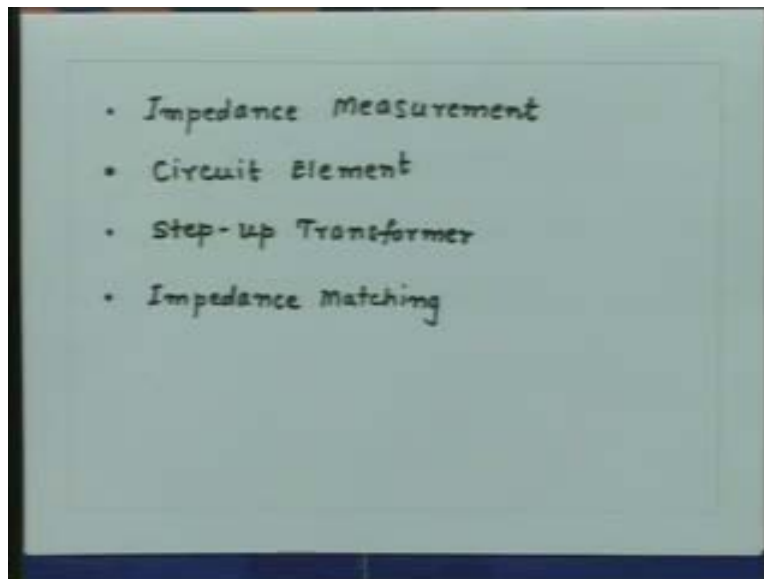


Transmission Lines & E M. Waves
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Lecture - 11

We are discussing applications of transmission lines. We saw the application of transmission line or measuring the unknown impedance at higher frequencies. As pointed out a higher frequency the measurement of phase is not very reliable. So one you would like to estimate the phase indirectly, without doing the direct measurement of the phase, say as we saw by using transmission lines just by measuring the magnitudes of the voltages on transmission lines, we can find the voltage standing waves, their parameters and from that we can indirectly estimate the phase of the unknown impedance. The another application which we have investigating, e is the transmission line as circuits elements.

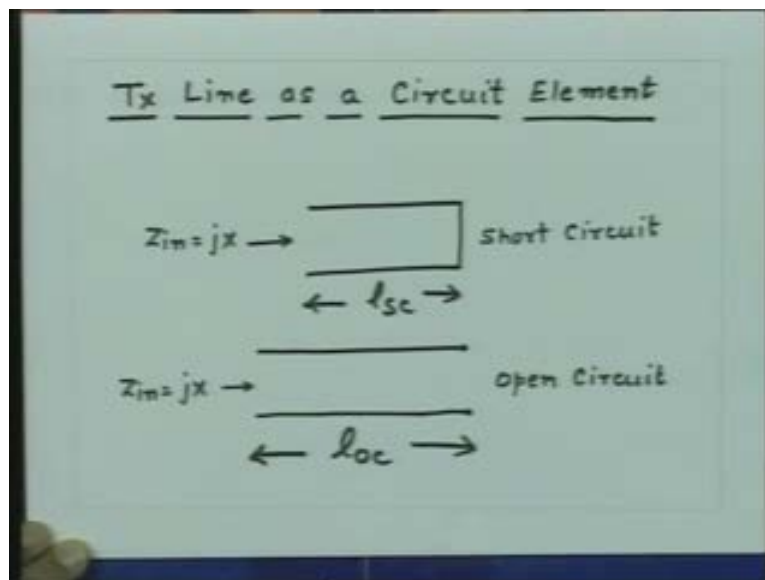
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We saw that if you take a section of a transmission line which is either open circuit or short circuit then, these sections of transmission lines can be use as a reactive element in the higher frequency circuits. So last time we saw the characteristics of the short circuit and open circuit

section of transmission lines and you also saw that these sections, if the length is multiples of $\lambda/4$ then these section can be used as a resonance circuit in the higher frequencies. So essentially by taking a length of a transmission line which is either $\lambda/4$ or $\lambda/2$, we can realize a parallel or serial resonant circuit. We saw that if the input impedance of this line is close to 0 then the circuit behave like a series resonant circuit, if the input impedance of the line is close to infinity then the line behaves like a parallel resonance circuit and once the circuit behave like a resonant circuit than the natural question to ask is what is the quality factor of the resonant circuit and that is the thing essentially we are investigating now, let if I consider this section of a transmission line as a resonance circuit then what is the quality factor of this resonance circuit.

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As mentioned earlier, there whenever we talk about resonant circuit the quality factor is related to the losses in the resonant circuit. However, up till now we have consider the line to be lossless. So by definition for a lossless line the quality factor is always infinitive. Now since we are investigating the quality factor no matter, how small the loss is in the transmission line. We have to account for that loss then and then only we will be able to get the correct answer for the quality factor.

So we investigate now, the input impedance of the transmission line, when the length is multiples of λ by 4 and from there then we find out the quality factor of the resonance sections of transmission line have been said that let us consider, now a specific case that the length of the line is λ by 4 either the line can be open circuited or short circuited. So the length which you call last time is LSC which is the length of a short circuited line, this length is now λ by 4. Similarly, LOC is also λ by 4, so since the length is λ by 4 and as you have seen then the impedance, normalize impedance inverse itself every distance of λ by 4. This open circuit will appear like a short circuit in the line, the short circuit will appear like a open circuit of the distance of λ by 4.

