

Audio System Engineering
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Lecture - 09
Spherical Waves Propagation

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Spherical Waves

$$\nabla^2 p = \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2}$$

$$x = r \sin \theta \cos \phi$$

$$y = r \sin \theta \sin \phi$$

$$z = r \cos \theta$$

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} + \frac{1}{r^2 \sin^2 \theta} \frac{\partial}{\partial \theta} (\sin \theta \frac{\partial}{\partial \theta}) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2}$$

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} \quad \rightarrow \quad \nabla^2 p = \frac{1}{r} \frac{\partial^2 (rp)}{\partial r^2}$$

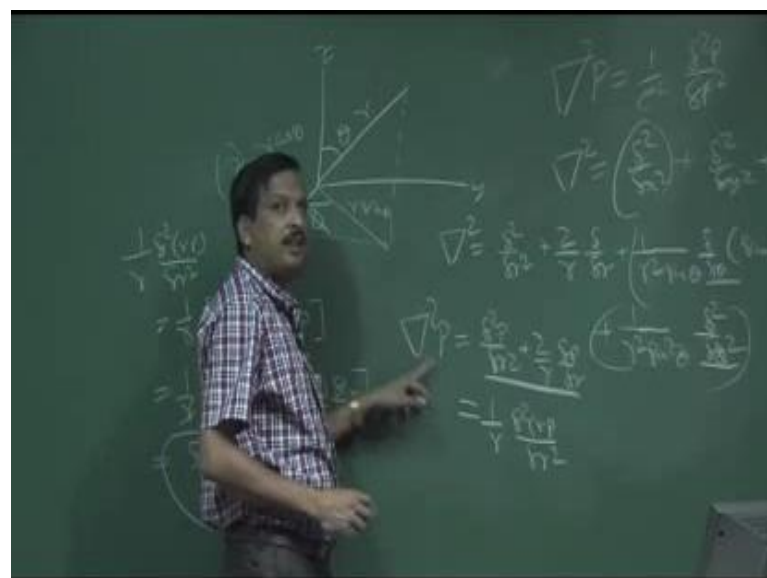
$$\frac{1}{r} \frac{\partial}{\partial r} \left(\frac{\partial (rp)}{\partial r} \right) = \frac{1}{r} \frac{\partial}{\partial r} (r \frac{\partial p}{\partial r} + p) = \frac{1}{r} \left(r \frac{\partial^2 p}{\partial r^2} + \frac{\partial p}{\partial r} + \frac{\partial p}{\partial r} \right) = \frac{\partial^2 p}{\partial r^2} + \frac{2}{r} \frac{\partial p}{\partial r}$$

Conservation of energy and the relationship $l = \frac{p^2}{2\rho_0 c}$ lead to that the pressure amplitude might fall off as $1/r$, so that the quantity rp would have amplitude independent of r .

$$\frac{\partial^2 (rp)}{\partial r^2} = \frac{1}{c^2} \frac{\partial^2 (rp)}{\partial t^2}$$

So, in last class, we are discussing about that spherical waves propagation.

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So, we are saying that let us the instead of linear waves, let us sound is propagated like a spherical sphere. So, I can say that let us this is the source of the sound, and it is propagated like a sphere. So, instead of linearly propagate, it propagate in all direction, in all three-dimensional space in all direction like a sphere, spherical wave propagation. Now, what is the difference, now I know that in Cartesian coordinate x, y, z we have consider. Now, in here, I have to converted that x, y, z to r, θ and ϕ , same three coordinates; instead of x, y, z , I convert that r, θ and ϕ . So, in that case, what we will do is we just write draw the coordinate x, y, z and we put a point here, which is r . So, this point is θ , so this is r , projection of r here, $r \cos \theta$.

