

**Introduction to Classical Mechanics**  
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**Lecture 10**

**Kinetic term in generalized coordinates**

So as I was saying, we will start looking at the kinetic term, when we are using generalized coordinates, we will see what are the differences compare to using Cartesian coordinates.

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Kinetic term

Cartesian Coordinates.  
 $T = \sum_i \frac{1}{2} m_i \dot{\vec{r}}_i^2$   
 $T \rightarrow$  quadratic in  $\dot{\vec{r}}_i$   
 $\rightarrow$  no dependence on  $\vec{r}_i$   
 $\rightarrow \dot{\vec{r}}_i \cdot \dot{\vec{r}}_j$  if  $i \neq j$   
 $\rightarrow$  No cross terms

Generalized Coordinates  
 $T = \frac{1}{2} m (\dot{\theta}^2 + \dot{r}^2)$  ← example

In general  
 $T = T_0 + T_1 + T_2$   
 $T_0 =$  independent of  $\dot{q}_\alpha$   
 $T_1 =$  linear in generalized velocities  $\dot{q}_\alpha$   
 $T_2 =$  quadratic in  $\dot{q}_\alpha$

So, first what do we see when we see, when we work with Cartesian coordinates. Actually it will be nice to have color, let see if I can do it quickly, I can, good, good, I think it is done, Cartesian. So, what do you see when you are using Cartesian coordinates, you get T as half mi, let me put r dot, half mi ri dot square, that is what it is. So, this is of course quadratic in velocities, note that there is no term which involves the product of velocity of particle i with a particle j, there is no such term. So they are all, T is really diagonal in the velocities.

Also note that there is no factor in here, which depends on the coordinates itself. So, T does not depend on the coordinate but only on the velocities, this is another thing and also note that there is no linear term, it is all just quadratic. So, when you are using generalized coordinates, T is quadratic in r dots, note that there is no dependence on r and also for example there could have been r square dependence which is absent and also there is no, what I wanted to say, there is no ri

dot  $r_j$  dot kind of a term here, where  $i$  is not equal to  $j$ . You have only the terms where  $i$  and  $j$  are equal, so there are no cross terms, t e r m, no cross terms are present.

Let see what the situation is when we start using generalized coordinates. Let see if I can make, I do not know what I am making it colorful but let see. Okay good. Now I do not know how to get the green back, that is sad. Yeah here is the green, let us look at the generalized coordinates now. Yeah we have already looked at  $T$  in at least few contexts and you remember when we were looking at particle which was moving in two dimensions  $XY$  plane and we were using polar coordinates.

At that time, we wrote the kinetic energy of that particle in polar coordinates and it was  $\frac{1}{2} m \dot{\theta}^2 + \dot{r}^2$ . So, in this example itself you see that this condition is, this not condition, this thing is not true. Generalized coordinate  $r$  can appear or in principle, in some different context, some other coordinates could also appear. So, it not just the generalized velocities,  $\dot{\theta}$  and  $\dot{r}$  appear, but also coordinate itself can appear which is the case here.

Here it is still of the form where cross terms are not present, so you still do not have a term like  $\dot{\theta} \dot{r}$ . But as I will show that this is also not a general statement, in principle they can be present, that is one thing and also here you do not have a term which is just involving the velocity of one of the coordinates and as you are going to see very soon that is also not a general statement. So let us see, let see what the general expression would be, this was one specific example which we covered in one of the previous videos.

So, in general the kinetic energy can be written as the sum of three terms, which I will call as  $T_0$  plus  $T_1$  plus  $T_2$ , where  $T_0$  is independent of  $\dot{q}$ , independent of the generalized velocities. So let me put  $\alpha$  and  $\alpha$  will run from 1 to whatever degrees of freedom you have. So, it will be independent of all the generalized velocities and that is what  $T_0$  signifies.  $T_1$  is the term which will be linear in generalized velocities, linear in generalized velocities, velocities, which is  $\sum \alpha \dot{q}_\alpha$ . Then  $T_2$  will be quadratic in  $\dot{q}$ . That is what we will prove now.