So this line, if the length is λ by 4, will appear like a parallel resonance circuit whereas this if the length is λ by 4, will appear like a series resonance circuit because the impedance will be very close to 0. So now you specifically we just consider the sections of transmission line which are having length of λ by 4 and find out this input impedance and then, go to the calculation of the quality factor. In general, since we are considering loss now, the input impedance of the open or short circuited line has to be written in the total propagation constant γ and not only the phase constant.

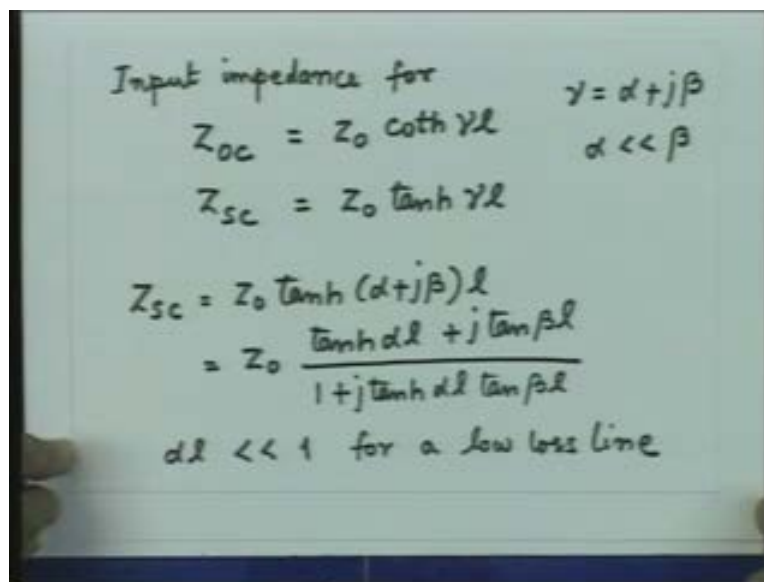
So if you go back to your original impedance relationship, if you have derive as substitute z equal to either 0 or infinity depending upon whether a short circuit or open circuit, I get the input impedance. So this is the input impedance for a line of length L , when the line has loss. So let us call for a open circuit deadline, let us say this is denoted by Z_{oc} . So that is equal to $Z_0 \cot$ hyperbolic of γL and for Z_{sc} the short circuited line that will be $z_0 \tan$ hyperbolic of γl , where if you remember γ is equal to $\alpha + j\beta$, the line is low offline.

So α is much much smaller than β , as we have consider earlier. When the line is Loc, the propagation \cot and now is $\alpha + j\beta$ and it is not only $j\beta$ is still assume that the characteristic impedance that this line is almost real. So, although we have set now the propagation constant is complex it is not purely imaginary, the z_0 be still consider to be almost here. So now, what again we are saying is consider any of these cases, the open circuit line or a short circuit line and under the assumption that α is much much less than β right and the

length is either multiples of lambda by 4 or any arbitrary length find out under these approximation, what will be the input impedance. We know if you consider a short circuit line and if the length is lambda by 4, if you appeal ideal like open circuit.

Similarly, if you take an open circuited line and take a length lambda by 4, the input impedance will be like a short circuit. However, in the presence of the loss that will not be true, so the impedance will neither be infinitive nor it will be 0 and that is what we are trying to calculate now that if, the loss is present what will be the input impedance of these 2 lines. So let us consider these Z_{sc} and substitute for gamma alpha plus j beta, so from here I can get Z_{sc} that is equal to $Z_0 \tanh$ hyperbolic of alpha plus j beta into l. I can expand this, so that will be $Z_0 \tanh$ hyperbolic alpha l plus j tan beta l divided by 1 plus tan hyperbolic alpha l j tan beta l. So, in general for a length of line, the input impedance, the short circuited line, the input impedance essentially given by this.

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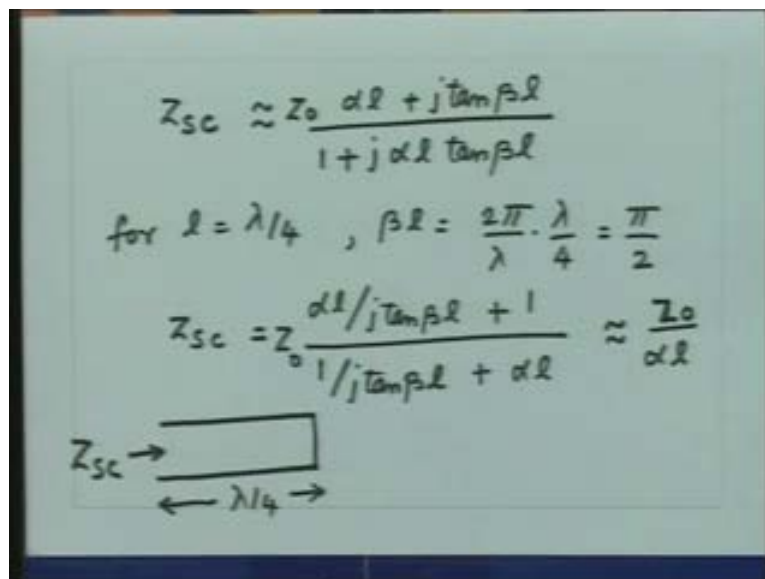
Input impedance for $\gamma = \alpha + j\beta$
 $\alpha \ll \beta$
 $Z_{oc} = Z_0 \coth \gamma l$
 $Z_{sc} = Z_0 \tanh \gamma l$
 $Z_{sc} = Z_0 \tanh (\alpha + j\beta) l$
 $= Z_0 \frac{\tanh \alpha l + j \tan \beta l}{1 + j \tanh \alpha l \tan \beta l}$
 $\alpha l \ll 1$ for a low loss line

Now since alpha is small for a length of lambda by 4 section of transmission line or a length which is smaller than lambda by 2, this quantity alpha l is much much smaller than 1. So we have alpha l much much smaller than 1, for a low loss line. Once, I get this then I can make an

approximation that $\tanh \alpha l$ will be approximately equal to αl , if I substitute that then the input impedance of a short circuited line, Z_{sc} will be approximately equal to αl plus $j \tan \beta l$ divided by $1 + j \alpha l \tan \beta l$.

