

Power System Analysis
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Lecture - 26
Load Flow Studies

And will come back.

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which is commonly referred to as the Surge Impedance.
 Its value varies between 250 Ω and 400 Ω in case of overhead transmission lines and between 40 Ω and 60 Ω in case of underground cables.

From eqns (54) and (58), we see,

$$v = \frac{\omega}{\omega \sqrt{LC}} = \frac{1}{\sqrt{LC}}$$

$\rightarrow \therefore v = \frac{1}{\sqrt{LC}} \dots (60)$

and

$$\lambda = \frac{2\pi}{\beta} = \frac{2\pi}{\omega \sqrt{LC}} = \frac{2\pi}{2\pi f \sqrt{LC}}$$

$\rightarrow \therefore \lambda = \frac{1}{f \sqrt{LC}} \dots (61)$

$\therefore \beta = \omega \sqrt{LC}$
 [from eqn (58)]

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$$\frac{dx}{dt} = \frac{\omega}{\beta} \dots (55)$$

Thus, the velocity of propagation is given by

$$v = \frac{\omega}{\beta} = \frac{2\pi f}{\beta} \dots (56)$$

A complete voltage cycle along the line corresponds to a change of 2π radian in the angular argument βx . The corresponding line length ($x = \lambda$) is defined as the wavelength. If β is expressed in rad/m ,

$$\beta \lambda = 2\pi$$

$\therefore \lambda = \frac{2\pi}{\beta} \dots (57)$

So, now, from equation 56 and 58 right we will get this equation 56 is this one, $2 \pi f$ upon beta.

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When line losses are neglected, i.e., $g=0$ and $r=0$, then real part of the propagation constant $\alpha=0$.

From eqn. (27)

$$\rightarrow \gamma = \alpha + j\beta = \sqrt{ZY} = \sqrt{(r+j\omega L)(g+j\omega C)}$$

$$\therefore \alpha + j\beta = \sqrt{j^2 \omega^2 LC} = j\omega \sqrt{LC}$$

$$\therefore \alpha = 0,$$

$$\rightarrow \beta = \omega \sqrt{LC} \quad \dots (58)$$

From eqn. (29), the characteristic impedance,

$$Z_c = \sqrt{\frac{Z}{Y}} = \sqrt{\frac{r+j\omega L}{g+j\omega C}} = \sqrt{\frac{L}{C}} \quad \dots (59)$$

And equation 58 we have seen beta is equal to omega root over l c right. So, therefore, v is equal to omega by beta actually and it is beta is equal to omega root over l c. So, v is equal to 1 upon root over l c right therefore, this is equation we have marked v is equal to 1 upon root over l c 60 equation 60 right. Now lambda is equal to we have seen beta lambda is equal to 2 pi. So, lambda is equal to 2 pi by beta and it is 2 pi and beta is equal to omega root over Lc. So, omega 2 pi f. So, it is 2 pi upon 2 pi f root over Lc. So, 2 pi 2 pi will be cancel and lambda is equal to 1 upon f in to root over l c this is the wave length, this is equation this is the equation of the wave length and this is equation 61 right; now will make some approximation.

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Now for a single phase line,

$$L = \frac{\mu_0}{2\pi} \ln\left(\frac{D}{r'}\right) \quad \text{and} \quad C = \frac{2\pi\epsilon_0}{\ln\left(\frac{D}{r'}\right)}$$
$$\therefore LC = \frac{\mu_0 \epsilon_0 \ln\left(\frac{D}{r'}\right)}{\ln\left(\frac{D}{r'}\right)}$$

Approximating $\ln\left(\frac{D}{r'}\right) \approx \ln\left(\frac{D}{r}\right)$

$$\therefore LC = \mu_0 \epsilon_0 \dots (62)$$

Substituting the expression LC from eqn

So, now for a single phase line while we are conducting inductance and capacitance, we got L is equal to μ_0 by 2π $\ln D$ upon r dash right. This we got this we are for the inductance calculation we got this right and capacitance C is equal to 2π ϵ_0 $\ln d$ upon r this 2 we got. Now multiply this 2 L and C , if you multiply LC is equal to is a μ_0 ϵ_0 $\ln d$ upon r dash divided by $\ln D$ r , D by r right.

Now will make an approximation that will assume will approximating it that $\ln D$ upon r dash the natural of D upon r dash is equal to $\ln D$ upon r this is an approximation you are making it an approximation right. If it is so, if this 2 are equal approximately if you so.


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$$L = \frac{\mu_0}{2\pi} \ln\left(\frac{D}{r'}\right) \quad \text{and} \quad C = \frac{2\pi \epsilon_0}{\ln\left(\frac{D}{r'}\right)}$$
$$\therefore LC = \frac{\mu_0 \epsilon_0 \ln\left(\frac{D}{r'}\right)}{\ln\left(\frac{D}{r'}\right)}$$

Approximating $\ln\left(\frac{D}{r'}\right) \approx \ln\left(\frac{D}{r}\right)$

$$\therefore LC = \mu_0 \epsilon_0 \quad \dots (62)$$

Substituting the expression LC from eqn. (62) into eqn. (60) and (61), we get,



Then LC actually is equal to $\mu_0 \epsilon_0$ this is equation 62 right. So, if substituting this expression of LC from equation 62 I mean this one in to equation 60 and 61 right; that means, if you substitute this equation in equation 60, that is v is equal to $1/\sqrt{LC}$ and in 61 by λ is equal to $1/f\sqrt{LC}$, you substitute here right that your LC is equal to $\mu_0 \epsilon_0$ so v is equal to $1/\sqrt{\mu_0 \epsilon_0}$.