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$$T = T_0 + \sum_{\alpha} T_{1\alpha} \dot{q}_\alpha + \sum_{\alpha\beta} T_{2\alpha\beta} \dot{q}_\alpha \dot{q}_\beta \quad \checkmark$$

Proof:  $T = \sum_i m_i \dot{\vec{r}}_i^2$

$\vec{r}_i = \vec{r}_i(q_1, \dots, q_{3N-k}, t)$

$\frac{d\vec{r}_i}{dt} = \frac{\partial \vec{r}_i}{\partial t} + \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial q_\alpha} \dot{q}_\alpha$

$\dot{\vec{r}}_i^2 = \left( \frac{\partial \vec{r}_i}{\partial t} + \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial q_\alpha} \dot{q}_\alpha \right) \cdot \left( \frac{\partial \vec{r}_i}{\partial t} + \sum_{\beta} \frac{\partial \vec{r}_i}{\partial q_\beta} \dot{q}_\beta \right)$

$= \left( \frac{\partial \vec{r}_i}{\partial t} \right)^2 + \underbrace{\frac{\partial \vec{r}_i}{\partial t} \cdot \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial q_\alpha} \dot{q}_\alpha + \frac{\partial \vec{r}_i}{\partial t} \cdot \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial q_\alpha} \dot{q}_\alpha}_{2 \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial t} \cdot \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial q_\alpha} \dot{q}_\alpha} + \sum_{\alpha} \sum_{\beta} \frac{\partial \vec{r}_i}{\partial q_\alpha} \cdot \frac{\partial \vec{r}_i}{\partial q_\beta} \dot{q}_\alpha \dot{q}_\beta$

Kinetic term

Cartesian Coordinates  
 $T = \sum_i \frac{1}{2} m_i \dot{\vec{r}}_i^2$   
 $T \rightarrow$  quadratic in  $\dot{\vec{r}}_i$   
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Generalized Coordinates  
 $T = \frac{1}{2} m (\dot{\theta}^2 + \dot{z}^2)$   $\leftarrow$  example

In general  
 $T = T_0 + T_1 + T_2$   
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 $T_1 =$  linear in generalized velocities  $\dot{q}_\alpha$   
 $T_2 =$  quadratic in  $\dot{q}_\alpha$

So, what I want to prove is this that, let us see, yeah, so I will introduce a notation and make what I said just now more explicit, that it will, T will have a form T<sub>0</sub> plus I said that this will be T<sub>1</sub> would be linear in velocities, so T<sub>1</sub> would be having a form like q<sub>α</sub> dot times some coefficients or some functions, let me call it T<sub>1α</sub>, I distinguish this T<sub>1α</sub> from T<sub>1</sub> because T<sub>1</sub> has no index and this guy has. So, it is clearly it is not the same quantity. So, you have T<sub>1α</sub> and there will be a sum over all the alpha. That is good.

Then as I said there will be terms which are quadratic, so you could have q<sub>α</sub> dot, q<sub>β</sub> dot and T<sub>2αβ</sub> summation over all alpha and beta. So, whatever I said here is more explicitly

written here and note that I am saying that it is not  $q_\alpha \dot{q}_\alpha$ , I am saying  $q_\alpha \dot{q}_\beta$  because cross terms may be present and this is what we want to prove now and proof is simple.

Proof, so what do I have to do, just remember that  $\frac{1}{2} m_i \dot{r}_i^2$  sum over all the particles, good. Now your  $r_i$  will be related to your generalized coordinates by the following,  $r_i = q_1$  to whatever  $q$  you have, let us say  $3N - k$ , total number of degrees of freedom and it could also possibly depend on time. Your constraints of, constraints on the system may be changing with time, so your transformation from  $r$  to the generalized coordinates would involve time as well. That is good.

All I have to do is take this, take the time derivative, so that I get an  $\dot{r}$  and take the  $\dot{r}^2$  and put in here, that is all I have to do which is not difficult. So,  $dr$  just a second, hold on please, so the velocity is  $dr_i / dt$  which is  $\partial r_i / \partial t$ , so you take the partial derivative and then you take the partial derivative with respect to all the coordinates,  $q_\alpha$  and you have  $q_\alpha \dot{q}_\alpha$  and you should sum over all the alphas. That is good.

