

Statistical Mechanics
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Lecture - 64
Discontinuity in the Specific Heat of a Bose Gas - Part 02

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Handwritten notes on a whiteboard:

$$\left. \frac{\partial N}{\partial \mu} \right|_{T, V} = \frac{C_\alpha (\alpha-1) \Gamma(\alpha-1)}{\beta_c^{\alpha-1}} \zeta(\alpha-1)$$

where $\beta_c^{\alpha-1} = \beta_c^{\alpha-1}$ and $\alpha-1 > 1$.

$$\left. \frac{\partial N}{\partial T} \right|_{\mu, V} = ?$$

$$N = C_\alpha \int_0^\infty d\epsilon \frac{g(\epsilon)}{z^{-1} e^{\beta \epsilon} - 1} = \frac{C_\alpha}{\beta^\alpha} \int_0^\infty d(\beta \epsilon) \frac{(\beta \epsilon)^{\alpha-1}}{z^{-1} e^{\beta \epsilon} - 1}$$

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So, we start off with the expression for N, which was $C_\alpha \int d\epsilon g(\epsilon) z^{-1} e^{-\beta \epsilon} - 1$. And we know how to convert this, because this $g(\epsilon)$ goes as $\epsilon^{\alpha-1}$.

So, $d(\beta \epsilon)$ times $\beta^{\alpha-1}$ divided by $z^{-1} e^{\beta \epsilon} - 1$ divided by β^α . Now, this is a

repeat, this is a repetitive statement, but it is just to ensure that if you have forgotten you can start off from the scratch.

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$$\begin{aligned}
 N &= C_\alpha \int_0^\infty d\epsilon \frac{\beta(\epsilon)}{z^{-1} e^{\beta\epsilon} - 1} = \frac{C_\alpha}{\beta^\alpha} \int_0^\infty d(\beta\epsilon) \frac{(\beta\epsilon)^{\alpha-1}}{z^{-1} e^{\beta\epsilon} - 1} \\
 &= \frac{C_\alpha}{\beta^\alpha} \int_0^\infty dx \frac{x^{\alpha-1}}{z^{-1} e^x - 1} = \frac{C_\alpha}{\beta^\alpha} \frac{\Gamma(\alpha)}{\Gamma(\alpha)} \int_0^\infty dx \frac{x^{\alpha-1}}{z^{-1} e^x - 1} \\
 N &= \frac{C_\alpha}{\beta^\alpha} \Gamma(\alpha) \mathcal{F}_\alpha^+(z) \quad z = 1 \\
 \frac{\partial N}{\partial T} &= \frac{\partial}{\partial T} \frac{C_\alpha}{\beta^\alpha} \Gamma(\alpha) \mathcal{F}_\alpha^+(z=1) = \frac{\partial}{\partial T} \frac{C_\alpha}{\beta^\alpha} \zeta(\alpha)
 \end{aligned}$$



So, $C_\alpha \beta^\alpha$ and then you see that, this integral is now 0 to infinity, 0 to infinity, 0 to infinity $dx x$ to the power $\alpha - 1$ $z^{-1} e$ to the power $x - 1$. So, that this gives me C_α over β^α and I need $\Gamma(\alpha)$ divided by $\Gamma(\alpha)$, because I need $\alpha - 1$ factorial to relate it to the Bose integrals that we have defined, $dx x$ to the power $\alpha - 1$ $z^{-1} e$ to the power $x - 1$. And this is C_α for β^α times $\Gamma(\alpha)$ $\mathcal{F}_\alpha^+(z)$.

Now, I want to take a derivative with respect to temperature, chemical potential being held constant, but this also has to be evaluated very close to T_C . So, to make my life simple, I mean I can simply say that look, I can replace T_C everywhere, but then my derivative is

going to be identically 0. So, what we will do is we will replace this, by saying that I have z equal to 1.

