

Statistical Mechanics
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Lecture - 63
Discontinuity in the Specific Heat of a Bose Gas - Part 01

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Discontinuity in specific heat of boson gas

$g(\epsilon) \sim \epsilon^{\alpha-1}$ $\alpha > 2$ there is a discontinuity in the specific heat.

Harmonically trapped boson gas $g(\epsilon) \sim \epsilon^2$ so that $\alpha = 3$.

U for $T < T_c$ as well as for $T > T_c$

At exactly $T = T_c$:

$\frac{U}{Nk_B} = 12 \frac{S(4)}{S(3)}$ at $T = T_c$ if we approach from below T_c



We want to discuss the Discontinuity in Specific Heat of a Boson Gas. We know that the general density of state goes as epsilon raised to the power alpha minus 1. For alpha greater than 2, there is a discontinuity in the specific heat and this we have seen. We looked at the case of a harmonically trapped gas Boson Gas. And here, we worked out the density of state to go and we saw that $g(\epsilon) \sim \epsilon^3$ behaves as ϵ^3 ; so that, we immediately identified alpha is equal to 3.

Consequently, if we calculated the specific heat both using by determining the internal at an energy U or E ; we will use this notion interchangeably for T less than T_c as well as for T greater than T_c . And we saw that the 2 results are not the same at exactly T is equal to T_c . U over Nk_B was sorry $12 \text{ times } \zeta(4) \text{ over } \zeta(3)$ at T is equal to T_c .

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Specific heat.

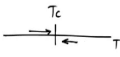
Harmonically trapped Boson Gas $g(\epsilon) \sim \epsilon^2$ so that $d=3$.

U for $T < T_c$ as well as for $T > T_c$

At exactly $T = T_c$.

$$\frac{C}{Nk_B} = 12 \frac{\zeta(4)}{\zeta(3)}$$


at $T = T_c$ if we approach from below T_c ($T < T_c$)



$$\frac{C}{Nk_B} = 12 \frac{\zeta(4)}{\zeta(3)} - 9 \frac{\zeta(3)}{\zeta(2)}$$

at $T = T_c$ if we approach T_c from above ($T > T_c$)

$$\frac{\Delta C}{Nk_B} = -9 \frac{\zeta(3)}{\zeta(2)}$$





If we approach from below T_c . In contrast U over Nk_B , was $12 \text{ over } \zeta(4) \text{ divided by } \zeta(3)$ minus $9 \text{ over } \zeta(3) \text{ divided by } \zeta(2)$ at T equal to T_c , if we approach T_c from above; which means from temperatures T greater than T_c and this means temperature below T less than T_c .

So, you have the critical temperature over here and you can approach the critical temperature from temperatures below T_c as well as from above. And you see that there is oops this is not

going to be U, this is going to be the specific heat if I call. And consequently, there is a difference in specific heat which is minus 9 zeta of 3 divided by zeta of 2 in this case.

Now, the question is we want to try to figure out what exactly this discontinuity this difference in amount is going to be for any general alpha.

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$$E(T, \mu)$$

$$dE = \left(\frac{\partial E}{\partial T}\right)_{\mu} dT + \left(\frac{\partial E}{\partial \mu}\right)_{T} d\mu$$

$$\boxed{\left(\frac{\partial E}{\partial T}\right)_{N} = \left(\frac{\partial E}{\partial T}\right)_{\mu} + \left(\frac{\partial E}{\partial \mu}\right)_{T} \left(\frac{\partial \mu}{\partial T}\right)_{N}}$$

$$\Delta C = \left(\frac{\partial E}{\partial \mu}\right)_{T} \left(\frac{\partial \mu}{\partial T}\right)_{N} \rightarrow \text{evaluated at } T = T_c.$$

$$E = \int_0^{\infty} d\epsilon g(\epsilon) \frac{\epsilon}{z^{-1} e^{\beta \epsilon} - 1} = C_{\alpha} \int_0^{\infty} d\epsilon \epsilon^{\alpha-1} \frac{\epsilon}{z^{-1} e^{\beta \epsilon} - 1}$$



For that, our starting point is to write down E as a function of T comma mu so that; the differential d E is going to be del E del T mu held constant d T plus del E del mu T held constant d mu. Consequently, I have del E del T; this is the specific heat that I am interested in is going to be del E with the N held constant del E del T mu held constant plus del E del mu T held constant del mu del N sorry, del mu del T N held constant.