Now, if I make a projection in x, y plane of this same arc, I get $r \sin \theta$. So, the projection of $r \sin \theta$ on x -axis gives me the x . So, I can write x is nothing but a $r \sin \theta \cos \phi$ if this angle is ϕ . Similarly, y is nothing but a $r \sin \theta \sin \phi$. And z is nothing but a $r \cos \theta$. So, you know what is the wave equation we know; now wave equation is nothing but a divergence square P is equal to one by C square $\nabla^2 P$ by $\nabla^2 t$ square, this is the linear wave equation. So, I can say what is the meaning of this thing, this thing is nothing but a ∇^2 square by $\nabla^2 x$ square plus ∇^2 square by $\nabla^2 y$ square plus ∇^2 square by $\nabla^2 z$ square. So, this is that things.

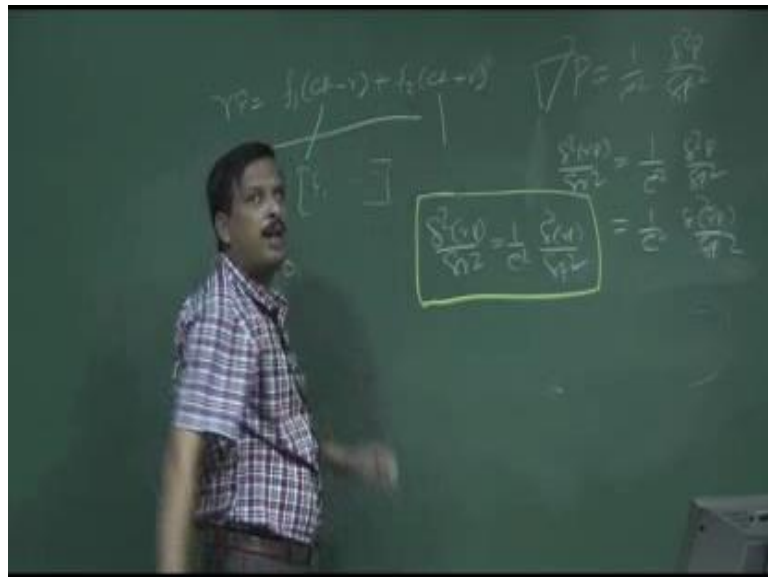
Now, I instead of x, y, z this term can be calculated from this $x, \nabla^2 x \nabla^2 x \nabla^2 x \nabla^2 x \nabla^2 x \nabla^2 x$ square I can calculate, ∇^2 square $\nabla^2 y$ square I can calculate, and ∇^2 square $\nabla^2 z$ square I can calculate. So, eventually, this Laplacian operator will become in θ, r and ϕ domain that will be ∇^2 square by $\nabla^2 r$ square first term plus 2 by $r \nabla^2$ by $\nabla^2 r$ plus one by r square $\sin^2 \theta$ into ∇^2 by $\nabla^2 \theta$ into $\sin^2 \theta$ ∇^2 $\nabla^2 \theta$ plus one by r square $\sin^2 \theta$ into ∇^2 square by $\nabla^2 \phi$ square. So, this is prove of this thing, this is details mathematically prove you can calculate in that case.

So, now if this is the case, so instead of this, I have to get this whole equation. Now, think I said sound is propagated or acoustic wave is propagated in a spherical in nature. So, in every angle, every direction, it is symmetry. If something is propagated in spherical nature; that means, I am saying it propagate in symmetric in all direction so if it is propagate symmetric in all direction, so the variation against the θ and ϕ will we not consider, nothing will be there. So, I can say the variation against θ and ϕ , no terminology will be there, so I can say these whole terminology will be vanished if I

write $\nabla^2 P$ then will be only $\nabla^2 p$ by $\nabla^2 r$ square plus 2 by r ∇p by ∇r .
 So, $\nabla^2 p$ by $\nabla^2 r$ square plus 2 by r ∇p by ∇r .

