

Introduction to Quantum Field Theory - II (Theory of Scalar Fields)
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Lecture - 34
Renormalization Group Equation 1

(Refer Slide Time: 00:15)

Renormalization Group Equations

Base Lagrangian : λ, m 2 parameters

Renormalized λ_R, m_R, μ $d = 4 - 2\epsilon$


μ has no physical meaning

$\lambda_R(\mu), m_R(\mu) \rightarrow \lambda = \lambda(\lambda_R(\mu), m_R(\mu), \mu)$
 $m = m(\lambda_R(\mu), m_R(\mu), \mu)$

How should λ_R & m_R change with μ ?

$\lambda = Z_\lambda \mu^{2\epsilon} \lambda_R$

$\div = Z_\lambda(\lambda_R(\mu), m_R(\mu), \mu, \epsilon) \mu^{2\epsilon} \lambda_R(\mu)$



So, today we will start renormalization group equations which are differential equations that govern how Greens functions are asymmetric elements or some coupling constants evolve as you change the renormalization parameter μ . So, that is what falls under this heading and I am going to first start by talking about the simplest thing which is the coupling constant. How can we will do it in some detail and let us see what we learn from there.

So, let me remind you what we have done so far. So, we have seen that the bear lagrangian density is written in terms of 2 parameters λ and M these are bear parameters but if you write the renormalize lagrangian density then instead of 2 you have three λ_R, M_R and μ where μ is the renormalization scale and you remember that we introduced this because the coupling constant became dimension full in d dimensions where these four - 2 Epsilon.

Let me also remind you that we are using $4 - 2\epsilon$ and good now if you calculate any physical quantities in principle you can keep Epsilon fixed. So, that nothing is divergent. So, you can think of a theory being studied in d dimensions and then λ and M are finite

they are singular rule in the abstract going to zero limit but to make our arguments it is better to think in terms of Epsilon finite non-zero.

So, these are constants. So, any observable that you calculate using this original lagrangian in terms of Λ and M will not depend on μ right because it depends only on Λ and M but if you calculate the same objects using these three observables now these are these are these three variables Λ M R at sorry using these three then it looks like you have an additional parameter.

So, instead of 2 you have three and as I said last time this allows you to change the value of μ because μ is not fixed by any physical scale in the theory even when you are scattering particles and let us say you have certain incoming momenta and certain outgoing momenta there is no relation of μ to any of those momenta this is completely arbitrary this has no physical relevance no physical meaning.

So, let me write the down μ has no physical meaning. So, the way you work is you fix some value of μ choose some value of Λ R and M R and that will fix the values of Λ and M and that defines your theory. So, all physical observables are then defined once you have made a choice of μ M R and Λ R because that fixes M and Λ for you for example the physical masses are or the probabilities of certain scatterings or any schedules.

And as we are this is last time that I can choose to change the value of μ without altering the theory provided I also change Λ R and M R such that it compensates the changes in μ . So, we said that Λ R should be written as a function of μ and R should be written as a function of μ says that we keep these parameters fixed. So, this fixes the theory. So, Λ the bear the bear parameter here is a function of Λ R of μ M R of μ and μ it could depend explicitly on μ .

And also implicitly through the coupling constant and the renormalized mass and similarly the where parameter M would depend on the renormalized coupling constant renormalize mass parameter and μ the dependence will be typically both explicit and implicit through these variables. So, now let us look at the coupling constant first NC what this requirement implies for the coupling constant meaning how Λ R changes when you change μ that is the question we want to ask.

So, how should Λ_R and M_R change with μ we want to know this how exactly that behaviour is. So, what is Λ ? Λ the bare parameter is $Z\Lambda$ that is the wave function sorry that is the renormalization constant that we had introduced and then to make Λ_R to be dimensionless we had absorbed the dimension of Λ here. So, we had introduced a parameter μ and raised to the tube to ϵ now I should be writing it more explicitly.

