

WAVE OPTICS
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Lecture - 45: Fresnel's half period zone (Cont.)

Hello, students, welcome to the wave optics course. Today, we have lecture number 45 and we will discuss Fresnel's half-period zone and a few more aspects of that. So today we have lecture number 45. So we have already started the discussion of Fresnel's half-period zone. So let me do that once again as a reminder that, suppose we have a wavefront that is propagating and this wavefront can be divided according to Fresnel's theory with few zones like this. Such that if I have a point here P, the distance between the nth zone and the central point, if I write this R naught, it is R naught plus n lambda by 2. That is the basic structure. So what we are doing is that if I want to find out the intensity at this point P, instead of considering the entire wavefront, we can divide this wavefront into a few zones, and these zones are situated or are located in such a way that if I join the nth zone to point P, and if I join the first zone, the central point to P, if it is r naught, if the r naught distance, that is the perpendicular distance from the plane wavefront to the point P, the point of observation P, this is r naught, then it is r naught plus n into lambda by 2. So, with this notation, it is easy to show. So, let me draw these zones once again. So, that we discussed in the last class, but I am doing this once again. So, suppose these are the concentric circles forming the zones, and each circle should have some area. I am shading the alternative zones so that one can understand how this is formed. A This is the way the zones are formed here and this is the distance P and this is R naught. So it is easy to show that the area of each zone is almost the same.

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Lec No = 45

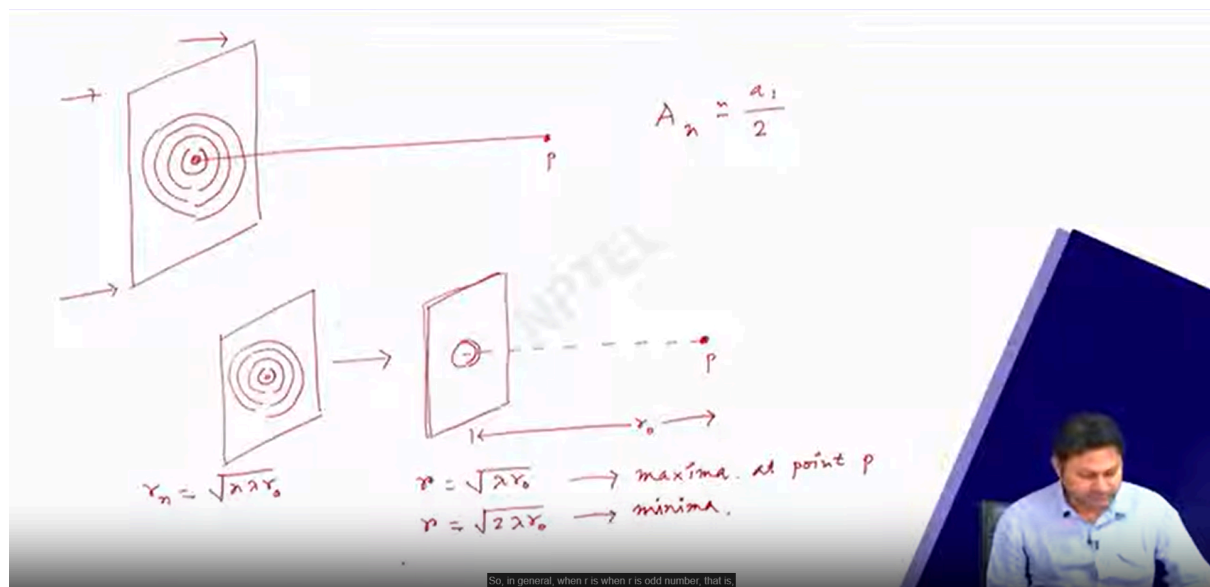
$A_m = \pi \lambda r_0$
 $r_n^2 + r_0^2 = (r_0 + n \frac{\lambda}{2})^2$
 $r_n^2 + r_0^2 = r_0^2 + n r_0 \lambda + \frac{\lambda^2 n^2}{4}$
 $r_n^2 \approx n r_0 \lambda$
 $r_n = \sqrt{n r_0 \lambda}$

