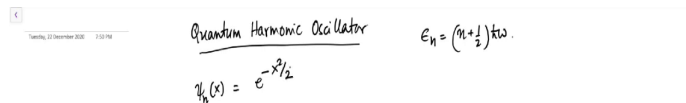


Statistical Mechanics
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Lecture - 44
Single Particle Quantum Partition Function Harmonic Oscillator – Part 1

Welcome back. So, we are interested in we looked at the density matrix for a free particle both in the momentum as well as in the coordinate presentation.

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Quantum Harmonic Oscillator

$$E_n = \left(n + \frac{1}{2}\right) \hbar \omega$$
$$\psi_n(x) = e^{-x^2/2}$$



Now, we are interested in looking at the harmonic oscillator, single particle harmonic oscillator. So, I will write a Quantum Harmonic Oscillator. As you know that in this particular case the energy eigenvalues are discrete. So, I have n plus half $\hbar \omega$. And the energy eigen functions are given by e to the power minus x square over 2.

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HARMONIC OSCILLATOR

$$H = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 \hat{q}^2 \quad \psi = E\psi \quad \hat{p} = -i\hbar \frac{d}{dx}$$

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + \frac{1}{2}m\omega^2 \hat{q}^2 \psi = E\psi \quad \psi_n(q)$$

$$\psi_n(q) = \left(\frac{m\omega}{\hbar\pi}\right)^{1/4} e^{-x^2/2} H_n(x) \quad \text{Hermite Polynomial.}$$

$$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} e^{-x^2}$$

$$= \frac{1}{\sqrt{\pi}} e^{x^2} \int_{-\infty}^{\infty} du (-2iu)^n e^{-u^2 + 2iux}$$



So, we will before we write this the Hamiltonian of the system is half p i p square over twice m plus half m omega square q square. This gives me discrete energy eigenvalue. Once I solve the equation as E of psi the operator p hat over here is ih bar d dx. So, that if I substitute this Hamiltonian, this becomes minus h square over twice m d 2 psi d x 2 plus half m omega square q square dq 2 because the coordinate in our case is q is going to be E times psi.

And we solve this equation to get the energy eigen functions psi n of q. And psi n of q is a little complicated is m omega h bar raised to the power one-fourth e to the power minus x square over 2, and then you have H n of x. This function is known as the Hermite polynomial, where H n of x is minus 1 raised to the power n e to the power x square d n dx n of e to the power minus x square.

So, once you operate the another derivative on this, finally, you are going to get a polynomial. Interestingly, this also has an integral representation and that is what we want to write down, we want to write down this as the integral presentation is square root pi e to the power x square minus infinity to plus infinity d u minus twice i u raised to the power n e to the power minus u square plus twice i v times x, sorry twice i u times x, because I am integrating over u there should not be any other variable.

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$$\begin{aligned}
 & \left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m \omega^2 x^2 \right) \psi = E \psi \quad \psi_n(x) \\
 & \psi_n(x) = \left(\frac{m\omega}{\hbar\pi} \right)^{1/4} e^{-x^2/2} H_n(x) \quad \text{Hermite Polynomial.} \quad x = \sqrt{\frac{\hbar m \omega}{\hbar}} x \\
 & \psi_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} e^{-x^2} \\
 & = \frac{1}{\sqrt{\pi}} e^{x^2} \int_{-\infty}^{\infty} du (-2iu)^n e^{-u^2 + 2iux} \\
 & \psi_n(x) = \left(\frac{m\omega}{\hbar\pi} \right)^{1/4} \frac{e^{-x^2/2}}{\sqrt{\pi}} \int_{-\infty}^{\infty} du (-2iu)^n e^{-u^2 + 2iux}
 \end{aligned}$$



So, that psi by the way x here is square root m omega over h bar times q. So, psi n of q then becomes m omega h bar pi raised to the power one-fourth e to the power minus x square by 2 e to the power x square, square root of pi minus infinity to plus infinity du twice i u raised to the power n e to the power minus u square plus twice i u x.