Now, if the line is short circuited and if the length is $\lambda/4$, this l is equal to $\lambda/4$. So for l equal to $\lambda/4$, βl will be 2π by λ into $\lambda/4$ that is equal to $\pi/2$. If I substitute now, the length of the line equal to $\lambda/4$, this quantity will be infinity this quantity will be infinity. So I can take this $\tan \beta l$ common, no infinity. So what will I see that Z_{sc} will be equal to if I take $j \tan \beta l$ are common, this will be αl divided by $j \tan \beta l$ plus 1 divided by $1 + j \alpha l \tan \beta l$ plus αl , substituting now for βl equal to $\pi/2$ that means $\tan \beta l$ is infinity, this quantity will be 0 , this quantity will be 0 .

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The image shows a handwritten derivation of the input impedance Z_{sc} for a short-circuited transmission line of length $l = \lambda/4$. The derivation starts with the general formula for the input impedance of a short-circuited line:

$$Z_{sc} \approx Z_0 \frac{\alpha l + j \tan \beta l}{1 + j \alpha l \tan \beta l}$$

For $l = \lambda/4$, the phase constant βl is calculated as:

$$\beta l = \frac{2\pi}{\lambda} \cdot \frac{\lambda}{4} = \frac{\pi}{2}$$

Substituting $\beta l = \pi/2$ into the formula, the $\tan \beta l$ term becomes infinity. The expression simplifies to:

$$Z_{sc} = Z_0 \frac{\alpha l / j \tan \beta l + 1}{1 / j \tan \beta l + \alpha l} \approx \frac{Z_0}{\alpha l}$$

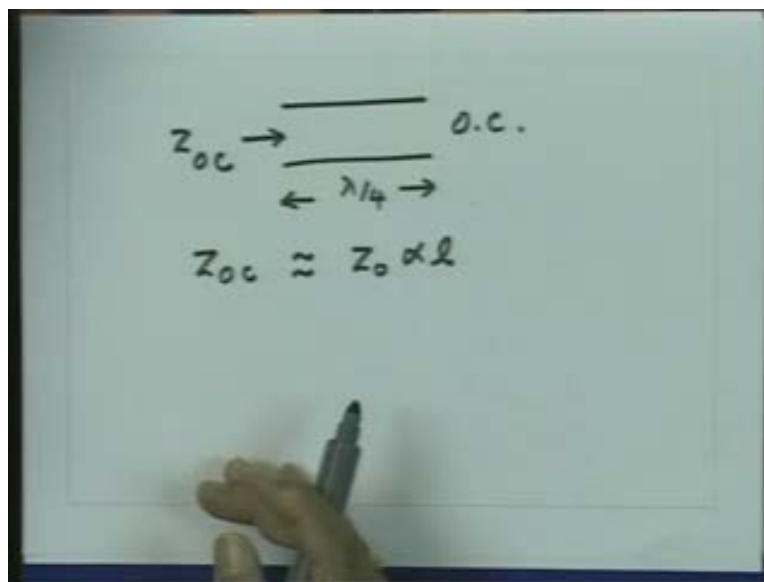
Below the equations, a circuit diagram is shown. It consists of a rectangular box representing the transmission line. An arrow labeled Z_{sc} points into the left side of the box. A double-headed arrow below the box indicates its length is $\lambda/4$.

So I will get Z_{sc} is approximately equal to $1/\alpha l$. So the input impedance of a short circuited line, if the length is $\lambda/4$ is a short circuit, this length is $\lambda/4$, ideally for this line it should have appear open circuit at the input, these impedance Z_{sc} . Now I do not see open circuit but I see an impedance which is very large because αl is very small quantity for a low loss line but this impedance is not infinity. So short circuited section of a transmission

line of length $\lambda/4$ does not in practice appear open circuit but it has some impedance measure at the input terminal which is essentially this.

Similarly, if I done the similar exercise for the Z_{oc} and it did the expansion and all those similar things by done then, the input impedance of open circuited line would appear to be αl multiply by z_0 not anyway here the you have to put z_0 not here. So this quantity should be multiplied by z_0 not, this is z_0 not, this is z_0 not. So doing the same exercise for the open circuit deadline this is $\lambda/4$, this is open circuit. so the impedance which we measure here, here is now Z_{oc} .

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So by doing the same exercise as we have done for the short circuit deadline the Z_{oc} will be approximately equal to αl . So for a short circuited section of a transmission line, the input impedance is z_0 divided by αl for a open circuited section of a transmission line, the impedance is z_0 multiplied by αl . So what we now note here is that the input impedance of a resonance section of a transmission line is ideally infinitive or 0 depending upon the resonance circuit is a series circuit or a parallel resonance circuit. However, in practice V_c , large impedance for a parallel resonant circuit and small impedance for a series resonance circuit, must you get

this either I can use this information that if I vary now the frequency around the resonance frequency, how these impedance changes from there I can get the impedance variation at the function of frequency or what is called frequency response of your circuit and from where I can calculate the quality factor of the circuit.

