

Solid State Devices
Dr. S. Karmalkar
Department of Electronics and Communication Engineering
Indian Institute of Technology, Madras
Lecture - 28
Bipolar Junction Transistor (Contd...)

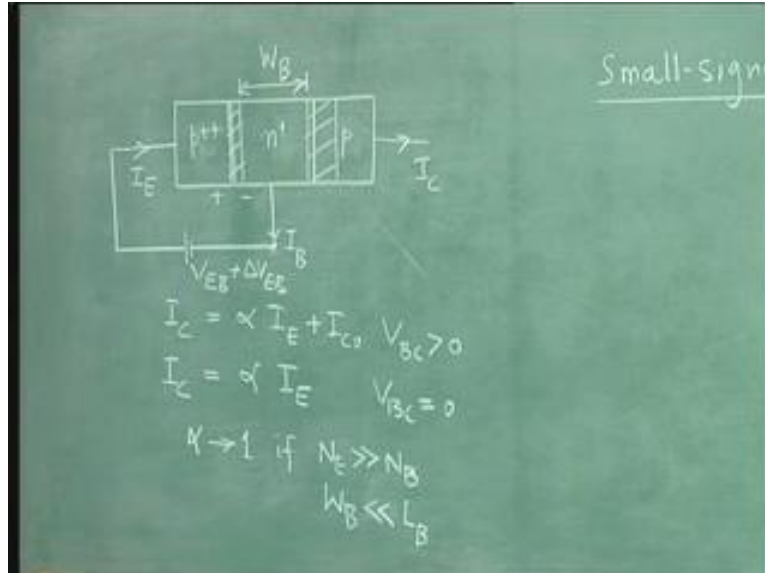
This is the 28th lecture of this course and the 3rd lecture on bipolar junction transistors. In the previous lecture we basically explained what transistor action is. We said that transistor action is the transfer of current from one junction which is forward bias to a nearby junction. Now this transfer is complete if this nearby junction which is called the collector junction is very close to the emitter junction as defined by this particular relation W_B is much less than L_B . W_B is the width of the base region here that is the distance between the two depletion edges.

And secondly, another condition for the transfer to be complete is that the doping N_E that is the doping in the emitter should be much greater than the doping in the base, this is the relation. So, if these two conditions are satisfied then you will have almost complete transfer of the current from the emitter junction to the collector junction or I_C will be very close to I_E . I_C is equal to α times I_E where α tends to 1. Now what is the condition for this transfer on the bias across the collector junction? What we said is that is this bias is 0 that is the collector base is sorted or if the collector base junction is reverse bias then in both these cases the transfer is almost complete. But if the collector junction is forward bias, that is, if it is like this the p-region positive with respect to the n-region then the transfer is seriously affected and the transistor action therefore is seriously affected. So we must ensure that if you want complete transfer of the current from the emitter to collector this junction should either be zero bias or reverse bias.

For the zero bias case the relation is I_C is equal to α times I_E so this is when V_{BC} is equal to 0. V_{BC} is the voltage between base and collector. And I_C is equal to α times I_E plus a small current if V_{BC} is greater than 0. That is, the base is positive with respect to collector or collector is reverse bias. So in both these cases you have the complete transfer. This is the only small difference in the reverse bias and zero bias case, this current I_C here.

Now in this lecture we will see how the transistor action can be used for the purpose of amplifying small signals. So the topic in this lecture is small signal amplification. Now let us look at the biasing arrangement for this purpose. The biasing arrangement is as follows. Small signal means the emitter base junction is forward biased and this forward bias is incremented by a small value δV_{EB} .

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Now as a result of this forward bias change there will be a change in the collector current ΔI_C . Of course there will be changes in emitter and base currents as well. So, for example, the base current would become I_B plus ΔI_B and the emitter current would become I_E plus ΔI_E . So they are the changes in response to the change in the controlling voltage that is the emitter base voltage. So, the first step of amplification is to relate this change in collector current to the emitter base voltage.