Now, the first term is going to be the same whether you are approaching from below temperatures below T_c or whether you are approaching from temperatures above T_c . Therefore, the difference in specific heat is going to be $\frac{\partial E}{\partial \mu}$ temperature held constant $\frac{\partial \mu}{\partial T}$ held constant and this quantity evaluated at T is equal to T_c . Our purpose is to figure out what these things are going to be.

So, the energy we again start is integration $d\epsilon$ $g(\epsilon)$. So, here again it is, because starting from scratch.

So, that; things are not very complicated for you to figure out is going to be this. And which I have as C_α integration $d\epsilon$ ϵ^α from 0 to infinity $\frac{\epsilon^\alpha}{z^{-1}e^\epsilon - 1}$, then I have an ϵ^α divided by $z^{-1}e^\epsilon - 1$.

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$$E = \int_0^\infty d\epsilon g(\epsilon) \frac{\epsilon}{z^{-1}e^\epsilon - 1} = C_\alpha \int_0^\infty d\epsilon \epsilon^{\alpha-1} \frac{\epsilon}{z^{-1}e^\epsilon - 1}$$

$$= C_\alpha \int_0^\infty d\epsilon \epsilon^\alpha \frac{1}{z^{-1}e^\epsilon - 1}$$

$$\left. \frac{\partial E}{\partial \mu} \right|_T = \frac{\partial E}{\partial z} \frac{dz}{d\mu} = \beta \left. \frac{\partial E}{\partial z} \right|_T \quad \begin{matrix} z = e^{\beta\mu} \\ dz = e^{\beta\mu} \cdot \beta \\ d\mu = \frac{dz}{\beta z} \end{matrix}$$

$$\left. \frac{\partial E}{\partial z} \right|_T = C_\alpha \int_0^\infty d\epsilon \epsilon^\alpha \frac{d}{dz} \left(\frac{1}{z^{-1}e^\epsilon - 1} \right) \quad x = \beta\epsilon$$

$$\frac{d}{dz} \left(\frac{1}{z^{-1}e^\epsilon - 1} \right) = -\frac{1}{z} \frac{d}{dx} \left(\frac{1}{z^{-1}e^x - 1} \right)$$



Which I have $C \alpha 0 \text{ to infinity } d \epsilon$, ϵ to the power $\alpha 1 \text{ over } Z \text{ inverse } e$ to the power $X \text{ minus } 1$. So that, $\frac{dE}{d\mu}$ temperature held constant I note that is going to be $\frac{dE}{d\mu}$.

So, $\frac{dE}{dZ}$ of e times $\frac{dZ}{d\mu}$ and $\frac{dZ}{d\mu}$ since Z is equal to e to the power $\beta \mu$ $\frac{dZ}{d\mu}$ is going to be e to the power $\beta \mu$ times β . So, I have β times $\frac{dE}{dZ}$ temperature held constant and we go back to the expression that we have written down over here. So, that $\frac{dE}{dZ}$ temperature held constant is going to be $C \alpha 0 \text{ to infinity } d \epsilon$ to the power $\alpha d \frac{dZ}{d\mu}$ of $1 \text{ over } Z \text{ inverse } e$ to the power $X \text{ minus } 1$.

Now, life is a little bit simpler, because I have seen that $\frac{d}{dZ}$ of $1 \text{ over } Z \text{ inverse } e$ to the power $X \text{ minus } 1$ is just going to be $\text{minus } 1 \text{ over } Z$ $\frac{d}{dX}$ of $1 \text{ over } Z \text{ inverse } e$ to the power $X \text{ minus } 1$.