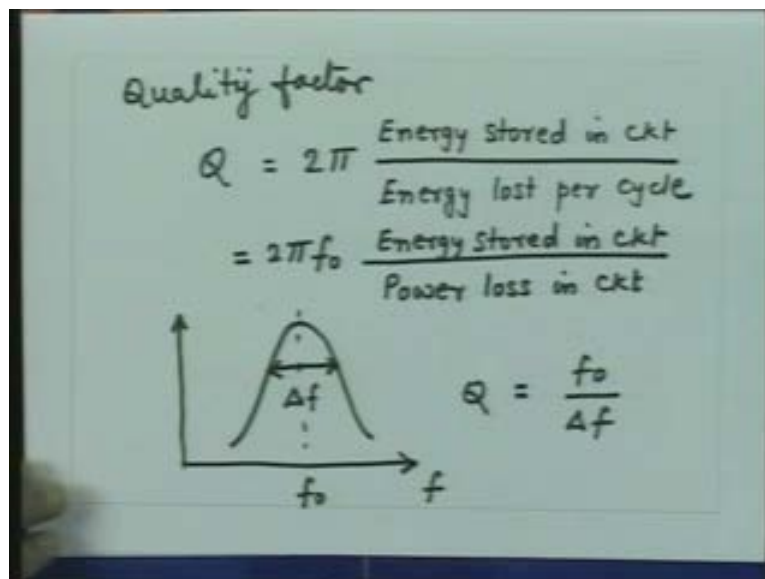
So quality factor can be calculated in 2 ways, one is from the frequency response of the circuit. Secondly, it can be calculated from the basic definition of the quality factor. So let us see what is the quality factor by definition. So the quality factor Q is defined 2π multiplied by the energy stored in the circuit divided by the energy loss per cycle, loss per cycle. If the resonance frequency is f not, the energy loss per cycle is there are f not cycle per second. So I can write down the same expression for the quality factor as 2π into the resonance frequency of the circuit multiplied by the energy stored in the circuit divided by the energy loss per second which is nothing but the power loss in the circuit. So this is divided by the power loss in circuit.

So to calculate the quality factor, if I find out what is the energy stored in the section of a transmission line. If I calculate what is the power loss in the transmission line then I can find out what is the quality factor of that section of a transmission line. Alternatively, as we mention that the quality factor is related to the frequency response. So, if I have a plot the variation of current or voltage, when a voltage or current sources applied to the input section of a transmission line and measure the 3 db bandwidth of the response, the center frequency divided by the 3 db bandwidth of the frequency response give you by the quality factor.

So if I have a frequency response of the circuit which is for this is my impedance at the function of frequency, I get to a frequency response which will look typically like that depending upon whether I am using series or parallel resonance circuit, this will be the voltage response of the current response but at resonant frequency I get the maximum response and other deviate from the resonant frequency, the response drops. So I can measure the 3 db frequency for this response where the amplitude reduces to $1/\sqrt{2}$ of its maximum value, if I take the amplitude response.

So this is um the Δf which is the 3 db bandwidth then, the quality factor for the circuit is equal to f_0 divided by Δf . So, now the quality factor can be calculated from the impedance calculation because the impedance can give me the frequency response. Alternatively, take a section of a transmission line find the voltage distribution on the transmission line and the current distribution on transmission line, find out what is the energy stored, find out what is the power loss any transmission line and use definition to calculate the quality factor. Let us go to this first method which is more fundamental method of calculating the quality factor for transmission line.

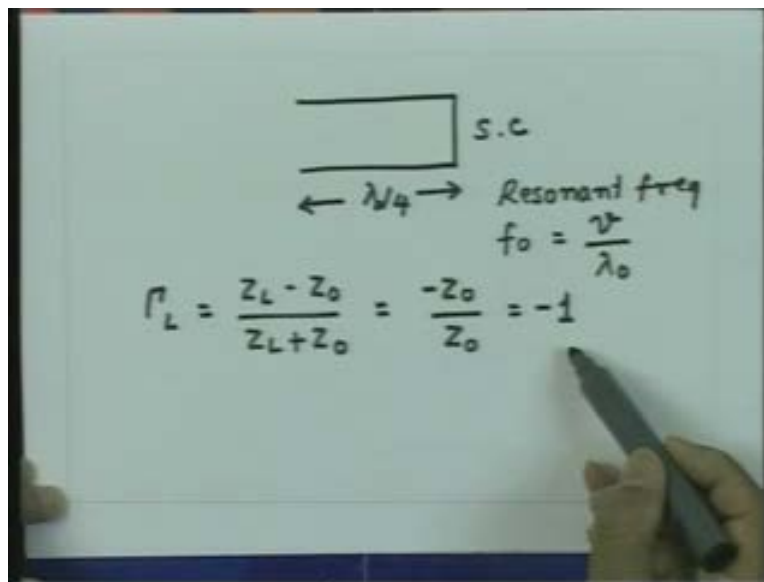
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So let us say I consider specifically a section of a transmission line which is of $\lambda/4$ length and let us say the line is short circuited. So let us take a specific case here, if the line is short circuited this length is $\lambda/4$ and since, I am not talking about the resonance frequency, the length cannot be $\lambda/4$ at all frequencies. So let us say this is having a specific frequency $\lambda_0/4$, where λ_0 corresponds to the wavelength of the resonant frequency. So this length is $\lambda_0/4$, so the resonant frequency of the circuit f_0 is nothing but the velocity of the wave divided by the length $\lambda_0/4$.

