

FOUNDATIONS OF QUANTUM THEORY: NON-RELATIVISTIC APPROACH

Dr. Sandeep K. Goyal
Department of Physical Sciences
IISER Mohali
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Composite Systems: Pure States, Schmidt Decomposition, Operators Acting on Composite Systems-Part 01

Hello everyone. Today's topic of discussion is composite systems. Here we will talk about tensor product space, how the entanglement comes into the picture, how we find the reduced density matrices, how do we perform measurements on a composite systems and things like that. This is one of the most important topics in foundation of quantum mechanics because this gives us the foundations of entanglement and how to deal with it. And entanglement is the resource which makes the quantum advantage possible in quantum computation and quantum key distribution and quantum communication in general.

A quantum system with more than one particle is called composite system. But there is a small confusion in this statement. That is, whenever we talk about a quantum system, we say an ensemble of particles. So, there are many, many particles. Like if you remember, when we were talking about the measurement, we had n identical particles going in the measurement setups.

And then from measuring upon those, we are finding the probabilities of different measurement outcomes. So, in that way, whenever we deal with a quantum system, we always talk about a large number of systems. Then what is so special about a composite system? So, here, when we say more than one particle, it means the single unit of quantum system, when we say, like we were saying earlier, qubit. So, qubit is one single unit, which has one particle in it.

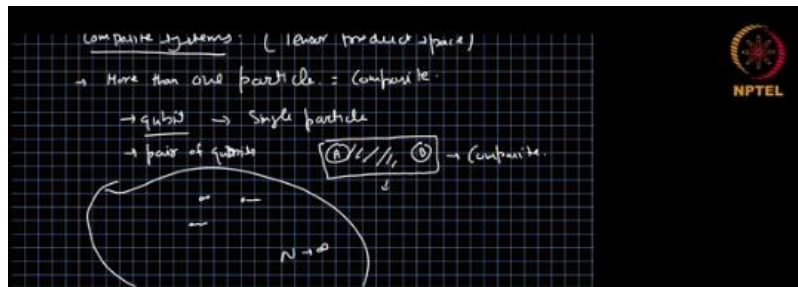
But if we have pairs of qubits coming into the picture always, this is a single particle we are seeing single particle even if we have an ensemble of particles, they are coming one by one to us, so, it's always a single particle but repeated copies multiple times. Now if we have a pair of qubits and we consider that as one quantum system, so, we have qubit A and we have qubit B and these two together we call it a quantum system. They need

not be together in the sense that they need not be present spatially, physically at the same place they can be at different places but this qubit A and qubit B belong together okay, so, in that way we will talk about the state of this two qubit system together and then this becomes our composite system. Of course, one pair of qubit is of no use in quantum mechanics, so we will have an ensemble of a pair of qubits, of course, again we are just representing them for the sake of representation but these pair of qubit need not be at the same location and they're like, n tending to infinity or very large number of those pairs.

For example, in a typical quantum communication protocol we have a source which is producing pair of qubits and it is sending those pairs at different locations. If they are in at the same location then they are not useful for communication, so we are sending them at a different location and we are tagging them with time, so, the pair produced at time t and t plus delta and t plus two delta and so on and so forth. So, we tag each pair of qubits with time. So, that becomes the identification. So, the two photos or two qubits we are getting in this way, they belong together.

So, all the processing is done on them together. Together in the sense like their measurements will be correlated or their operations will be correlated in some sense. So, this becomes a composite system. This is two qubits, we can have three qubits also, we can have four qubits, we can have not just qubits, we can have larger systems, n -dimensional systems, d -dimensional systems. So, all of them are composite systems.