So, I will write all the arguments of $Z\Lambda$. So, Z of Λ will be in general a function of Λ_R which in turn is a function of μ should be a will be a function of M_R which again is a function of μ the renormalization scale μ can also appear explicitly and of course you will have ϵ also appearing in the argument that is how we wrote down Λ earlier all I have done is just made all the arguments of Z of Λ exclusive.

So, what I have done is I have written all the quantities that can in principle appear without leaving out anything unless I have an argument why a certain parameter will not appear I should write all the things I know when I renormalized I have to subtract Infinities poles will be there. So, there is absolute dependence of course when you calculate fermion and diagrams they will depend on Λ_R and M_R and you have already seen $\log \mu$ appearing explicitly.

So, there is an explicit new dependence as well. So, I have listed down everything both the implicit dependence on μ and explicit dependence on here good but see left hand side doesn't know about μ right hand side does. So, if I change μ Λ should not change but right hand side even though individual pieces $Z\Lambda\mu^{-2\epsilon}$ and Λ_R individually they will change.

But they all should change together in a manner that the dependence completely gets cancelled because the left hand side is independent and that is all I am going to utilize now no great wisdom just this much.

(Refer Slide Time: 09:23)

$$\lambda = z_\lambda \mu^{2\epsilon} \lambda_R \quad \mu \frac{d\lambda}{d\mu} = 0$$

$$0 = \mu \frac{d\lambda}{d\mu} = \mu \frac{dz_\lambda}{d\mu} \mu^{2\epsilon} \lambda_R + 2\epsilon z_\lambda \mu^{2\epsilon} \lambda_R + z_\lambda \mu^{2\epsilon} \mu \frac{d\lambda_R}{d\mu}$$

Divide by $z_\lambda \mu^{2\epsilon}$


$$\lambda_R \frac{1}{z_\lambda} \left(\mu \frac{dz_\lambda}{d\mu} \right) + 2\epsilon \lambda_R + \mu \frac{d\lambda_R}{d\mu} = 0$$

Define β -function:

$$\tilde{\beta} = \mu \frac{d\lambda_R}{d\mu} \quad ; \quad \tilde{\beta} \rightarrow \text{finite because } \lambda_R \rightarrow \text{finite in } \epsilon \rightarrow 0 \text{ limit.}$$

$$\tilde{\beta} = -2\epsilon \lambda_R - \lambda_R \frac{1}{z_\lambda} \left(\mu \frac{dz_\lambda}{d\mu} \right)$$

Since $\tilde{\beta} \rightarrow \text{finite}$ $\frac{1}{z_\lambda} \left(\mu \frac{dz_\lambda}{d\mu} \right) \rightarrow \text{also finite in } \epsilon \rightarrow 0 \text{ limit.}$



So, because Lambda is a bare object an independent of mu is this recording must be. So, d Lambda over d mu should be 0. That is the simple statement that Lambda does not depend on mu. Now this is zero I can multiply mu here you cannot stop me from doing that. So, I will just multiply that the statement is still true and it is customary to introduce mu here as this Factor now what it means for the individual factors.

So, I will just take the derivative on the right hand side of this equation. So, 0 is Mu d Lambda or d mu that is the expression. So, I am just differentiating that maybe I will write it here again to oh fine. So, the first factor is z you remember that. So, mu dz Lambda over d mu and the other factors were mu to the 2 Epsilon and Lambda R at it here Lambda is Z Lambda u to the 2 Epsilon and Lambda R.

I am differentiating this then the second term is when I differentiate this mu to the 2 Epsilon it gives me a 2 Epsilon times mu to the 2 Epsilon - 1 but then I have a mu multiplying here. So, that makes it again mu to the 2 Epsilon. So, I get and finally a differentiate Lambda R. So, I get Z of Lambda mu to the 2 Epsilon and mu d Lambda R over d B that is what I get now let us divide by Z Lambda mu to the 2 Epsilon if I do.