Now, we can get the voltage and current expression on the transmission line with the load short circuited and when the load is short circuit, the reflection coefficient at the load point is minus 1 we have seen earlier. So reflection coefficient load Γ_L that is z_L minus z_0 divided by z_L plus z_0 and z_L is now short circuit, so z_L is 0. So that is equal to minus z_0 , upon z_0 that is equal to minus 1. So the reflection coefficient for short circuit load is minus 1 which is saying that the amplitude of reflection coefficient is 1 and it has a phase change of 180 degrees.

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We will make use of this information later when we go to the application of transmission line for the voltage or current standing of transformers. But at this point, once I know the reflection coefficient I can write down the voltage and current equation and I can find out, what is the variation of the voltage and current on the transmission line. Going to the basic equations of the voltage and current, the voltage now on the transmission line can be given as v^+ plus $e^{-j\beta l}$ to the power minus $j\beta l$ but for reflection coefficient of minus 1 at the load, this quantity Γ_L is nothing but v^- divided by v^+ .

So, if I substitute for v^- is equal to minus v^+ plus the voltage on the line is essentially given by this. I can combine this, I can take v^+ common, this is $e^{-j\beta l}$ minus $e^{-j\beta l}$.

the power minus $j\beta l$ which is nothing but $j2$ into sign of βl . So the voltage on transmission line can be $j2V^+ \sin \beta l$, since there is nothing very special about this parameter V^+ plus I can combine $2V^+ \sin$ to define another quantity which is V_0 . So the voltage magnitude on the transmission line is magnitude of V_0 , some voltage multiplied by \sin of βl . Similarly, for the current if I substitute for the reflection coefficient equal to minus 1, the current on the transmission line will be given by this which again by combining these 2 terms, I will get a V_0 divided by Z_0 multiplied by cosine of βl .

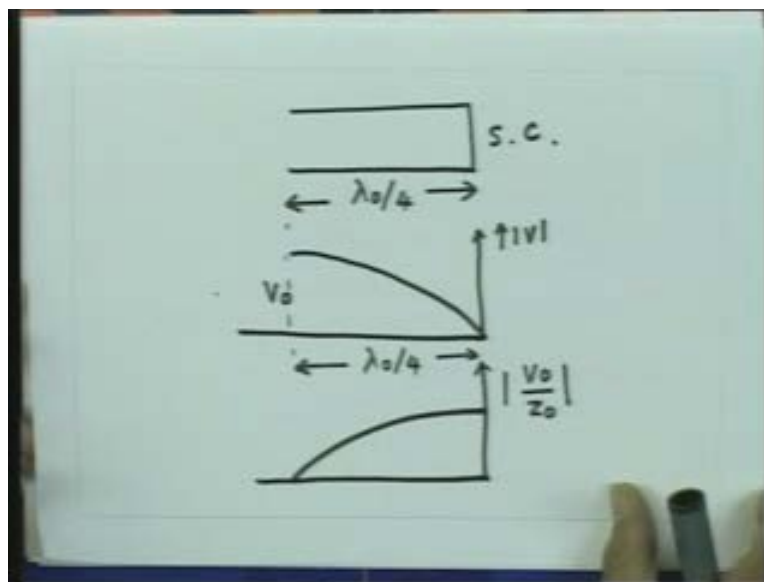
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$$\begin{aligned}
 V &= V^+ e^{j\beta l} + V^- e^{-j\beta l} \\
 &= V^+ e^{j\beta l} - V^+ e^{-j\beta l} \quad \Gamma_L = -1 \\
 &= j2V^+ \sin \beta l \\
 |V(l)| &= |V_0 \sin \beta l| \quad \text{where } V_0 = 2V^+ \\
 I &= \frac{V^+}{Z_0} e^{j\beta l} + \frac{V^-}{Z_0} e^{-j\beta l} \\
 &= \frac{2V^+}{Z_0} \cos \beta l \\
 |I(l)| &= \left| \frac{V_0 \cos \beta l}{Z_0} \right|
 \end{aligned}$$

So if I plot now this voltages on the line, say this is now, $\lambda/4$ line, this is short circuit. So the voltage is 0 at this point because its variation is the sign of βl . So at l equal to 0 which is said, if the load the voltage is 0 and when I go to a distance of $\lambda/4$, this quantity βl becomes $\pi/2$ that is quantity 1. So I see the voltage which is V_0 , so I get a voltage variation which is a 1 cycle variation from lower point to the input point. So this voltage here is V_0 and this is, this point is 0. So this is the variation of magnitude of magnitude of voltage on the transmission line. Similarly, if I plot the current the current has a variation which is cosine variation.

So if l equal to 0, this quantity is V_0 up on z naught and when l is $\lambda/4$ that time this quantity will be 0. So I will get a current variation which will go like that, so this is mode of V plus divided, this length is $\lambda/4$. So I know now on the section of the transmission line, the voltage and current variation and once, I know the voltage and current variation then I can find out what is the capacity when inductive energy stored at different locations on transmission lines. So if I take a small section of transmission line, the capacitance of the small section of transmission line is C multiplied by the length of the line dl , as we discussed in the very first lecture that for a small section of a transmission line, we define the capacitance as the capacitance per unit length multiplied by the length of the section of the transmission line.

