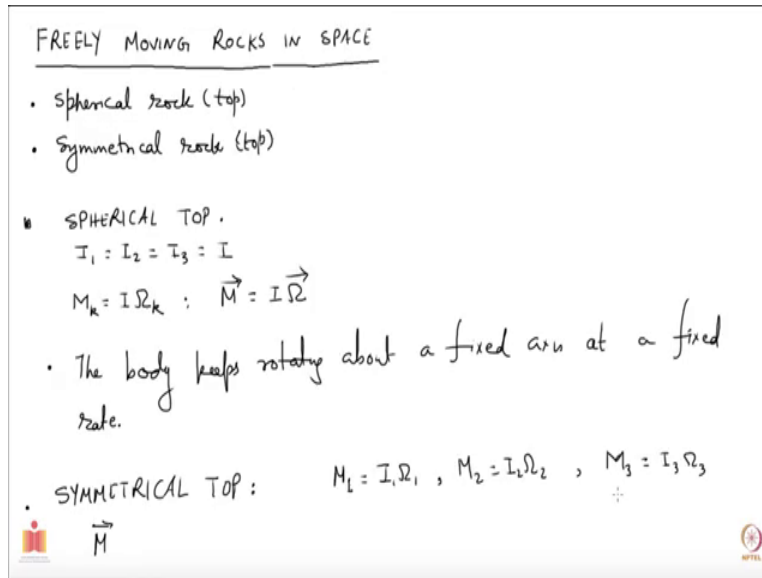


Introduction to Classical Mechanics
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Lecture 46
Motion of a free symmetric top

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In this video we will start looking at the motion of rocks in space and these rocks will be considered as free that is there will be no force acting on them and as we have seen that these rigid bodies or rocks are generically referred as top and they come in different varieties, they could be spherical, completely spherical symmetric, so you say that is spherically symmetric rock as a spherical top.

If your rock is symmetric about some axis, you say you have a symmetric top and if it does not have any symmetry then you say asymmetric top. So let us imagine some rock is moving around in space and it is spherically symmetric and then we want to know how it will be moving about and as before, in the last video we would not be looking at the motion of its centre of mass. So, we are only looking at the non-trivial part, the part due to rotation of this rock.

So, first case of spherical rock is very easy to look at. So let me write down here, spherical rock. So, in this case as you know already all the principal moments of inertia are equal, let us say the common value is I and then if you want to know the intrinsic angular momentum M , let us look

at the K th momentum, K could be 1, 2 or 3, then this would be $I \omega_K$, that is what it is which means that intrinsic angular momentum is parallel to the angular velocity.

Because there are no forces on the top, your angular momentum is not going to change and which also implies that the intrinsic angular momentum is not going to change because the part which describes the motion of the centre of mass that will also be conserved. So, that would mean that the intrinsic angular momentum which we denote by M is also going to be constant in time, which means that the angular velocity ω is also going to be constant.

So, the direction, the axis about which the body which is spherically symmetric is going to rotate is not going to change with time. In general you know that ω vector is going to change with time, the axis of rotation is going to change with time, but for a spherical top because of its highly symmetric construction it is not going to change, which means the picture is this. The body keeps spinning or rotating about a fixed axis at a fixed rate.

So, that is the simple motion that it is going to execute. Let us now reduce the degree of symmetry that you have and let us say we do not have a spherically symmetric top but you have a symmetrical top, meaning your rock which you are looking at has only one axis of symmetry, so I want to look at symmetrical top and you can again imagine rock which is moving freely deep in space, it has been thrown in some manner and now you just want to know how it moves around.

So, let us see what we can say. If you look at the intrinsic angular momentum then M , so I am imagining that I am decomposing the angular momentum of this body, whenever I say angular momentum I mean intrinsic angular momentum, so this M . Because there are no forces this is going to be conserved. So, let us say you think of a moment, one particular instant and at that instant you are going to decompose this M into its components about body axis.