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$$\begin{aligned} \left(\frac{\partial E}{\partial z}\right)_T &= C_\alpha \int_0^\infty d\epsilon \epsilon^\alpha \frac{d}{dz} \left(\frac{1}{z^{-1} e^{\beta\epsilon} - 1} \right) \quad x = \beta\epsilon \\ &= C_\alpha \int_0^\infty d(\beta\epsilon) \frac{(\beta\epsilon)^\alpha}{\beta^\alpha} \frac{d}{dz} \left(\frac{1}{z^{-1} e^{\beta\epsilon} - 1} \right) \\ &= \frac{C_\alpha}{\beta^{\alpha+1}} \int_0^\infty dx x^\alpha \frac{d}{dz} \left(\frac{1}{z^{-1} e^x - 1} \right) \end{aligned}$$



So that, this relation then, I have is going to be C alpha integration 0 to infinity. Now, look, I have used e to the power X and I have used e to the power X which is not the right thing to do. So, we have beta epsilon over here, did we keep it as beta epsilon yeah. So, unnecessarily, I have complete.

Anyway, so since this is beta X is equal to beta epsilon I can straightforward substitute this as d beta epsilon 1 over beta, beta epsilon raised to the power alpha beta to the power alpha d Z of 1 over Z inverse e to the power beta epsilon minus 1 and this becomes C alpha raised to the power beta alpha plus 1 0 to infinity dX X to the power alpha d Z of 1 over Z inverse e to the power X minus 1.

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$$\begin{aligned}
 &= \frac{C_\alpha}{\beta^{\alpha+1}} \int_0^\infty dx \, x^\alpha \frac{d}{dz} \left(\frac{1}{z^{-1}e^x - 1} \right) \\
 &= \frac{C_\alpha}{\beta^{\alpha+1}} \int_0^\infty dx \, x^\alpha \left(\frac{-1}{z} \right) \frac{d}{dx} \left(\frac{1}{z^{-1}e^x - 1} \right) \\
 &= -\frac{1}{z} \frac{C_\alpha}{\beta^{\alpha+1}} \int_0^\infty dx \, x^\alpha \frac{d}{dx} \left(\frac{1}{z^{-1}e^x - 1} \right) \\
 &= \left[\frac{1}{z^{-1}e^x - 1} x^\alpha \right]_0^\infty - \int_0^\infty dx \, \frac{\alpha x^{\alpha-1}}{(z^{-1}e^x - 1)}
 \end{aligned}$$



So, that I have C alpha beta raised to the power alpha plus 1 0 to infinity dX X to the power alpha and d dz is minus 1 over Z d dX of 1 over Z inverse e to the power X minus 1. So, this becomes minus 1 over Z C alpha beta alpha plus 1 0 to infinity dX X to the power alpha d dX of 1 over z inverse e to the power X minus 1.

Now, we are familiar with such integrals. We have done it several times in the past and the trick is to integrate by parts taking this as the first function. If I take this as the first function, then it follows that I have 1 over Z inverse e to the power X minus 1 X to the power alpha 0 to infinity minus 0 to infinity dx; a derivative of this is alpha X to the power minus X to the power alpha minus 1 and then, I have Z inverse e to the power X minus 1.

This vanishes in both the limits when X is equal to 0 you have it identically 0. When X is equal to infinity the exponential blows up and you have this also vanishes.

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$$\begin{aligned}
 &= \ominus \frac{1}{z} \frac{C_\alpha}{\beta^{\alpha+1}} \left[\frac{1}{z^\alpha e^X - 1} \Big|_0^\infty - \int_0^\infty dx \frac{X^\alpha}{(z^\alpha e^X - 1)} \right] \\
 &= + \frac{\alpha}{z} \frac{C_\alpha}{\beta^{\alpha+1}} \int_0^\infty dx \frac{X^{\alpha-1}}{z^\alpha e^X - 1} \quad N = \int_0^\infty dx \frac{1}{z^\alpha e^X - 1} \\
 &= \frac{\alpha}{z \beta} \boxed{\frac{C_\alpha}{\beta^\alpha} \int_0^\infty dx \frac{X^{\alpha-1}}{z^\alpha e^X - 1}} \quad N = \frac{C_\alpha}{\beta^\alpha} \int_0^\infty dx \frac{X^{\alpha-1}}{z^\alpha e^X - 1} \\
 &\qquad \qquad \qquad N
 \end{aligned}$$