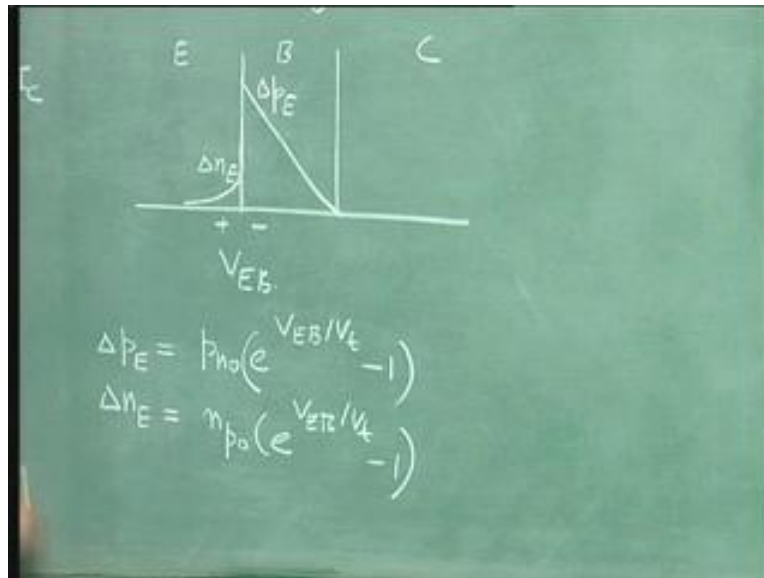
Let us see what is this change ΔI_C in the response to ΔV_{EB} . We can then build up other changes also that is ΔI_B and ΔI_E . And then we will see how this change in collector current can help us to achieve the amplification. Now to begin with we will assume that the collector base junction is zero bias. So here this is shorted so what we will do now is, if you start with the equation I_C is equal to α times I_E then we know that ΔI_C is equal to α times ΔI_E . This is the first step.

What is ΔI_E in response to ΔV_{EB} is what we need to see. It can be easily shown that the relation between ΔI_E and ΔV_{EB} is nothing but the diode relation that is the relation between the current and the voltage across a diode and this can be shown easily as follows. Let us draw the minority carrier distributions in the emitter or base and collector. We are assuming that the base width is very small as compared to the diffusion length L_B . So this distribution is almost a straight line as we have pointed out in the previous lecture and this distribution is exponential. So this is emitter this is base and this is collector this is Δp_E and this is Δn_E this is p-type region, this is n-type region and this is again p-type region. So this is the excess minority carrier distribution when the emitter base junction is forward bias and the collector junction is zero bias. We have drawn this in the previous lecture. Using this distribution let us explain the relation between the emitter current and the emitter base voltage.

Now, Δp_E according with the law of the junction can be written as Δp_E is equal to $p_{n0} e^{(V_{EB} / V_t - 1)}$ where the p_{n0} is the equilibrium concentration of minority

carriers in the base region. We can similarly write Δn_E is equal to $n_{p0}(e^{V_{EB}/V_t} - 1)$ where this is equilibrium concentration of minority carriers in the emitter region. Now, we have been using the symbol n_{p0} also for the electron concentration in the collector under equilibrium. That is minority carrier concentration in the collector because collector is also p-type like the emitter. Now we may therefore think that there can be some confusion.

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Depending on the context it will be clear whether n_{p0} corresponds to emitter or the collector. For example, here we are talking about Δn_E that is the excess concentration of electrons in the emitter at the depletion edge. Therefore obviously n_{p0} should correspond to emitter. So, depending on the context it will become clear whether n_{p0} is related to the emitter or to the collector. Now, proceeding further how can we write the emitter current in terms of these concentrations? We again follow the PN junction theory and we can write I_E is equal to I_{Ep} plus I_{En} so I_E is the current across this particular emitter junction. I_{Ep} is due to these holes which are injected from emitter into base. And I_{En} is due to these electrons which have been injected into the emitter from the base region.

So we can write I_{Ep} which is q into the diffusion coefficient of the carriers in the base into Δp_E by W_B , Δp_E by W_B is nothing but D_p by Dx in the base that is the slope of this line so W_B is this width. Now here we have not shown the depletion edges or rather the depletion layer in the emitter and the collector to avoid complication in the diagram. It is understood that this concentration corresponds to the concentration at the depletion edge. And similarly these other concentration also corresponds to concentration at the depletion edge. So q into D_B into Δp_E by W_B is the current density because of diffusion where D_B is the diffusion coefficient of holes.

Let us explain this nomenclature. We are always going to consider the diffusion coefficient of minority carriers when we talk about the currents in a PN junction.

So we have only one suffix there which shows a region in which the minority carriers are being considered. So, D_B would imply the diffusion coefficient of minority carriers in the base. Now we should multiply this by the area of the emitter to get the current. So this is I_{Ep} in terms of Δp_E that is this current. Similarly we can write I_{En} as area of the emitter into q into diffusion coefficient in the emitter into Δn_E by L_E where L_E is the diffusion length of electrons in the emitter, this an exponential decay so you take the diffusion length of the electrons there. Now, we can combine these relations for Δp_E and Δn_E and then we can write the expression for I_E as I_E is equal to A_E into $q[D_B$ into p_{n0} by W_B plus D_E into n_{p0} by L_E] (e^{V_{EB}/V_t} minus 1) because this term will be common for both these terms. Now one can easily recognize that this particular form is nothing but the reverse saturation current of the emitter junction and we will therefore represent this as I_{E0} . This is nothing but the diode law I_E as a function of V_{EB} the exponential and this is a reverse saturation current.