So, let us say at a given moment the principal axis of the body are oriented in some manner and at that time you want to decompose M along those directions, so that is what we are doing. So, M_1 will be $I_1 \omega_1$, M_2 will be $I_2 \omega_2$ that is what we have seen already, it is the general thing and the third component will $I_3 \omega_3$.

Now, what I am going to do is look at where x_3 points, so x_3 is the, by x_3 I am denoting the third coordinate which is attached with the body and I look at the vector M . So, vector M is going to be fixed in space and at that instant I am looking at x_3 direction, so this direction of M and the direction of x_3 , they define a plane, any two lines define a plane or any two directions you can take and that will define a plane, so that is the thing which I am going to utilize now.

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$\vec{\Omega} = \vec{\Omega}_{pr} + \vec{\Omega}_a$
 $\Omega_a + \Omega_{pr} \cos \theta = \Omega_3$ (1)
 $\Omega_a = \Omega_3 - \Omega_{pr} \cos \theta$ (1)
 $\Omega_{pr} \sin \theta = \Omega_1$ (2)
 Substitution
 $M_k = I_k \Omega_k$
 $\Omega_{pr} = M/I_1$: Axis of the top rotates at a uniform rate about the direction of angular momentum.

FREELY MOVING ROCKS IN SPACE

- Spherical rock (top)
- Symmetrical rock (top)
- SPHERICAL TOP.
 $I_1 = I_2 = I_3 = I$
 $M_k = I \Omega_k$: $\vec{M} = I \vec{\Omega}$
- The body keeps rotating about a fixed axis at a fixed rate.
- SYMMETRICAL TOP : $M_1 = I_1 \Omega_1$, $M_2 = I_2 \Omega_2$, $M_3 = I_3 \Omega_3$
 $\vec{M} = (M_1, 0, M_3)$ $\Rightarrow \Omega_2 = 0$

So, let us see here, I have drawn already this setup, do not look at x_1 , this I will explain. I wish I had not drawn this already but just to save a little bit of time I have done that. So here is the

direction M which is not going to change with time, at the moment I am looking at things, let us say x_3 axis is in this direction and I say that the plane which is defined by M and x_3 is the plane of your screen, so that is the plane we have. So your M and x_3 are in your screen.

Now how about x_1 and x_2 , well they have to be perpendicular to x_3 , but because your body has symmetry about x_3 you can take x_1 and x_3 to be in any directions, that is what we discussed earlier also. So what I am going to do is I will choose to put x_1 direction in the plane of your screen and x_2 will be of course perpendicular to the screen, so this choice I can make because of the symmetry of this object.

So, here we are, all the x_1 , x_3 and M they are all in the same plane at this instant. At the next instant of course things will change, but at the instant we are looking at everything is in same plane. So if I do so, then M_2 is 0 because M is in this plane, M is in x_1 , x_3 plane so it does not have a component M_2 , so there is no M_2 component which means M_2 is 0 at this instant, which means ω_2 is 0.

This is good which means, ω has only, ω_1 and ω_3 as the components and ω_2 is 0 which means your vector ω lies in x_1 , x_3 plane. It is ω_1 , 0, ω_3 . This is a vector which is lying in x_1 and x_3 plane, so I want to draw that and of course as you know it will not point along M , so let me draw that, let us see if I can do it or should I, there is a, yes. Looks like it is going to work, perfect. I think this is good. Let me use some colour, maybe no colour, fine no colour. Let us use red.

So, this is the vector ω and ω is also going to be in this plane. That is nice, let me cross it. This is my vector and that is fine. Now, I want to first tell that the angle which axis of the top, this x_3 , x_3 is the direction of the axis of the top, the symmetry axis. This makes an angle θ at this instant, let us say this is angle θ that is good. Now, let me draw a line parallel to M , so this should be vertical, let me maybe I should make a nice, now that works. So this direction which I am about to draw is parallel to intrinsic the angular momentum.