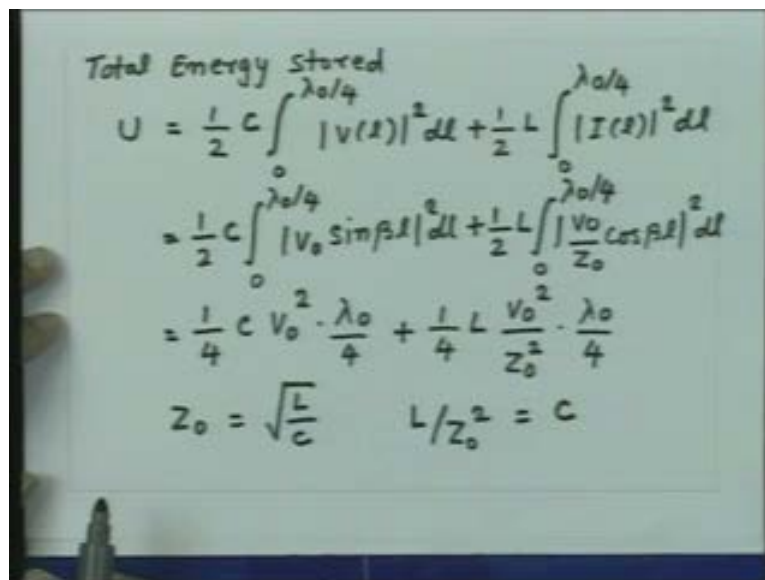
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So for a section of the transmission line of length dl , the capacitance is C into dl . So the capacitive energy stored in the section of a transmission line is half C into dl into the voltage square at their location. Similarly, I can get the inductive energy look it in the infinitive section of a transmission line. So once I get that the total energy stored in the transmission line can be written as, so this is the total energy stored, let us called that u that will be equal to half C , total energy will be the integrated version over the length of the line.

So this is from 0 to $\lambda_0/4$ mod of v into l square into dl plus half l integral 0 to $\lambda_0/4$ mod of v into l square into dl , substituting for voltage and current, this is \sin of βl this is \cos of βl , this is half into c integral 0 to $\lambda_0/4$ mod of v into $\sin^2 \beta l$ dl plus half l integral 0 to $\lambda_0/4$ mod of v into $\cos^2 \beta l$ dl . The integrals are very straight forward, you can calculate the value of these integrals and that will be equal to $\frac{1}{4} c v_0^2 \lambda_0$ plus $\frac{1}{4} L v_0^2 \lambda_0$ upon z_0^2 square into $\lambda_0/4$. Now since z_0 is square root of L upon c , the quantity L upon z_0^2 square, this quantity is nothing but c . So, what you find is that these 2 terms are equal, what does means is that the energy stored in this resonant circuit is equally distributed into the capacity energy and the inductive energy.

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Handwritten derivation of total energy stored in a transmission line:

$$\begin{aligned}
 \text{Total Energy stored} \\
 U &= \frac{1}{2} C \int_0^{\lambda_0/4} |V(z)|^2 dz + \frac{1}{2} L \int_0^{\lambda_0/4} |I(z)|^2 dz \\
 &= \frac{1}{2} C \int_0^{\lambda_0/4} |V_0 \sin \beta z|^2 dz + \frac{1}{2} L \int_0^{\lambda_0/4} \left| \frac{V_0}{Z_0} \cos \beta z \right|^2 dz \\
 &= \frac{1}{4} C V_0^2 \cdot \frac{\lambda_0}{4} + \frac{1}{4} L \frac{V_0^2}{Z_0^2} \cdot \frac{\lambda_0}{4} \\
 Z_0 &= \sqrt{\frac{L}{C}} \quad L/Z_0^2 = C
 \end{aligned}$$

So this step is represented the energy stored in the inductance of the line, this term shows the energy stored in the capacitance of the line and from here, we see these 2 quantities are equal that means the energy stored in the inductance and capacitance of line are equal. If I substitute for this then the total energy stored will be nothing but twice of this term, what twice of this term because these 2 terms are equal. So I get the total energy stored now U that is equal to half into $c v_0^2 \lambda_0$ to $\lambda_0/4$, $\lambda_0/4$ divided by 4. This is the total energy stored in

the transmission line. For calculation of the quality factor now we require, the power loss in the transmission line. One way of doing that is again, you go to the primary parameters that the resistance and the conductance of transmission line and get the voltage and current on the transmission line again integrated to the transmission line, to find out the $I^2 R$ losses, but we can do something different here to get the loss and that is if I take a section of its transmission line here, the energy source is connected to the line here, see if I find out that what is power loss at this location, this power loss will be nothing but the power loss in this line because the line is short circuited to the other end that means the load is not consuming any power, see by at all that any power is supply to this.

This power will be only equal to the loss which are taken place inside the section of this line. So without going into the primary constant of transmission line just by calculating the input impedance of this line, we can find out what will be the equivalent resistance to be the power is supplied and from there we can calculate the loss of power in this section of transmission line. Precisely, this is what we do now, so we say that since the line is short circuited though input impedance which will seen here will be Z_0 divided by αL .

So the input impedance of this line, you are seen earlier this is Z_{sc} is approximately equal to Z_0 divided by α into l , where l is equal to $\lambda/4$. So, now the power loss in the transmission line, say we call it P_{loss} that will be equal to V_{naught}^2 but that is the voltage which you are seeing at the input terminals here. So it is as you I am having a voltage source of V_{naught} here which is supplying power to a resistance whose value is this. So, I get the power loss which is V_{naught}^2 divided by the R which is Z_{naught} divided by αl .