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The image shows a handwritten derivation on a green chalkboard. At the top, the equation $I_E = I_{Ep} + I_{En}$ is written. Below it, arrows point from I_{Ep} and I_{En} to their respective terms in a combined equation. The first term is $\frac{A_E q D_B \Delta p_E}{W_B}$ and the second term is $\frac{A_E q D_E \Delta n_E}{L_E}$. These are then combined into a single equation: $I_E = A_E q \left[\frac{D_B p_{n0}}{W_B} + \frac{D_E n_{p0}}{L_E} \right] (e^{V_{EB}/V_t} - 1)$. A bracket under the terms in the square brackets identifies this as I_{E0} , the reverse saturation current.

So, as compared to the PN junction theory we have discussed earlier the only difference is that here for this term instead of the diffusion length we are having the width of the particular region. And this is because this is like a short region and therefore the diffusion length in this region is being replaced by W_B that is the width of that region because as we have said in our transistor W_B is much less than L_B for the device to act like a good transistor or to have the efficient transistor action. That is why L_B is being replaced by W_B here and this is the difference. Now we can write this formula I_E is equal to I_{E0} (e^{V_{EB}/V_t} minus 1).

And coming back to our relations here we need to obtain ΔI_C which is given by α times ΔI_E in response to ΔV_{EB} . So now we can get ΔI_E in response to ΔV_B using this formula. So we can write ΔI_E is equal to I_{E0} (e^{V_{EB}/V_t} minus 1) by V_t into ΔV_{EB} where in if V_{EB} by V_t is more than about three times V_t which will practically be the

case (eV_{EB} by V_t) is much greater than 1 and therefore $I_{E0} e(V_{EB}$ by $V_t)$ is nothing but I_E itself. So we can write this as approximately equal to I_E into ΔV_{EB} by V_t .

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$$I_E = I_{E0} (e^{V_{EB}/V_t} - 1)$$

$$\Delta I_E = \frac{I_{E0} e^{V_{EB}/V_t}}{V_t} \Delta V_{EB}$$

$$\approx \frac{I_E \Delta V_{EB}}{V_t}$$

The increment in the emitter current is proportional to the increment in the emitter base voltage and the proportionality constant is I_E by V_t . So substituting this relation in this formula here we can write ΔI_C is equal to α times I_E by V_t into ΔV_{EB} where in α times is I_E nothing but I_C . This can be further simplified to I_C by V_t into ΔV_{EB} . So ΔI_C is nothing but I_C by V_t into ΔV_{EB} . So increment in the collector current is proportional to the increment in the emitter base voltage. This is a consequence of the exponential dependence of the emitter current on the emitter base voltage. And since most of the emitter current is transfer to the collector the collector current also depend on the exponentially on emitter base voltage. Therefore when you differentiate or when you take the increments you end up getting a linear relation between the increment in the collector current and the increment in the emitter base voltage. This term I_C by V_t has dimensions of 1 by resistance or conductance and therefore we can represent this using a symbol g_m and we can write this as $g_m \Delta V_{EB}$. This g_m is called the transconductance. That is the relation between the increment in the collector current and the increment in the emitter base voltage.

Now, let us see how this increment in the collector current can be used for purposes of amplification. Suppose we pass this current to a resistor R and we try to find out the voltage change across this resistor as compared to the change in the emitter base voltage. Let us try to relate the change in the emitter base voltage to the change in voltage across this resistor R because of the increment in collector ΔI_C . Now we shall call the increment in the voltage across the resistance as ΔV_0 that is we shall assume that the voltage across the resistance is the output voltage. The V_{EB} emitter base voltage is the input voltage and voltage across this resistor is the output voltage. So we will denote this as ΔV_0 and we write ΔV_0 is equal to R into ΔI_C .

Now, expressing ΔI_C in terms of ΔV_{EB} we can write the relation ΔV_0 by ΔV_{EB} as R into g_m where ΔI_C is nothing but g_m . So ΔV_0 by ΔV_{EB} is simply $R(g_m)$ where g_m is I_C by V_t . Now I_C corresponds to the voltage V_{EB} . So, when you make an increment in the emitter base voltage what you find is that there is an increment in the voltage across the resistor which depends on g_m and R . Now, if R into g_m is greater than 1 then we find that we have amplification because ΔV_0 is more than ΔV_{EB} . The change in the output voltage is more than the change in the input voltage.