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$$v = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \quad \dots (63)$$


and

$$\lambda = \frac{1}{f\sqrt{\mu_0 \epsilon_0}} \quad \dots (64)$$

Substituting for $\mu_0 = 4\pi \times 10^{-7} \text{ H/m}$ and $\epsilon_0 = 8.854 \times 10^{-12} \text{ F/m}$

$v \approx 3 \times 10^8 \text{ m/sec}$, i.e. approximately velocity of light at 60 Hz,

the wave length $\lambda \approx 5000 \text{ km}$.



If you if you do so, then μ will become 1 upon root over $\mu_0 \epsilon_0$ and λ will become 1 upon f root over $\mu_0 \epsilon_0$ right. So, you know the μ_0 value 4π in to 10 to the power minus 7 henry per meter, and you know that ϵ_0 also 8.854 in to 10 to the power minus your 12 farad per meter right. So, approximately v will be 3 in to 10 to the power 8 meter per second, which is approximately the velocity of light at frequency 60 hertz right. So, this is that your velocity and wave length λ will become 5000 kilo meter. If you put in that second equation here right λ will become 5000 kilo meter right this is more or less some ideas regarding your velocity your velocity of propagation and your wave length λ right.

But this is approximate this values are approximate where because you have taken l and d upon r dash approximately is equal to l and d r upon r right, but this will give a good idea about this thing about your v and λ right.

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For a lossless line $\gamma = j\beta$ and the hyperbolic functions

$$\cosh(\gamma x) = \cosh(j\beta x) = \cos(\beta x) \text{ and}$$

$$\sinh(\gamma x) = \sinh(j\beta x) = j\sin\beta x.$$

The equations for the rms voltage and current along the line, given by eqn. (38) and (39) become,

$$V(x) = \cos(\beta x) V_R + jZ_c \sin(\beta x) I_R \quad \dots (65)$$

$$I(x) = j\frac{1}{Z_c} \sin(\beta x) V_R + \cos(\beta x) I_R \quad \dots (66)$$

At the sending end $x = l$, $V(x) = V(l) = V_S$ and $I(x) = I(l) = I_S$,

$$\therefore V_S = \cos(\beta l) V_R + jZ_c \sin(\beta l) I_R \quad \dots (67)$$

$$I_S = j\frac{1}{Z_c} \sin(\beta l) V_R + \cos(\beta l) I_R \quad \dots (68)$$

Now, for a loss less line we have seen for a loss less line γ is equal to velocity of propagation α plus j beta. So, α is 0. So, for a loss less line γ is equal to j beta right and the hyperbolic function \cosh hyperbolic γx actually will become that if you substitute γ is equal to j beta for a loss less line, it will be \cos hyperbolic j beta x is equal to nothing, but \cos beta x right and \sinh hyperbolic γx is equal to \sin hyperbolic j beta x is equal to j \sin beta x right.

So, is the simply $j \sin \beta x$. So, if you take those expression and put this value definitely, you will get the cos hyperbolic $j \beta x$ is equal to $\cos \beta x$ details are given while developing long transmission line that is why not showing it here I understandable right; the equations for the r m s voltage and current along the line is given by equation your 38 and 39 so equation 38 actually this and 39 this 2 equations rewriting here right. $V(x)$ is equal to cos in that case instead of cos hyperbolic x , we are putting $\cos \beta x$. For a loss less line just now we have written here we have written here cos hyperbolic x for a loss less cos hyperbolic γx is equal to basically $\cos \beta x$, because α is equal to 0 and γ is simply $j \beta$ and cos hyperbolic $j \beta x$ is equal to $\cos \beta x$.

So, instead of cos hyperbolic we are writing $\cos \beta x$ V_R plus Z_c , $j Z_c$ rather $j Z_c \sin \beta x$ I_R ; because sin hyperbolic sin γx hyperbolic γx is equal to sin hyperbolic $j \beta x$ is equal to $j \sin \beta x$. Similarly here also right this is all this you just replace those hyperbolic term you will get $I(x)$ is equal to that is equation 38 and 39 in equation 38 you put that because in 39 you put for $I(x)$ right, then you will get $j \frac{1}{Z_c} \sin \beta x$ upon V_R plus $\cos \beta x$ I_R this is equation 66 this 165 and this 166 right.

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$\cosh(\gamma x) = \cosh(j\beta x) = \cos(\beta x)$ and
 $\sinh(\gamma x) = \sinh(j\beta x) = j \sin \beta x$,
 The equations for the rms voltage and current along the line, given by eqn. (38) and (39) become,

$$V(x) = \cos(\beta x) V_R + j Z_c \sin(\beta x) I_R \quad \dots (65)$$

$$I(x) = j \frac{1}{Z_c} \sin(\beta x) V_R + \cos(\beta x) I_R \quad \dots (66)$$
 At the sending end $x = l$, $V(x) = V(l) = V_S$,
 and $I(x) = I(l) = I_S$.

$$\therefore V_S = \cos(\beta l) V_R + j Z_c \sin(\beta l) I_R \quad \dots (67)$$