$$\left(\frac{\partial F}{\partial \mu} \right)_T = \beta \left(\frac{\partial F}{\partial z} \right)_T = \beta \frac{\alpha}{z \beta} N$$



So, in the 2 limits the first term vanishes. So, that the alpha I can bring it in front and I have minus alpha over $Z C_\alpha \beta^{\alpha+1}$. Also, note that I have a minus over here and here, I have minus 1 by $C_\alpha \beta^{\alpha+1}$. So that, this minus and this minus gives me a plus I have 0 to infinity $dX X^{\alpha-1} Z^{-\alpha} e^{-X}$. Let us rewrite this expression as this. Alpha divided by Z times beta.

So, that; I have $C_\alpha \beta^{\alpha+1} \int_0^\infty dx X^{\alpha-1} Z^{-\alpha} e^{-X}$. And if you recall, then this is exactly the particle number N . So, particle number N was integration $d\epsilon \epsilon^{-\alpha} e^{-\epsilon}$.

epsilon minus 1 and if you substitute for epsilon to the power alpha minus 1, you are going to get C alpha divided by beta alpha dX X to the power alpha minus 1 Z inverse e to the power X minus 1 0 to infinity.

So, that del E del mu temperature held constant is beta times del E del Z temperature held constant which is alpha over sorry beta times alpha over Z beta times N. So, this looks slightly suspicious, because in the denominator I have a Z. So, we go back and see where we have done the mistake and it is over here.

We did the derivative correctly. So, that this is going to be beta times Z we have to include a beta times Z here. So, that this expression is beta times Z and this is going to be alpha times N.

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$$\epsilon_T \left[\frac{1}{N} \right]$$

$$\left(\frac{\partial E}{\partial \mu} \right)_T = \beta Z \left(\frac{\partial E}{\partial Z} \right)_T = \beta Z \frac{\alpha}{Z \beta} N = \alpha N$$

$$\boxed{\left(\frac{\partial E}{\partial \mu} \right)_T = \alpha N}$$

$$\left(\frac{\partial \mu}{\partial T} \right)_N \left(\frac{\partial N}{\partial T} \right)_\mu \frac{\partial T}{\partial T}$$



So, that $\partial E / \partial \mu$ comes out to be a very very simple and elegant expression $\alpha \mu$. Now, in this expression I have calculated the first term. So, this is done I have to look at this quantity $\partial \mu / \partial T$ N held constant. So, $\partial \mu / \partial T$ N held constant; to evaluate this I am going to use the cyclic identity that we learned in thermodynamics, that is essentially $\partial N / \partial T$ μ held constant and then, I have ∂T .

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$$\left(\frac{\partial \mu}{\partial T}\right)_N \left(\frac{\partial T}{\partial \mu}\right)_\mu \left(\frac{\partial \mu}{\partial T}\right)_T = -1$$

$$\left(\frac{\partial \mu}{\partial T}\right)_N = - \frac{\left(\frac{\partial N}{\partial T}\right)_\mu}{\left(\frac{\partial N}{\partial \mu}\right)_T}$$

$$N = \frac{C_\alpha}{\beta^\alpha} \int_0^\infty dx \frac{x^{\alpha-1}}{z^{-1} e^x - 1}$$

$$\left(\frac{\partial N}{\partial \mu}\right)_T = \left(\frac{\partial N}{\partial z}\right)_T \left(\frac{dz}{d\mu}\right) = \beta z \left(\frac{\partial N}{\partial z}\right)_T$$



So, no this is not right this is going to be $\partial T / \partial N$ μ held constant $\partial N / \partial \mu$ T held constant is going to be minus 1. Therefore, I have $\partial \mu / \partial T$ N held constant is going to be minus $\partial N / \partial T$ μ held constant divided by $\partial N / \partial \mu$ T held constant. So, these are the two partial derivatives which I have to evaluate, but I know the expression for N . N was C_α raised to the power $\beta \alpha$ 0 to infinity dX X to the power α minus 1 Z inverse e to the power X minus 1 $\partial N / \partial T$.