Let me, put just this looks perpendicular. So if I or maybe I just fill it up and good. So this one, I am going to call this vector from here to there as ω_{pr} , pr for precision, you will see why I have used this subscript. So, anyhow it is just the projection of ω along the direction of M

and then this remaining part, so this vector from here to there. Let me again bring this. This vector is what I am going to call as ω_A . I call it ω_A , because this is along the axis of symmetry. a means, subscript a axis of symmetry.

So, I have written my angular velocity vector as sum of these two vectors, where ω_A points in the direction of the axis of symmetry and the other one points or is along the direction of M , so that is how I decompose it. That is good. Now let me see, so let us draw perpendicular here, so I project ω along this direction which will be ω_3 , from here to there, it is ω_3 .

Third component of the angular velocity and then this anyway is ω_a and this piece from here to there is $\omega_p \cos \theta$ because this angle is also θ . So this will be $\omega_p \cos \theta$, $\omega_p \cos \theta$. That is good now let us have understanding of this. So, first thing is, let me write down a few questions now, I think I have written everything. So, first is that my ω is $\omega_p \cos \theta + \omega_a$, that is the decomposition of this vector we have done. So let me nothing, I think I have said already.

So, here ω has been written as a sum of ω_p which is parallel to M and ω_a , which is along the symmetry axis of the top. So, from the figure it is very clear that you can write $\omega_a + \omega_p \cos \theta$, this is ω_3 , so here you have ω_3 from here to there and this is sum of ω_A and $\omega_p \cos \theta$, that is what I have written, so let me call this as equation 1, very good and from this I can conclude that ω_a is $\omega_3 - \omega_p \cos \theta$.

Just taken one term on the other side and also your $\omega_p \sin \theta$ is ω_1 . You see if you take this ω and project it along x_1 , this is what you get. It is equivalent to projecting your, it is equivalent to projecting your ω_p along x_1 axis. So, your $\omega_p \sin \theta$ is equal to ω_1 . That is correct. Now I know my M , angular momentum and I would like to determine these ω s and what I am interested is in ω_a and ω_p .

So, I am going to use my ω_1 as M_1 / I_1 and ω_3 as M_3 / I_3 ; that is what I am going to use. So substituting that I get, let me write substituting M_1 , let us say M_k , $I_k \omega_k$, there is no summation over k . I get the following, I get you can easily check. That is trivial, that

omega precision is M over I_1 . That is one result I have. So, this says that your top is rotating at a uniform rate about the direction of the angular momentum.

See the M is not going to change with time, it is going to always point in the same direction and you are looking at a vector ω which points along M which means that when you are looking at ω , it corresponds to the motion of the rigid body as a rotation about this axis, ω as the axis, so that is why it corresponds to rotation about the direction of angular momentum and that is what we call as precession.

So, let me write it down, this corresponds to the axis of the top at a uniform rate about the direction of angular momentum, about the direction of angular momentum and then the remaining thing let me talk about.

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$$\begin{aligned} \Omega_a &= \Omega_3 - \dot{\theta} \cos \theta \\ &= \frac{M \cos \theta}{I_3} - \frac{M}{I_1} \cos \theta \\ \Omega_a &= M \left(\frac{1}{I_3} - \frac{1}{I_1} \right) \cos \theta. \quad \checkmark \\ T &= \frac{1}{2} I_1 \Omega_1^2 + \frac{1}{2} I_2 \Omega_2^2 + \frac{1}{2} I_3 \Omega_3^2 \\ &\left(\begin{array}{l} \Omega_2 = 0 \\ \frac{1}{2} I_1 \Omega_1^2 + \frac{1}{2} I_3 \Omega_3^2 \end{array} \right) \\ &= \frac{1}{2} M^2 \left[\frac{\sin^2 \theta}{I_1} + \frac{\cos^2 \theta}{I_3} \right] \\ \theta &\Rightarrow \text{const.} \end{aligned}$$

$\vec{\Omega} = \vec{\Omega}_{pr} + \vec{\Omega}_a$
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 $\Omega_a = \Omega_3 - \Omega_{pr} \cos \theta$ (1)
 $\Omega_{pr} \sin \theta = \Omega_L$ (2)

Substitution
 $\Omega_{pr} = M/I_1$: Axis of the top rotates at a uniform rate about the direction of angular momentum.