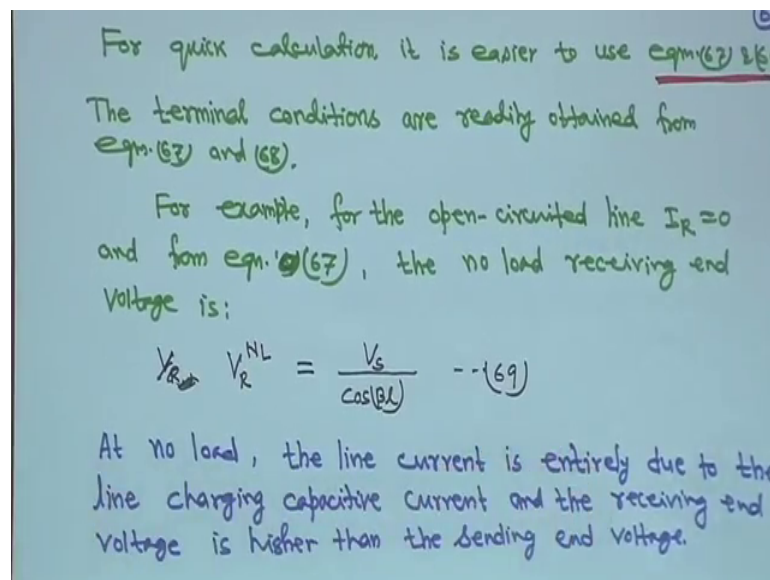
$$I_S = j \frac{1}{Z_c} \sin(\beta l) V_R + \cos(\beta l) I_R \quad \dots (68)$$

Now, at the sending end because you are made it in the distance from the receiving end to sending end, so at the sending end x is equal to l . So, $V(x)$ equal to $V(l)$ is equal to sending end voltage v_s right similarly $I(x)$ is equal to $I(l)$ is equal to I_s right this one. So,

sending end voltage can be written as $\cos \beta l V_R + j Z_c \sin \beta l I_R$ this is 67 and I_s is equal to $j \omega C \sin \beta l V_R + \cos \beta l I_R$ right.

So, actually for quick calculations for a b c d parameters this equation is I think less laborious I mean easier to compute right this equation 67, but 68, but remember this is for a loss less line right. Now that is why I have written for quick calculation it is easier to use 67 and equation 67 and 68 right.

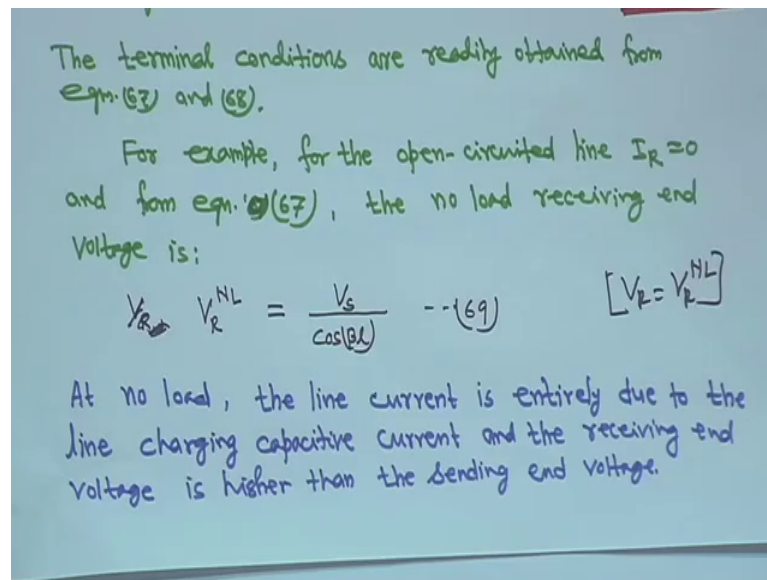
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Now, terminal conditions are readily obtained from equation 67 and 68; that means, all this terminal conditions you can easily get it from equation 67 and 68 right. So, for example, for the open circuit line that receiving end current is 0, $I_R = 0$ when line is open circuit right and from equation 67 right.

That means, from this equation, equation 67 this equation you put I_R is equal to 0 in equation 67 here you put this is for open circuit right if you do so, the no load receiving end voltage will become at that time at the line is open circuit. So, this I mean condition no load condition is that while line is open circuit when I_R is equal to 0 right in this equation at no load when I_R is equal to 0, V_R is equal to V_R^{NL} right.

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So, when you are telling that when I_R is equal to 0; here I am somewhere I am writing that V_R is equal to V_R^{NL} right. So, I am not writing here, but understandable that when open circuit line I_R is equal to 0 and V_R is equal to V_R^{NL} that will be easier to understand therefore, this equation that equation 67, it will become that V_s is equal to $\cos(\beta l)$ and V_R is equal to V_R^{NL} because I_R is 0 therefore, V_R^{NL} is equal to V_s by $\cos(\beta l)$ right. So, that is what I have writing V_R^{NL} is equal to V_s upon $\cos(\beta l)$ right that is equation 69. So, what does it signify? That at no load condition that receiving end voltage right is greater than the sending an voltage because $\cos(\beta l)$ is less than 1 right.

This is some time this is called Ferranti effect and this is happening because of your charging admittance in the line right. So, no load condition this is happening. So, that is why V_R^{NL} is equal to V_s upon $\cos(\beta l)$ this approximate, but in the same your long line expression also same logic you can use and you will find that V_R^{NL} will be greater than the sending end voltage at no load right.