So, the easier one is del N del mu T held constant and that is going to be again del N del Z T held constant dZ d mu which is going to be tell me sorry, dZ d mu is again beta times Z del N del Z T held constant.

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$$N = \frac{C_\alpha}{\beta^\alpha} \int_0^\infty dx \frac{x^{\alpha-1}}{z^{-1}e^x - 1}$$

$$\left(\frac{\partial N}{\partial \mu}\right)_T = \left(\frac{\partial N}{\partial z}\right)_T \left(\frac{dz}{d\mu}\right) = \beta z \left(\frac{\partial N}{\partial z}\right)_T \quad \frac{d}{dz} \frac{1}{z^{-1}e^x - 1} = -\frac{1}{z} \frac{d}{dx} \frac{1}{z^{-1}e^x - 1}$$

$$\left(\frac{\partial N}{\partial z}\right)_T = \frac{C_\alpha}{\beta^\alpha} \int_0^\infty dx x^{\alpha-1} \frac{d}{dz} \left(\frac{1}{z^{-1}e^x - 1} \right)$$

$$= \frac{C_\alpha}{\beta^\alpha} \int_0^\infty dx x^{\alpha-1} \left(-\frac{1}{z}\right) \frac{d}{dx} \frac{1}{z^{-1}e^x - 1}$$



So, let us look take this expression for N and open figure out the derivative del N del Z T held constant. This is going to be C alpha beta raised to the power alpha 0 to infinity d of x, X to the power alpha minus 1 d dZ of 1 over Z inverse e to the power X minus 1. Again, I know the standard trick how to do this. I know that d dZ of 1 over Z inverse e to the power X minus 1 is minus 1 by Z d dX of 1 over Z inverse e to the power X minus 1.

So, that I have C alpha divided by beta alpha this is 0 to infinity d of X x to the power alpha minus 1 minus 1 over Z d dX of 1 over Z inverse e to the power X minus 1.

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$$\begin{aligned}
 &= \frac{C\alpha}{\beta^\alpha Z} \int_0^\infty dx \ x \left(-\frac{1}{Z}\right) \frac{d}{dx} \frac{1}{Z^{-1}e^x - 1} \\
 &= -\frac{C\alpha}{\beta^\alpha Z} \int_0^\infty dx \ x^{\alpha-1} \frac{d}{dx} \frac{1}{Z^{-1}e^x - 1} \\
 &= \left(-\frac{1}{Z}\right) \frac{C\alpha}{\beta^\alpha} \left[\int_0^\infty \frac{1}{Z^{-1}e^x - 1} dx - \int_0^\infty \frac{(x-1)}{Z^{-1}e^x - 1} dx \right] \\
 &= \frac{1}{Z} \frac{C\alpha}{\beta^\alpha} \int_0^\infty dx \ x^{\alpha-2} \frac{1}{Z^{-1}e^x - 1}
 \end{aligned}$$



The minus 1 by Z I can bring out. C alpha beta raise to the power alpha divided by time Z 0 to infinity dx x to the power alpha minus 1 d dx of 1 over Z inverse e to the power x minus 1. Now, this integral again I know how to handle and it we do just integration by parts.

So, minus we will just separate this out as minus 1 over Z C alpha beta alpha and then, this becomes my first function; so that, you have Z inverse e to the power x minus 1 x to the power alpha minus 1 0 to infinity minus integration 0 to infinity dX; the derivative of the second function which is going to be alpha minus 1 X to the power alpha minus 2 and the integral of the first function which is just going to be Z inverse e to the power X minus 1 looks very nice and note that this vanishes in both the limits.

