \sin \theta_0}{mg}\right)$. The second line shows the expansion for small k : $= \frac{m}{k} \cdot \frac{k v_0 \sin \theta_0}{mg} = \frac{v_0 \sin \theta_0}{g}$.

Because T is equal to m over k log of, 1 plus $k v_0 \sin \theta_0$ over mg . In the limit of very small k can be written as m over k expanding log, I get $k v_0 \sin \theta_0$ over mg this cancels, this cancels and I get my familiar answer, $v_0 \sin \theta_0$ over g .

After doing this exercise, let us go on to find the trajectory. For trajectory, what I will do is, find x and y separately and plot them. Recall that without any drag force, one could eliminate T and find x and y , a direct relationship between x and y . This is not possible in this case so, going back to our expressions.

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A photograph of a chalkboard showing two equations. The first is $\dot{x}(t) = v_0 \cos \theta_0 e^{-k/m t}$. The second is $x(t) = \frac{m v_0 \cos \theta_0}{k} (1 - e^{-k/m t})$.

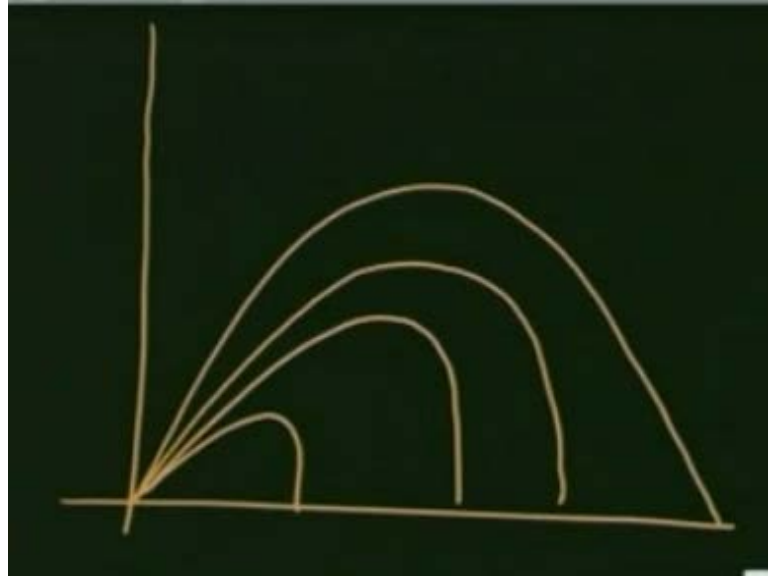
We had \dot{x} is equal to $v_0 \cos \theta_0 e^{-k/m t}$, which gives me x is equal to $\frac{m v_0 \cos \theta_0}{k} (1 - e^{-k/m t})$. This is my expression for x , you are also familiar with this from the earlier part of the lecture.

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The image shows a chalkboard with two equations. The top equation is the derivative of y with respect to time, $\dot{y}(t) = -\frac{mg}{k} (1 - e^{-k/m t}) + v_0 \sin \theta_0 e^{-k/m t}$. The bottom equation is the integrated form of y(t), $y(t) = \frac{m v_0 \sin \theta_0}{k} (1 - e^{-k/m t}) - \frac{m g}{k} (t - \frac{m}{k} (1 - e^{-k/m t}))$.

