

Introduction to Classical Mechanics
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Lecture 49
First Integral of a Heavy Symmetric Top

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Heavy symmetric top with one point fixed



$$L = \frac{I_1'}{2} (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + \frac{I_3}{2} (\dot{\psi} + \dot{\phi} \cos \theta)^2 - \mu g l \cos \theta$$

θ & ϕ are cyclic. [First integrals]

$$\frac{\partial L}{\partial \dot{\phi}} = M_2 = \dot{\phi} (I_1' \sin^2 \theta + I_3 \cos^2 \theta) + \dot{\psi} I_3 \cos \theta. \quad (1)$$

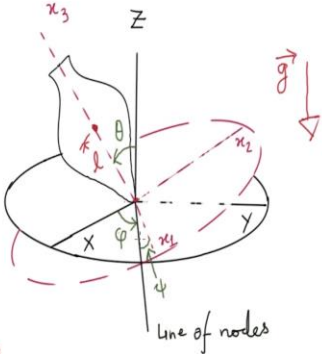


$$\frac{\partial L}{\partial \dot{\psi}} = M_3 = \dot{\phi} I_3 \cos \theta + \dot{\psi} I_3 \quad (2)$$

Conservative force \Rightarrow Energy is conserved

$$E = \frac{I_1'}{2} (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + \frac{I_3}{2} \underbrace{(\dot{\psi} + \dot{\phi} \cos \theta)^2}_{\omega_3^2} + \mu g l \cos \theta \quad (3)$$



Heavy Symmetrical top with lowest point fixed.

- Origin is chosen to be at the fixed point
- Euler angles completely specify the orientation
- One of the principal axis is the axis of symmetry - x_3

So, let us continue our discussion of Heavy Symmetric Top with one point fixed that is what we were looking at last time. Where that point is the lowest point of the top. So, the Lagrangian we wrote down last time was this the Lagrangian of the system is I_1' prime over 2 and then you have θ dot square plus ϕ dot square sin square θ , these are your euler angles and their

derivatives and you have then $I_3 \dot{\psi}^2 + I_1 \dot{\phi}^2 \cos^2 \theta$ minus the potential energy term which is $\mu g l \cos \theta$, where μ denotes a total mass of this top.

So, we quickly note that the coordinates ϕ and ψ are cyclic you see ϕ and ψ only appear as derivatives not otherwise. So, our coordinates ϕ and ψ are cyclic and you know that they imply conservation of their corresponding conjugate momentum, which means that p_ϕ which is $\frac{\partial L}{\partial \dot{\phi}}$ which is p_ϕ , which is also same as let us go back and see what is $\dot{\phi}$, see ϕ remember ϕ was the first rotation we talked about and it happens about the Z axis.

So, this corresponds to the angular momentum M_z . So, you take your angular momentum and project it along Z that is what is $\frac{\partial L}{\partial \dot{\phi}}$ and this comes out to be what? $\frac{\partial L}{\partial \dot{\phi}}$, so this is and you are not going to contribute, so this term will contribute and this term will contribute.

So, if I take the derivatives I get this, $\dot{\phi} I_1 \sin^2 \theta$, so there is the first this contribution from here that 2 cancels with this 2 here plus $I_3 \cos^2 \theta$ that is from the square term and then from the cross term we get $\dot{\psi} I_3 \cos \theta$ that is let us say our equation 1.

Now, ψ is also cyclic. So, I write down $\frac{\partial L}{\partial \dot{\psi}}$ will be a constant. So, here M_z is a constant, now $\frac{\partial L}{\partial \dot{\psi}}$ will be constant and what is $\frac{\partial L}{\partial \dot{\psi}}$ let us go back, $\dot{\psi}$. So, this rotation ψ was happening about the x_3 axis remember that the final rotation. It was about the final location of your how do you say anyway, it was about x_3 axis. This was the final rotation. So, I call it the M_3 . So, we remember we are denoting by 1, 2 and 3 these coordinates x_1, x_2 and x_3 and by capital X, Y and Z the fixed coordinates and that is why I am putting a 3.