Now, if I say this term can be written as 1 by r $\nabla^2 r p$ divided by $\nabla^2 r$ square. I can prove it, 1 by r $\nabla^2 r p$ divided by $\nabla^2 r$ square. What is this? It is nothing but a 1 by r , ∇ by ∇r into r ∇p by ∇r plus t . Then I took those things also. So, 1 by r into r $\nabla^2 p$ by $\nabla^2 r$ square plus ∇p by ∇r plus ∇p by ∇r 1 by r and 1 by r . So, r , r cancel, it is nothing but a $\nabla^2 p$ by $\nabla^2 r$ square plus 2 by r ∇p by ∇r . So, these things and these things are same.

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So, I can write $\nabla^2 P$ is nothing but a 1 by r $\nabla^2 r p$ divided by $\nabla^2 r$ square.
 So, my acoustics wave equation, I just put these values in here. In case of spherical wave propagation, it becomes 1 by r $\nabla^2 r p$ divided by $\nabla^2 r$ square is equal to 1 by C square $\nabla^2 p$ del p square. Now, if I say P is my function, the pressure amplitude, so I said the pressure amplitude, independent pressure amplitude is independent of r , and can I say that. The pressure amplitude P is independent of r , can I say that. So, if I say that spherical wave propagation also preserve the conservation of energy. So, now let us take a small sphere with a radius lets r 1 , and the intensity of the wave is I . So, energy at this surface of the sphere, if all direction I will be same, so I multiply by the $4 \pi r$ 1 square.

Similarly, if I take another at this distance a sphere, whose radius is r_2 ? So, energy here is $4\pi r_2^2 I$. So, now, if I say the conservation of the energy, energy let us there is not loss in acoustic wave propagation; consider there is no other kinds of losses. So, energy in here and energy in here will be the same. So, $I_1 4\pi r_1^2$ and $I_2 4\pi r_2^2$ will be the same. Since, r_1 and r_2 is not same, these two is has to be same, then I has to be change.

So, what is I , I is nothing but a P^2 by $2\rho_0 c$. So, P cannot be a parameter, which is independent of r . So, instead of P , I consider rP is a parameter, rP is pressure amplitude, amplitude which is independent of r . So, instead of P , here I consider rP is the function. So, I can write rP divided by Δt^2 . So, this $1/r$ will be not there. So, I can say in case of spherical wave propagation, wave equation become $\Delta^2 rP$ divided by Δr^2 is equal to $1/C^2 \Delta^2 rP$ divided by Δt^2 . So, this is the spherical wave propagation equation.

So, in that case, what is the solution in spherical wave propagation? So, I can say rP is a consistent of two functions of wave, one is called forward wave or you can say converging, and another is called divergence. Forward, forward means it is going out; it is diverging wave, in case of sphere. And if it is backward waves, so it is converging wave. So, I can write lets this rP , the two function one is f_1 into $Ct - r$ and plus f_2 $Ct + r$. So, this is a converging wave, and this is a diverging wave. So, waves are diverted is a forward wave; and once it is backward wave means over converted. So, it is a diverging wave, and this is a converging wave.

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$$p = \frac{1}{r} f_1(ct-r) + \frac{1}{r} f_2(ct+r) \quad \text{For all } r > 0$$

A spherical wave diverging from origin

A spherical wave converging at origin

At $r=0$ the solution fails

- For out going waves solution fails at origin because some source of sound is required to supply the energy carried away but wave equation does not contain any term representing this energy source. → So its means that the medium must be excluded from some volume of space including the origin which will be occupied by the sound source
- For incoming waves energy is being focused at the origin and the small amplitude approximations will fail. → this will manifest itself in a nonlinear wave equation and strong acoustic losses.

Now, what is P, P is nothing but a $\frac{1}{r}$ that function F_1 plus F_2 . So, now if it is one by r, so at r equal to 0, this concept does not valid; diverging wave, converging wave at r equal to 0, P no solution is valid. So, in that case, what will happen, so solution fail? So, in case of solution fail, for outgoing wave solution fail at origin required some source of sound is required to supply the energy carried away, but wave equation does not contain any term representing this energy source. So, it means that medium must excluded from some volume from the media some volume must be excluded from the media, from the medium, so that space at origin, at origin that space, which acquired by the sound source or acoustic source.