And based on the result, what I do is this

So the area, for example, of the mth zone is simply, so yesterday we calculated that the area of the mth zone is roughly pi into lambda into R naught, that is the area of the individual

zone, and it is approximately the same for all the cases because this right-hand side doesn't depend on m . We can also calculate the radius of these individual zones, which we didn't calculate precisely. So today we will do that. So let me draw that. So suppose this is the tip, this is the n th zone I am drawing, and this is the central point. So, according to, this is the n th zone, and this is point P, this is R_0 , and if I join this line from here to here, that value is, say $r_0 + n\lambda/2$. So that is how we define it. So, if I want to find out the radius of the n th zone, then we can also calculate that with very easy steps. And that is, we have an equation. So, suppose the radius here is R_n . So I have simply $R_n^2 + R_0^2$ is equal to $R_0 + n\lambda/2$, a whole square of that or $R_n^2 + R_0^2$ square is $R_0^2 + nR_0\lambda + \lambda^2/4$. Now this R_0^2 , R_0^2 is going to cancel out and essentially what I get is R_n^2 is nearly equal to $nR_0\lambda$ or in other words R_n because R_0 in general is very, very large compared to the wavelength λ . So that is why I can neglect this λ^2 term and essentially R_n is equal to $\sqrt{nR_0\lambda}$. That is an interesting result we have that this is the way the radius is defined for the n th zone. Based on the result, what I do is this: Before that, let me remind you what we had, and that was if I want to find out the intensity and we get a remarkable result here that this is the wavefront that is containing, that is a propagating wavefront, the same figure that I had used. So this is a propagating wave front containing zones like, this my drawing is not with that period scale, obviously I'm drawing this by hand but you should understand the concept and then what happened that if I want to find out what is the intensity or the amplitude at this point P, what we find that the total disturbance that we get due to the entire wavefront that is a n is simply equal to a 1 divided by 2, that is the contribution of the entire wave front is summed up at the point P, and we find it is nothing but whatever the contribution we have for the first zone that is this zone divided by 2, that is a remarkable result. Now what we can do is this, suppose we have a wavefront that is coming, and we have an aperture here.

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This is the wavefront that is coming with this kind of zone, and it is moving, and we also have an aperture here this is the system where we have an aperture, the light is passing

through this aperture, and this is p and suppose here to here this distance is as before r_0 . So now, if I increase the radius of this aperture, what happens is that the p point will be gradually illuminated, and it will be eliminated most when the aperture r will be equal to r_0 because here we know that r_0 is equal to $\sqrt{n \lambda r_0}$ that is the result just derived. So when this aperture radius is exactly equal to $\sqrt{\lambda r_0}$, that means I am only allowing the first zone to pass through this aperture, then we get a maxima here in the point P . So, we get a maxima. Now, interestingly, if I increase this aperture to such that this aperture is allowing not one zone, but two zones, then what happens? This aperture will have a value like $2 \sqrt{\lambda r_0}$. That means I am increasing the aperture to $\sqrt{2} \sqrt{\lambda r_0}$, then here only one zone is allowed to pass through this aperture, but here in the next case, two zones are allowed to pass through, and we know that when there are two zones that are passing through even number of zones, that is passing through the contribution of the amplitude is such that they are canceling each other. And if that is the case, here we will get a minimum. That is a very interesting kind of result, which means if I increase the aperture in general, it suggests that we should get more intensity at point P , but here it is not the case, and that is happening entirely because of the wave nature of these things. So we get a maxima at point P . And if I increase more, then we are going to get a minimum. In a similar way, if I increase more, then we are going to get maxima and minima. So, in general, when r is an odd number, that is, $2n + 1$ multiplied, this is the aperture with aperture radius we are talking about circular aperture. For example, $\sqrt{\lambda r_0}$ then we get maximum and when r is an even number this integer is an even number that is $2n \sqrt{\lambda r_0}$, then we have a minimum here where n can go to $0, 1, 2$, etcetera. So, this is the way we can see that even if we increase the aperture, the value at some point doesn't increase the intensity but it is decreasing and it is changing in an alternative way. We will show how it is changing. In fact, another important thing is if I change the location point of the distance, then we will see a similar kind of change. So, for example, before that, let me also make a few notes. So note 1 is that if I increase the aperture, I discuss this issue; another issue is if instead of having an aperture, I have an opaque region, then what happens? (Refer slide time: 22:32)

$$r = \sqrt{(2n+1)\lambda r_0} \quad \text{Maxima.} \quad (n = 0, 1, 2, \dots)$$

$$r = \sqrt{2n\lambda r_0} \quad \text{Minima.} \quad (n = 1, 2, 3, \dots)$$

Note 1 → If a circular opaque object is placed in such a way that it covers the 1st HPZ. The resultant intensity will be $(\frac{a_2}{2})^2$ and so on.

$$I \propto \left(\frac{a_1}{2}\right)^2$$

We can see that this value will go to change. Accordingly also another thing I like to note and that is note 2, say, so this is the way from that is coming here