So, here I have used the fact that the conjugate momentum corresponding to this rotation is angular momentum. So, that us your M_3 . Let us, take a derivative with respect to $\dot{\psi}$. So, these two do not contribute. This does not contribute only this term contributes and of course, there are two terms one from $\dot{\psi}^2$ and another from the cross term and you get $\dot{\phi} I_3 \cos \theta + \dot{\psi} I_3$.

Let us, call it equation number 2. So, these are our two first integrals we get. We also have one first integral because this is a conservative system and total energy will be conserved. So, we have one more, so conservation of energy conservative force implies energy conservation total energy conservation.

So, one more first integral and that is the total energy E is conserved which we have already seen but I will write $\dot{\theta}^2 + \dot{\psi}^2 \sin^2 \theta$ these are your kinetic term which I am writing right now $\frac{I_3}{2} \dot{\psi}^2 + \dot{\theta}^2 \cos^2 \theta$ this is your kinetic and then you have the potential energy $\mu g l \cos \theta$, that is our equation number 3.

Now, look at this part, this is, you remember what this was? This was the kinetic energy was half $I_1 \omega_1^2$ plus then you had $I_2 \omega_2^2$ and then the third was half $I_3 \omega_3^2$. So, if you recall, this is just the expression of ω_3^2 there is a reason why I am writing it like this here.

So, these are our first integrals which we can now look at because they are simpler than second order differential equations. So, what I can do immediately is look at this equation 1 and 2, they both involve $\dot{\psi}$ and $\dot{\theta}$ there is no nothing else with the derivative. So, I can solve these two because my M_z and M_3 are constants because they are conserved quantities. I can solve these two for $\dot{\psi}$ and $\dot{\theta}$ and that will be in terms of θ . So, let me write it down.

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Using (1) & (2) we can write



$$\dot{\psi} = \frac{M_z - M_3 \cos \theta}{I_1 \sin^2 \theta}; \quad \dot{\psi} = \frac{M_3}{I_3} - \left(\frac{M_z - M_3 \cos \theta}{I_1 \sin^2 \theta} \right) \cos \theta \quad (5)$$

$M_3 = I_3 \Omega_3 \Rightarrow \Omega_3$ is a constant

Define $E' = E - \frac{I_3}{2} \Omega_3^2$

$$E' = \frac{I_1}{2} \dot{\theta}^2 + U_{\text{eff}}(\theta) \quad ; \quad U_{\text{eff}}(\theta) = \frac{(M_z - M_3 \cos \theta)^2}{2 I_1 \sin^2 \theta} + \mu g l \cos \theta$$

↑ one dimensional problem
we should be interested
in the turning points.

Heavy symmetric top with one point fixed

$$L = \frac{I_1'}{2} (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + \frac{I_3}{2} (\dot{\psi} + \dot{\phi} \cos \theta)^2 - \mu g l \cos \theta$$

ϕ & ψ are cyclic. [First integrals]

$$\frac{\partial L}{\partial \dot{\phi}} = M_2 = \dot{\phi} (I_1' \sin^2 \theta + I_3 \cos^2 \theta) + \dot{\psi} I_3 \cos \theta \quad (1)$$

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Conservative force \Rightarrow Energy is conserved

$$E = \frac{I_1'}{2} (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + \frac{I_3}{2} \underbrace{(\dot{\psi} + \dot{\phi} \cos \theta)^2}_{\Omega_3^2} + \mu g l \cos \theta \quad (3)$$

So, using 1 and 2 we can write $\dot{\phi}$ is equal to $M_2 \cos \theta - M_3$ divided by $I_1' \sin^2 \theta$ and then you will also have $\dot{\psi}$, maybe I should give a little bit more gap, $\sin^2 \theta$ and your $\dot{\psi}$ will be M_3 / I_3 , this you can check and $M_2 \cos \theta - M_3$ over $I_1' \sin^2 \theta$ and then the entire thing is multiplied with the $\cos \theta$, maybe I should put a bracket, this looks fine.