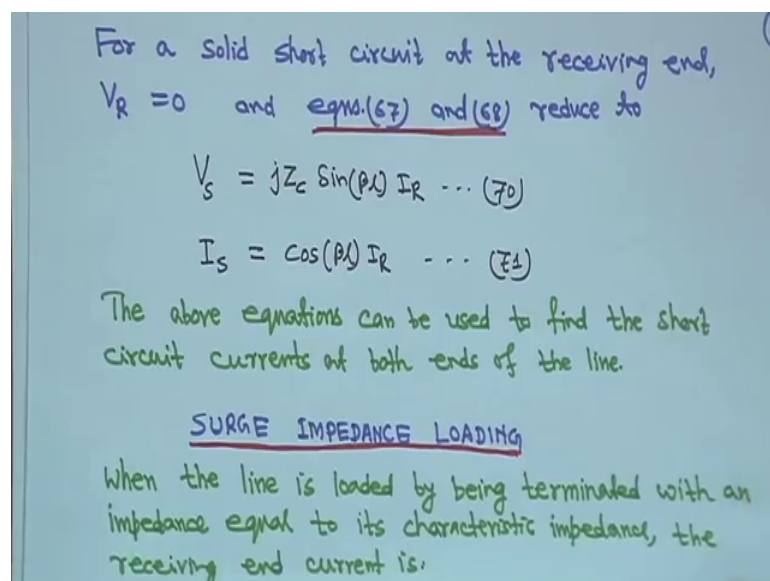
So, that is why at no load the line current is entirely due to the charging capacity capacitive current, and the receiving end voltage is higher than the sending end voltage because when line is no load there is no I said line no line is there is no load connected here right, but charging capacitance are there. So, because of that your what you call this that this receiving no load voltage is becoming higher than the sending end voltage. So,

suppose a transmission line this a question to you, suppose if transmission line is your this thing what you call not carrying any power under no load condition right and it is not carrying any power.

Suppose a man standing on the ground with bare feet say with no proper protection or insulation and taking a conductive knot and touching that conductor, what will what will happen to that parcel. I repeat this suppose a line transmission line high tension transmission line was carrying power, but after that that load is disconnected. So, now, this condition whatever will putting that it is that I R is 0 right it is not, but a man standing on the ground right with bare feet say right and taking a conductive rod and touch that conductor that is the line right.

What will happen to that parcel this is a question to you right. Later you can send this when you listening this you send your answer to me right, but here I am not giving that answer right.

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So, this one next for a solid short circuit right for a solid short circuit at the receiving end solid means there is no fault impedance fault will come later right much later. So, for a solid short circuit at the receiving end the V_R is equal to 0 right and in equation 67 and 68 that is your this equation 67 and 68 this 2 equation you put that condition that for V_R is equal to 0 right. If you do so, then you will get sending end voltage is equal to $j Z_c \sin \beta l I_R$ this is equation 70 and I_S is equal to $\cos \beta l I_R$ this is equation 71. The

above equation actually can views to find the short circuit current at both ends of the line right.

Next is that your surge impedance loading when the line is loaded by being terminated with an impedance equal to its characteristic impedance that is Z_c is equal to Z_c is equal to $\sqrt{L/Y}$ right the receiving end current you can define; that means, when a line is loaded by being terminated right with an impedance equal to its characteristic impedance, the receiving end current you be receiving end voltage is V_R and it is and that load.

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Handwritten notes on a blue background showing equations for surge impedance loading (SIL). The text is written in green and red ink. A small number '67' is written in the top right corner.

$$I_R = \frac{V_R}{Z_c} \dots (72)$$

For a lossless line Z_c is purely resistive. The load corresponding to the surge impedance at rated voltage is known as the surge impedance loading (SIL), given by

$$\rightarrow SIL = 3 V_R I_R^* = \frac{3 |V_R|^2}{Z_c} \dots (73)$$

Since $V_R = \frac{V_{L, rated}}{\sqrt{3}}$, SIL in MW becomes

$$\rightarrow SIL = \frac{(kV_{L, rated})^2}{Z_c} \text{ MW} \dots (74)$$

We learn that line is loaded by being terminated with an impedance equal to your characteristic impedance that is Z_c ; that means I_R will be V_R upon Z_c this is equation 72. The receiving end voltage V_R and line is terminated with in your what you call with an impedance equal to Z_c .

So, I_R will become V_R upon Z_c this is equation 72 right. Now for a loss less line Z_c is purely resistive because you have seen Z_c is equal to $\sqrt{L/Y}$. So, the load the corresponding the surge impedance at rated voltage is known as the surge impedance loading that is SIL right. So, this load. So, for Z_c is purely the load corresponding to the surge impedance at rated voltage is known as surge impedance loading the loading they SIL and is given by SIL is equal to $3 V_R I_R$ conjugate, this we have seen also I have shown you separately how this $V_R I_R$ conjugate or V_R conjugate

I R it come ins right. So, this is $3 V_R I_R$ conjugate is equal to this I R is equal to V_R upon Z_c .