Let us put some typical values and see how much can ΔV_0 by ΔV_{EB} be in practice. Let us assume these values which are typical I_C is equal to 1 mA supposing we set up 1 mA of current in the transistor. Let us take room temperature so V_t is 0.026 volts and let us assume a resistance of 1 kilo ohm. R into g_m is given by 1000 ohms into 1 mA upon 0.026 volts. So ohms into ampere by volts this cancels giving a dimensionless quantity 10 to the power minus 3 and this cancels so 1 by 0.026 that is 1000 by 26. So R times g_m is equal to 1000 by 26 so this is close to 1000 by 25 that is about 40 maybe less than 40 and that maybe around 38 or something. We want just an approximate figure so it is about 40. So you find that the change in the voltage across a resistance is forty times the change in the voltage across the emitter base junction. This is what is meant by amplification.

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Small-signal amplification

$$\Delta V_0 = R \Delta I_C$$

$$\frac{\Delta V_0}{\Delta V_{EB}} = R g_m \quad \leftarrow I_C / V_t$$

$$R g_m = \frac{1000}{26}$$

$$\frac{\Delta V_0}{\Delta V_{EB}} \approx 40$$

$I_C = 1 \text{ mA}$ $T = 300 \text{ K}$
 $V_t = 0.026 \text{ V}$
 $R = 1 \text{ K}$

$$R g_m = \frac{1000 \times 1 \times 10^{-3}}{0.026} \approx 40$$

This is what small signal amplification is. So, ΔV_{EB} is the input small signal and ΔV_0 is the output small signal and the ratio between these two small signal voltages is 40 so you are getting voltage amplification here. We emphasize that this amplification is for small signal. So we are not taking the ratio between V_0 and V_{EB} , this V_0 is the voltage across the resistance are because of this current I_C plus ΔI_C and V_B is the DC voltage so we are not taking the ratios of the DC current the total current and total voltage. We are taking the ratios of the increment in the voltage across resistor and the increment in the emitter base voltage. This is therefore an incremental picture. This is what is meant

by small signal amplification. Now there is one catch here and that is, if you pass this current I_C through the resistor it develops a voltage drop which is like this. Now ΔV_0 only correspond to I_C so strictly speaking this is V_0 plus ΔV_0 . Since we are only concerned about the ΔV_0 in response to ΔI_C we showed as ΔV_0 . So this is the voltage drop because of I_C plus ΔI_C . Now obviously this voltage drop is going to come here, it appears like this across the junction.

The moment you put a resistance here in this collector to base lead actually your collector base junction has got forward bias by this much amount. And if there is a forward bias across the collector base junction then your transistor action is destroyed. We cannot write I_C is equal to αI_E where αI_E is very close to 1. Now, that being the case, actually this formulae we have derive are not valid for this particular circuit unless we do something to bring this voltage back to zero bias. If you can bring this voltage across the collector base junction, that is, back to zero bias then all that we have discussed is valid. Now this is what is important for using the transistor as a small signal amplifier.

One simple way of doing that is you include a battery whose polarity is opposite to that of the battery here. That would mean we must include a battery which is positive on the base side and negative on the collector side. Now you can see that when you go like this then this voltage will compensate this voltage and if you choose this battery to be exactly equal to V_0 plus ΔV_0 then this voltage will return to 0. Let us call this voltage V_{cc} that is c stands for collector. So this is actually a collector voltage between collector and base. Now if V_{cc} is exactly equal to V_0 plus ΔV_0 then the collector base voltage is zero and then this formula is valid.

Now can you always maintain the V_{cc} is equal to V_0 plus ΔV_0 ?

Obviously this is not possible because ΔV_{EB} means you are changing the emitter base voltage and this is going to change the ΔI_C and ΔV_0 . So, as you go on changing your emitter base voltage your ΔV_0 will change so you have to keep changing your V_{cc} in conformity with the emitter base voltage. So, if V_{EB} for example is a sinusoid your V_{cc} will also have to be a sinusoid of appropriate amplitude because it must compensate this voltage. Now it is obviously not possible, you cannot have a battery whose voltage is going on changing with time, this is practically not possible so what you do?

Now this is where we make use of another result that we obtained in the previous lecture. That is, even if the collector base junction is reverse bias the transistor action is not seriously affected and we can still write I_C is equal to αI_E except that there is going to be a small current I_{C0} . But when we take increments ΔI_C in response to ΔI_E because of ΔV_{EB} there will not be any significant change in I_{C0} . This current is a reverse current of the collector base junction which is reverse bias.