So this is equal to here V_0^2 square divided by Z_{naught} into α into L which is equal to $\lambda/4$ in this case, this is $\lambda/4$. Once, I know the energy stored in the circuit at the power loss in the circuit then, I can get in the quality factor Q that is equal to 2ϕ into f_0 into energy stored in the circuit. So this is half c , $V_{naught}^2 \lambda/4$, V_0^2 square upon $Z_{naught} \alpha$ into $\lambda/4$.

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Handwritten equations on a whiteboard:

$$U = \frac{1}{2} C V_0^2 \frac{\lambda_0}{4}$$

Power loss

$$P_{\text{loss}} = \frac{V_0^2}{(Z_0/\alpha l)} = \frac{V_0^2}{Z_0} \alpha \cdot \frac{\lambda_0}{4}$$

Quality Factor

$$Q = 2\pi f_0 \frac{\frac{1}{2} C V_0^2 \lambda_0/4}{\frac{V_0^2}{Z_0} \alpha \cdot \lambda_0/4} = \frac{2\pi f_0 Z_0 C}{2\alpha}$$

So this quantity cancels V_0^2 cancels, so essentially we get this equal to $2\pi f_0 Z_0 C$ divided by 2α . Now we can do some small manipulation because you know this $Z_0 C$ square root divided by C if I write this quantity here Z_0 into C , so Z_0 into C will be square root of L upon C into C that is equal to square root of LC and the quantity here $2\pi f_0$ is nothing but ω . So the numerator now in the quality factor Q is ω square root LC divided by 2α , where ω is nothing but 2π into the frequency. Now if you were called this quantity ω square root LC is nothing but the phase constant of the transmission line β .

So this quality factor Q is β divided by 2α , same expression you can obtain as I mentioned by finding out the frequency response of the section of a transmission lines. So if I measure the impedance and measure the variation of the impedance as a function of frequency, I can find out the cdb bandwidth of the circuit and then from their using the definition that the quality factor is the resonant frequency divide by the 3 db bandwidth, I will get the same expression for the quality factor. So now the quality factor is now related to the phase and the amplitude constant or the phase and attenuation constant of the transmission line and for a low loss line, since β is much much larger than α , this quantity is very large.

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Handwritten equations on a whiteboard:

$$Z_0 C = \sqrt{\frac{L}{C}} \cdot C = \sqrt{LC}$$
$$Q = \frac{\omega \sqrt{LC}}{2\alpha} \quad \omega \equiv 2\pi f_0$$
$$Q = \frac{\beta}{2\alpha}$$

So for low loss line beta is much much greater than alpha, so Q is much much greater than 1, on thus a very useful property that when you go to high frequency, a section of a transmission line can really give me a very high quality factor circuit and since, as you mention the quality factor is relatively to the 3 db bandwidth higher the quality factor, smaller is the 3 db bandwidth or more tune in the circuit is. So the frequency selectivity of a circuit is very good, if the quality factor of the circuit is very large.

So for a low loss section of a transmission lines in the quality factor is large and typically, you can get of few 100 mega, the quality factor could be as higher as few 100 or 1000, you can really get a very good frequency selectivity by the sections of transmission lines used as a resonant circuit that is the reason when we go to the high frequencies, the sections of transmission lines serve very good purpose for realization of the resonant circuits. So this is one of the very important application of transmission line that is in the high frequency circuit. The next application of the transmission line is the voltage or current stepping up transformer.

So third application which we have here is voltage or current, step up transformer. Let us again consider a quarter wavelength section of a transmission line which is shorted at one end and open

circuited other. So let us say, I have a line and let me draw little big a picture here. So this is short circuited, this end of the line is open circuited and this length is $\lambda/4$. Let us say, by some means, some voltage is induced on some point on this section of transmission line.

So let us say at some location here some voltage is induced, if I consider a voltage source connected to this at this location, it will see as if there are 2 sections of transmission line connected in parallel, this section and this section. Now since, the standing wave is still or not set up and just voltage just getting induced inside a transmission line, this voltage does not see any other impedance but the characteristic impedance of the line. So as soon as the voltage tries to get induced inside this structure it sends 2 traveling waves on both sides of the center transmission line because it will see as if, the energy supply to the characteristic impedance on this 2 transmission line.

So essentially 2 equal amplitude waves are sent from this location. This wave travels up to this point, when it reaches here, this wave sees the short circuit that means you see the reflection coefficient of minus 1. So this wave here is reflected completely from here with a reflection coefficient of minus 1 that is magnitude is unity and the phase is 180 degrees. So here, I have a reflection coefficient Γ that is equal to minus 1, say it is equal to $1 \angle \pi$. So this wave which was traveling here rightwards after reflection undergoes the phase change of π , this wave now traverses all the way up to open circuit, I am considering only this wave, say it wave comes here it gets reflected from the short circuit point travels all the way up to the open circuit point. At the open circuit again the reflection coefficient is plus 1, its magnitude is 1 its angle is 0.