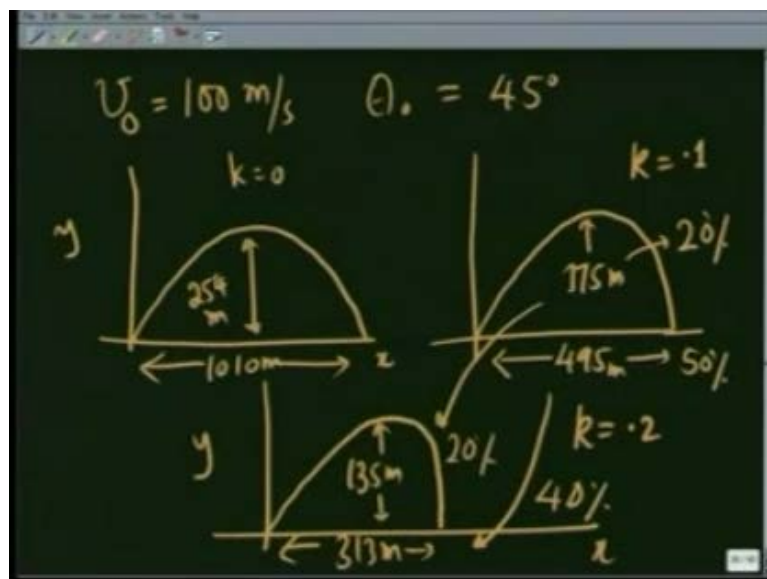
Similarly, I have \dot{y} which is nothing but $-\frac{mg}{k} (1 - e^{-k/m t}) + v_0 \sin \theta_0 e^{-k/m t}$. Integrating this I get an expression for y as a function of time and it comes out to be $\frac{m v_0 \sin \theta_0}{k} (1 - e^{-k/m t}) - \frac{m g}{k} (t - \frac{m}{k} (1 - e^{-k/m t}))$. You see because t appears in linear form as well as this form, that is why I cannot eliminate it to get a direct expression between x and y . However, for given time I can calculate x and y and plot the trajectory.

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If I do that for different case, and a fix v_{naught} and $\sin \theta_{\text{naught}}$, I find that for k equal to 0 it is a perfect parabola. However, as k is introduced that means, drag force comes into play, its height and range both start changing and all have got some numbers here, which I will show you how for different parameters it changes. As an example of how the trajectories are affected for different values of that coefficient or by this drag.

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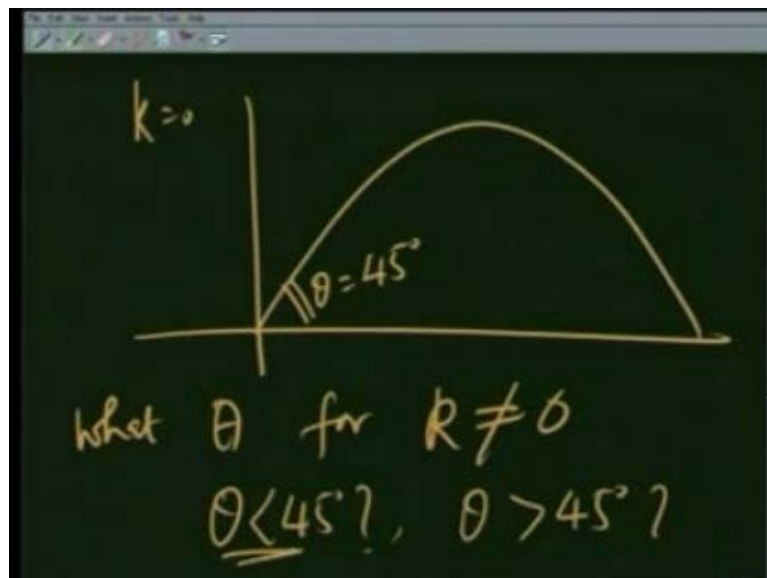
I have taken a sample problem of a particle being thrown in the speed of 100 meters per second at an angle θ_0 equal to 45 degrees, and I will plot for three different case, the

trajectories which I have calculated. I take k is equal to 0, which is the case of no drag or no viscosity, in that case you will familiar as this comes out as a parabola. Its height is about 254 meters, and the range comes to about 1 kilometer 1010 meters. Now, I take k to be point 1, the trajectory changes slightly. It goes and comes down, but not as a parabola, but something like this, its height now reduces to 175 meters and range comes out to be 495 meters.

Notice, that the reduction in height is about 20 percent whereas, range as gone down by about 50 percent. If I go further and make k equals 0.2, the trajectory becomes even more like this with height now becoming 135 meters, and the range becoming 313 meters. Again notice, that the reduction in range is much more which is about 40 percent whereas, the height goes down by about 20 percent.

So, you see how introducing drag or this frictional force due to fluid or the gas affects the motion of a trajectory. This is a very important problem when we consider a motion of particles some motion of say bombs or motion of artillery pieces being thrown, being projected in a battle field. An interesting question you may ask at this point is, you all know from your 12th grade that.