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$$\psi_n(x) = \left(\frac{m\omega}{\hbar\pi}\right)^{1/4} \frac{e^{-x^2/2}}{\sqrt{\pi}} \int_{-\infty}^{\infty} du (2iu)^n e^{-u^2}$$

$$\rho_{mn} = \rho_n \delta_{m,n} \quad \rho_n = \frac{1}{Z} e^{-\beta \epsilon_n} = \frac{1}{Z} e^{-\beta \hbar \omega (n+1/2)}$$

$$\text{Tr}(\hat{\rho}) = 1 \Rightarrow \frac{1}{Z} \sum_n e^{-\beta \hbar \omega (n+1/2)} = 1$$

$$\frac{1}{Z} e^{-\beta \hbar \omega / 2} \left[1 + e^{-\beta \hbar \omega} + e^{-2\beta \hbar \omega} + \dots \right] = 1$$

$$\frac{1}{1 - e^{-\beta \hbar \omega}}$$



And this gives me $m \omega \hbar \pi$ raised to the power one-fourth e to the power x square by $2 \sqrt{\pi}$ minus infinity to plus infinity du twice minus of twice iu raised to the power n minus e to the power minus u square plus twice iu . We will later on use this relation.

Now, if you work in the energy basis, then of course your density matrix is very very simple ρ_{mn} is $\rho_n \delta_{m,n}$, where ρ_n is going to be e to the power minus beta of ϵ_n 1 over Z , and you have 1 over Z e to the power minus beta $\hbar \omega n$ plus half. If I now use that the trace of the density matrix must be 1.

This implies sum over n 1 by Z e to the power minus beta $\hbar \omega n$ plus half is going to be 1, so that you have minus beta $\hbar \omega$ 1 over Z . And if you open the sum, then you will see you are going to have a series and infinite series minus beta $\hbar \omega$ 1 plus e to the

power minus beta h bar omega twice beta h bar omega higher order terms is going to be 1. This is simply a geometric series whose sum is minus beta h bar omega.

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$$Z(T, V, 1) = \frac{e^{-\beta \hbar \omega / 2}}{1 - e^{-\beta \hbar \omega}} = \frac{1}{e^{\beta \hbar \omega / 2} [1 - e^{-\beta \hbar \omega}]} = \frac{1}{e^{\beta \hbar \omega / 2} - e^{-\beta \hbar \omega / 2}}$$

$$\boxed{Z = \frac{1}{2 \sinh(\beta \hbar \omega / 2)}}$$

$$\langle q^2 \rangle = \sum_{n, n'} \langle q | n' \rangle \langle n' | \hat{q}^2 | n \rangle \langle n | q \rangle$$

$$\langle q | n' \rangle = e^{-\beta \epsilon_n} \delta_{n, n'} \langle n | q \rangle$$



So, that the partition function $Z(T, V, 1)$ for a single harmonic oscillator is going to be minus beta h bar omega by 2 divided by 1 minus beta h bar omega. You can write down this as e to the power and beta h bar omega by 2 $1 - e$ to the power beta h bar omega which gives 1 over e to the power beta h bar omega by 2 minus beta h bar omega by 2.

And this gives you we can immediately see that this as the form e to the power x minus x which is nothing but the sin hyperbolic function. So, you have an answer twice of sin hyperbolic beta h bar omega over 2. So, this is your partition function – micro-canonical partition function.

Now, the idea is to construct the density matrix in the coordinate representation. And this is where it becomes a little bit more involved, but it is very straightforward mathematics that is nothing extremely difficult which you cannot do. So, we write down this as ρ , n prime q . This quantity I know this is the density matrix in the energy presentation. And therefore, this becomes $\frac{1}{Z} e^{-\beta \hat{H}}$ n by Z n comma n prime and q q prime n prime.