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Handwritten derivation on a chalkboard:

$$3V_R I_R^* = 3V_R \cdot \left(\frac{V_R}{Z_c}\right)^* = 3 \frac{V_R \cdot V_R^*}{Z_c} = \frac{3V_R^2}{Z_c}$$

$$I_R = \frac{V_R}{Z_c} \dots (72)$$

For a lossless line Z_c is purely resistive. The load corresponding to the surge impedance of rated voltage is known as the Surge impedance loading (SIL) given by

So, this is actually $3 V_R I_R$ conjugate right and I R is equal to V_R upon Z_c . So, it is V_R and your V_R upon Z_c conjugate right multiplied by three. This is $3 V_R$, V_R conjugate and Z_c because Z_c is a real quantity. So, V_R in to V_R conjugate it is $3 V_R$ square divided by Z_c right. So, that is why you are writing that this one is equal to SIL is equal to 3 in to voltage magnitude V_R square upon Z_c . This is equation your what you call this thing and V_R is your phase voltage that is why this is multiplied by 3 right.

So, since V_R is equal to V_L rated upon root 3 that is V line rated voltage you are taking here in general. So, SIL in megawatt, it becomes SIL is equal to KVL , KVL rated whole square upon Z_c megawatt what we are doing is V_R is equal to V_L rated upon root 3 we know, but this rated voltage we are actually making it to kilo volt that is why I am writing KVL rated then suppose it is suppose KVL rated means suppose 220.

Then 220 square upon Z_c whatever will be there this is kilo volt right this is kilo volt and that is square and this Z_c is the ohmic value ohm. So, this is actually megawatt right because it is kilo volt square and this is actually equation 74.

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Substituting $I_R = \frac{V_R}{Z_c}$ [Eqn. (72)] in ~~eqn (65)~~ (68)

in eqn (65)

$$V(x) = \cos(\beta x) V_R + j Z_c \sin(\beta x) I_R$$

$$\therefore V(x) = \cos(\beta x) V_R + j Z_c \sin(\beta x) \times \frac{V_R}{Z_c} = (\cos(\beta x) + j \sin(\beta x)) V_R$$

$$\rightarrow V(x) = V_R \angle \beta x \dots (75)$$

and $\frac{V_R}{Z_c} = I_R$ in eqn (66)

$$I(x) = j \frac{1}{Z_c} \sin(\beta x) V_R + \cos(\beta x) I_R$$

$$\therefore I(x) = j \sin(\beta x) I_R + \cos(\beta x) I_R$$

$$I(x) = (\cos(\beta x) + j \sin(\beta x)) I_R$$

$$\rightarrow I(x) = I_R \angle \beta x \dots (76)$$

Now, substituting you this is I R is equal to V R upon Z c that is that I that I told you right in equation 65. So, in equation this is actually equation 65 I have rewritten here right. So, cosine beta x V R plus Z c sin beta x I R. Now this I R is equal to V R upon Z c you substitute here I R is equal to V R upon Z c, then Z c Z c will be cancel I have substituted here that this is V R upon Z c. So, Z c Z c will be cancel and it will become actually cosine beta x plus j sin beta x in to V R that means, V x is equal to V R angle beta x this is equation 75.

Similarly that your equation your what you call that in equation see this V R same thing here writing it other way that V R upon Z c is equal to I R you substitute it in equation 66.

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in eqn(65)

$$V(x) = \cos(\beta x) V_R + j Z_c \sin(\beta x) I_R$$

$$\therefore V(x) = \cos(\beta x) V_R + j Z_c \sin(\beta x) \times \frac{V_R}{Z_c} = \{\cos(\beta x) + j \sin(\beta x)\}$$

→ $V(x) = V_R | \beta x \dots (75)$

and $\frac{V_R}{Z_c} = I_R$ in eqn(66)

$$I(x) = j \frac{1}{Z_c} \sin(\beta x) V_R + \cos(\beta x) I_R$$

$$\therefore I(x) = j \sin(\beta x) I_R + \cos(\beta x) I_R$$

$$\therefore I(x) = \{\cos(\beta x) + j \sin(\beta x)\} I_R$$

→ $I(x) = I_R | \beta x \dots (76)$

So, equation 66 I am writing here that $I(x)$ is equal to $\frac{1}{Z_c} \sin(\beta x) V_R + \cos(\beta x) I_R$. So, $\frac{V_R}{Z_c}$ is equal to I_R you substitute. If you do so, that skipping this line same thing just rearranging $I(x)$ is equal to $\cos(\beta x) I_R + j \sin(\beta x) I_R$ in to I_R ; that means, $I(x)$ is equal to $I_R \cos(\beta x)$; that means, this equation 76, $V(x)$ is equal to $V_R \cos(\beta x)$ and here in 76 $I(x)$ is equal to $I_R \cos(\beta x)$.

That means, whatever may be the value of x right that V_R and I_R at every point of x that $v(x)$ and V_R that value remain same right that whatever at every whatever may be the value at any point of x $V(x)$ is equal to V_R , V_R is total independent of x and here also I_R total independent of x . So, $I(x)$ is equal to this I_R and V_R will not change even x is change right at very at every point this $v(x)$ and $I(x)$ both remain same right.

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Eqs (75) and (76) show that in a lossless line under surge impedance loading, the voltage and current at any point along the line are constant in magnitude and are equal to their sending end values.

Since Z_c has no reactive component, there is no reactive power in the line, i.e., $Q_s = Q_r = 0$.