So, that del N del T is going to be say, del del T of C alpha gamma alpha raised to the power beta alpha F alpha plus 1 evaluated at z equal to 1 extremely close to T C and this becomes del del T of C alpha gamma alpha raised to the power beta sorry, divided by beta raised to the power alpha and then I have zeta of alpha, the Riemann zeta value.

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$$\begin{aligned}
 &= C_\alpha \Gamma(\alpha) \zeta(\alpha) \frac{\partial}{\partial T} \left(\frac{1}{\beta^\alpha} \right) \\
 &= C_\alpha \Gamma(\alpha) \zeta(\alpha) \frac{\partial}{\partial \beta} \left(\frac{1}{\beta^\alpha} \right) \frac{\partial \beta}{\partial T} \\
 &= C_\alpha \Gamma(\alpha) \zeta(\alpha) \left(\frac{\partial}{\partial \beta} \frac{1}{\beta^\alpha} \right) \cdot \left(\frac{\partial \beta}{\partial T} \right) \\
 &= \frac{\alpha}{T} \frac{C_\alpha \Gamma(\alpha) \zeta(\alpha)}{\beta^\alpha} \quad \beta \rightarrow \beta_c \quad \begin{aligned} \frac{\partial \beta^{-\alpha}}{\partial \beta} &= -\alpha \beta^{-(\alpha+1)} \\ \frac{\partial \beta}{\partial T} &= \frac{\partial}{\partial T} \left(\frac{1}{k_B T} \right) \\ &= -\frac{1}{k_B T^2} \\ &= -\frac{1}{(k_B T)} \frac{1}{T} \\ &= -\frac{\beta}{T} \end{aligned} \\
 &= \frac{\alpha}{T} \frac{C_\alpha \Gamma(\alpha) \zeta(\alpha)}{\beta_c^\alpha}
 \end{aligned}$$



This quantity then becomes C alpha gamma alpha times zeta of alpha and then, I have derivative of del del T of 1 over beta raised to the power alpha, which is C alpha gamma alpha zeta of alpha del del beta of 1 over beta to the power alpha times del beta del T. This is zeta alpha.

Now, $\frac{d\beta}{d\alpha}$ is sorry, $\frac{d}{d\alpha} \beta^\alpha$ is $\alpha \beta^{\alpha-1}$. And $\frac{d\beta}{dT}$ is going to be $\frac{1}{k_B T^2}$, which is going to be $-\frac{1}{k_B T^2}$ and which I can write, as $-\frac{1}{k_B T} \times \frac{1}{T}$. So, that this becomes $-\alpha \beta^{\alpha-1} \times \frac{1}{k_B T^2}$.

So, I have $-\alpha \beta^{\alpha-1} \times \frac{1}{k_B T^2}$ times $-\beta^\alpha$ divided by T and you see, that 1 factor of β cancels with this 1 factor sitting over here and the two minus 2 negative signs gives me a plus.

So, this becomes $\frac{\alpha \beta^\alpha}{k_B T^3}$ and here, in the place of β raised to the power α , since I am extremely close to $T C$, I am now, going to substitute this as βC . So, β now, goes to βC only after you have taken the derivative not before that. So, that $\frac{\alpha (T C)^\alpha}{k_B T^3}$.

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$$\begin{aligned}
 &= C_\alpha \Gamma(\alpha) S(\alpha) \left(\frac{\beta^\alpha}{\beta^{\alpha+1}} \right) \cdot \left(\frac{\partial \beta}{\partial T} \right) && \text{or } \frac{\partial \beta}{\partial T} \\
 &= \frac{\alpha}{T} \frac{C_\alpha \Gamma(\alpha) S(\alpha)}{\beta^\alpha} && \boxed{\beta \rightarrow \beta_C} && = -\frac{1}{k_B T^2} \\
 & && T \rightarrow T_C && = -\frac{1}{(k_B T)} \frac{1}{T} \\
 & && && = -\frac{\beta}{T} \\
 &= \frac{\alpha}{T} \frac{C_\alpha \Gamma(\alpha) S(\alpha)}{\beta_C^\alpha} && \boxed{N = \frac{C_\alpha \Gamma(\alpha) S(\alpha)}{\beta_C^\alpha}} \\
 \left. \frac{\partial N}{\partial T} \right|_{\mu} &= \frac{\alpha N}{T}
 \end{aligned}$$