So, then I get. So, there is a factor of Lambda r. So, this mu 2 the 2 Epsilon goes away from each of the terms here in this one I will get. So, I am writing this Lambda R first Lambda R then Z of Lambda will come in the denominator one of one over Z Lambda mu d z Lambda over do you know Plus 2 Epsilon I have divided by Z Lambda. So, this goes away mu to the 2 Epsilon goes away and I am left with Lambda R and in this term.

Again Z Lambda goes away μ to the 2ϵ goes away and you are left with $\mu \frac{d\Lambda}{d\mu}$ R over $d\mu$ is equal to 0 and it is $\frac{d\Lambda}{d\mu} R$ over $d\mu$ that is what you want to calculate you want to know how ΛR changes as you change μ . So, that thing is here which you are searching for. So, let us define this as β tilde it is called beta function define beta function which is β is equal to $\mu \frac{d\Lambda}{d\mu} R$ over d .

So, from the definition itself it is clear because ΛR is a finite object it is a finite number as ϵ goes to zero and μ is anywhere parameter this β tilde is a finite object β tilde is finite because ΛR is finite we are not expecting the derivative to blow up. So, this group is finite. So, substituting this in this equation I will get. So, this is β . So, β tilde I take these 2 terms on the other side it is other side it is $-2\epsilon \Lambda R - \Lambda R$ over $Z \Lambda \mu \frac{d\Lambda}{d\mu}$ over d .

Now β tilde or the beta function is finite in ϵ I am going to zero limit this object is finite in the ϵ going to zero limit β tilde moves to zero. So, that goes to zero ΛR is anyway finite and that implies that 1 over $Z \Lambda \mu \frac{d\Lambda}{d\mu}$ this factor is also finite. So, since we total lies finite 1 over $Z \Lambda \mu \frac{d\Lambda}{d\mu}$ is also finite. Now I will look at the let us suppose that we are doing the calculation in the \overline{MS} or \overline{MS} scheme.

And you remember the advantage of using \overline{MS} or \overline{MS} scheme is that $Z \Lambda$ and $Z M$ is the renormalization constant for the mass parameter they do not depend on the scale μ and this is an all order statement and that further implies that $Z \Lambda$ and $Z M$ they do not depend on $M R$ also.

(Refer Slide Time: 16:37)

In MS (MS)

$\rightarrow Z_\lambda, Z_m$ do not depend on μ explicitly.

\rightarrow " " " " " " " " m_R .


This is why MS scheme is called mass-independent renormalization scheme.

$$Z_\lambda = 1 + \lambda_R \frac{a_{1,1}}{\epsilon} + \lambda_R^2 \left(\frac{a_{2,2}}{\epsilon^2} + \frac{a_{2,1}}{\epsilon} \right) + \dots + \lambda_R^n \left(\frac{a_{n,n}}{\epsilon^n} + \dots + \frac{a_{n,1}}{\epsilon} \right) + \dots$$

$$Z_m = 1 + \lambda_R \frac{b_{1,1}}{\epsilon} + \lambda_R^2 \left(\frac{b_{2,2}}{\epsilon^2} + \frac{b_{2,1}}{\epsilon} \right) + \dots + \lambda_R^n \left(\frac{b_{n,n}}{\epsilon^n} + \dots + \frac{b_{n,1}}{\epsilon} \right) + \dots$$

a_{ij} & b_{ij} are constants.

$\ln 4\pi - \gamma_E$
 $\left(\frac{\mu}{m_R}\right)$
 $1 + f(m)$



So, let me say this I will be using now MS or MS bar scheme I will call them together as MS minimal subtraction because the difference between MS and MS bar is purely a constant. It does not involve any mu the difference is just log mu - sorry log of 4 Pi - solid gamma you have seen this log of 4 Pi - all of gamma. So, in MS you subtract only the poles and in MS bar you subtract together with the pole this Factor sorry these 2 terms which always accompany the poles that is a statement at one Loop but similar statement is through it hair loops also.