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$$\begin{aligned}
 \rho &= \sum_{n, n'} \langle q | n' \rangle \langle n' | \hat{\rho} | n \rangle \langle n | q \rangle \\
 &= \sum_{n, n'} \langle q | n' \rangle \frac{e^{-\beta E_{n'}}}{Z} \delta_{n, n'} \langle n | q \rangle \\
 &= \sum_n \langle q | n \rangle \frac{e^{-\beta E_n}}{Z} \langle n | q \rangle \\
 &= \sum_n \psi_n^*(q) \frac{e^{-\beta E_n}}{Z} \psi_n(q) \\
 &= \sum_n \left(\frac{m\omega}{\hbar\pi} \right)^{1/4} \frac{1}{\sqrt{2^n n!}} \frac{e^{x'^2/2}}{\sqrt{\pi}} \int_{-\infty}^{\infty} du (-2iu)^n e^{-u^2 + 2iux'} \frac{e^{-\beta(\frac{m\omega}{2} + \frac{1}{2})\hbar u}}{Z} \left(\frac{m\omega}{\hbar\pi} \right)^{1/4} \frac{1}{\sqrt{2^n n!}} \frac{e^{x'^2/2}}{\sqrt{\pi}} \int_{-\infty}^{\infty} dv (-2iv)^n e^{-v^2 + 2ivx}
 \end{aligned}$$



Now, clearly there is sum over n and n prime and I can carry forward I can take care of one of the sums over n prime to take care of this delta function, so that this becomes sum over n q prime n e to the power minus beta epsa n Z n of q . Let me see whether here I have missed out a factor that one should include. Where is that?

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$$\begin{aligned}
 \psi_n(x) &= (-1)^n e^{x^2} \frac{d^n}{dx^n} e^{-x^2} \\
 &= \frac{1}{\sqrt{\pi}} \frac{e^{x^2}}{\sqrt{2^n n!}} \int_{-\infty}^{\infty} du (-2iu)^n e^{-u^2 + 2iux} \\
 \psi_n(x) &= \left(\frac{m\omega}{\hbar\pi}\right)^{1/4} \frac{e^{-x^2/2}}{\sqrt{\pi} \sqrt{2^n n!}} \int_{-\infty}^{\infty} du (-2iu)^n e^{-u^2 + 2iux} \\
 \psi_n(x) &= \left(\frac{m\omega}{\hbar\pi}\right)^{1/4} \frac{e^{x^2/2}}{\sqrt{\pi} \sqrt{2^n n!}} \int_{-\infty}^{\infty} du (2iu)^n e^{-u^2 + 2iux} \\
 \int_{m,n} &= \int_n \delta_{m,n} \quad \int_n = \frac{1}{2} e^{-\beta} \Gamma_n = \frac{1}{2} e^{-\beta} \Gamma(n + \frac{1}{2})
 \end{aligned}$$



Yeah there has to be a 2 to the power n and n factorial with square root, there has to be a square 2 two to the power n, n factorial, square root 2 to the power n, n factorial this is the thing which I have missed not in h n x, but in psi n x not in h n x.

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Quantum Harmonic Oscillator

$$H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 q^2 \rightarrow E_n = (n + \frac{1}{2})\hbar\omega$$

$$H\psi = E\psi$$

$$\left(-\frac{\hbar^2}{2m} \frac{d^2}{dq^2} + \frac{1}{2}m\omega^2 q^2 \right) \psi = E\psi$$



$$\psi_n(q) = \left(\frac{m\omega}{\hbar\pi} \right)^{1/4} \frac{e^{-x^2/2}}{\sqrt{2^n n!}} H_n(x)$$

Hermite Polynomial.

$$x = \sqrt{\frac{m\omega}{\hbar}} q$$

$$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} e^{-x^2}$$

$$= \frac{1}{\sqrt{\pi}} \frac{e^{x^2}}{x^2} \int_{-\infty}^{\infty} du (-2iu)^n e^{-u^2 + 2iux}$$

$$\int_{-\infty}^{\infty} du (-2iu)^n e^{-u^2 + 2iux}$$



So, in psi n of q, yeah, here it has to be square root n factorial. And this comes from the normalization of this wave function right. So, this is sum over n this quantity is psi star n of q prime e to the power minus beta epsa n divided by Z, and this quantity is psi n of q. Such psi n of q as we have seen was m omega h bar pi 1 fourth square root 2 to the power n n factorial e to the power x square by 2 over square root pi.

And then I have this integral presentation of the Hermite polynomial which I use over here minus u square plus was there two factor, then yeah plus twice iu of x. So, then we write on this sum over n m omega over h bar pi one-fourth square root 2 to the power n, n factorial e to the power x prime square over 2 because this is q prime now. So, this becomes x prime square over 2 square root pi minus infinity to plus infinity du minus twice iu raised to the power n e to the power minus u square plus twice iux prime.