Recall that this C_α , $\Gamma(\alpha)$ divided by β C raised to the power α times ζ α ; the condition for the condensation to happen at β_C is given by this. So, you realize that this quantity that you see over here is nothing but N .

So, I have αN divided by T_C . So, here also I have T_C , because I have replaced β by β_C , which effectively also means that I have replaced T by T_C . So, $\left. \frac{\partial N}{\partial T} \right|_{\mu}$ being held constant is given by this. So, now, we have expression for all the partial derivatives we need we have this value, we have this value.

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$$\rightarrow \left(\frac{\partial \mu}{\partial T} \right)_N = - \frac{\left(\frac{\partial N}{\partial T} \right)_\mu}{\left(\frac{\partial N}{\partial \mu} \right)_T}$$

$$N = \frac{C_x}{\beta^x} \int_0^\infty dx \frac{x^{x-1}}{z^{-1} e^x - 1}$$

$$\frac{\left(\frac{\partial N}{\partial \mu} \right)_T}{\left(\frac{\partial N}{\partial T} \right)_T} = \frac{\left(\frac{\partial N}{\partial z} \right)_T \frac{dz}{d\mu}}{\left(\frac{\partial N}{\partial T} \right)_T} = \left(\beta z \right) \frac{\partial N}{\partial z} \Big|_T$$

$$\frac{\partial N}{\partial z} \Big|_T = \frac{C_x}{\beta^x} \int_0^\infty dx x^{x-1} \frac{d}{dz} \left(\frac{1}{z^{-1} e^x - 1} \right)$$

$$\frac{d}{dz} \frac{1}{z^{-1} e^x - 1} = -\frac{1}{z} \frac{d}{dz} \frac{1}{z^{-1} e^x - 1}$$



And the quantity that we are looking for, is essentially $\left(\frac{\partial \mu}{\partial T} \right)_N$, which is $\left(\frac{\partial N}{\partial T} \right)_\mu$ held constant, $\left(\frac{\partial N}{\partial \mu} \right)_T$ held constant with a minus sign in front of it.

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$$\begin{aligned}
 \left. \frac{\partial \mu}{\partial T} \right|_N &= - \frac{\partial N}{\partial T} \bigg|_{\mu} = - \frac{\alpha N}{T_C} \left[\frac{C_\alpha (\alpha-1) \Gamma(\alpha-1) \zeta(\alpha-1)}{\beta_C^{\alpha-1}} \right]^{-1} \\
 &= - \frac{\alpha N}{T_C} \frac{\beta_C^{\alpha-1}}{C_\alpha (\alpha-1) \Gamma(\alpha-1) \zeta(\alpha-1)} \\
 &= - \frac{\alpha N k_B}{k_B T_C} \frac{\beta_C^{\alpha-1}}{C_\alpha (\alpha-1) \Gamma(\alpha-1) \zeta(\alpha-1)}
 \end{aligned}$$



So, that I have $\left. \frac{\partial \mu}{\partial T} \right|_N$ held constant is minus $\left. \frac{\partial N}{\partial T} \right|_{\mu}$ held constant divided by $\left. \frac{\partial N}{\partial \mu} \right|_T$ held constant. And this is minus αN over T_C and then we go back and look up the expression for $\left. \frac{\partial N}{\partial \mu} \right|_T$ evaluated at T is equal to T_C plus as this.