So that, I have the minus and this minus makes it a plus. So, I have 1 over Z C alpha beta alpha 0 to infinity; there is a alpha minus 1 dX X to the power alpha minus 2 Z inverse e to

the power X minus 1. Let us go back and see I originally want $\frac{\partial N}{\partial \mu}$ at T held constant and that is just β times Z .

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$$\begin{aligned}
 &= \frac{1}{Z} \frac{C_\alpha (\alpha-1)}{\beta^\alpha} \int_0^\infty dx \frac{x^{\alpha-2}}{z^{-1} e^x - 1} \\
 \left. \frac{\partial N}{\partial \mu} \right|_T &= \beta Z \left. \frac{\partial N}{\partial z} \right|_T = \beta Z \cdot \frac{1}{Z} \frac{C_\alpha (\alpha-1)}{\beta^\alpha} \int_0^\infty dx \frac{x^{\alpha-2}}{z^{-1} e^x - 1} \\
 &= \frac{C_\alpha (\alpha-1)}{\beta^{\alpha-1}} \int_0^\infty dx \frac{x^{\alpha-2}}{z^{-1} e^x - 1} \\
 \left. \frac{\partial N}{\partial \mu} \right|_{T=1} &= \frac{C_\alpha (\alpha-1)}{\beta^{\alpha-1}} \int_0^\infty dx \frac{x^{\alpha-2}}{e^x - 1} \quad Z=1
 \end{aligned}$$



So that; $\frac{\partial N}{\partial \mu}$ at T held constant is βZ times $\frac{\partial N}{\partial Z}$ at T held constant which is going to be βZ times $\frac{1}{Z} \frac{C_\alpha (\alpha-1)}{\beta^\alpha}$. I have $\alpha - 1$ and then, I have 0 to ∞ dx x to the power $\alpha - 2$ divided by $Z^{-1} e^x - 1$ the not the b .

So, the cancellation means that Z cancels out and β cancels with one factor of β raised to the power α . So, that this becomes $\beta^{\alpha-1}$ and the answer that I have is $C_\alpha (\alpha-1)$ divided by $\beta^{\alpha-1}$ raised to the power $\alpha - 1$ 0 to ∞ dx x to the power $\alpha - 2$ $Z^{-1} e^x - 1$. Now, I have to

evaluate this derivative. If you recall we originally started off with this derivative and this is the derivative that I have to evaluate at T c.

So, which means; I can replace Z is equal to 1. Since, I am evaluating this Z very close to T c. So, that this becomes del N del mu T is equal to T c is C alpha alpha minus 1. Well, you can also have T c plus 1 T c plus. So, that you are approaching from the high temperature side beta C raised to the power alpha minus 1. Here, Z is equal to 1.

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$$\begin{aligned} \left. \frac{\partial N}{\partial \mu} \right|_{T=T_c^+} &= \frac{C_N (\alpha-1)}{\beta^{\alpha-1}} \int_0^{\infty} dx \frac{x^{\alpha-2}}{e^x - 1} \\ &= \frac{C_N (\alpha-1)}{\beta^{\alpha-1}} \int_0^{\infty} dx \frac{x^{\alpha-2}}{e^x - 1} \quad Z=1 \\ &= \frac{C_N (\alpha-1)}{\beta^{\alpha-1}} \Gamma(\alpha-1) \end{aligned}$$



And you have integral 0 to infinity d of X x to the power alpha minus 2 divided by e to the power X minus 1 which is going to be C alpha alpha minus 1 beta C alpha minus 1 and this you see is gamma of alpha minus 1 this is the gamma function. So, I have gamma of alpha minus 1 is del N del mu evaluated at T equal to T c plus.