If a circular opaque is placed in such a way that it covers the first half period zone, I write in short hpz then the result in intensity will be the result and intensity will be A^2 by 2, we

already calculated earlier, that what should be the resultant intensity square. If everything is open. Then what happened? We have a resulting intensity a_1 by 2 whole squares. Now what I am doing, is the way the wavefront is moving. Let me show the figure. And we place the opaque object here. So, this is the way the other rays are contributed and come to this point. So the intensity of these things is a_1 amplitude, for this, it is a_2 , for this it is a_3 , and so on. So the resultant amplitude, if I cover this, then the intensity will be proportional to the next one which is A_2 divided by 2 squares. Now, if I cover another in such a way that two zones are now covered, then the intensity will be A_3 by 2 whole squares, and so on. So, it will be this and so on. So, in the previous case, we show that we have an aperture and in this case, we have a opaque object that we are placing in front of the wavefront in such a way that it covers the entire first zone and then the intensity I am calculating at this point P and if I increase this point we can see that this value will go to change. Accordingly, also another thing I like to note is note 2, say, so this is the way front that is coming here and over here I have several points, say, p_1 , p_2 , I am moving towards or away from this whatever the aperture we have. Suppose at a distance p on this axis what happened that the aperture transmitted, only the first half period zone. So at this point, if I come forward, then the value of R_0 will change and since the value of R_0 will reduce, we have reconstruction of the zone because the zone area and the zone radius both depend on this value R naught. So, if I move from this point to this point, this value will change. So, if the distance is such that only the first half period zone of the wave is coming here at this point, that should be the value of I_p , at this point, the intensity that will be proportional to a 1 square. In a similar way if I wider the aperture for the same value, if I come and we need to wider that these things but for a wider aperture or for a near point of P_1 on the axis in such a way that lets the axis in the aperture it transmit only the first two half period zone, instead of one. Then what happened? The intensity I_{P1} will be, now proportional to a_1 minus a_2 square. So that is for another point, if I move to another point, the zone will be constructed accordingly because, as I mentioned, the zone construction depends on what is the distance from this point. So I am moving over these points, and accordingly, this intensity is going to change, and you can see that this value is very close to 0. Again, if I move to another point so, this is I_{P1} , I_{P2} , and I_{P3} that is proportional to A_1 minus A_2 plus A_3 , and this is again a high value because A_1 , A_2 will cancel out, and this is a high value. So, the point is if I move along P_1 , P_2 , P_3 , the intensity variation should also come as a, change as a periodic way. Sometimes it is less, sometimes it is large and that is also an interesting aspect of the Fresnel's zone. So if I plot this stuff, so two cases are there one, if I increase this aperture and we keep fixed the detector here which is capturing intensity but we find that there will be a variation of the intensity and this intensity variation is to some extent period. I am just drawing a very rough thing. So, the intensity will go down and up and down and up and down and up something like this. If the radius of the aperture r is changing, if r is increasing the interesting fact is this the intensity at point p , the interesting point is minimal. That means even though we are moving, we are increasing the aperture the intensity is reducing. In a similar way we get a similar kind of figure if I just plot how the intensity is changing, if I change my r naught, that is if I change this distance then also we are supposed to get a variation something like this. So how is this variation, what is the functional form of this variation? It shows that like a sinusoidal kind of curve, but there is also something more to it. So we're going to discuss this in the later part of the lecture. But

here also we have, if you go towards the aperture or go away from this aperture, we will get not only the maxima but also the minima, these periodic variations are there. So today I like to discuss these two aspects of Fresnel's zones. I don't have much time to discuss further an instrument called the zone construction or an instrument through which we can prepare this. So, in the next class, I will try to discuss these zones' construction do detailed calculations, and show how the relationship is there with distance and wavelength and we will find a nice-looking expression as we get in the case of lenses. It acts like a lens where light can be confined to a particular point. So with that note, I like to conclude here. Thank you very much for your attention and see you in the next class.

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$N \propto z^2$
 $I_{P_1} \propto a_1^2$
 $I_{P_2} \propto (a_1 - a_2)^2 \rightarrow \text{Very close to zero.}$
 $I_{P_3} \propto (a_1 - a_2 + a_3) \rightarrow \text{High value.}$

I_0
 r
 I_P
 r_0

So, in the next class, I will try to discuss about these zones construction and do the detailed calculations and show that how the relationship is there with distance and wavelength and we will find a nice looking expression like we get in the case of lens.