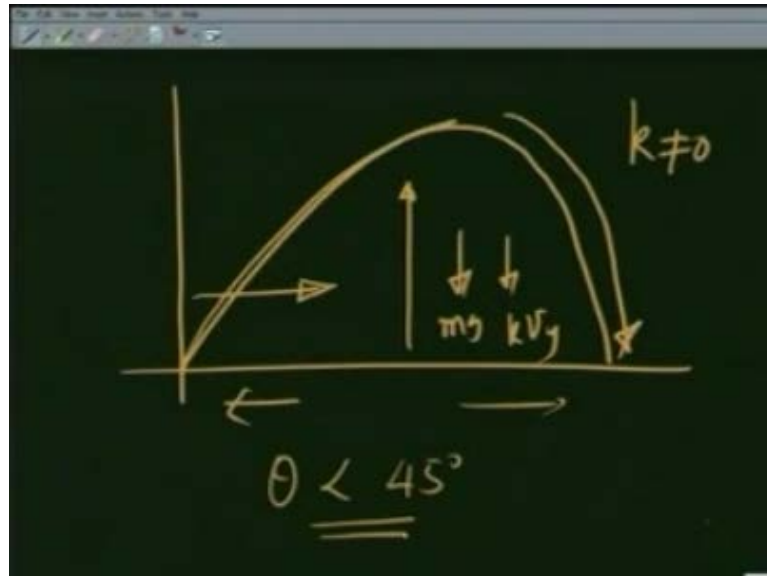
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If k is equal to 0, the maximum range which is attained is when θ is equal to 45 degrees. What θ for k not equal to 0? Is θ less than 45 degrees? Is θ more than

45 degrees? The answer to this question is the theta should be slightly less than 45 degrees.

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Because you saw that the height is not affected so much by introducing of k as the range. So, what I want to do is give slightly more of x component to the velocity, if I want to fire the projectile further and further. The other way to look at the same problem is that, when the particle is moving is up, it spends more time because it is being pull down by mg , it is also being pull down by kV_y and spends more time in going the same height then, when it is coming down.

So, in that time, you want to cover as much distances as possible because x is, x component of the velocity is also going down. So, to attain maximum range, I should fire at theta less than 45 degrees. So, far I have considered the drag force which is proportional to the velocity, and in the opposite direction, and shown you through some examples, how it affects the motion. This is the simplest possible drag force.

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Speed increases

$$\vec{F} = -k_1 \vec{v} - k_2 v^2 \hat{v}$$

As the speed increases more dependence on velocity may come in the form of F may be equal to minus $k v$, minus, let me write this k_1 , $k_2 v^2$, left side a unit vector \hat{v} so that, it tells you this is opposite to the velocity vector, and higher order terms, this being the simplest. When these terms come into play, problem becomes non-linear.

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$$m \ddot{\vec{r}} = \vec{F}_{\text{applied}} - k_1 \dot{\vec{r}} - k_2 (\dot{\vec{r}})^2 \hat{\dot{\vec{r}}}$$

Because now you have $m \ddot{\vec{r}} = \vec{F}_{\text{applied}} - k_1 \dot{\vec{r}} - k_2 (\dot{\vec{r}})^2 \hat{\dot{\vec{r}}}$, And these are not easy to solve. In most of the cases, you have to apply numerical methods to get the final solution. However, there are

some cases where you can apply some techniques and get solutions. I will solve one example to illustrate this, and that would complete our lecture on drag forces.

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The image shows handwritten notes on a blackboard. At the top, the drag force is given as $\vec{F}_{\text{drag}} = -k v^2 \hat{v}$. Below this, there are two diagrams illustrating the forces on a particle. The first diagram shows a particle moving upwards with initial velocity v_i and time t , with a height h reached. The second diagram shows a particle moving downwards with final velocity v_f . To the right of these diagrams, two equations of motion are written and circled: $m\ddot{y} = -mg - kv^2$ for upward motion and $m\ddot{y} = -mg + kv^2$ for downward motion.

The example I take is that of F_{drag} being proportional to v square. And obviously, it is in direction opposite to its velocity. And I want to find, if I throw a particle up with some initial speed, v initial, how long does it take to reach up? What is the height it can go up to? And what is the speed when it comes down? v final.