So, I have already told you that Z Lambda Z M do not depend on mu explicitly and because Z is 1 plus order Lambda term and one is dimensionless and because mu does not appear you cannot construct any dimensionless ratios meaning you because you do not have mu there is no way you can construct such an object which is dimensionless. So, there is no way Mr can appear because 1 plus if you were to construct some function which has M R it will be dimension full.

The only way it could have been dimensionless is if you had mu also available to you then you could have divided mu or M R mu by M R and that ratio will be dimensionalized but that is not there. So, it also does not depend on mu and for this reason these are these schemes are called mass independent renormalization schemes because there is no mass scale appearing in these schemes. So, this is why MS schemes scheme is called mass independent renormalization scheme.

So, now let me tell you how Z_Λ would look like and Z_M would look like if you were to calculate to had orders in perturbation theory. So, the general form of Z_Λ would be this and you have already seen at one Loop what it looks like. So, it starts with one because you do not require renormalization. If you are at the tree level tree level does not require any renormalization there are no infinities.

So, Z_Λ is one then typically you will get corrections at order ΛR . So, that is ΛR some of the coefficients may vanish but in general you will get non-zero coefficients now you have seen that you get a pole simple Pole or one over Epsilon pole. So, you are going to get something here which is proportional to one over Epsilon times some constant what you have in the numerator is a constant it does not depend on $\mu M R$ and scales neither the external momenta they are irrelevant.

So, what you get is a constant and let me call that constant as a one, 1 the first one here stands for the order of Λ and the second one here stands tells you that you have a pole simple Pole right the Epsilon power is one that is what it is telling you then if you go to order Λ^2 meaning you are going higher order in perturbation Theory then you will encounter 2 loops then you will have in addition to a single pole you will have a double pole also.

So, the highest order pole will be one over Epsilon Square it will be more Divergent and again the coefficient will be a constant and I will call it A_2 that 2 first 2 stands for the second order in ΛR , 2 and this 2 stands. For the second order of the pole so, this is a order 2 pole and that is why this 2 here and you will also have a single pole. So, you will have a $2, 1$ over Epsilon this 2 refers to again the fact that you are at ΛR^2 and this one tells you that you are looking at a simple pole 1 over Epsilon to the one.

And the general structure is that it will be ΛR to the n a_n . So, you will get order n Pole at order ΛR to the n then you will have lower poles and this n Z finally a simple pole at order n . So, at every order you are going to get poles off. So, at order n you are going to get poles of order $n - 1$ $n - 2$ up to order 1 a simple pole. So, all these will be present. So, that is the structure of Z_Λ .

So let me also write down what Z M would look like it will be identical except for the difference in coefficient. So, let us call let us call them b b 1 b 2 and so on so forth. So, that is the structure and as I said a ij and maybe I should use a ij and b ij are constants. So, that is the general structure of these renormalization constants in MSP or M is working the same thing. Now I can of course rewrite it slightly differently.

(Refer Slide Time: 24:19)


$$Z_\lambda = 1 + \frac{1}{\epsilon} (\lambda_1 a_{1,1} + \lambda_2^2 a_{2,1} + \dots + \lambda_n^n a_{n,1}) + \frac{1}{\epsilon^2} (\dots) + \dots$$

$$Z_\lambda = 1 + \frac{A_1}{\epsilon} + \frac{A_2}{\epsilon^2} + \frac{A_3}{\epsilon^3} + \dots$$

$$Z_m = 1 + \frac{B_1}{\epsilon} + \frac{B_2}{\epsilon^2} + \frac{B_3}{\epsilon^3} + \dots$$

Now we will see the $\frac{1}{Z_\lambda} \left(\mu \frac{dZ_\lambda}{d\mu} \right)$ does not depend on μ explicitly.