So, I have; so, this and then I have the expression $C_\alpha (\alpha-1) \Gamma(\alpha-1) \zeta(\alpha-1) \beta_C^{\alpha-1}$ divided by $\beta_C^{\alpha-1}$ whole minus 1. So, that this is minus αN over T_C and now, I want to see so, I have $C_\alpha (\alpha-1) \Gamma(\alpha-1) \zeta(\alpha-1) \beta_C^{\alpha-1}$ and then, I have $\beta_C^{\alpha-1}$ raised to the power $\alpha-1$.

So, let us bring in a k_B factor in the numerator and write down this as $k_B T C^\alpha$ raised to the power $\alpha - 1$ divided by $C^\alpha \Gamma(\alpha - 1) \zeta(\alpha - 1)$.

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$$\begin{aligned}
 &= -\frac{\alpha N k_B}{k_B T C} \frac{C^{\alpha-1}}{C^\alpha \Gamma(\alpha-1) \zeta(\alpha-1)} = -\frac{\alpha N k_B C^{\alpha-1}}{C^\alpha \Gamma(\alpha-1) \zeta(\alpha-1)} \\
 &= -\frac{\alpha N k_B C^{\alpha-1}}{C^\alpha} \frac{1}{\Gamma(\alpha-1) \zeta(\alpha-1)} \\
 &= -\frac{\alpha k_B}{C} \frac{\Gamma(\alpha) \zeta(\alpha)}{\Gamma(\alpha)} \frac{1}{\Gamma(\alpha-1) \zeta(\alpha-1)} \\
 \therefore &= -\frac{\alpha k_B}{C} \frac{\Gamma(\alpha) \zeta(\alpha)}{\Gamma(\alpha-1) \zeta(\alpha-1)}
 \end{aligned}$$



So, that this becomes minus $\alpha N k_B$ times $C^{\alpha-1}$ divided by $C^\alpha \Gamma(\alpha - 1) \zeta(\alpha - 1)$. Looks complicated, but if we are going to simplify this now, the $C^{\alpha-1}$ factor in the numerator cancels with the C^α in the denominator.

So, I have $N k_B$ times $C^{\alpha-1}$ divided by C^α and then, I have $\alpha - 1$ gamma of $\alpha - 1$ zeta of $\alpha - 1$. One over this, but go back and I know that N is $C^\alpha \Gamma(\alpha) \zeta(\alpha)$. So, let

us substitute. So, minus alpha C alpha gamma alpha divided times zeta of alpha beta C raised to the power alpha is N.

And then I have beta C raised to the power alpha C alpha 1 over alpha minus 1, gamma alpha minus 1, zeta of alpha minus one. So, things cancel out very nicely, C alpha C alpha cancels out, except that I have missed out the factor K B, which should have been sitting over here. So, I have alpha K B beta C beta C cancels out. So, that I have minus alpha K B over alpha minus 1 and then I have gamma of alpha divided by gamma alpha minus 1, I have zeta of alpha divided by zeta of alpha minus 1.

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$$\begin{aligned}
 & k_B T C \frac{C_\alpha^{(\alpha-1)} \Gamma(\alpha-1) \zeta(\alpha-1)}{C_\alpha^{(\alpha-1)} \Gamma(\alpha-1) \zeta(\alpha-1)} \\
 &= -\alpha \frac{N k_B \beta C^\alpha}{C_\alpha} \frac{1}{(\alpha-1) \Gamma(\alpha-1) \zeta(\alpha-1)} \\
 &= -\alpha k_B \frac{\beta C^\alpha \Gamma(\alpha) \zeta(\alpha)}{\beta C^\alpha} \frac{1}{(\alpha-1) \Gamma(\alpha-1) \zeta(\alpha-1)} \quad E(T, \mu) \\
 & \left(\frac{\partial \mu}{\partial T} \right)_N = -\alpha \frac{k_B}{(\alpha-1)} \frac{\Gamma(\alpha)}{\Gamma(\alpha-1)} \frac{\zeta(\alpha)}{\zeta(\alpha-1)} \\
 & \Delta C \\
 & dE = \frac{\partial E}{\partial T} dT + \frac{\partial E}{\partial \mu} d\mu \\
 & \frac{\partial E}{\partial T} = \frac{\partial E}{\partial T} + \frac{\partial E}{\partial \mu} \frac{\partial \mu}{\partial T} \\
 & \Delta C = \frac{\partial E}{\partial T} \frac{\partial \mu}{\partial T}
 \end{aligned}$$