This indicates that for SIL, the reactive losses in the line inductance are exactly offset by reactive power supplied by the shunt capacitance or $\omega L |I_x|^2 = \omega C |V_R|^2$. From this we find that $Z_c = \frac{V_R}{I_x} = \sqrt{L/C}$, which verifies the result in eqn (53)

That means it shows that in a lossless line under surge impedance loading, the voltage and current at any point along the line are constant in magnitude and are equal to their sending end values. If you put here x is equal to l ; that means, V_x means V_l is equal to this is the this will be the V_s is equal to $V_R \cos \beta l$ and this will become I_x is equal to rather I_s is equal to $I_R \cos \beta l$ right; that means, and in magnitude and are equal to their sending end values right since Z_c .

So, any point you take right. So, the Z_c has no reactive component we have seen that Z_c is equal to $\sqrt{L/C}$. So, there is no reactive component; that means, in the line Q_s is equal to Q_r is equal to 0 right. Under this condition this indicates for surge impedance loading the reactive losses in the your line inductance, are exactly off set by the reactive power supplied by the shunt capacitance right that is your charging capacitance. So, if it is true then your $\omega L |I_x|^2$ will become $\omega C |V_R|^2$ square right if this is true.

So, then your if you use this condition right. So, you will get Z_c is equal to V_R upon I_R if this 2 are equal then it will be $\sqrt{L/C}$; that means, that your what I will say right whatever your this thing that this inductance there the line inductance exactly off set by the power supplied by the shunt capacitance; that means, your this one that line inductance is $x L$ is equal to $L \omega$.

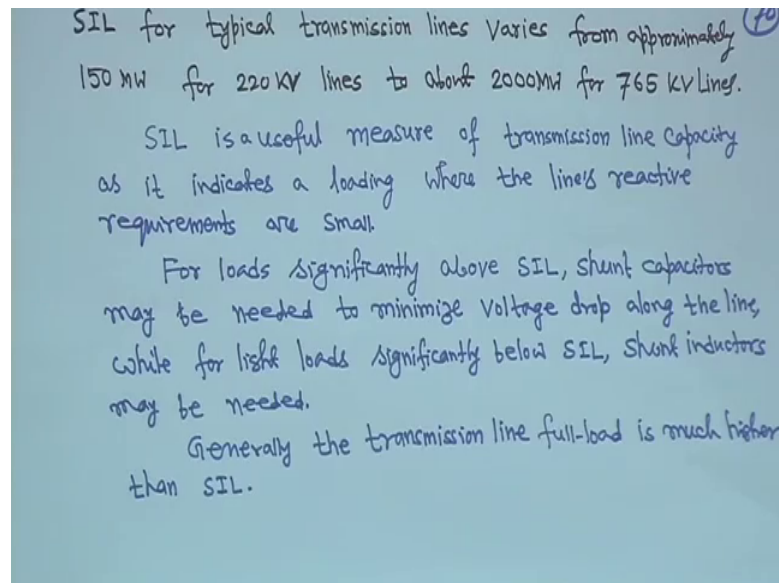
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$$X_L = \omega L$$
$$|Z_L|^2 X_L = \omega L |I_R|^2$$
$$X_C = \frac{1}{\omega C}$$
$$\frac{V_R^2}{X_C} = \omega C V_R^2$$

So, that is right; that means, you are what you call I R square in to your I R square x L is equal to l omega I R square.

So, that is why we are writing actually omega L I R square right and that charging admittance is omega c V R square right. So, in that case the power supplied or this reactive power supplied by the charging admittance actually, it is your V R square by x c right and you know x c is equal to 1 upon omega c. So, you substitute here you will become omega c V R square. So, that is why this total your reactive losses supplied by your charging capacitance. So, that is why omega l I square, is equal to omega c V R square from which from which we will get it that Z c is equal to V R upon I R is equal to root over L by c which verifies the result in equation 59.

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SIL for typical transmission lines varies from approximately 150 MW for 220 kV lines to about 2000 MW for 765 kV lines. SIL is a useful measure of transmission line capacity as it indicates a loading where the line's reactive requirements are small. For loads significantly above SIL, shunt capacitors may be needed to minimize voltage drop along the line, while for light loads significantly below SIL, shunt inductors may be needed. Generally the transmission line full-load is much higher than SIL.

So; that means, when you have your this thing j this kind of condition that there is no reactive power in the line right then what does it mean for surge impedance loading. So that means, for surge impedance first you see for surge impedance loading for a typical transmission line varies approximately from 150 megawatt for 220 k v lines this approximate values, to about 2000 megawatt for 765 kilo volt lines, but generally this line actually of when its operated at the full load or this, that it operates much about the surge surge impedance loading. So, SIL is useful measure of transmission line capacity as it indicates a loading where the lines reactive requirements are small right.

For load significantly above SIL shunt capacitor may be needed to minimize the voltage drop along the line right, generally line operates above this limit right. While for light load significantly below SIL you need basically shunt inductors right. So, if the if it is above the surge impedance loading, then you need shunt capacitor and if it is below the surge impedance shunt impedance loading, then you need shunt inductor right. So, generally the transmission line full load is much higher than the shunt what you call that surge impedance loading.

So, whatever now whatever this thing we have studied for this long transmission line that one is that surge impedance loading, then you have studied velocity of propagation then that wave length lambda right and that physical significance of the your surge impedance loading and the terminal conditions and then we saw that under no load condition that

charging this thing what you call receiving end voltage is higher than the sending end voltage this is typical phenomena for a transmission line right this is this happen you do your charging admittance right.