$$\frac{1}{Z_\lambda} \left(\mu \frac{dZ_\lambda}{d\mu} \right) = \frac{1}{Z_\lambda} \left[\mu \frac{\partial}{\partial \mu} + \tilde{\beta} \frac{\partial}{\partial \lambda_k} + d_f \frac{d m_k}{d\mu} \frac{\partial}{\partial m_k} \right] Z_\lambda$$

$$= \tilde{\beta} \frac{1}{Z_\lambda} \frac{\partial Z_\lambda}{\partial \lambda_k}$$


And let us instead of writing it as an expansion in Lambda I will write it as an expansion in the pole the order of the poles and then we are going to see something very interesting. So, I can write Z Lambda as 1 plus so let us collect all the order one over Epsilon terms. So, it will be what it will be Lambda R times a 1,1 that is the coefficient of 1 over Epsilon coming from one Loop the lowest order diagram which is proportional to Lambda R.

Then you will have contributions coming from Lambda R square terms from higher order calculations from order Lambda R square calculation with description and. So, forth yeah I am just regrouping I am not doing anything major and then you will also have other terms like one over Epsilon Square terms and you can similarly write what should be here plus one over Epsilon cube and so forth.

So, the structure will be that Z Lambda I can write as 1 plus A 1 over Epsilon. So, this Epsilon and A 1 is this piece. So, even depends on Lambda R it is a polynomial in Lambda R of I mean in principle you have this goes to this keeps continuing right you I should not have stopped here because you can calculate up to infinite orders. So, this is a polynomial then you

will have one over Epsilon Square terms it is called the coefficient to be A^2 and A^2 is again a function of ΛR .

And the dependence is like this Plus and so forth and the same is true for Z of M where A^1 is for example listing. Now before proceeding further let me first show you that β tilde or the β function does not depend on μ explicitly in the \overline{MS} scheme again that is what I am going to utilize later. So, let me do that. So, from here it is not. So, much clear that β tilde does not depend on μ explicitly right.

So, this term as far as this term is concerned it is clear that dependence on μ is through μR . So, that is implicit this factor is implicit but then you have a factor of μ here. So, it is not. So, obvious that it does not depend on μ explicitly right the Z 's do not depend on you explicitly because you are in \overline{MS} scheme and I should argue that this entire Factor also here at the bottom does not depend on you explicitly do we see that.

So, 1 over $Z \Lambda$ times $\mu \frac{dZ \Lambda}{d\mu}$ this is 1 over $Z \Lambda$ and then let us write the total derivative as the partial derivatives. So, $\mu \frac{\partial}{\partial \mu} + \frac{\partial}{\partial \Lambda R} \times \nu \frac{\partial \Lambda R}{\partial \mu} + \frac{\partial}{\partial \mu} \Lambda R$ and that is what is called β tilde Plus $\frac{\partial}{\partial M R} \frac{d M R}{d \mu}$ acting on Z of Λ and as I have said already that $Z \Lambda$ is independent of $M R$.

So, this derivative acting on this gives you zero similarly this derivative acting on $Z \Lambda$ gives you 0 because there is no explicit dependence on μ . So, that also gives you 0. So, what we are left which is only this term. So, you get β tilde $\frac{\Delta Z \Lambda}{\Delta \Lambda R} \times \frac{1}{Z \Lambda}$. So, that is what you get now let me put that back in this equation this let me box it hmmm it is definitely not bad I want also maybe this equation.

So, yeah so I will substitute that thing in here I have just written down $\frac{1}{Z \Lambda} \mu \frac{dZ \Lambda}{d\mu}$ as β tilde $\frac{1}{Z \Lambda} \frac{\Delta Z \Lambda}{\Delta \Lambda R}$. Now the derivative is with respect to ΛR instead of μ . So, I will go I am going to put this in here. So, what do I get? So, I will write it again.