So, this is going to be my del mu del T N held fixed. Now, recall the difference in specific heat, we wanted to evaluate and here we said that this is E; we started off with E as a function of T comma mu. So, that dE was del E del T mu held constant times dT plus del E del mu T

held constant times d mu and then we said look, I am interested in evaluating this derivative del E del T N held constant.

So, that that is going to be del E del T mu held constant plus del E del mu T held constant, del mu del T N held constant and this is the part, which I have evaluated. Now, I want to figure out well we already have figured out del E del T, I del E del mu T held constant and the difference in the specific heat is just this term is going to be del e del mu T held constant del mu del T N held constant.

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N

$$\left(\frac{\partial E}{\partial \mu}\right)_T = \beta z \left(\frac{\partial E}{\partial z}\right)_T = \beta z \frac{\alpha N}{z \beta} = \alpha N$$

$$\left(\frac{\partial E}{\partial \mu}\right)_T = \alpha N$$

$$\left(\frac{\partial \mu}{\partial T}\right)_N \left(\frac{\partial T}{\partial \mu}\right)_N \left(\frac{\partial N}{\partial \mu}\right)_T = -1$$

$$\rightarrow \left(\frac{\partial \mu}{\partial T}\right)_N = - \frac{\left(\frac{\partial N}{\partial T}\right)_\mu}{\left(\frac{\partial N}{\partial \mu}\right)_T}$$

So, we were looking for the term now, del mu del T and we have del E del mu is equal to alpha N.

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$$\left. \frac{\partial \mu}{\partial T} \right|_N = - \frac{\alpha k_B}{(\alpha-1)} \frac{\Gamma(\alpha)}{\Gamma(\alpha-1)} \frac{\zeta(\alpha)}{\zeta(\alpha-1)}$$

$$\Delta C = - \frac{\alpha^2 N k_B}{(\alpha-1)} \frac{(\alpha-1)!}{(\alpha-2)!} \frac{\zeta(\alpha)}{\zeta(\alpha-1)}$$

$$\left. \frac{\partial E}{\partial \mu} \right|_T = \alpha N$$

$$\Delta C = - \frac{\alpha^2 N k_B}{(\alpha-1)} \frac{(\alpha-1)!}{(\alpha-2)!} \frac{\zeta(\alpha)}{\zeta(\alpha-1)}$$

$$= - \alpha^2 N k_B \frac{\zeta(\alpha)}{\zeta(\alpha-1)}$$

$$\frac{(\alpha-1)!}{(\alpha-2)!} = (\alpha-1)$$

$$\frac{\partial E}{\partial T} = \frac{\partial E}{\partial T} + \frac{\partial E}{\partial \mu} \frac{\partial \mu}{\partial T}$$

$$\Delta C = \frac{\partial E}{\partial T} - \frac{\partial E}{\partial \mu} \frac{\partial \mu}{\partial T}$$



So, I have therefore, delta C as minus alpha square N K B divided by alpha minus 1. Gamma alpha is alpha minus 1 factorial and gamma alpha minus 1 is alpha minus 2 factorial, zeta alpha divided by zeta of alpha minus 1. Here we have used the result, that del e del mu T held constant plus alpha n.

So, that delta C becomes minus alpha square N K B; now, you immediately see that if you take the ratio alpha minus 1 factorial divided by alpha minus 2 factorial and this ratio is nothing but alpha minus 1. So, that this alpha minus 1 is going to cancel with the denominator and then I have zeta alpha divided by zeta of alpha minus 1. So, things looks very nice and I have minus alpha square N K B times zeta alpha divided by zeta alpha minus 1.