So, these are the things we have studied, but will take a small example of this your what you call velocity of propagation then surge impedance after that will go to for 3 phase your what you call power flow through transmission line right.

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Example-7 (7)

A three phase, 50 Hz, 400 kV transmission line is 300 km long. The line inductance is 0.97 mH/km per phase and capacitance is 0.0115 μ F/km per phase. Assume a lossless line. Determine the line phase constant β , Z_c , v and λ .

Soln.

$\beta = \omega \sqrt{LC}$
 $\therefore \beta = 2\pi \times 50 \sqrt{0.97 \times 10^{-3} \times 0.0115 \times 10^{-9}}$
 $\therefore \beta = 0.00105 \text{ rad/km.}$

Surge impedance,
 $Z_c = \sqrt{\frac{L}{C}} = \sqrt{\frac{0.97 \times 10^{-3}}{0.0115 \times 10^{-9}}} = 290.43 \Omega$

Velocity of propagation
 $v = \frac{1}{\sqrt{LC}} = \frac{1}{\sqrt{0.97 \times 10^{-3} \times 0.0115 \times 10^{-9}}}$
 $\therefore v = 2.994 \times 10^5 \text{ km/sec}$

Line wavelength is
 $\lambda = \frac{v}{f} = \frac{2.994 \times 10^5}{50}$
 $\therefore \lambda = 4988 \text{ km}$

And so, this one we take a small example right that is 3 phase 50 hertz 400 k v transmission line is 300 kilo meter long. The line inductance is 0.97 mili Henry per kilo meter right per phase and capacitance is 0.0115 micro farad per kilo meter per phase

So, assume here you assume a loss less line you have to compute beta, the line beta that is the line phase constant Z_c that characteristic impedance or surge impedance v velocity of propagation and lambda the wave length right. So, this thing solution is very simple you know beta is equal to omega root over LC it is 50 has system L is given C is given all you substitute you will get beta is equal to 0.00105 radian per kilo meter right.

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A three phase, 50 Hz, 400 kV transmission line is given. The line inductance is 0.97 mH/km per phase and capacitance is 0.0115 μ F/km per phase. Assume a lossless line. Determine the line phase constant β , Z_c , v and λ .

Soln.

$$\beta = \omega \sqrt{LC}$$

$$\therefore \beta = 2\pi \times 50 \sqrt{0.97 \times 10^{-3} \times 0.0115 \times 10^{-6}}$$

$$\therefore \beta = 0.00105 \text{ rad/km.}$$

Surge impedance,

$$Z_c = \sqrt{\frac{L}{C}} = \sqrt{\frac{0.97 \times 10^{-3}}{0.0115 \times 10^{-6}}} = 290.43 \Omega$$

Velocity of propagation

$$v = \frac{1}{\sqrt{LC}} = \frac{1}{\sqrt{0.97 \times 10^{-3} \times 0.0115 \times 10^{-6}}}$$

$$\therefore v = 2.994 \times 10^5 \text{ km/sec}$$

Line wavelength is

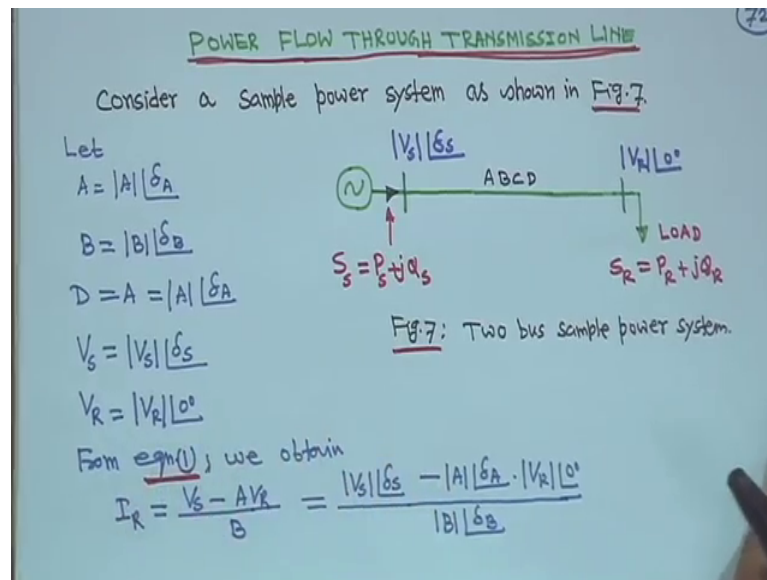
$$\lambda = \frac{v}{f} = \frac{2.994 \times 10^5}{50}$$

$$\therefore \lambda = 4990 \text{ km.}$$

And then surge impedance is $Z_c = \sqrt{L/C}$, L is given C is given. So, it is 290.43 ohm right and velocity of propagation v is equal to $1/\sqrt{LC}$. So, L is known C is known you substitute you will get 2.994 in to 10 to the power 5 kilo meter per second right and line wave length λ is equal to v upon f right. So, is equal to 2.994 we have seen no λ is equal to you are a 1 upon what you call f root over 1 upon 1 see that is basically v upon f , because v is competent here.

So, you put the value of b divided by f . So, approximately 4990 990 kilo meter right this is you are the small example for this one, but after this will go for power flow through transmission line right.

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So, in terms of your A B C D parameter; so this one I will get it for you for the receiving end one and sending one I will give you the your final expression, but you can you can derive this right.