Once a reverse bias is more than three times V_t any change in the character base reverse bias does not change this I_{C0} . And therefore we can still write ΔI_C is equal to α times ΔI_E when we take increments because ΔI_{C0} is 0 and if that is the case this formula will still hold even if this is reverse biased. Now, what this means is that we can choose a V_{cc} whose value is such that it compensates this voltage whenever it is maximum. So this V_{cc} is equal to V_0 plus ΔV_0 the maximum value of this. And when

this voltage reduces because of change in the emitter base voltage then the junction will get reverse bias to a certain extent. For example, let us take V_{CC} is equal to V_0 plus delta V_0 . Supposing this is what we have chosen and when your delta V_{EB} increment this maximum then this is the voltage and at that point this is zero bias.

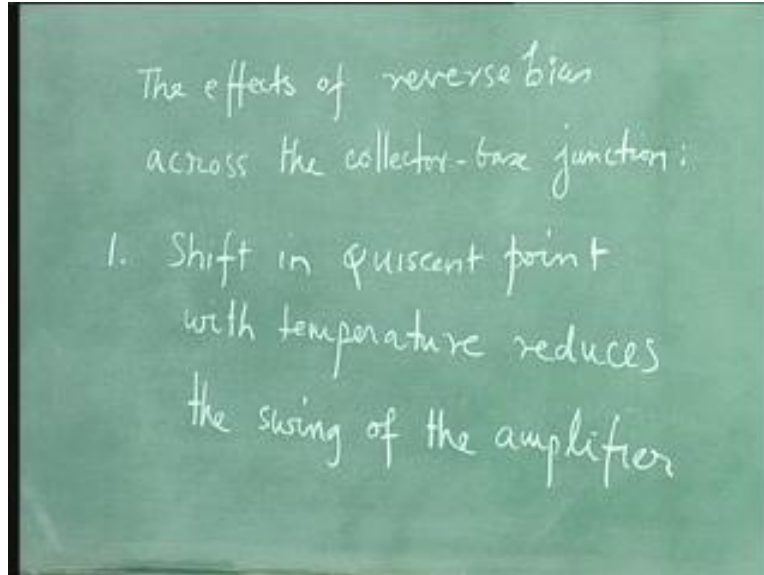
Now suppose this voltage becomes V_{EB} and therefore this current becomes I_C and therefore this becomes V_0 then you can see that a bias equal to delta V_0 will appear here but it would be reverse bias because this voltage is more than this voltage, the bias with this polarity and under that condition this equation will be valid but for incremental purposes same old equation we considered will be valid and therefore our entire analysis will hold. Therefore the moral of the whole story is that you will have to include a power supply in this lead whose magnitude is equal to the maximum voltage drop across the resistor then your emitter base voltage is changing. Now this is the circuit that will behave and it will give you amplification. Is there any effect of reverse bias across a collector on the amplification? We said that as far as this equation is concerned the collector current is equal to alpha times I_E plus I_{C0} .

The I_{C0} is the only addition as far as DC is concerned. Now, as far as AC is concerned, that is when you make increments in the currents is the amplification not affected by the presence of a change in reverse bias across the collector base junction? So let us examine this issue in little detail. The effects of reverse bias across the collector base junction. There are two main effects we will show. The first is the following. Now because of the presence of I_{C0} the difficulty is that if the temperature goes on changing which is what can happen in practice then this I_C will be changing with temperature even if you keep your I_C constant by maintaining V_{EB} constant.

Supposing you consider the situation when there are no increments so you have set up a V_{EB} here and as a result you have a current set up in the transistor. Now, for some reason the temperature starts changing, this happens in practice. What are the reasons because of which temperature can change? One reason is that the transistor is dissipating power and this power is dissipated as heat. After all you have voltages and currents in the transistor and obviously the voltage into current that much power is dissipated in the device and that power is dissipated as heat therefore the temperature of the transistor can raise.

Similarly, ambient temperature can rise. The temperature of the room in which you are setting up this device as an amplifier can rise. So because of such reasons temperature can change and in such a case this I_{C0} will change rapidly with temperature. We know that I_{C0} is the reverse saturation current of a PN junction and as we have explained it doubles approximately for around every 8 to 10 degree C rise in temperature for silicon diodes. So, for silicon transistors also similarly I_{C0} will change rapidly with temperature. Now as a result of this the bias point of the transistor under DC conditions can shift which is also called as the quiescent point. So it is the shift in quiescent point with temperature.