And then you have e to the power minus $\beta \epsilon \pi n$ which is going to be n plus half \hbar bar ω divided by Z , and then you have ψ of q which again gives you $m \omega$ over \hbar bar π raised to the power $1/4$ 1 over square root 2 the power n , n factorial e to the power x square by 2 , because this is now q not q prime square root of π minus infinity to plus infinity dv minus twice iv raised to the power n e to the power minus v square plus twice ivx . I look horribly complicated.

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$$\begin{aligned}
 &= \frac{1}{h} \left(\frac{m\omega}{\hbar\pi} \right)^{1/2} \int_{-\infty}^{\infty} \psi^* \psi \, dx \\
 &= e^{-\beta\hbar\omega/2} \left(\frac{m\omega}{\hbar\pi} \right)^{1/2} \frac{1}{\sqrt{2}} \frac{1}{\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-(x^2+v^2)/2} e^{(u^2+v^2-2iux'-2ivx)} \sum_n \frac{1}{2^n n!} e^{-\beta\hbar\omega n} (-4uv)^n \\
 &= e^{-\beta\hbar\omega/2} \left(\frac{m\omega}{\hbar\pi} \right)^{1/2} \frac{1}{\sqrt{2}} \frac{1}{\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-(u^2+v^2-2iux'-2ivx)} e^{-2uv} \sum_n \frac{1}{n!} (-2uv)^n e^{-\beta\hbar\omega n} \\
 &= e^{-\beta\hbar\omega/2} \left(\frac{m\omega}{\hbar\pi} \right)^{1/2} \frac{1}{\sqrt{2}} \frac{1}{\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-(u^2+v^2-2iux'-2ivx)} e^{-2uv} \sum_n \frac{1}{n!} (-2uv)^n e^{-\beta\hbar\omega n} \\
 &= e^{-\beta\hbar\omega/2} \left(\frac{m\omega}{\hbar\pi} \right)^{1/2} \frac{1}{\sqrt{2}} \frac{1}{\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-(u^2+v^2-2iux'-2ivx)} e^{-2uv} \left[1 - 2uv e^{-\beta\hbar\omega} + \frac{(2uv)^2}{2!} e^{-2\beta\hbar\omega} + \dots \right] \\
 &= e^{-\beta\hbar\omega/2} \left(\frac{m\omega}{\hbar\pi} \right)^{1/2} \frac{1}{\sqrt{2}} \frac{1}{\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-(u^2+v^2-2iux'-2ivx)} e^{-2uv} e^{-\beta\hbar\omega} \\
 &= e^{-\beta\hbar\omega/2} \left(\frac{m\omega}{\hbar\pi} \right)^{1/2} \frac{1}{\sqrt{2}} \frac{1}{\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-(u^2+v^2-2iux'-2ivx)} e^{-2uv} e^{-\beta\hbar\omega}
 \end{aligned}$$



But things are going to simplify very soon. This together with this gives me $m \omega$ \hbar bar π half, then I have this square root π and this square root π that gives me 1 by π , and I have a Z gives me 1 by Z right. The exponential factors that you see over here is not being integrated over neither is being sum.

So, it is $x^2 + x'^2$ by 2. Let us take the integrals over here du and then I have dv . Now, I take care of the sum n . But before that I also note that this factor that you see over here gives me e to the power minus $\beta \hbar \omega$. So, I will bring out e to the power minus $\beta \hbar \omega$. And hopefully, you have taken care of everything. Let us see.

Then I have this and this which is e to the power minus $u^2 + v^2$ minus twice uv prime minus twice uv times sum over n $1/2$ to the power n , $n!$ because you have a square root here and the square root here, the product of these two gives you 2 to be power $n/n!$, and then you have e to the power minus $\beta \hbar \omega$; not only that now you also have this term and this term.

So, this term together that minus and minus gives you plus, 2, 2 gives you 4. But you have an i^2 , so that this is minus 4 uv raised to the power n . Let us concentrate on this sum, this sum I can write down as $1/n!$ minus $4uv$ by 2. The 2 I can bring inside the bracket raised to the power n $\beta \hbar \omega$, which is sum over n $1/n!$ the 2, 2 gets cancelled out. And I have minus twice uv raised to the power n e to the power $\beta \hbar \omega$.