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$$\begin{aligned}
 \left. \frac{\partial N}{\partial \mu} \right|_T &= \beta z \left. \frac{\partial N}{\partial z} \right|_T = \beta z \cdot \frac{1}{z} \cdot \frac{C_\alpha (\alpha-1)}{\beta^{\alpha-1}} \int_0^\infty dx \frac{x^{\alpha-2}}{z^{-1} e^x - 1} \\
 &= \frac{C_\alpha (\alpha-1)}{\beta^{\alpha-1}} \int_0^\infty dx \frac{x^{\alpha-2}}{z^{-1} e^x - 1} \quad \int_0^\infty f_m(z) \\
 &= \frac{C_\alpha (\alpha-1)}{\beta^{\alpha-1}} \frac{1}{(\alpha-2)!} \int_0^\infty dx \frac{x^{\alpha-2}}{z^{-1} e^x - 1} \\
 \left. \frac{\partial N}{\partial \mu} \right|_T &= \frac{C_\alpha (\alpha-1)}{\beta^{\alpha-1}} \Gamma(\alpha-1) \int_0^\infty f_m(z) \quad (\alpha-2)! = \Gamma(\alpha-1)
 \end{aligned}$$



So, this expression now, I can write down as C alpha alpha minus 1 beta raised to the power alpha minus 1. And I see that this I can write down in terms of the integrals f of m eta well f of m plus of Z by identifying the fact that I have a missing 1 over alpha minus 2 factorial over here.

So, for that I introduce 1 over alpha minus 2 factorial in the numerator and alpha minus 2 factorial in the in sorry 1 minus alpha minus 2 1 by alpha minus 2 factorial in the denominator and alpha minus 2 factorial in the numerator and I write down this as dX X to the power alpha minus 2 divided by Z inverse e to the power X minus 1.

This quantity alpha minus 2 factorial I know is gamma of alpha minus 1. So, that I have C alpha alpha minus 1 divided by beta alpha minus 1 and then I have gamma of alpha minus 1 f of alpha minus 1 plus of Z. Now, I need to evaluate this. So, this is the derivative. Let us write

down, del N del mu T held constant. Now, this is the derivative that I am going to use, but this derivative I am going to evaluate very close to T c.

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$$\begin{aligned}
 &= \frac{C_\alpha (\alpha-1)}{\beta^{\alpha-1}} \frac{1}{(\alpha-1)!} \int_0^{\alpha-1} dx \frac{x^{\alpha-2}}{z^{\alpha-1} e^z - 1} \\
 \left. \frac{\partial N}{\partial \mu} \right|_T &= \frac{C_\alpha (\alpha-1) \Gamma(\alpha-1)}{\beta^{\alpha-1}} \underbrace{\int_{d-1}^{\alpha} f_{d-1}(z)}_{\Gamma(\alpha-1)} \quad (\alpha-1)! = \Gamma(\alpha-1) \\
 \left. \frac{\partial N}{\partial \mu} \right|_{T=T_c} &= \frac{C_\alpha (\alpha-1) \Gamma(\alpha-1)}{\beta_c^{\alpha-1}} \underbrace{\int_{d-1}^{\alpha} f_{d-1}(z)}_{\Gamma(\alpha-1)} \quad \beta^{\alpha-1} \rightarrow \beta_c^{\alpha-1} \\
 &= \frac{C_\alpha (\alpha-1) \Gamma(\alpha-1) \zeta(\alpha-1)}{\beta_c^{\alpha-1}} \quad \alpha-1 > 1
 \end{aligned}$$



So, I have del N del mu. What I require is del N del mu T is equal to T c plus extremely close to T c. For that; since I want this. I look at this expression on top and I immediately say that oh I am going to put Z equal to 1 over here. So, that this becomes C alpha alpha minus 1 gamma of alpha minus 1 raised to the power beta C alpha minus 1.

So, the 2 substitutions I made. I made Z equal to 1 that is a not the 2 substitutions I put 2 values Z equal to 1 and beta to the power alpha minus 1 goes to beta C raised to the power alpha minus 1 and this becomes alpha minus 1 of 1. So, this function is now, C alpha alpha minus 1 gamma of alpha minus 1 divided by beta C of alpha minus 1 times zeta of alpha minus 1. Of course, we are considering alpha great alpha minus 1 is greater than 1 right.