(Refer Slide Time: 32:33)

$$\begin{aligned} \tilde{\beta} &= -2\epsilon\lambda_R - \lambda_R \frac{1}{Z_\lambda} \mu \frac{dz_\lambda}{d\mu} \\ &= -2\epsilon\lambda_R - \lambda_R \tilde{\beta} \cdot \frac{1}{Z_\lambda} \cdot \frac{dz_\lambda}{d\lambda_R} \end{aligned}$$

\Rightarrow $\tilde{\beta}$ does not depend explicitly on μ .
 • || || || • || on m_R

$$\tilde{\beta}(\lambda_R, m_R, \epsilon) \equiv \tilde{\beta}(\lambda_R, \epsilon) \quad \text{in MS Scheme}$$

Now, we will see that only single pole terms contribute to β -function.



So, we had there we go so, that equation was beta tilde is equal to $-2\epsilon\lambda_R - \lambda_R \frac{1}{Z_\lambda} \mu \frac{dz_\lambda}{d\mu}$ and I have found that $\frac{1}{Z_\lambda} \mu \frac{dz_\lambda}{d\mu}$ is beta tilde times this factor. So, it becomes this is beta tilde times $\frac{1}{Z_\lambda} \frac{dz_\lambda}{d\lambda_R}$.

Now if you look at this equation there is mu dependence which is implicit here through λ_R again implicit through Z_λ then you have a beta tilde Z_λ does not depend on mu explicitly. So, the derivative with respect to λ_R also gives you a function that does not depend explicitly on mu and this of course $\frac{1}{Z_\lambda}$ does not depend explicitly on mu.

So, when you solve this equation you will get for beta tilde function that does not depend explicitly on mu. So, from here you conclude that beta tilde does not depend on mu plus settling on another thing for the same argument from the same argument you also see that beta tilde does not depend on m_R right. Because m_R does not appear anywhere here because the only way it could have entered is through Z_λ but Z_λ and MS scheme does not depend on m_R .

So, beta tilde also does not depend on m_R good. So, yeah so what do I have. So, I started with beta tilde as a function of λ_R, m_R and mu and now I have argued that this is and of course would also be a function of epsilon but now I have concluded that beta tilde is only a

function of Λ_R it does not depend on μ and M_R is μ explicitly and there is no depend on M_R excellent.

Now yeah now what I will show you is that if you look at the beta function and beta function is important because that is telling you how Λ_R changes as you change μ . So, that you know your a bear coupling constant does not change remember we always we are always trying to ensure that when we change μ . We do not start changing the theory and that is what is ensured by making sure that the bear parameters do not change.

And that has led us to the beta function and beta function is telling you how Λ_R the renormalized coupling constant changes when you change μ . So, as of now I have concluded that beta function is independent of μ and M_R and is dependent only on Λ_R and ϵ and also I can calculate that because here beta tilde can be calculated once you know Z of Λ because you just take these derivatives.

And Z Λ how do you get? Z Λ are the counter terms that you I mean this coupling const sorry renormalization constants that you get by subtracting the Infinities which is arranged by the counter terms. So, when you do a loop calculation to one Loop or 2 Loop or whatever Loop order that will give you the expression of Z Λ because Z Λ or Z M these are the things which you adjust to remove the poles.

So, a loop calculation will give you that Λ take the derivative and that will give you beta Λ beta tilde that is how you obtain the beta function. So, now I want to show you that the beta function is receives contributions only from the single pole terms. So, what I mean to say is keep this. So, if you look at this expression which is going to give you beta tilde upon taking these derivatives of Z it is it looks like that all the poles will contribute right because Z Λ has poles simple poles.

Poles of order 2 meaning one over ϵ Square terms poles of order three and even our ϵ Cube terms as I wrote somewhere here. So, Z Λ has all kinds of terms right on our ϵ^{-1} over ϵ^2 one over ϵ^3 what I am going to show you is that these terms these coefficients A_3 A_2 they do not appear in the beta function what appears is only A_1 .