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I want to emphasize what is meant by the quiescent point. Quiescent point means the collector current and the collector base voltage when there is no signal when there is no disturbance. The word quiescent means calm. This means the condition when there is no ΔV_{EB} . When there is no signal you have a certain collector current here and then you have a certain collector base voltage. Both these I_C and V_{CB} is referred to as the quiescent point so what is happening is that when your temperature increases the I_{C0} changes rapidly as a result of which I_C changes because of which V_0 changes because of which the V_{cc} minus V_0 which is the voltage drop here goes on changing.

What is the effect of that?

The effect is that your swing capacity of the amplifier is reduced. That is the maximum signal amplitude that can be amplified is reduced. What is the maximum signal that you can amplify? This depends on the difference between V_{cc} and V_0 when there is no signal. If the I_C is more this voltage drop or V_0 is more the difference between this V_{cc} and V_0 is less. Therefore you can only increase V_0 until it becomes equal to V_{cc} . If V_0 becomes more than V_{cc} then this will be forward bias and the transistor action will be destroyed. So the difference between V_{cc} and V_0 that is the difference between the voltage drop across the resistor and the collector power supply represents the maximum voltage that you can get at the output as result of the signal amplification. So signal voltage that you can get at the output depends on the difference between V_{cc} and voltage drop across R which depends on I_C .

If I_C is more and that difference is less then the maximum swing of the amplifier, the maximum input voltage that it can amplify or the maximum output signal voltage you can get is reduced. So shift in quiescent point temperature reduces the swing of the amplifier. This is one effect of the reverse bias across the collector base junction. This is happening because this I_{C0} cannot be control by any means. This is dependent on temperature and this cannot be controlled. And even if I want to maintain I_E constant I cannot control this

I_{C0} . So this lack of control of this current that is coming about because of reverse bias is the cause of the shifting quiescent point with temperature.

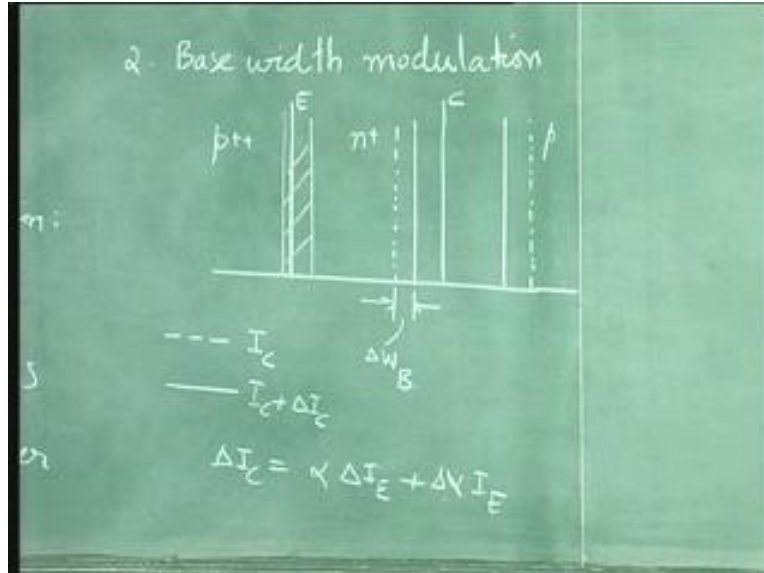
The next important thing that happens is the reduction in amplification because of what is called base width modulation. Let us understand what is this base width modulation. The base width modulation is the change in W_B with change in V_{EB} or I_C because of change in the depletion width here. The voltage drop across this junction depends on the difference in the voltage of the power supply and the voltage drop across R . When you introduce a signal V_{EB} changes, I_C changes, this voltage changes and this voltage here is changing. So when the amplification is taking place and signal amplification is taking place collector base voltage is changing with time. Therefore the depletion width is also changing with time.

Now because of a change in depletion width the W_B is also changing. So this variation in W_B in the presence of signal is what is called base width modulation. Let us see how the base width modulation can reduce the amplification. When the collector current is at the maximum the difference in these voltages is the minimum here. Therefore the depletion width across the collector junction is small. When the depletion width is small the base width is large. Let us show this effect by exploding this particular portion of the diagram. So we will exaggerate the collector base depletion width for one. This is the collector junction, this is the emitter junction, this is the collector base depletion width. This is emitter p plus plus, n plus and p. So this is the collector depletion width for one reverse bias across the collector base junction.