So, once I have this expression now; all I am left to evaluate now is $\frac{\partial N}{\partial T}$ μ held constant.

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$$\begin{aligned} \frac{\partial N}{\partial T} \Big|_{\mu} &= ? \\ N &= \frac{C_{\alpha}}{\beta^{\alpha}} \int_0^{\infty} dx \frac{x^{\alpha-1}}{e^x - 1} \\ &= \frac{C_{\alpha}}{\beta^{\alpha}} \Gamma(\alpha) f_{\alpha-1}^+(z) \\ N &= \frac{C_{\alpha}}{\beta^{\alpha}} \Gamma(\alpha) \zeta(\alpha) \\ \frac{\partial N}{\partial T} \Big|_{\mu} &= C_{\alpha} \Gamma(\alpha) \zeta(\alpha) \frac{\partial}{\partial T} \left(\frac{1}{\beta^{\alpha}} \right) \end{aligned}$$



So, I have to evaluate $\frac{\partial N}{\partial T}$ μ held constant. Now, N is again given by $C_{\alpha} \beta^{\alpha}$ raised to the power α 0 to infinity dx x to the power $\alpha - 1$ divided by Z inverse e to the power $x - 1$. And this derivative again, I am going to use very very close to T c . So, that this becomes $C_{\alpha} \beta^{\alpha}$ raised to the power sorry β raised to the power α . I am going to bring in a factor of $\Gamma(\alpha)$ and this equation becomes f of $\alpha - 1$ plus of set.

And close to T_c ; I have N as $C \alpha \beta^\alpha \Gamma(\alpha) \zeta(\alpha)$ of α very close to T_c ; so that, $\frac{\partial N}{\partial T} \mu$ held constant becomes $C \alpha \Gamma(\alpha) \zeta(\alpha)$ of α over β^α .

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$$\begin{aligned}
 & \beta^{-\alpha} \int_0^{\infty} z^{\alpha-1} e^{-z} dz \\
 &= \frac{C \alpha}{\beta^\alpha} \Gamma(\alpha) \zeta(\alpha) \\
 N &= \frac{C \alpha}{\beta^\alpha} \Gamma(\alpha) \zeta(\alpha-1) \\
 \left. \frac{\partial N}{\partial T} \right|_{\mu} &= C \alpha \Gamma(\alpha) \zeta(\alpha-1) \frac{\partial}{\partial T} \left(\frac{1}{\beta^\alpha} \right) \\
 &= C \alpha \Gamma(\alpha) \zeta(\alpha-1)
 \end{aligned}$$



This I rewrite as $C \alpha \Gamma(\alpha) \zeta(\alpha)$. I think I have made a mistake over here. So, this is going to be $\alpha \zeta(\alpha-1)$. This is also going to be $\zeta(\alpha-1)$ and this is going to be $\zeta(\alpha-1)$. I hope in the earlier one I have written it correctly. This is $\zeta(\alpha-1)$ yeah. So, this is correct. So, this was originally correct this has to be $\zeta(\alpha)$.

(Refer Slide Time: 28:30)

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$$\frac{\partial N}{\partial T} = ?$$

$$N = \frac{C_\alpha}{\beta^\alpha} \int_0^\infty dx \frac{x^{\alpha-1}}{e^x - 1}$$

$$= \frac{C_\alpha}{\beta^\alpha} \Gamma(\alpha) f_{\alpha}^+(z)$$

$$N = \frac{C_\alpha}{\beta^\alpha} \Gamma(\alpha) \zeta(\alpha)$$

$$\frac{\partial N}{\partial T} = C_\alpha \Gamma(\alpha) \zeta(\alpha-1) \frac{\partial}{\partial T} \left(\frac{1}{\beta^\alpha} \right)$$



And I had X to the power alpha minus 1. So, this has to be f of alpha.