Let us see what all answers can we get. If I write the equation of motion since, this is a 1 dimensional case, let me take y to be going up, I can write $m y$ double dot is equal to minus mg minus $k v$ square, this is when the particle is going up. Notice in this case, I have to be careful when considering the motion of particle coming down. When the particle is coming down, I would have $m y$ double dot is equal to minus mg plus $k v$ square because in that case the force, drag force would be acting in upward direction. That is why I took this particular example. So, let us see how do we go about solving this?

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$$m\ddot{y} = -mg - k\dot{y}^2$$
$$\ddot{y} = \frac{d}{dt}(\dot{y}) = \frac{d}{dy}(\dot{y}) \frac{dy}{dt}$$
$$= \frac{1}{2} \frac{d}{dy}(\dot{y}^2)$$
$$\frac{1}{2} \frac{d}{dy}(\dot{y}^2) = -\frac{mg}{m} - \frac{k}{m}\dot{y}^2$$

While the particle is moving up I write m y double dot equals minus mg , minus k y dot square, and I go back to my technique of writing y double dot as d over dt , y dot which is same as d over dy , y dot, dy dt which is same as d over dy of, y dot square $1/2$. And therefore, I can write my equation as $1/2$ d over dy , y dot square is equal to minus mg over m , minus k over m y dot square.

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$$\frac{1}{2} \frac{d}{dy}(\dot{y}^2) = -g - \frac{k}{m}\dot{y}^2$$
$$\boxed{\frac{1}{2} \frac{d}{dy}(\dot{y}^2) + \frac{k}{m}\dot{y}^2 = -g}$$
$$\dot{y}^2 = C e^{-\frac{2k}{m}y} - \frac{mg}{k}$$

And I have $1/2$ d over dy , y dot square is equal to minus g , minus y dot square, $1/2$ d over dy y dot square, plus y dot square equals minus g . I can solve for y dot square

directly in terms of y . You see I have eliminated time. So, in this problem it will not be possible for me to calculate how long does it take before the particle comes to stop, but I can certainly calculate, how high the particle goes before it comes to stop.

You can directly see that, the solution I have forgot k over m here. So, there is going to be a k over m here, \dot{y}^2 is going to be of the form $C e^{-2k/m y} - \frac{mg}{k}$. You can directly substitute this and see if I substitute \dot{y}^2 with $\frac{mg}{k}$, it gives me minus g and this satisfies the homogenous part. Now, two initial condition.

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$$\begin{aligned}\dot{y}^2 &= C e^{-\frac{2k}{m}y} - \frac{mg}{k} \\ \dot{y}^2 (y=0) &= v_i^2 \\ v_i^2 &= C - \frac{mg}{k} \\ \Rightarrow \dot{y}^2 &= \left(v_i^2 + \frac{mg}{k}\right) e^{-\frac{2k}{m}y} - \frac{mg}{k}\end{aligned}$$

So, I have \dot{y}^2 equals some constant $e^{-2k/m y}$, minus mg over k , and I know the initial velocity \dot{y}^2 , when y is 0 is equal to v initial square. And therefore, v initial square is going to be equal to C , minus mg over k and that gives me \dot{y}^2 is equal to v_i^2 , plus mg over k , $e^{-2k/m y}$, minus mg over k .

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$$\dot{y}^2 = 0 \text{ at } y = h$$
$$\left(v_i^2 + \frac{mg}{k} \right) e^{-\frac{2k}{m}h} = \frac{mg}{k}$$
$$e^{\frac{2k}{m}h} = \left(1 + \frac{v_i^2 k}{mg} \right)$$

To find how high the particle went, I take y dot square to be 0 at y equal to h and therefore, I get v_i square, plus mg over k , e raise to minus $2k$ over m h , equals mg over k . And that tells me that, e raise to $2k$ over m h is equal to 1 , plus v_i square k upon mg .