$\text{III: } (\vec{\Omega} \times \vec{r})^2 = (\vec{\Omega} \times \vec{r})_i (\vec{\Omega} \times \vec{r})_i$
 $= \epsilon_{ijk} \Omega_j r_k \epsilon_{ilm} \Omega_l r_m$
 $= \Omega_j (\epsilon_{ijk} \epsilon_{ilm} r_k r_m) \Omega_l$
 $= \Omega_j I'_{jk} \Omega_l$

$\epsilon_{ijk} \epsilon_{ilm} = \delta_{jl} \delta_{km} - \delta_{jm} \delta_{kl}$

$I'_{jl} = (\delta_{jl} \delta_{km} - \delta_{jm} \delta_{kl}) r_k r_m$
 $= \delta_{jl} r_k^2 - r_j r_l \leftarrow \text{Inertia tensor.}$
 $I_{jkl} = \sum_m I'_{jkl}$
 $T = \frac{1}{2} \mu V^2 + \frac{1}{2} \Omega_j I_{jkl} \Omega_l$; $L = T - U$

So my omega a is, where is it? I think I wrote it before is omega 3 minus omega p cos theta, so this is the equation I am using, so let me write it again. Omega Pp cost of theta, and what is omega 3. Omega 3 is your M3 over I3, but your M3 is M cos theta minus, omega p we have already found M over I1 and the cos theta factor is here anyway.

So, your omega a is M1 over I3 minus 1 over I1 cos of theta. So, that is the omega a which is in the direction of x3, so it corresponds to the motion of your top about the x3 axis which is basically the angular velocity of the top spinning about its axis; that is the spin of the top about its axis that is what your omega a gives and omega pr is giving you the precession of x3 about M.

Now I want to establish that the $\cos \theta$ is a constant in this problem. So, for that I want to utilize, I would like to utilize the fact that the kinetic energy T is a constant because there are no forces on this body and I write T as $\frac{1}{2} \omega_2^2 + \frac{1}{2} I_3 \omega_3^2$. I believe I have done this already, let me check, we were looking at correct.

So here you see we wrote down the kinetic energy as sum of these two terms and you had ω_j , $I_j \omega_j$, all you have to do is now go to the principal axis, so once you are in the principal axis your I will be diagonal and it will have components I_1, I_2, I_3 and this will become $I_1 \omega_1^2 + I_2 \omega_2^2 + I_3 \omega_3^2$ and each of the terms will be multiplied with the half. That is what I am using now, that is what I am writing now.

And the instant at which I am looking at everything, ω_2 is 0 remember that. Your vector ω lies in the x_1, x_3 plane. So, ω_2 is 0 which means your T is $\frac{1}{2} I_1 \omega_1^2 + \frac{1}{2} I_3 \omega_3^2$ and if I again substitute what ω_1 and ω_3 are in terms of angular momentum, you get $\frac{1}{2} M^2 \sin^2 \theta / I_1 + \frac{1}{2} M^2 \cos^2 \theta / I_3$. T is constant, M is constant, you conclude that your θ is constant.

Which means that the body is going to make a fixed angle with respect to M , the x_3 is going to make a fixed angle with respect to M . This θ is not going to change. The rock is going to spin about its axis with angular velocity ω_A that we have determined already here and it is also the axis, the x_3 axis is also going to precess about M and x_3 is going to move at a uniform rate about M , it will keep going like this and the angle will not change.

That is the general motion of a free rigid body which is symmetric. But here you note that the direction of M and ω were not same. If they were same, if they were in the same direction then what would happen, I will leave for that you to figure out, should be easy. That is good and maybe one more thing.

I will also leave you to figure out what happens if your I_3 is smaller than I_1 or if your I_3 is bigger than I_1 because there will be a difference in \sin in these two cases. So, that is another thing I would urge you to think about. So, there is some motion of a symmetric rigid body in free space and we will continue more on rigid bodies in the next video.