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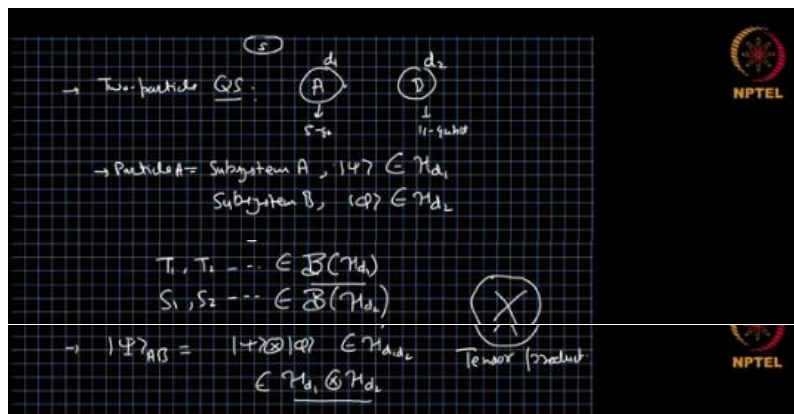


So, here we are making distinction between single particle and quantum system and more than one particle quantum system. So, more than one particle quantum system is the composite system. So, now let us take a very simple case to understand the whole scenario. Let us talk about two particle quantum systems. So, there is a party one or A and party B they can be of like five qubits and it can be of 11 qubits or some such number like we don't care just that we have divided the whole system into two parts, part A and part B. Dimension of part A is d_1 and dimension of part B is d_2 , they can be same they

can be different they can be as low as two. The part one or particle one or subsystem one sub subsystem A, as the state of that, let us call psi and they are from the Hilbert space Hd1 where d1 stands for the dimension of the Hilbert space and we are assuming the dimensions. So, the subsystem B, the states, let us call they have phi states and they belong to the Hilbert space Hd2.

We can have operators, let us say T1, T2 are the operators acting on Hilbert space Hd1. So, they belong to the set of operators acting on Hd1. Okay, let us represent the set of operators acting on Hd1 by this B of Hd1 and S1 S2 are the operators acting on subsystem B, so, they belong to set of operators acting on Hd2. Now our aim is to write the joint state of A and B because we want to represent these two subsystems together, to make a bigger quantum system which is the bipartite quantum system, so, there are two parties, so, we can, we generally use the name bipartite quantum system. So, we want to write the joint state of the two, so we have psi for state A, system A and phi for system B. We can use some symbol here and the most favorite system symbol is this, let me zoom in, this is a circle with a cross and this is called tensor product, okay, symbol for tensor product, we will see what are the properties of this.

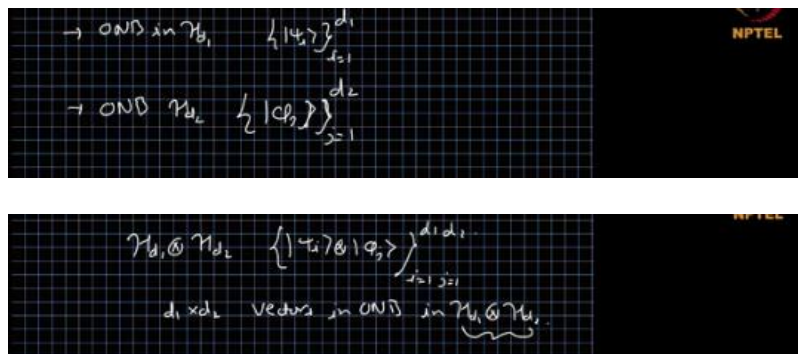
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So, whenever we write the state of two systems, where the state of subsystem A is psi and subsystem B is phi, then the joint state is psi tensor phi. And they belong to Hilbert space Hd1 tensor Hd2. Okay this is a symbolic way but actually it belongs to Hilbert space H of d1 d2, okay, so we'll see that, let us say there's an orthonormal basis in Hd1 and let us represent it with psi i where i is from 1 to d1. This is an orthonormal basis in Hd1.