This is another collector depletion width for another reverse bias across the collector base junction a higher reverse bias. So, when the collector current is more V_{CB} is less so let us say this dotted line corresponds to the current I_C then the solid line depletion-region will correspond to I_C plus ΔI_C because when the collector current is more this voltage is more and the difference in this voltage which is this voltage is less so the depletion width is less. So I_C plus ΔI_C is solid line, this and I_C is the dotted line. Therefore this is the difference in the width of the base region. So we have not shown the emitter depletion layer here but you can show that. The change in this is small so we are not showing the change here. So this difference from here to here is the base width so this is the ΔW_B .

Now this is the base width modulation effect. As a result of this base width modulation your amplification will reduce. How do we show that? We know that the change in the base width is going to affect the α of the transistor. So, in this particular equation ΔI_C is equal to $\alpha \Delta I_E$ is not exactly correct when the base width modulation is present because this assumes that α is constant when you make the change. So the correct equation in the presence of base width modulation is ΔI_C is equal to $\alpha \Delta I_E$ plus $\Delta \alpha I_E$. Now, depending on this $\Delta \alpha$ your ΔI_C in presence of base width modulation can be more or less. For example, if this $\Delta \alpha$ is negative then this term will subtract from this and therefore in the presence of base width modulation your increment in the collector current will be less because of increment in the emitter base voltage.

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We will exactly show what happens, delta alpha is negative when you increase the collector current the alpha decreases and delta alpha is negative. How do we show that? For this purpose we must derive an equation for alpha. This equation can be derived as follows. The equation for alpha is gamma into b where gamma is injection efficiency and b is base transport factor where gamma is $1 / (1 + I_{En} / I_{Ep})$ and B is $1 - I_r / I_{Ep}$. The base transport factor b is nothing but I_{Cp} the hole current reaching the collector divided by the hole current injected from the emitter. And that hole current reaching the collector is nothing but the current injected from the base minus the recombination taking place in the base. So, from there we get this relation. Now you can write equations for these terms using the equations for minority carrier concentration.

Now I_{En} / I_{Ep} will depend on Δn_E and Δp_E . We can write this as Δn_E into q into D_E by L_E . I shift this L_E here; I_{Ep} is $q \Delta p_E$ into D_B by W_B . Of course you also have the emitter areas coming both in numerator and denominator. Now this q cancels and A_E cancels and Δn_E by Δp_E we can write in terms of the doping levels. So Δn_E by Δp_E will be in the reverse ratio of the doping in the emitter and base. That is, Δp is inversely proportional to doping here. So doping in the base will come in the numerator and doping in the emitter will come in the denominator. So I_{En} / I_{Ep} is, we can remove this and write it the other way. So, by transferring that information here this is equal to $1 / (1 + N_B D_E W_B / N_E D_B L_E)$.

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Equation for α .

$$\alpha = \frac{I_c}{I_e} = \frac{I_{Eh}}{I_{Ep} + I_{Er}}$$

$$= \frac{1}{1 + \frac{I_{Er}}{I_{Ep}}} \left(1 - \frac{I_{Er}}{I_{Ep}} \right)$$

$$= \frac{1}{1 + \frac{N_B D_E W_B}{N_E D_B L_E}} \left(1 - \frac{I_{Er}}{I_{Ep}} \right)$$


$$\frac{I_{Eh}}{I_{Ep}} = \frac{N_B D_E W_B}{N_E D_B L_E}$$

Similarly we can write I_r by I_{Ep} . I_r is the recombination current which is related to this area so we can write I_r by I_{Ep} here. Now this area is half of Δp_E into W_B this is the area under the triangle this difference is W_B the base width.

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$\left(1 - \frac{I_{Er}}{I_{Ep}} \right)$



$$\frac{I_{Er}}{I_{Ep}} = \frac{\left(\frac{1}{2} \Delta p_E W_B A_E / \tau_B \right) W_B}{\frac{1}{2} \Delta p_E D_B A_E}$$

$$= \frac{W_B^2}{2 D_B \tau_B} = \left(\frac{W_B^2}{2 L_B^2} \right)$$

Now this area is multiplied by the charge q multiplied by the area of the emitter A_E . So A_E into W_B is the volume and the recombination current will depend on this stored charge divided by the lifetime in this region. Let us call the lifetime in the base region as τ_B so charge by lifetime this is from the law of the junction, it is a recombination in this region.

So divided by I_{EP} we can rewrite that equation which we wrote for I_{EP} as $Q \Delta p_E D_B$ into W_B because the slope of this line is ΔP by W_B so this is ΔP by W_B .