Now, in case of diverging wave, it is fail explanation is there. For incoming wave, energy is being focused at the origin and the small amplitude approximation will fail. This will manifested itself in a non-linear wave equation and strong acoustic losses. So, this is the physical meaning of if r equal to 0.

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Now, let us r is not equal to 0, and we proceed. Then I can say if this is my wave equation, this is the wave equation, and then what is the solution of general solution of this wave equation. So, general solution is that $r p$ is equal to $A e$ to the power $J \omega t$ minus $K r$, K is the wave constant and r is the distance from the source. So, in that case, P is nothing but a A by $r e$ to the power $J \omega t$ minus $K r$. So, this is the P . Now, if this is P , what is velocity potential ϕ ? So, it is written that $\rho_0 \frac{\partial \phi}{\partial t}$ is equal to minus P . We said $\frac{\partial \phi}{\partial t}$, we always write $J \omega$, so it is $\rho_0 J \omega \phi$ is equal to minus P . So, ϕ is equal to minus P divided by $J \omega \rho_0$, velocity potential.

Now, what is the particle velocity u is nothing but a divergence of ϕ . So, in case of spherical coordinates, divergence, $\frac{d \phi}{d r}$ only. So, if I put that $\frac{d \phi}{d r}$ in case of P , what I will get, let us put that $\frac{d \phi}{d r}$, so ϕ is equal to minus P by $J \omega \rho_0 r$. So, I can write minus A by $r e$ to the power $J \omega t$ minus $K r$ divided by $J \omega \rho_0 r$. Now, I take the $\frac{\partial \phi}{\partial r}$. What I will get, first order differential equation, so I will get I just writing $1 - \frac{J}{K r}$ into $\frac{P}{\rho_0 C}$. If you do the differentiation, you will get this thing; $\frac{\partial \phi}{\partial r}$ is nothing but a $1 - \frac{J}{K r} P$ by $\rho_0 C$, which is nothing but u .

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Now, see that in case of linear wave equation that pressure wave P and ϕ are in the same phase, but here u is a function of some complex thing. So, if I want to calculate the impedance z , so z is nothing but a P by u ; z is nothing but a P by u . So, if I calculate the z , what I will get, I will get, so P will be cancel, P by $1 - jkr$ into P by $\rho_0 c$, so P, P cancel. So, it is nothing but a $\rho_0 c$ divided $1 - jkr$. So, I can write it is nothing but a $\rho_0 c$ into kr by $kr - j$.

Let us consider in case of linear wave, z is nothing but a $\rho_0 c$. Here I am saying there is another complex stock is there. Let us consider this J , create an angle so if I say representation of the J , if this is θ then if I represent this complex number in a trigonometry form, so it will kr , it will be one and it will be root over of $1 + kr^2$, if this angle is θ . This is kr ; this is one. So, I can write this complex in term of amplitude and θ , so I can write let us a minus Jb , so I can write root over a square plus b^2 is the amplitude and e to the power $J\theta$. And θ will be $\tan^{-1} b/a$, let's this will be minus, this is minus, so I write minus θ . So, instead of plus θ , I write θ then it will be plus.

So, instead of one by $kr - j$, I can write $\rho_0 c$ into kr by amplitude $\sqrt{1 + kr^2}$ into e to the power minus $J\theta$. Same thing, I can write $\rho_0 c$ kr by $\sqrt{1 + kr^2}$ into e to the power $J\theta$, minus would be plus, when it go up. Now, what is kr by $\sqrt{1 + kr^2}$? So, we have drawn that triangle like this way. This is