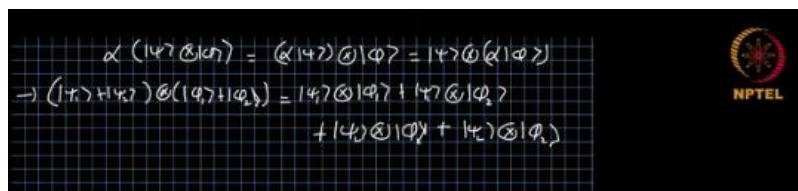
Similarly, we have orthonormal basis in H_{d2} and we represent it by ϕ_j . j is from 1 to d_2 . Okay, so now the basis for the $H_{d1} \otimes H_{d2}$ will be $\psi_i \otimes \phi_j$ where i is from 1 to d_1 and j is from 1 to d_2 . So, in that way there are d_1 times d_2 vectors in the basis. Hence the dimension of this space is d_1 times d_2 so that's why we are writing space $d_1 d_2$. Now let us discuss about the properties of the tensor products or some properties in general.

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So, if we multiply the product this tensor of ψ tensor ϕ with a scalar α then it is same as α multiplied with ψ tensor ϕ and this is same as ψ tensor $\alpha \phi$ so it's linear in both the arguments ψ and ϕ if we have $\psi_1 + \psi_2$ tensor $\phi_1 + \phi_2$, then it will be it can be opened up like a normal product this is just keeping the order, ψ_1 tensor ϕ_2 plus ψ_2 tensor ϕ_1 plus ψ_2 tensor ϕ_2 . This is another property the tensor product. We can have inner product. Let us say we have two vectors, ψ_1 tensor ϕ_1 and ψ_2 tensor ϕ_2 . And we want to calculate the inner product.

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That will be ψ_1 tensor ϕ_1 , ψ_2 tensor ϕ_2 . And that will be, so, the subsystem A states will have inner product separately and subsystem B states will have inner product separately. So, it will be ψ_1 , ψ_2 , tensor ϕ_1 , ϕ_2 . And since this is scalar and this

is scalar, the tensor product does not act on scalar, so it will be just product, $\psi_1, \psi_2,$
 $\phi_1, \phi_2.$

This is the definition of inner product and this can be understood in the following sense that, subsystem A, if we have two states in subsystem A, ψ_1 and $\psi_2,$ and we have states in subsystem B, ϕ_1 and $\phi_2,$ and they are not correlated in any sense, till now we have not talked about correlations, then they don't feel the existence of the other system. So, A does not feel the existence of subsystem B, and B does not feel the existence of A. So in their labs, the inner product will be just as if there was nothing else. So, for A it will be just $\psi_1 \psi_2$ and for B it will be $\phi_1 \phi_2.$ When we have correlation which we will discuss shortly, when we have the correlations then these inner products might be slightly more complicated, but it will satisfy this definition and then little bit of generalization of that.

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$$\begin{aligned}
 & (\psi_1 \otimes \phi_1, \psi_2 \otimes \phi_2) \\
 &= \langle \psi_1 | \psi_2 \rangle \langle \phi_1 | \phi_2 \rangle \\
 &= \langle \psi_1 | \psi_2 \rangle \langle \phi_1 | \phi_2 \rangle \\
 &= \langle \psi_1 | \psi_2 \rangle \langle \phi_1 | \phi_2 \rangle
 \end{aligned}$$

Now, what we mean by the tensor product is, let us say ψ is the vector, $s_1 s_2$ up to $s_{d1},$ some vector we are not talking about the normalization or anything is just a vector, a matrix vector, so, and ϕ is another vector, let us say with elements $p_1 p_2$ up to p_{d2} where d_1 and d_2 are the dimensions of the respective vectors then ψ tensor ϕ will be s_1 times ϕ s_2 times ϕ s_3 times ϕ and so on, so it will be a vector of $d_1 d_2$ dimension d_1 times d_2 dimension and the elements will be written like this, so, it's just a shorthand we are writing if we expand it fully then it will be ψ tensor ϕ will be $s_1 p_1 s_1 p_2$ up to $s_1 p_{d2}$ there's no line here $s_2 p_1 s_2 p_2$ up to $s_2 p_{d2}$ then it will be $s_{d1} p_1 s_{d1} p_2$ up to $s_{d1} p_{d2}.$ These will be the elements of the vector the tensor product vector ψ A B, it's such a long vector that it is not coming in one screen, so that the size of this vector will be d_1 times d_2 and it belongs to a Hilbert space dimension $d_1 d_2.$ So, this is the like working definition of defining definition of tensor product.