So, let me call this equation 3 and 4 this is 3, this is 4, no 4 and 5, equation 4 I think there is a way I can do it. So, this is 4 and this is 5 that is good. Now, note that your M_3 is just $I_3 \omega_3$ third component of angular momentum is $I_3 \omega_3$ because these are your principal axis and because your M_3 is a constant which we have seen here, M_3 is a constant because $\dot{\psi}$ is cyclic your ω_3 is also constant, that is good.

Now, what I do is let me go back to E. So, you have E equal to this thing plus half $I_3 \omega_3^2$, now ω_3 is constant this is constant, so I take this entire piece to the left and combined with E to define a new quantity, which I call E prime. So, I define E prime is equal to E minus $I_3 / 2 \omega_3^2$. It makes sense to define this because these are constants, E is constant that is the total energy and this entire quantity is a constant.

So, with this I have now E prime is equal to what? Let us go back this term plus that term. So, this term gives you a $\dot{\theta}^2$ term, plus something which involves only θ and something involves only θ and in this $\dot{\phi}$ I am going to substitute this thing, because I have already determined $\dot{\phi}$ in terms of θ .

So, I get when I substitute in E' the following, I get $2I_1' \dot{\theta}^2$ that is this thing and then $\dot{\phi}$ is expressed in terms of θ this is any way involving θ , this is gone to the left side, this involves only θ . So I define a new quantity, which is a function of only θ and I call it $U_{\text{effective}}(\theta)$, where $U_{\text{effective}}(\theta)$ is defined by $Mz - M^3 \cos^2 \theta$ this whole thing is squared and we should divide it by $2I_1' \sin^2 \theta$ make sure that there are no mistakes I think this is correct $\mu g l \cos \theta$, that is nice.

And it is nice because it looks very familiar, does it look familiar this equation. Well you have encountered the same thing when we were looking at a two body problem. Remember the equation for R coordinate. The radial separation. It was $E = \frac{1}{2} \mu \dot{R}^2 + U_{\text{effective}}(R)$. So, instead of R you have θ here, but as far as the mathematics is concerned and other logics is concerned is the same there is no difference.

So, we can do what we did there. So, we are happy. We have a one dimensional problem as far as it θ variable is concerned. So, what we would like to know here is, let us recall what we were asking when we had this equations in the central problem or two body problem. We wanted to know where the turning points are, where $\dot{\theta}$ goes to 0 and depending on the turning points we knew where between what limits are the particle is going to move.

So, similarly here once we have found the turning points we will know that the θ has to vary within those limits of θ which we determined by finding the turning points. So, that is what we will do. So, let me just write out. We should be interested in now in the turning points that is good. But before we do all that, let us do a little bit more of work. Now, I will make a substitution I will define u as a new variable and that will be equal to $\cos \theta$.

So, what I am going to do is I am going to take my E' which is here and wherever I have $\cos \theta$ I am going to substitute u , see this will become u , this will become u , this will become $1 - u^2$. So, as far as $U_{\text{effective}}$ is concerned you will immediately have it in terms of u and $\dot{\theta}^2$ will also give you a nice expression of u which I will show.

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Define $u = \cos \theta$ (9)

$$\Rightarrow \dot{\theta}^2 = \frac{\dot{u}^2}{1-u^2} \quad (10)$$

$$U_{\text{eff}}(u) = \frac{(M_2 - M_3 u)^2}{2I_1'(1-u^2)} + \mu g l u \quad (11)$$