Of course the area of the emitter also comes there. Now we can see that we can cancel this q , we can cancel this Δp_E , can cancel this A_E and then this is equal to W_B square by $2 D_B$ into τ_{EB} which is nothing but W_B square by $2 L_B$ square and we got W_B square because of this W_B into W_B . So I_r by I_{EP} is W_B square by $2 L_B$ square. So we can write this here. This is the expression for α and it shows how α depends on base width W_B . You can clearly see that if W_B is much less than L_B this quantity is negligible and therefore it becomes close to 1.

Similarly, if W_B is much less than W_B and also if the base doping is much less than emitter doping this quantity is very small and again this tends to 1 and that is how α tends to 1 for N_B much less than N_E and W_B much less than W_B . Now what is important to see from here is that, as your W_B increases your α this term will increase and also this hole term will become less and therefore your α is going to reduce. So, as W_B increases α reduces. Therefore because of increase in W_B the increment in α will be negative. So when you increase ΔI_C what you have seen is that the collector base voltage reduces which means the depletion width reduces and ΔW_B therefore is increasing. So, base width modulation is increased in the base width when your collector current is increasing. And increase in the base width causes α to fall. Therefore $\Delta \alpha$ in response to ΔI_C will be negative.

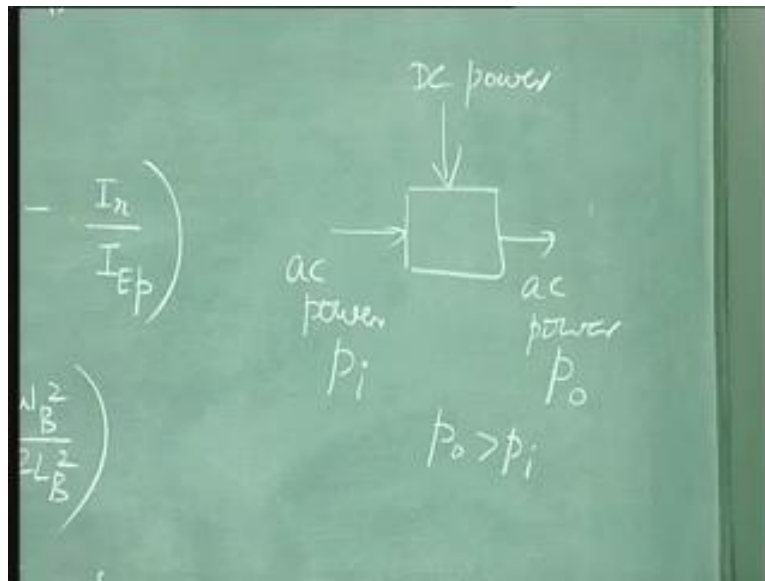
So we can write this as $\alpha \Delta I_E$ minus modulus of $\Delta \alpha$ into I_E . Now this is a positive quantity subtracting from this quantity and therefore it clearly shows ΔI_C is negative. Therefore the increment in the collector current is reduced, it is not negative it is still positive because this is more than this but the increment in the collector current is reduced. Therefore your amplification is affected it is reduced because amplification depends on ΔI_C . This explains how base width modulation reduces amplification.

We can write this statement as W_B increases α falls therefore the base width modulation we will abbreviate as BW_M base width modulation. Therefore BW_M reduces amplification. These are the two effects of the reverse bias. So, base width modulation affects amplification to some extent and also it results in change in the quiescent point width temperature which affects the swing. But for these two effects your gain of the transistor is quite large as we have shown that the voltage gain is large and therefore the device works very well as amplification.

Now please note that it is a small signal amplifier and we are only amplifying small signals. So you are super imposing a small signal over a DC voltage and you are amplifying that small signal or disturbance. Now, if there is amplification then it means there is a power gain. So how can there be gain in power? Your output signal power is more than input signal power. How can this happen because conservation of energy should be there. It is to be understood here that there is an increase in the ac power. If you see the input power of the small signal here and compare that with the small signal output power the output small signal power is more than input small signal power. But this

output extra small signal power is coming from the power supply. So what the amplifier is doing is it is converting DC power into ac power that is why it is called an active device. Transistor is called an active device because it converts DC power into ac power. The diagram is something like this. You have a DC power as input you also have an ac power as input. Now you have ac power output.

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Now this small signal ac power p_o and if this is p_i p_o is greater than p_i . But there is this DC power that is coming in. So you are supplying energy from the power supply. And one should not think that you are getting something out of nothing. You are getting ac power amplification but definitely you are supplying DC power and only in the presence of DC power the amplification can take place. So unless you have DC conditions maintained you have V_{EB} , I_C and V_{CC} you cannot have this small signal amplification. With this we complete the important application of the transistor action namely small signal amplification. We have explained how you can get the voltage gain. But you can also show how you can get a current gain, and we will see this in the next class.