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$$\rightarrow |\psi\rangle = \begin{bmatrix} a_1 \\ a_2 \\ \vdots \\ a_n \end{bmatrix} = |\phi\rangle = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_m \end{bmatrix} \quad |\psi \otimes \phi\rangle = \begin{bmatrix} a_1 \phi_1 \\ a_2 \phi_2 \\ \vdots \\ a_n \phi_n \end{bmatrix}$$

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$$|\psi \otimes \phi\rangle = \begin{bmatrix} a_1 \phi_1 \\ a_2 \phi_2 \\ \vdots \\ a_n \phi_n \\ \vdots \\ a_m \phi_1 \\ a_m \phi_2 \\ \vdots \\ a_m \phi_m \end{bmatrix}$$

Now, if we have ψ_1 and ψ_2 states. So, these states belong to H_{d1} tensor H_{d2} and it's a valid Hilbert space it's a valid quantum system and stuff like that if it is a bipartite quantum system can be treated as a single quantum system then by the exams of quantum mechanics the first example of quantum mechanics the system can be in a state which is superposition of any of its states, so, a system can be $\sum_{ij} \alpha_{ij} \psi_i \otimes \phi_j$. This is the most general state of the composite system, where ψ_i and ϕ_j we have taken as the orthonormal basis or any state which exists in the Hilbert space. Now, can we always write these kinds of states as some state ψ tensor some state ϕ ? This is the question we should ask now before we start talking about the quantum correlations.

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$$\begin{aligned} & \rightarrow |\psi_1\rangle \text{ and } |\psi_2\rangle \in H_{d1} \quad |\phi_1\rangle, |\phi_2\rangle \in H_{d2} \\ & |\psi_1 \otimes \phi_1\rangle, |\psi_1 \otimes \phi_2\rangle, |\psi_2 \otimes \phi_1\rangle, |\psi_2 \otimes \phi_2\rangle \in H_{d1 \otimes d2} \\ & \rightarrow |\Psi\rangle = \sum_{ij} \alpha_{ij} |\psi_i \otimes \phi_j\rangle \end{aligned}$$

So, whenever we can write a state as a tensor product, ψ tensor ϕ , then it means we have individual identity of subsystem A and subsystem B. So, it's as if we have two qubits, for example, and one qubit is in lab A and other qubit is in lab B. We prepare the qubit A in state ψ and prepare the qubit B in state ϕ . But the state of the composite system, more general state is of this form. Now question is, can we always write it as ψ tensor ϕ ? Now let us see that we are saying that ψ is the orthonormal basis,

In H_{d1} , this is the orthonormal basis and ϕ_j is the orthonormal basis in H_{d2} . It means we can write ψ as sum over i $s_i \psi_i$ and we can write ϕ as sum over j $p_j \phi_j$. What we are saying is if the composite state $\psi \otimes \phi$ can be written as a product ψ tensor ϕ tensor product ψ tensor ϕ in for some choice of ψ , some particular specific choice of i and ϕ then ψ is some state from H_{d1} , so we can write it as a superposition of its orthonormal basis. We don't know yet what is s_i , but it can be written in this form for some suitable choice of s_i , the coefficients. Similarly, the state ϕ can also be written as a superposition of the orthonormal basis with ϕ_j with some unknown coefficient p_j . Now, this product, ψ tensor ϕ , will be sum over ij , $s_i p_j \psi_i \otimes \phi_j$. Right? But we are given that this state $\psi \otimes \phi$ is sum over ij , $\alpha_{ij} \psi_i \otimes \phi_j$.

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The image shows handwritten mathematical derivations on a grid background. At the top, it defines a state $|\psi\rangle_{AB}$ as a tensor product $|\psi\rangle_A \otimes |\phi\rangle_B$. Below this, it expresses $|\psi\rangle_A$ as a sum over basis states $|\psi_i\rangle$ with coefficients s_i , and $|\phi\rangle_B$ as a sum over basis states $|\phi_j\rangle$ with coefficients p_j . The final result shows the expansion of the tensor product state into a double sum over i and j of $s_i p_j |\psi_i\rangle \otimes |\phi_j\rangle$, which is equated to a sum over ij of $\alpha_{ij} |\psi_i\rangle \otimes |\phi_j\rangle$. An NPTEL logo is visible in the top right corner of the slide.