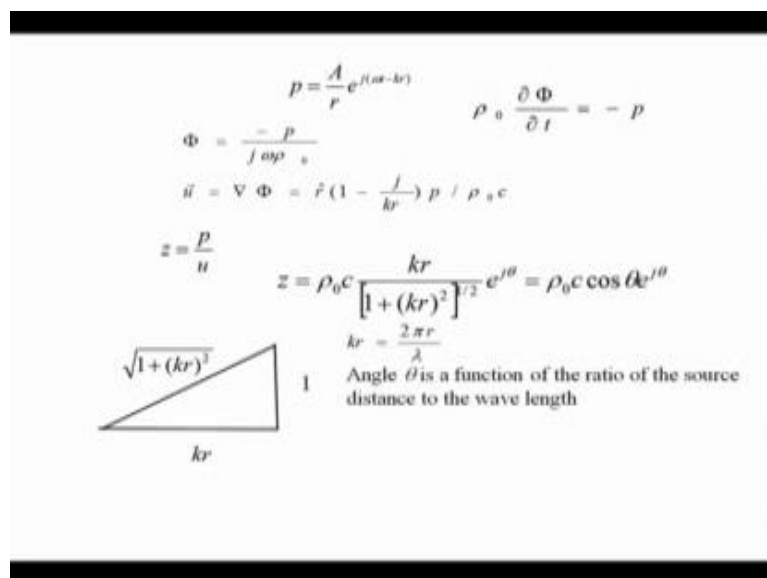
theta, this is one, this is K r, and this is 1 plus K r square. So, what is cos theta, cos theta is nothing but a base divided by 1 plus K r square, in that case, I can write this is nothing but a rho 0 C cos theta into e to the power J theta. So, rho 0 C cos theta into e to the power J theta, if I say the amplitude part of the z is not rho 0 C.

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So, in case of spherical wave propagation, the impedance - acoustic impedance is not rho 0 C, instead of rho 0 C, it is a rho 0 C cos theta. What is the theta? Theta is the angle between the pressure and velocity.

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Now, if this is theta, what is theta, I can write $\cot \theta$ is equal to $K r$ by 1. So, $\cot \theta$ is nothing but a $K r$. Now, what is $K r$, what is $K r$, r is the distance from the source, what is the K , K is the ω sub c , so I can write $K r$ is nothing but a $2 \pi r$ divided by λ . So, I can say the theta is a function of the ratio between the distance from the source and wavelength of the acoustic source, acoustic wave. So, distance from the source and wavelength of the acoustic source.

Now, let us the spherical wave is started here and going like this way. If this r is small, very small compared to this λ , what will happen? If r is very small, that means, the distance travelled by the source, the distance travelled by the acoustic wave from the source is very small compared to the wavelength of that acoustic wave. So, compared to the wavelength of the acoustic wave, if the distance is very small, very small then the complex the theta, value of the theta will be large, theta will be increase, theta will be very large.

Now, if the distance is very far away, let us from the source compare to λ is r is very large, that means, the acoustics wave is travelled far away from the source in that case, theta will be very small, because $\cot \theta$, so theta will be very small. So, in that case, theta will be very small means angle between the pressure wave and particle velocity or velocity wave will be almost NIL. So, in that case, z becomes ρ_0 , sorry ρ_0 C , \cos zero one. So, it is far away from the source, it is looks like a plane wave; although it is a spherical wave, it is look like a plane wave.

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$$z = \rho_0 c \frac{kr}{[1+(kr)^2]^{1/2}} e^{j\theta} = \rho_0 c \cos \theta e^{j\theta}$$