And what is A_1 ? A_1 gets a contribution from one order Lambda term lowest order calculation will give you a $1/\epsilon$ second order calculation will give you a $2/\epsilon^2$ and so forth but the result is going to depend only on one over Epsilon sorry the coefficients of only one over Epsilon will contribute. Meaning if you do a calculation at one Loop 2 loop 3 Loop and so forth what is relevant for beta function is just the coefficient of simple pole terms from one loop from 2 loop from 3 Loop etcetera.

You do not have to include you will not be including this kind of terms that is what I am going to show you now. So, now we will see only single pole terms contribute to Beta function. So, let us see that um.

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The slide contains the following handwritten equations and notes:

$$\mu \frac{dZ_\lambda}{d\mu} = \tilde{\beta} \frac{dZ_\lambda}{d\lambda_R} = \tilde{\beta} \left(\frac{1}{\epsilon} \frac{dA_1}{d\lambda_R} + \frac{1}{\epsilon^2} \frac{dA_2}{d\lambda_R} + \dots \right) \quad (2)$$

Recall $\frac{1}{Z_\lambda} \mu \frac{dZ_\lambda}{d\mu} \rightarrow f(\lambda_R)$ (3)
 $f(\lambda_R) \rightarrow f(\lambda)$

$$f(\lambda_R) = \frac{\tilde{\beta}}{Z_\lambda} \left(\frac{1}{\epsilon} \right) \left[\frac{dA_1}{d\lambda_R} + \frac{1}{\epsilon} \frac{dA_2}{d\lambda_R} + \dots \right]$$

$$\tilde{\beta} = -2\epsilon \lambda_R - \lambda_R f(\lambda_R)$$

$$\left(1 + \frac{A_1}{\epsilon} + \frac{A_2}{\epsilon^2} + \dots \right) f(\lambda_R) = (-2\lambda_R - \frac{1}{\epsilon} \lambda_R f(\lambda_R)) \left(\frac{dA_1}{d\lambda_R} + \frac{1}{\epsilon} \frac{dA_2}{d\lambda_R} + \dots \right)$$

Compare powers of ϵ on both sides
 $f = -2\lambda_R \frac{dA_1}{d\lambda_R}$

NPTL

So, let us look at $\mu \frac{dZ_\lambda}{d\mu}$ this is we have seen just now that there is that this is $\mu \frac{dZ_\lambda}{d\mu}$ is beta tilde Z_λ over $\Delta \lambda_R$ I'm just dropping out one over Lambda from both the sides. So, it is beta tilde actually I can replace by total derivative it does not matter because there is no other variable in Z_λ lights just λ_R . So, let me in fact.

So, what does that give you it gives you beta tilde. Now we are taking the derivatives. So, Z_λ the first term is one differentiated with λ_R gives you 0. So, that goes away second term is $1/\epsilon$. So, the expression of Z_λ that I am taking is this one in this form not in this form not this one when they are both the same as there is no difference but this is how I have organized in terms of poles.

So, that gives you one over Epsilon and the derivative of the simple pole term the coefficient of the simple pole plus $\frac{1}{\epsilon^2} \frac{d}{d\lambda} R + \frac{1}{\epsilon^3}$ and so forth now also recall that we have already argued that $\frac{1}{Z} \mu \frac{dZ}{d\lambda}$ is finite and this thing I am going to call as $f(\lambda, R)$. So, $f(\lambda, R)$ is finite that is just a name given to this factor this object.

So, now you substitute this in one. So, what you get is. So, I am just which one did I call something else one earlier one let me call this 2 let me call this 3. So, I am just putting this here in this equation 2. So, you get $f(\lambda, R)$. So, I'm dividing the equation 2 by Z of λ Z of λ . So, that gives you $f(\lambda, R)$ sorry $f(\lambda, R)$ is equal to $\tilde{\beta}$ that is the $\tilde{\beta}$ here over Z of λ .