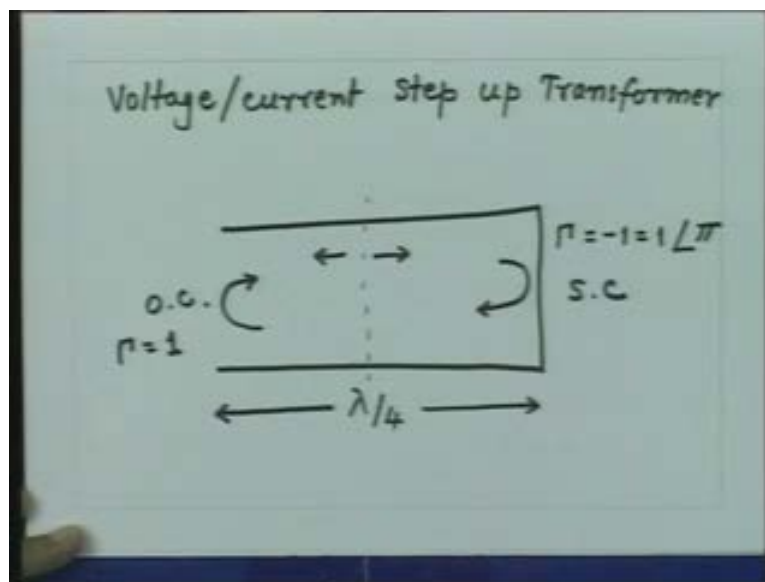
So again the entire wave from here gets reflected with no phase change and it reaches to this location. So by the time this wave has traveled a round trip from here to here, it has traveled a distance of $\lambda/2$ and has gone through an additional phase change of π because of this reflection coefficient here. Now a distance travel of $\lambda/2$ corresponds to a phase change of π plus a phase change of π which is taking place because of the reflection coefficient here.

So this wave which travels a round trip and reaches here has undergone a phase change of 2π

that means the wave which is getting induced at that this location, C is now a phase which is of the wave which is reflected, same as the wave which is getting excited. In other words, it is some kind of a positive feedback which is going on now. The induced voltage and the voltage which is gone on these after round drift, they add up. So the 2 voltages together now start travelling on this, when they again come after round trip, it is again in phase.

So the induced voltage again adds through this voltage, so the voltage essentially starts growing on the transmission line, exactly same thing happens to this wave also. This wave will go like that after reflection from here, it will not undergo any phase change but the entire r will be reflected. It will travel a distance of $\lambda/2$, this will be against your reflection coefficient of minus 1 here. So this wave also will have a regenerative process, so this wave will grow in amplitude, this wave will grow in amplitude exactly same wave and the standing wave of this line will grow because the 2 waves are equal in amplitude and they are going exactly the same way.

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So we will develop the fully standing wave of this line because of the small induced voltage at this location, one may wonder how for this going of the wave will go on. In fact if there are no losses in the transmission line, this growth of the wave will grow up to infinity. So the voltage

will go on growing for infinite amplitudes because there is no controlling parameter on the transmission line.

So even if a small voltage is induced on a transmission line by some **stay** coupling or something and if the section of transmission line is resonant then, the voltage which you see on this transmission line is now much much larger compared to the induced voltage because of this regenerative feedback. So this process, now can be used for your advantage that the voltage which you measure on transmission line could be much larger compared to the coupling voltage at this location or in other words, this resonant section of a line can be used as the voltage stepping of transformer. Precisely, this application we are talking about takes a voltage source which is a small source which is connected to this resonant section of a line and if I measure the voltage on the open circuit of the line, this voltage is much much larger compared to the coupling voltage on the line.

So if there is no loss of the transmission line ultimately in the steady state, the voltage will reach to infinity between the terminals of the line. However, in practice it does not happen because when the voltage and current start going on the transmission line, the ohmic losses also go on increasing and when the loss in the line because of the resistance and the conductance becomes equal to the energy supplied by the coupling source at that point that is the energy balance now and the growth of the standing wave on this section stops. But before this stage reaches the voltage and currents already reach to a substantially **large** value.

So even in practice, when the losses of the transmission line are small one may generate very large voltage and currents, voltage here and the current here, by a small coupling source connected to the section of a transmission line. So this application is the useful application whenever, we want to step up the voltages and currents at higher frequencies, one can show that the voltage which is going to grow that is related to the quality factor of the transmission line.

So again the higher quality factor means low losses on transmission line and that will give me the higher voltage on the transmission line. The phenomena which can be used for your advantage, for stepping up the voltages and currents could be harmful also in many cases. Consider an application, where a small energy source gets unwillingly coupled to the section of

transmission line and if the section of transmission line is of resonant length then, the voltages which will develop on this line will be much larger which your electronic circuit can handle.

So any small coupling of the energy to a resonant section of transmission line, may unknowingly develop the voltages which will be much larger compare to what a circuit can handle and it is possible that your electronic and circuit can get damage because there are very large voltages which are developed here. So while doing the design of the high frequency circuit you should be kept in mind that no forces get coupled to the sections of transmission line, especially if there of resonant length. Otherwise you will unknowingly, you will be developing very large voltages at some locations right in the circuit which might be a harmful to the electronics which are employed with high frequency.

So the phenomena of stating a voltage as in current could act as a **later end**, when unknowing voltage sources or current source get coupled to the section of whereas in transmission line. These are certain applications now which are for designing the higher frequency circuits and as I mention earlier, the sections of transmission lines are use for high frequency circuit designs.

So whenever, we have high frequency circuit design especially a micro of frequencies, we rarely see the components like capacitances and inductances, what we see are the active devices like transistor effects and the reactive components are completely realize by the sections of transmission lines. So most of the microwave circuits will appear as some arbitrary sections of transmission lines connected here and there across the active devices. So the sections of transmission line play a very important role in the design of the higher frequency circuits.