If you now write down this expand the whole thing, then the first term is 1, the second term is minus twice uv e to the power $\beta \hbar \omega$, then you have twice uv^2 by 2 factorial e to the power twice $\beta \hbar \omega$, and then we have higher order terms. And if you look at this carefully, then you see that this is nothing but the expansion of minus twice uv e to the power minus $\beta \hbar \omega$.

So that what you originally started off with now reduces to $\beta \hbar \omega$ over $\hbar \pi^{1/2}$ $1/Z$ and $1/\pi$, then you have $x^2 + x'^2$ by 2 you have minus infinity to plus infinity du minus infinity to plus infinity dv you have e to the power minus $u^2 + v^2$ minus twice uv prime minus twice uv , and then you have this which is e to the power minus 2 uv e to the power minus $\beta \hbar \omega$.

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$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\frac{1}{2}(u^2 + v^2 - 2iux' - 2ivx + 2uv e^{-\beta t \omega})} e^{-2uv e^{-\beta t \omega}}$$

$$N = \begin{pmatrix} u \\ v \end{pmatrix} \quad A = 2 \begin{pmatrix} 1 & e^{-\beta t \omega} \\ e^{-\beta t \omega} & 1 \end{pmatrix} \quad b = \begin{pmatrix} x' \\ x \end{pmatrix}$$

$$\frac{i}{2} W^T A N + i b^T W$$

$$W^T A N = 2 \begin{pmatrix} u & v \end{pmatrix} \begin{pmatrix} 1 & e^{-\beta t \omega} \\ e^{-\beta t \omega} & 1 \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}$$

$$= 2 \begin{pmatrix} u & v \end{pmatrix} \begin{pmatrix} u + v e^{-\beta t \omega} \\ u e^{-\beta t \omega} + v \end{pmatrix} = 2 \left[u^2 + v^2 + 2uv e^{-\beta t \omega} \right]$$



Now, we concentrate on this integral. So, we will rewrite this integral as $\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\frac{1}{2}(u^2 + v^2 - 2iux' - 2ivx + 2uv e^{-\beta t \omega})}$. And then I have $2uv e^{-\beta t \omega}$. This happens to be a multivariate Gaussian integral. This term that you see over here I can write down this as half of $b^T A^{-1} b$, sorry not $b^T A^{-1} b$.

Now, to evaluate this, you define a vector a column vector which is going to be your u and v , a matrix A which is going to be your $1 - e^{-\beta t \omega}$ and $1 + e^{-\beta t \omega}$, and another column vector which is going to be your x' and x . Once you have defined this you can recast this part that is underlined over here as half of $W^T A^{-1} W + i b^T W$ which is basically $b \cdot W$.

W transpose A W is going to give you u v 2 times this, the 2 comes from this part 1 e to the power minus beta h bar omega e to the power minus beta h bar omega 1 times u v. So, that if you now do this, this is going to be twice u v times u plus v e to the power minus beta h bar omega u e to the power minus beta h bar omega plus v.

And if you multiply now this one – the transpose, you are going to get u square plus v square plus twice uv minus beta h bar omega. So, that you see u square plus v square combined with this is given by this term. The two factor that you see cancels with this half factor.

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$$\begin{aligned}
 & \frac{1}{2} W^T A W + i b^T N \\
 W^T A W &= 2(u \ v) \begin{pmatrix} 1 & e^{-\beta t \omega} \\ e^{\beta t \omega} & 1 \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} \\
 &= 2(u \ v) \begin{pmatrix} u + v e^{-\beta t \omega} \\ u e^{\beta t \omega} + v \end{pmatrix} = 2 \left[u^2 + v^2 + 2uv e^{-\beta t \omega} \right] \\
 i b^T N &= 2i(x' \ x) \begin{pmatrix} u \\ v \end{pmatrix} = 2i x' u + 2i x v
 \end{aligned}$$



Similarly, b transpose W is going to be x prime x u v which is going to be x prime u plus xv. And b transpose has to be ok, so b there has to be a factor 2 here. So, i times b transpose is going to be 2 i, and I have 2 i; and I have 2 i and that is precisely what I have.