Now I take this which is  $\dot{r}_i$  and square it,  $\dot{r}_i$  and square it. So, I write this as  $\partial r_i / \partial t$  plus the same thing which you have above, then I should dot it with again the same piece  $\partial r_i / \partial t$ , remember  $i$  is the level for the particle,  $\sum_i \partial r_i / \partial t \cdot \sum_\alpha q_\alpha \dot{q}_\alpha$  summation over  $\alpha$ .

Now if you have not already noticed the mistake and if you have not noticed still despite might (( ))(12:47) that there is a mistake, you need to know one thing that this is going to give a trouble if I use the same index  $\alpha$  which is here and the same index  $\alpha$  here. I should be very careful and I should not use the same index on both the sides.

So it will be wiser to change the  $\alpha$  to  $\beta$ ,  $q_\beta$ ,  $\dot{q}_\beta$ ,  $\beta$  and I leave it to you to figure out what will be the problem if I use  $\alpha$  in this piece and  $\alpha$  in here also, you should know why you will get into trouble. That is good, now I should take the dot products or have been, yeah. Now this, so I take a dot of this with that which is easy,  $\dot{r}_i$  over  $\Delta t$  square. Then I will take a dot product of this piece with this term,  $\partial r_i / \partial t$ , this piece with that term and you will have another similar term coming from this one and this one and they are both same because the  $\alpha$  is dummy.

I am just summing over all possible alpha. So in these two products, this times that and this times this, I can use the same indices, maybe I will go slowly so that you can see plus let us multiply this with that one. So, I get  $\sum_{\alpha} \frac{\delta r_i}{\delta t}$ , that is correct dot summation over beta  $\sum_{\alpha} \frac{\delta r_i}{\delta q} \beta_{\alpha}$  and then you look at the product of this guy with this guy which is  $\sum_{\alpha} \frac{\delta r_i}{\delta t}$  coming from here dot summation over alpha  $\sum_{\alpha} \frac{\delta r_i}{\delta q} \alpha_{\alpha}$ .

This dot product is because of the two vectors here, that is good and then you have one more term plus product of these two, these two will give you term which is quadratic in the velocities. These ones are generating terms which are linear in velocities, there is the only one velocity involved here, one velocity involved here plus so I write down summation over alpha summation over beta  $\sum_{\alpha} \frac{\delta r_i}{\delta q} \alpha_{\alpha}$  dot because there is a dot product here,  $\sum_{\alpha} \frac{\delta r_i}{\delta q} \beta_{\alpha}$  and then you have  $\sum_{\alpha} \alpha_{\alpha} \sum_{\beta} \beta_{\beta}$ ,  $\sum_{\alpha} \alpha_{\alpha} \sum_{\beta} \beta_{\beta}$ .

Now what I was saying earlier was that this beta here can be replaced by alpha, it does not matter, you are just summing over all the coordinates and multiplying with the relative, corresponding velocities. So, I have replaced this beta by alpha, then these two become identical and the sum will have a factor of two which is, so these two together become  $2 \sum_{\alpha} \frac{\delta r_i}{\delta t}$  dot summation over alpha  $\sum_{\alpha} \frac{\delta r_i}{\delta q} \alpha_{\alpha}$  dot. So, this is what I was claiming here, first term no velocities are involved in here but this is a function of q and t, right now I am writing r as a function of q and t.


So, this is a function of q and t but there are no velocities involved in here. So, that is belongs to  $T_{naught}$ . Then this piece has only linear dependence on the generalize velocities. So, that is what is here and so you can read off from here after summing over all the masses and all the particles, your  $T_1$  alpha and this is the one which is going to give contribution to  $T_2$  alpha. So let us see, I will write this down what happened. So, here we have or maybe I can let see if I can do it there itself. Let us, yeah maybe not, let us be neat and clean.