$$z = \rho_0 c \frac{(kr)^2}{[1+(kr)^2]} + j \rho_0 c \frac{kr}{[1+(kr)^2]}$$

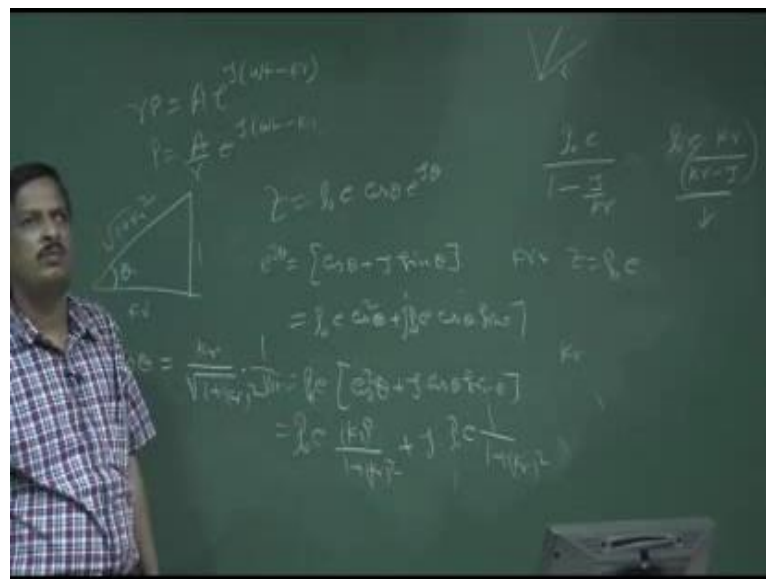
$$\cos \theta = \frac{kr}{\sqrt{1+(kr)^2}}$$

$$\sin \theta = \frac{1}{\sqrt{1+(kr)^2}}$$

Acoustic resistance Acoustic reactance

For large value of kr the $z = \rho_0 c$
 For small value of kr the $z = 0$

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Now, I can express z, in other term also. So, you have seen z is the complex function. So, it is a two parameter function. So, z is nothing but a rho 0 C cos theta e to the power J theta. What is e to the power J theta is nothing but a cos theta J sin theta. So, I can write z is nothing but a rho 0, C cos theta into cos theta, so cos square theta plus rho 0 C cos theta J sin theta. Or I can write rho 0 C, cos square theta plus J cos theta sin theta. So, what is the value of cos theta and sin theta. I said the cos theta is nothing k r by root over of 1 plus k r square. So, I can write rho 0 C k r square divided by 1 plus k r square plus J

$\rho_0 C$ into $\cos \theta$ $k r$ by $1 + k r$ square and $\sin \theta$ is $1 + k r$ square, so it is nothing but a $1 + k r$ square.

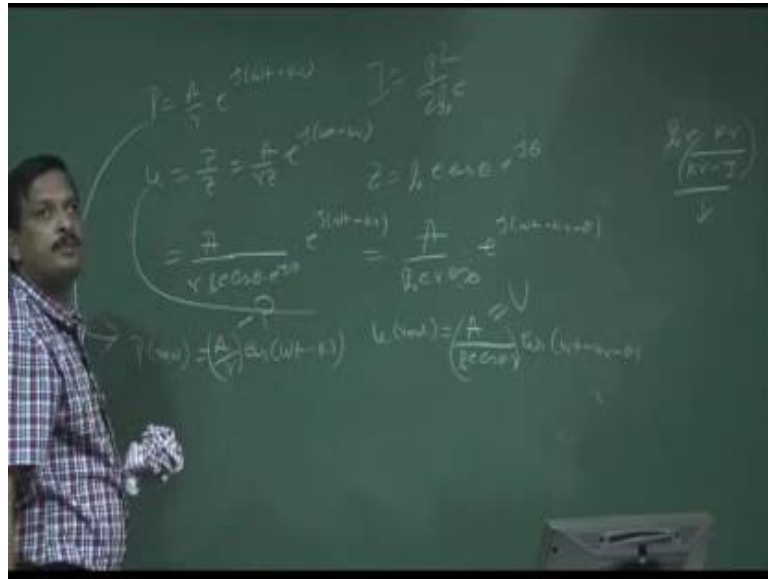
So, if you see a plus $J b$, so z is not a resistive not only the resistive term, there also some term, which is reactance. So, it is called acoustic resistance and it is the acoustic reactance. For large value of $k r$, if the $k r$ is very large, then you see this term is 0, and this term is one. So, it is z is nothing but a so at $k r$ is very large, z is nothing but a $\rho_0 C$. Now, if the $k r$ is very small, then z becomes 0, so that cannot be explained. So, the $k r$ is very small, z become 0.