So, you consider a sample power system as shown in figure this. So, this side is sending end side. So, power is taken P_s plus $j q_s$ this side is receiving end side. So, power is p_r plus $j q_r$ this is load. Sending end voltage is magnitude V_s angle δ_s receiving end and V_R angle 0° . So, it is actually your taken as say reference right and this line is represented by A B C D parameter, that is why I am writing here A B C D right and. So, we assume now first let A is equal to magnitude a angle delta a B is equal to magnitude B angle delta B all this thing, D is equal to A is equal to magnitude A angle delta A because throughout that we replace by this thing. Instead of D will write a because A is equal to D right. V_s is equal to right magnitude V_s angle delta s we have made it here V_R is equal to magnitude V_R angle 0° we are made it here.

So, from equation one; that means, very first equation right for this topic we have you know V_s is equal to $A V_R$ plus $B I_R$. So, from here it is coming I_R is equal to V_s minus $A V_R$ upon B right. So, $V_s - V_R A B$ all you substitute in terms of their magnitude and angle right therefore, angle V_s angle delta s. So, minus mode A angle delta A minus mode V_R angle 0° divided by mode B angle delta B right.

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$$\therefore I_R = \frac{|V_s|}{|B|} \angle \delta_s - \delta_B - \frac{|A| \cdot |V_R|}{|B|} \angle \delta_A - \delta_B \dots (77)$$

The receiving end complex power,

$$S_{R(3\phi)} = P_{R(3\phi)} + jQ_{R(3\phi)} = 3V_R I_R^* \dots (78)$$

Using eqn. (78) and (77), we get

$$S_{R(3\phi)} = 3 \cdot \frac{|V_s| |V_R|}{|B|} \angle \delta_B - \delta_s - 3 \cdot \frac{|A| \cdot |V_R|^2}{|B|} \angle \delta_B - \delta_A$$

or in terms of line-to-line voltage,

$$S_{R(3\phi)} = \frac{|V_{s,L-L}| |V_{R,L-L}|}{|B|} \angle \delta_B - \delta_s - \frac{|A| |V_{R,L-L}|^2}{|B|}$$

So, in that case this one you all this things you simplify this equation you simplify, then you will get I_R is equal to $\frac{|V_s|}{|B|} \angle \delta_s - \delta_B$, it will be angle δ_s minus δ_B minus mode A mode V_R by mode B angle δ_A minus δ_B this is equation 77. We are also same thing we have seen the receiving end complex power S_R it is receiving end power 3 phase that is why writing 3 phase.

Is equal to $P_{R(3\phi)} + jQ_{R(3\phi)}$ is equal to $3V_R I_R^*$ right this is equation 78. Now from equation 78 and 77 this I_R is given. So, you have to you find out what is I_R conjugate right. So, if you take the conjugate this is I_R if you take the conjugate one. So, V_s upon b will remain and as it is conjugate. So, this angle will become δ_B minus δ_s right that is why in this expression it is δ_B minus δ_s , and for this conjugate this angle also will be δ_B minus δ_A because of conjugate. So, this angle is δ_B minus δ_A . So, it is $3 \cdot \frac{|V_s| |V_R|}{|B|} \angle \delta_B - \delta_s$ minus $3 \cdot \frac{|A| |V_R|^2}{|B|} \angle \delta_B - \delta_A$ right.

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The receiving end complex power,

$$S_{R(3\phi)} = P_{R(3\phi)} + jQ_{R(3\phi)} = 3V_R I_R^* \quad \dots (78)$$

Using eqn. (78) and (77), we get

$$S_{R(3\phi)} = 3 \cdot \frac{|V_s| |V_R|}{|B|} \angle \delta_B - \delta_s - 3 \cdot \frac{|A| |V_A|^2}{|B|} \angle \delta_B - \delta_A \quad \dots (79)$$

or in terms of line-to-line voltage,

$$S_{R(3\phi)} = \frac{|V_{s,L-L}| |V_{R,L-L}|}{|B|} \angle \delta_B - \delta_s - \frac{|A| |V_{R,L-L}|^2}{|B|} \angle \delta_B - \delta_A \quad \dots (79)$$

So, in terms of the your what you call this line to line voltage right if you take this line to line voltage.

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$$3 |V_s| |V_R| = \frac{\sqrt{3} |V_s| \times \sqrt{3} |V_R|}{(|V_{s,L-L}|) \times (|V_{R,L-L}|)}$$

$$= \frac{(\sqrt{3} |V_s|)^2}{|V_{R,L-L}|}$$

$$= |V_{R,L-L}|$$

Then we can write this expression it is 3 V s V R; that means, it is 3 V s V R right this one you can write know it is root 3 in to V s into root 3 in to V R right the. So, this one actually will become I have taken here it is V s line to line right in to V R line to line. So, this term is taken that V s magnitude V s line to line V R line to line by mode B angle delta b my delta, delta B minus delta s right. So, this 3 actually written like this and

similarly this $\sqrt{3}VR$ square right. So, in that case also this this can be written as actually $\sqrt{3}VR$ your square because this 3 I take it inside that bracket. So, $\sqrt{3}VR$ square $\sqrt{3}VR$ square.

So, basically this one can be written as $\sqrt{3}VR$ line to line your square that is why it is written $\sqrt{3}VR$ line to line square upon mode B angle delta B minus delta A right. We will come back soon.

Thank you.