Now I will pull out one over Epsilon times what you have here is $\frac{d}{d\lambda} R + \frac{1}{\epsilon} \frac{d^2}{d\lambda^2} R$ plus other higher order pole terms. Now we can write $\tilde{\beta}$ using $f(\lambda, R)$. So, $\tilde{\beta}$ if you recall it was $-2\epsilon \lambda - \lambda$ times $f(\lambda, R)$ sorry $f(\lambda, R)$ right let me show you if you do not recall here. So, $\tilde{\beta}$ is $-2\epsilon \lambda - \lambda$ times this is what I have defined as $f(\lambda, R)$ which is finite $f(\lambda, R)$ sorry keep saying $f(\lambda, R)$.

So, that is good now hmm now I will take this Z of λ and multiply on the left hand side and Z of λ is $1 + \frac{1}{\epsilon} + \frac{A^2}{\epsilon^2}$ and so forth into $f(\lambda, R)$ that is $f(\lambda, R)$ that is equal to. So, this is a λ is on the other side Now $\tilde{\beta}$ where $\tilde{\beta}$ is $-2\epsilon \lambda - \lambda$ and I will divide by this Epsilon. So, it gives you $-2\lambda - \lambda$ it is λ times $f(\lambda, R)$.

And then because of this division by Epsilon it gives you a factor of $\frac{1}{\epsilon} \frac{d}{d\lambda} R + \frac{1}{\epsilon^2} \frac{d^2}{d\lambda^2} R$ and other higher order terms. Now let us compare on both sides powers of Epsilon. So, compare powers of Epsilon on both sides. So, what do you get let us look at the lowest order term which is I mean the finite term of $\frac{1}{\epsilon^0}$ or Epsilon to the zero you get 1 times $f(\lambda, R)$.

So, on the left you have f on the right you have $-2\lambda - \lambda$ times $\frac{d}{d\lambda} R$ any other term will be either one over Epsilon or one over Epsilon square and so forth. So, if I am just comparing Epsilon to the zero term on the left then it is f is equal to $-2\lambda - \lambda$

times dA_1 over $d\lambda_R$ remember what f is f is this piece and F is basically this piece this piece here get that factor here.

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$$f = -2\lambda_R \frac{dA_1}{d\lambda_R}$$

In MS scheme

$$\tilde{\beta} = -2\epsilon\lambda_R - \lambda_R \left(-2\lambda_R \frac{dA_1}{d\lambda_R} \right)$$

$$\tilde{\beta} = -2\epsilon\lambda_R + 2\lambda_R^2 \frac{dA_1}{d\lambda_R}$$

So, what do I get let me write again f is now $-2\lambda_R$ times dA_1 over $d\lambda_R$. So, in the \overline{MS} scheme or \overline{MS} bar also we get $\tilde{\beta}$ equal to $-2\epsilon\lambda_R - \lambda_R$ times this Factor right that factor is $f - 2\lambda_R$ check yeah there is a explicit factor of λ_R here λ_R times 1 over Z λ_R μ dz λ_R or $d\mu$ right that is what we have.

That is what we have calculated 1 over Z λ_R μ dz λ_R over $D\mu$. So, that gives you what was that $-2\lambda_R dA_1$ over $d\lambda_R$ so, $\tilde{\beta}$ is $-2\epsilon\lambda_R$ plus $2\lambda_R^2 dA_1$ over $f\lambda_R$ I hope I did not make any mistakes. So, what does that tell it tells that when you are looking at the beta function only A_1 contributes and what was A_1 A_1 was here A_1 was the coefficient of 1 over Epsilon term. So, you see all these other poles of higher order they have not contributed.

So, if you as I was saying earlier if you were to calculate a three loop or for loop whatever for the beta function only the coefficient of simple pole is relevant not the poles of higher orders good. So, we have now some understanding about the renormalization group of group equation for the coupling constant and we also understand what kind of terms contribute to the beta function. We will continue our discussion further in the next video.