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$$= -\alpha^2 N k_B \frac{\zeta(\alpha)}{\zeta(\alpha-1)}$$

$$\frac{\Delta C}{N k_B} = -\alpha^2 \frac{\zeta(\alpha)}{\zeta(\alpha-1)}$$

$$\frac{\Delta C}{N k_B} = -\alpha^2 \frac{\zeta(\alpha)}{\zeta(\alpha-1)} \quad \alpha = 3$$

$$\frac{\Delta C}{N k_B} = -9 \frac{\zeta(3)}{\zeta(2)}$$

→ Harmonically trapped Bose Gas

$\alpha = \frac{3}{2}$



So, that the difference the discontinuity in specific heat is delta C, we write down in a non dimensional way is going to be minus alpha square zeta alpha divided by zeta of alpha minus 1. So, once I have this, looks is minus alpha square zeta of alpha zeta of; so, there is a something problem with this pen alpha minus 1.

Let us validate alpha is equal to 3 and then you realize, that delta C over N K B is going to be minus 9 zeta 3 over zeta 2. This is exactly the expression that we got, when we looked at a harmonically trapped Bose gas and this was the difference in the specific heat. For alpha is equal to 3 half there is no discontinuity.

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$$\frac{\Delta C}{Nk_B} = -9 \frac{J(3)}{J(2)} \rightarrow \text{Harmonically trapped Bose Gas}$$

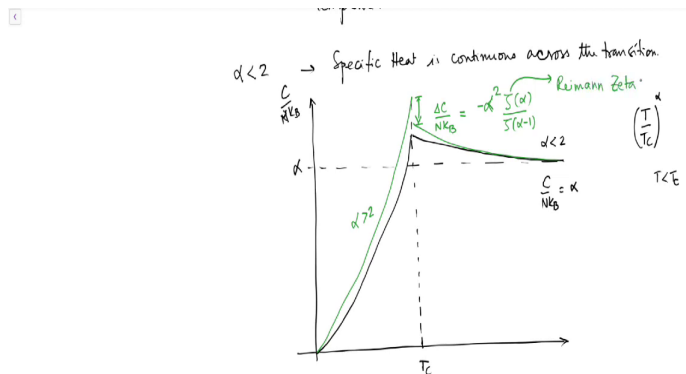
$d > 2 \rightarrow$ Specific Heat exhibits a discontinuity as a function of temperature. $d=3$

$d < 2 \rightarrow$ Specific Heat is continuous across the transition.



So, again let us reiterate that, that for alpha greater than 2, specific heat exhibits a discontinuity as a function of temperature, for alpha less than 2 specific heat is continuous across the transition.

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So, if you want to plot, C over $N K B$ and this is your transition temperature. So, let us draw a dotted line, that essentially marks this and let us draw the classical reiterates out like this, which is α , then we have the result that it must go as T over $T C$ raised to the power α ; for T less than $T C$ that is what we saw.

So, that it goes all the way up to here. For α less than 2, we also saw α less than 2 it is continuous not only that the specific heat at high temperatures, approaches the line the classical result C over $N K B$ is equal to α from top not from below. So, therefore, it is continuous and therefore, we can always write down like this way. In contrast; so, let us just to say that this is going to be α less than 2 and then I can do I can take the case, when α is it is a very poor schematics now.

So, let's it goes all the way up to here, for α greater than 2 again it goes behaves like this way, but there is a discontinuity and the discontinuity will somehow come like this way. Across the transition the function, the specific heat as a function of temperature exhibits a discontinuity and this is the amount of discontinuity that you get, which is ΔC over $N K_B$ and this expression is $-\alpha^2 \zeta(\alpha) / \zeta(\alpha - 1)$, where this is your Riemann zeta function.