Substituting (10) & (11) in (7) we get

$$\rightarrow E'(1-u^2) = \frac{I_1' \dot{u}^2}{2} + \frac{(M_2 - M_3 u)^2}{2I_1'} + \mu g l u(1-u^2).$$

$$\frac{2E'}{I_1'} \equiv \alpha, \quad \frac{M_2}{I_1'} \equiv a, \quad \frac{M_3}{I_1'} \equiv b, \quad \frac{2\mu g l}{I_1'} \equiv \beta$$

$$\dot{u}^2 = f(u)$$

$$f(u) = (\alpha - \beta u)(1-u^2) - (a - bu)^2.$$

Using (1) & (2) we can write

$$\dot{\phi} = \frac{M_2 - M_3 \cos \theta}{I_1' \sin^2 \theta}; \quad \dot{\psi} = \frac{M_3}{I_3} - \left(\frac{M_2 - M_3 \cos \theta}{I_1' \sin^2 \theta} \right) \cos \theta \quad (5)$$

$$M_3 = I_3 \Omega_3 \Rightarrow \Omega_3 \text{ is a constant}$$

Define $E' = E - \frac{I_3}{2} \Omega_3^2$

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↑ One dimensional problem
we should be interested
in the turning points.

So, let us define u equals \cos of θ , which gives you immediately \dot{u} , so maybe I can write it is obvious so you can immediately get from this that \dot{u}^2 all you have to do is just take a derivative and you get the following \dot{u}^2 is not this way I want to $\dot{\theta}^2$ from this you get $\dot{\theta}^2$ as \dot{u}^2 over $1 - u^2$, that is nice.

And my $U_{\text{effective}}$ becomes $U_{\text{effective}}$ of u is $M_2 - M_3 u$, $M_2 - M_3 u$ whole squared over $2 I_1' (1 - u^2)$, u squared over $2 I_1' (1 - u^2)$ plus $\mu g l u$, that is what our $U_{\text{effective}}$ would be and if you substitute this let us call it what 5, 6 and 7 maybe no.

Let us, call it 6, let us call this one 7 and this one 8 and 9, 10, 11. So, I am going to put 10 and 11 into 7 and 8 into 7.

So, what do I get? I get substituting, substituting 10 and 11 in 7, what do I get? Let us go back. So, you get the following $E' = 1 - u^2$, so I am taking $1 - u^2$ also and multiplying on the left hand side and this is equal to $\frac{1}{2} \mu^2 + Mz - \frac{1}{2} \mu^2$, please check that is correct. $1 - u^2$ I think, what was that? $1 - u^2$ that is because of this term getting multiplied $1 - u^2$, that is correct. So, that is nice.

Now, what I will do is, I will take the coefficient of u^2 and divided this coefficient throughout this equation. So, I am going to divide everything by $\frac{1}{2}$, I am going to divide everything by $\frac{1}{2}$, so that the coefficient of this becomes 1 that is what I am going to do and for that I mean, of course I will I mean I do that and I also define.

So, you see this E' will get divided by E' will carry a factor with it $\frac{1}{2}$ over $\frac{1}{2}$. Because will you divide it becomes $\frac{1}{2} \mu^2$ over $\frac{1}{2}$, so μ^2 over $\frac{1}{2}$ then you will have this one will become 1, the 2 when I multiply this entire equation by 2 to get rid of this there will be a 2 here which will cancel this 2 and then you have μ^2 divided here.

So, that will make μ^2 which I want to bring in this and it will be dividing Mz over $\frac{1}{2}$ and Mz , Mz over $\frac{1}{2}$, so I will have Mz over $\frac{1}{2}$ and I will have, first let me write Mz , Mz over $\frac{1}{2}$, then I will have Mz over $\frac{1}{2}$ and then the last term will have, what will it have it will have 2 will come in the numerator. So, $2 \mu^2$ over $\frac{1}{2}$ this will have this, that is correct.

So, I will define all these constants to be the following, this is very standard notation. This is α , this should have been called β , this is called Mz over $\frac{1}{2}$ is called a Mz over $\frac{1}{2}$ is called b . So, that is these are what I am defining. So, with this my equation becomes the following, so you can substitute this this definition into this equation and you will obtain u^2 , u^2 and the entire remaining thing will be some function of u where $f(u)$ is given by the following. It will be $\alpha - \beta u$ then you will have $1 - u^2$ that is correct then minus $a - b u^2$, that is what you will get.

We will continue with further discussion on this in the next video. So, please before watching this next video, please ensure that you have reproduce all these steps which are fairly easy and then we will be ready to analyse the motion of this top, which is quite interesting actually. See you then.