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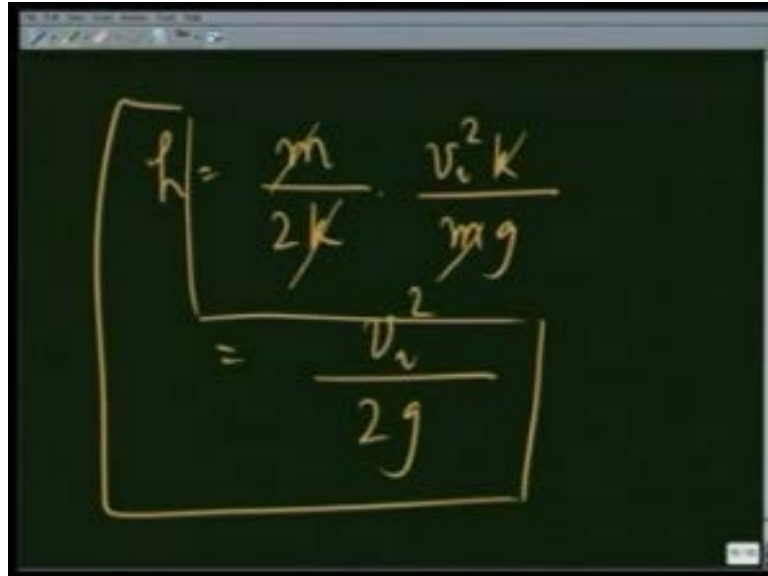
$$\frac{2k}{m}h = \ln \left(1 + \frac{v_i^2 k}{mg} \right)$$
$$h = \frac{m}{2k} \ln \left(1 + \frac{v_i^2 k}{mg} \right)$$

$k \rightarrow 0 \quad \frac{v_i^2 k}{mg}$

Or $2k$ over m h , equals \ln of 1 , plus v_i square k upon mg , or the height it went up to is equal to m over $2k$ \ln of 1 , plus v_i square k upon mg , this is the height at which it stops. Again, I will go back to limit of k going to 0 and see if my answer is correct. When the

limit k equal to 0 is taken, I can expand the log term and write this as v_i square k over mg .

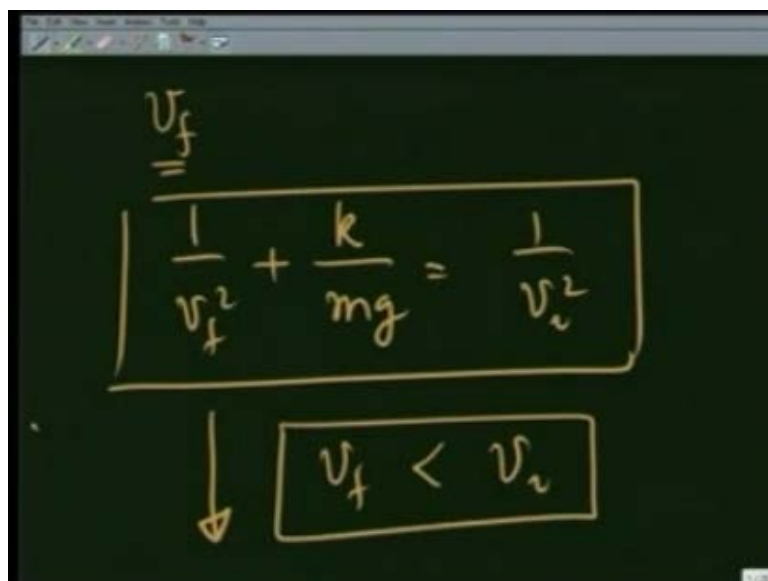
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A handwritten derivation on a chalkboard showing the calculation of height h . The first line is $h = \frac{m}{2k} \cdot \frac{v_i^2 k}{mg}$. The second line shows the result after simplification: $h = \frac{v_i^2}{2g}$.

And therefore, the height that I get is h equals m over $2k$, v_i square k over mg , which is our familiar result from your previous classes. So, I have calculated the height up to which the part will go, will go with the drag included. I leave it for you as an exercise that when it comes back.