Comparing these two, it turns out that α_{ij} must be $s_i p_j$. Now let us say we have matrix A such that its ij th element is α_{ij} . We have a vector S, such that its i th element is s_i and we have a vector P such that its j th element is p_j . Then what we are saying is from this equation, this equation is equivalent to A matrix being S vector times P vector transpose, okay, so, these two equations are identical, that α_{ij} being $s_i p_j$ and A matrix being S times P transpose, S times P transpose is the matrix.

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$$\alpha_{ij} = S_{i1} P_{1j} \leftrightarrow A = S P^T$$

$$A_{ij} = \alpha_{ij} ; \quad S_i = s_i ; \quad P_j = p_j$$

So, these two statements are identical. Now, what is the condition over α_{ij} for it to represent a state ψ ? The condition on α_{ij} is $\sum_{ij} |\alpha_{ij}|^2 = 1$. That is the normalization. Other than that, there is no restriction over the choice of α_{ij} .

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$$A_{ij} = \alpha_{ij} ; \quad S_i = s_i ; \quad P_j = p_j$$

$$\sum_{ij} |\alpha_{ij}|^2 = 1 \quad \text{Normalization}$$

$\rightarrow A$ is arbitrary $d_1 \times d_2$ complex.

$$\text{Rank}(A) = \min(d_1, d_2)$$

So, in that way, A is arbitrary. And the dimension of A is d_1 by d_2 . And it's complex. So, A is an arbitrary d_1 by d_2 complex matrix. So, the rank A will be minimum of d_1 and d_2 in general. There can be specific cases where it will be less than this but in general or there exist lot of matrices A for which rank is minimum of d_1 d_2 , so, the rank of A is minimum of d_1 d_2 d_1 and d_2 but what is the rank of matrix S times P transpose.

The rank of matrix S times P transpose is one, okay, it's a one one rank one matrix, so, it means ψ is not generally equal to $S P^T$ or what we can say is there exist enough states which cannot be written as $\psi = \psi \otimes \phi$ so, ψ need not always be $\psi \otimes \phi$ state. What I am saying is the state ψ can be written as a superposition of tensor product states but it need not be a tensor product state itself. So, the states, which can be written as tensor product $\psi \otimes \phi$, they are called product states. They are called separable states and they are called unentangled non-entangled states.

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$\text{Rank}(S P^T) = 1$
 $A \neq S P^T$
 $|\psi\rangle_{AB} \neq |\psi\rangle \otimes |\phi\rangle$
 $\rightarrow |\psi\rangle = \sum_{ij} d_{ij} |\psi_{ij}\rangle$
 $\neq |\psi\rangle \otimes |\phi\rangle$

States which cannot be written as side tensor phi, they are called entangled state. Here entanglement can be thought of as the correlation between subsystem A and B, which cannot be achieved by preparing the states locally in A and B. If we prepare the states of subsystem A and B locally, then it will be psi tensor phi. So, entanglement is the correlation which cannot be achieved by local operations, and classical communication, because when we have classical communication, that point will be clear later on, because till now we are only considering the pure states when we talk about the mixed state and classical communication also will come into picture. So, these states cannot be prepared locally.

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$\rightarrow |\psi\rangle = \sum_{ij} d_{ij} |\psi_{ij}\rangle$
 $\neq |\psi\rangle \otimes |\phi\rangle$
 $\rightarrow |\psi\rangle_{AB} = |\psi\rangle \otimes |\phi\rangle \rightarrow$ Product states
 \rightarrow Separable states
 \rightarrow Non-entangled states
 $|\psi\rangle_{AB} \neq |\psi\rangle \otimes |\phi\rangle \rightarrow$ Entangled states

We cannot have two qubits where we prepare their states independently and then put them together. They will not be entangled. So, two qubits have to be simultaneously prepared in some state to get entanglement. These things will be more clear when we talk about mixed states. Consider the state psi to be 0 tensor 0 plus 1 tensor 1 over root 2.