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$$\begin{aligned}
 p &= \frac{A}{r} e^{j(\omega t - kr)} \\
 u &= \frac{p}{z} = \frac{A}{zr} e^{j(\omega t - kr)} & z &= \rho_0 c \cos \theta e^{j\theta} \\
 u &= \frac{A}{\rho_0 c \cos \theta e^{j\theta} r} e^{j(\omega t - kr)} = \frac{A}{\rho_0 c \cos \theta} e^{j(\omega t - kr - \theta)} \\
 p(\text{real}) &= \frac{A}{r} \cos(\omega t - kr) = P \cos(\omega t - kr) & P &= \frac{A}{r} \\
 u(\text{real}) &= \frac{A}{\rho_0 c \cos \theta} \cos(\omega t - kr - \theta) = U \cos(\omega t - kr - \theta) & U &= \frac{A}{\rho_0 c \cos \theta} \\
 P &= U \rho_0 c \cos \theta \\
 I &= \frac{1}{T} \int_0^T P \cos(\omega t - kr) U \cos(\omega t - kr - \theta) dt = \frac{PU \cos \theta}{2} = \frac{P^2}{2 \rho_0 c}
 \end{aligned}$$

Now, that wave equation I have understand. Now, one thing how do we express the intensity of the sound. Most of the cases, most of the sound wave propagation we discuss about the spherical wave propagation. So, now we have to find out what should the expression of the intensity in case of spherical wave propagation. So, in case of plane wave propagation I is nothing but a P square by $2 \rho_0 C$. I have to say whether this I expression is remain same in case of spherical wave or it is change.

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So, let us I consider that P is nothing but a A by r e to the power J omega t minus k r. And in u is nothing but a P by z, I can write, so it is nothing but a A by r z, z is the impedance e to the power J omega t minus k r, where z is equal to rho 0 C cos theta e to the power J theta. So, I can write u is nothing but a A by r rho 0 C cos theta e to the power J theta into e to the power J omega t minus k r. So, I can write A by rho 0 C r cos theta into e to the power J omega t minus k r minus theta, e to the power theta you can term into here, so minus theta.

Now, this is u and this is P is this. So, what is the real part of the P? If I say the real part of the P, so P real – real part of the P is nothing but A by r cos omega t minus k r. Now, what is the real part of u, u real – real part is nothing but A by rho 0 C cos theta into r that term will be there into cos omega t minus k r minus theta. Let us this term as a capital U, and this term as a capital P. So, now if that is the case, then I can write P is equal to.

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So, I can write sorry I can write u is equal to small u and this is the capital U , and this is the P , P is equal to A by r , capital P is equal to A by r , and capital U is equal to A by $\rho_0 C \cos \theta$ into r . So, A by r is nothing but a P , capital P by $\rho_0 C \cos \theta$ (Refer Time: 30:05). So, I can write capital P is nothing but a capital U into $\rho_0 C \cos \theta$.

Now, what is I , what is the intensity of the wave? Intensity of the wave I is nothing but a average intensity, time average zero to T pressure into velocity. So, what is pressure, pressure is capital $P \cos \omega t - k r$ into capital $U \cos \omega t - k r - \theta$ into dt . Now, do this whole integration, if you do this integration, you will get, what will get P into $U \cos \theta$ divided by two. So, if I get that what is P , P is A by r or I can in term of P or U any one I can replace the U in terms of P , so U is nothing but a so P into P by two $\rho_0 C \cos \theta$ into $\cos \theta$; $\cos \theta$, $\cos \theta$ cancel. So, it is P^2 by two $\rho_0 C$. So, it is same as linear wave intensity equation, so it is same as linear wave intensity equation.

So, next or next class, I will discuss the physical meaning of acoustic intensity propagation. What is decibel, how decibel is measure, why it is decibel, and why intensity energy is different that is we will discuss in the next class.

Thank you.