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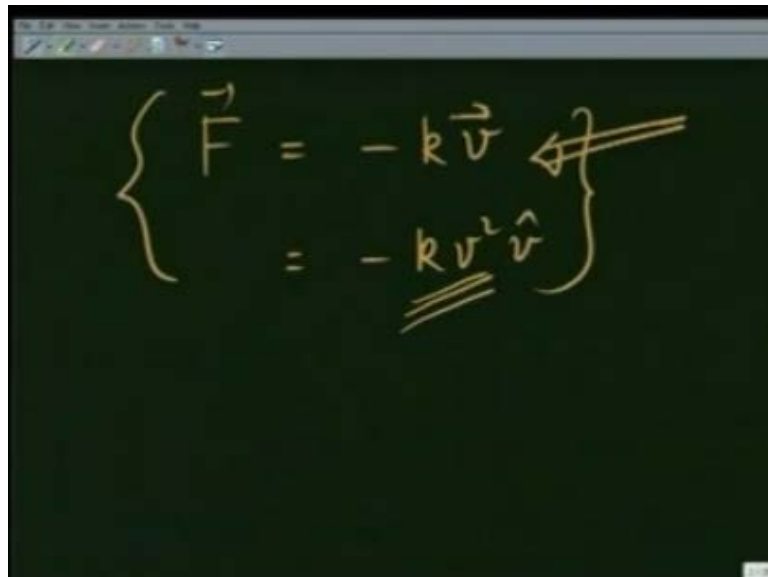


A handwritten derivation on a chalkboard showing the relationship between final velocity v_f and initial velocity v_i . The first line is $v_f = \frac{1}{\frac{1}{v_f^2} + \frac{k}{mg}} = \frac{1}{\frac{1}{v_i^2}}$. Below this, a downward arrow points to a boxed inequality: $v_f < v_i$.

The velocity v final then, the relationship between v final and v initial square is going to be 1 over v final square, plus k over mg is equal to 1 over v initial square. The word of caution though, when considering the motion of part of the coming down you should take proper account of the force of friction.

The force of friction changes direction, and I was showed in one of the previous slides, that you have to change the sign of the force. V final of courses is going to be smaller than v initial because in this friction, the particle has loss energy. So, what we have done in this lecture is, shown through some examples, how the drag force affects the motion of a particle.

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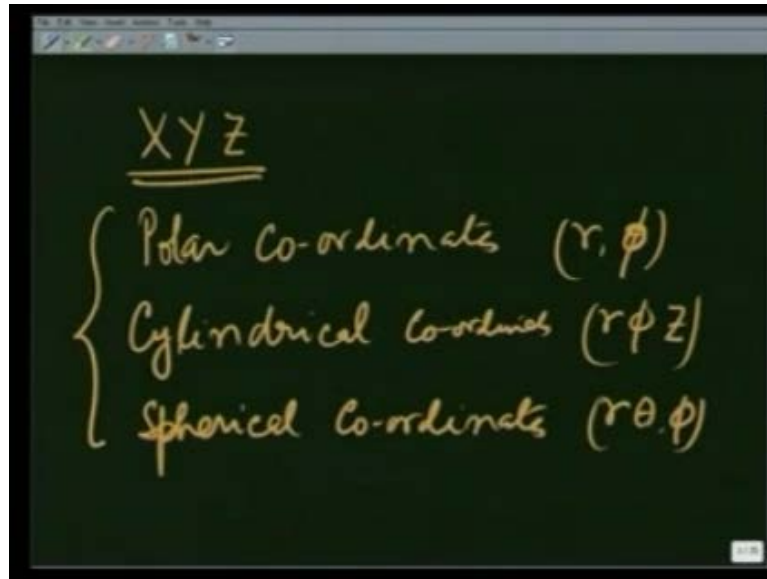


A photograph of a chalkboard with handwritten equations. The equations are enclosed in large curly braces. The top equation is $\vec{F} = -k\vec{v}$, with a double-headed arrow pointing to the right next to the velocity vector. The bottom equation is $= -k\underline{v^2}\hat{v}$, where v^2 is underlined with three lines. The entire set of equations is written in yellow chalk on a dark green background.

I have considered the simplest possible form where F is proportional to the velocity itself, and another form where F is proportional to the velocity square. This is encountered when velocities are small and suffices in most of the cases.

Another mathematical advantage or analytical advantage with this is, that most of the cases I can solve it analytically. If the speeds become higher then, higher orders of v 's power come in to play and then, one has to solve the problem numerically. With these we finish the first part of these lectures on engineering mechanics. In this part I made you familiar with different co-ordinate systems that we use besides.