So, here, alpha ij is, there are four alphas, alpha zero zero, alpha zero one, alpha one zero and alpha one one, there are four alphas alpha zero zero is one over root two, alpha zero one is zero, alpha one zero is zero and alpha one one is one over root two. So, we have two qubit system, d1 and d2 are two and we have this four coefficients. So, the matrix A becomes one over root two one zero zero one okay, so, this is the matrix A and the rank

of this matrix is two, hence, psi is entangled is an entangled state. This was a trivial example, let us take slightly more complicated example, let us say psi A B is zero, so, for shorthand known word we will be most often we will be representing psi tensor phi as psi phi. So, it will be 0, 0, 0, 1, 1, 0 and 1.

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So, let us take the coefficient of this plus, plus, plus and 1 over 2 for the normalization. So, in this case we have alpha 0 0 to be 1 over 2, alpha 0 1, alpha 1 0, alpha 1 1. So, here the matrix A is 1 over 2 1 1 1 1 now the rank of the matrix A is what, is the question. To find the rank of the matrix we see how many independent rows and columns we have. We have one column here and we have, we have the same repeating, so, we have only one independent column and similarly one independent row. So, the rank of A is one. In fact, if you see, this A can be written as a vector one one times, this is the ket vector this is the bra vector and there is a vector one over two outside. So, this can be calculated to be one one one one, okay. So, the rank of matrix A is one, this implies that the state psi is a product state. So, we can always find psi, a small psi and small phi such that psi can be written in this form.

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Now, in our earlier treatment, we saw that if A is S times P transpose, there are vectors, then it's, it's a, it's equivalent to capital psi to be psi tensor phi. Then there we, here we see that A can be written as S vector and P vector where S vector turns out to be 1 over

root 2, 1, 1, and this is same as the P vector. Hence, the state ψ can be written as $\sum_i \psi_i |i\rangle$, ψ_i are zeros and ones, so sum over i , so, this will be one over root two, zero plus one. Similarly, ϕ_i , everything is coming out of the same so it will also be one over root two zero plus one. So, you can see that if you take these ψ 's and ϕ 's then we get the same the state capital ψ which was given here. So, this was the example of separable states.

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$\text{Rank}(A) = 1$
 $A = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$
 $\Rightarrow |\psi\rangle = |1\rangle + |2\rangle$

$A = \sum p_i |i\rangle\langle i|$, $S = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \end{bmatrix} = \vec{P}$
 Hence, $|1\rangle = \sum_j S_j |j\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$
 $|\psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$

Now, let us take another example just to finish things off. Let us say again we have $0|0\rangle + 0|1\rangle + 1|0\rangle - 1|1\rangle$ over 2. This looks very much similar to the state considered in the previous example. But in this case, let us see α_{00} is 1 over 2, that is equal to α_{01} , that is equal to α_{10} , that is equal to minus of α_{11} . So, the matrix A becomes 1 over 2, 1, 1, 1, minus 1.

What is the rank of this matrix? Here again, we see that the one of the column here is 1 1 and other column is 1 minus 1. These are actually orthogonal to each other. So, they are two independent columns, similar two independent rows. So, the rank of the matrix A is 2, which is not equal to 1.

Hence, ψ is an entangled state. I hope with these examples the entanglement, like it's clear what what we mean by entanglement and how to find out whether it's entangled or not.

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$$\rightarrow |\Phi\rangle = \frac{|00\rangle + |01\rangle + |10\rangle - |11\rangle}{2}$$

$$d_{00} = \frac{1}{2} = d_{01} = d_{10} = -d_{11}$$

$$A = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \quad \text{Rank}(A) = ?$$

$$\text{Rank}(A) = 2 \neq 1$$

Hence $|\Phi\rangle \rightarrow$ Entangled state.