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$$T = \sum_i \frac{1}{2} m_i \left( \frac{\partial \vec{r}_i}{\partial t} \right)^2 \quad \leftarrow T_0$$

$$+ \underbrace{\sum_{\alpha} \left( 2 \sum_i \frac{1}{2} m_i \frac{\partial \vec{r}_i}{\partial t} \cdot \frac{\partial \vec{r}_i}{\partial q_{\alpha}} \right)}_{T_{1\alpha}} \dot{q}_{\alpha}$$

$$+ \sum_{\alpha \neq \beta} \underbrace{\left( \sum_i \frac{1}{2} m_i \frac{\partial \vec{r}_i}{\partial q_{\alpha}} \cdot \frac{\partial \vec{r}_i}{\partial q_{\beta}} \right)}_{T_{2\alpha\beta}} \dot{q}_{\alpha} \dot{q}_{\beta}$$



$$T = T_0 + \sum_{\alpha} T_{1\alpha} \dot{q}_{\alpha} + \sum_{\alpha \neq \beta} T_{2\alpha\beta} \dot{q}_{\alpha} \dot{q}_{\beta} \quad \checkmark$$


proof:  $T = \sum_i \frac{1}{2} m_i \dot{\vec{r}}_i^2$



$$\vec{r}_i = \vec{r}_i(q_1, \dots, q_{3N-k}, t)$$

$$\frac{d\vec{r}_i}{dt} = \frac{\partial \vec{r}_i}{\partial t} + \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial q_{\alpha}} \dot{q}_{\alpha}$$

$$\dot{\vec{r}}_i^2 = \left( \frac{\partial \vec{r}_i}{\partial t} + \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial q_{\alpha}} \dot{q}_{\alpha} \right) \cdot \left( \frac{\partial \vec{r}_i}{\partial t} + \sum_{\beta} \frac{\partial \vec{r}_i}{\partial q_{\beta}} \dot{q}_{\beta} \right)$$

$$= \underbrace{\left( \frac{\partial \vec{r}_i}{\partial t} \right)^2}_{T_0} + \underbrace{\frac{\partial \vec{r}_i}{\partial t} \cdot \sum_{\beta} \frac{\partial \vec{r}_i}{\partial q_{\beta}} \dot{q}_{\beta} + \frac{\partial \vec{r}_i}{\partial t} \cdot \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial q_{\alpha}} \dot{q}_{\alpha}}_{2 \sum_{\alpha} \frac{\partial \vec{r}_i}{\partial q_{\alpha}} \dot{q}_{\alpha}} + \sum_{\alpha \neq \beta} \frac{\partial \vec{r}_i}{\partial q_{\alpha}} \cdot \frac{\partial \vec{r}_i}{\partial q_{\beta}} \dot{q}_{\alpha} \dot{q}_{\beta}$$



So I will write T as, I am just now bringing in the half mi which was remaining and I have to sum over all the particles. So, if you look at what we had earlier, I had dt square. Let see, ri del over del t of this square, so that is the first term which is what your T0 is now. Let us look at the second term now, good, I will multiply this with half mi and sum over all the is. So, you will get summation over i 2 times half of mi delta ri over delta t dot delta ri over delta q alpha and this is the summation over alpha which I can take out and I can put the summation over i in here and put the summation over alpha here, this entire thing with q alpha dot.

So let us check,  $q_\alpha \dot{q}_\alpha$ , this summation I have pulled out and the sigma over  $i$ , I have pulled in. So, this is the piece which is, this piece, this entire thing is your  $T_1$  alpha, good plus the term which is quadratic in velocities and here you will have summation over alpha and beta then summation over  $i$  half  $m_i$  and we had  $\delta r_i$  over  $\delta q_\alpha q_\beta$ , very good,  $q_\alpha \dot{q}_\alpha q_\beta \dot{q}_\beta$  and this is what your, this piece is  $T_2$  alpha beta, let us see what did we call  $T_2$  alpha beta, yes.

So, I have proofed what I was claiming that you have all kinds of terms here and note that these coefficients  $T_1$  alpha and  $T_2$  alpha beta and  $T$  naught, they depend on generalized coordinates and time in general and there may be cases depending on what kind of coordinates you are using that some of the coordinates may not appear and some of velocities, generalized velocities may, will generalized velocities may not appear linearly but those will be specific to those cases. But that is the general result in front of you. We will stop this video here and we will meet next time.