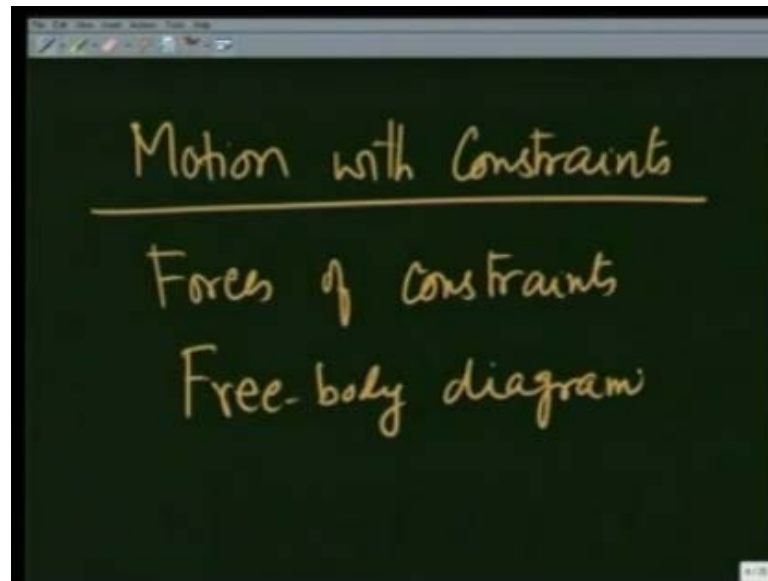
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The xyz co-ordinate system, cartesian co-ordinate systems. We considered the planar polar co-ordinates r and ϕ . We considered cylindrical co-ordinates r , ϕ and z . We considered spherical co-ordinates r , θ , and ϕ . And we used some of them while solving problems.

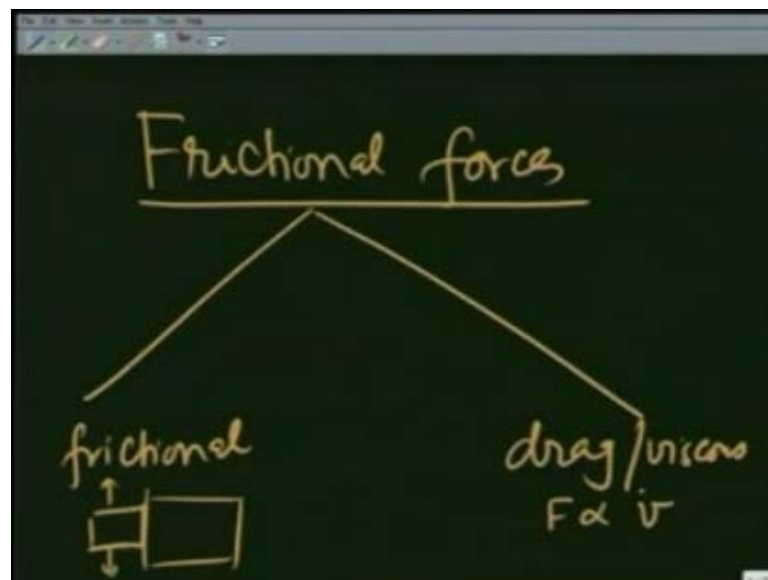
Not only these co-ordinates systems are useful in solving mechanics problems, you will find them very useful in solving problems in electrodynamics and wherever there are symmetries. After this introduction to co-ordinate systems, we considered the motion of particles single particles with constraints and learnt.

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What are forces of constraints? What is free body diagram? How to take constraints into account and how to solve problems when constraints are there? Following that, we considered a very special class of forces.

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The frictional forces and took them into account when solving dynamical problems. In these we considered regular frictional force that you are familiar with, when one body slides over the other body, solid body. And we also considered drag or viscous force. In their simplest possible forms, one where F is proportional to the velocity and also

proportional to the v square. In the coming few lectures, we are now going to get more sophisticated, considered the work done by forces, do work energy theorem and then, go to many particle systems where we will consider rigid body motion and things like those.