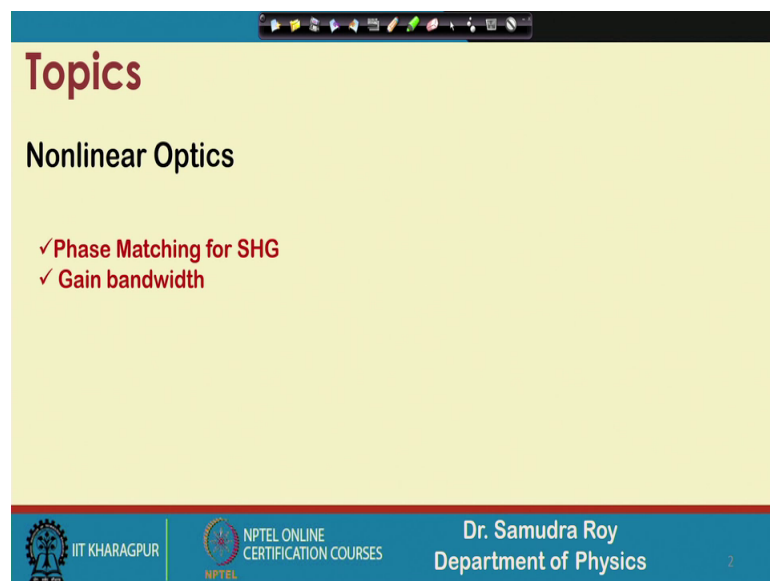


Introduction to Non-Linear Optics and its Applications
Prof. Samudra Roy
Department of Physics
Indian Institute of Technology, Kharagpur

Lecture – 19
Phase Matching of SHG, Gain Band Width Calculation

So, welcome student, to the class of Introduction to Non-Linear Optics and its Applications. So, this is lecture number – 19.

(Refer Slide Time: 00:28)



The slide is titled "Topics" and lists "Nonlinear Optics" as the main subject. Under "Nonlinear Optics", two sub-topics are listed with checkmarks: "Phase Matching for SHG" and "Gain bandwidth". The slide footer includes the IIT Kharagpur logo, the NPTEL Online Certification Courses logo, and the name and department of the lecturer, Dr. Samudra Roy, Department of Physics.

In the previous lecture, so, let us see what we have in this lecture. In the previous lecture, we learn about the second harmonic generation. The mathematics of the second harmonic generation we derived and now, today we will going to learn about the phase matching and then the gain band width of the efficiency curve. So, efficiency is something which also we discussed in the lass class.

So, let us find, what is the meaning of phase matching in second harmonic.

(Refer Slide Time: 00:28)

Evolution equations of E_1 and E_2

$$\frac{\partial E_2}{\partial z} = i \frac{\omega d}{c n_2} E_1^2 e^{-i \Delta k z}$$

$$\frac{\partial E_1}{\partial z} = i \frac{\omega d}{c n_1} E_1^* E_2 e^{i \Delta k z}$$

Conversion efficiency is very low so we can consider E_1 to be constant.

$P_1 \sim 1 \text{ watt}$
 $P_2 \sim 10 \text{ mW}$
 $\eta \propto \frac{P_2}{P_1}$
 $E_2 \ll E_1$
 $\eta \propto \frac{|E_2|^2}{|E_1|^2}$

$E^{(o)}$
 $E^{(2o)}$

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physi

Well, this is our expression that we have derived in the last class. So, once again let me remind you what is the meaning of this. So, this is the evolution of this term is the evolution of the second harmonic and then this term is evolution of the fundamental wave. The evolution of the second harmonic wave which is defined by E_2 is function of z , obviously, that is why the derivative with respect to z in the left hand side gives us something. Here, the function z is appearing we can see clearly, but also E_1 suppose to have function of z .

But, if we say that the efficiency is very low; that means, if I launch a fundamental wave of the order of say 1 watt and we are able to generate. So, this is E_1 of the order of 1 watts, say. Now, if we generate the amount of second harmonic if we generate this of the order of say 10 milli watt, then E_2 is very very less than E_1 and efficiency is efficiency η is proportional to mod of E_2 square divided by E_1 square which is essentially the ratio of the power.

Now, this quantity should be very small. So, here E_1 , I should write not E_1 because it is a unit of watt. So, better to write it is P_1 and P_2 . So, P_1 and P_2 is of the order of 1 watt and 10 milli watt then the efficiency which is proportional to this quantity or efficiency which is of the order of which is proportional to say P_2 divided by P_1 will be very very small. When the efficiency is very small then we can consider this E_1 as a constant because E_1 not going to change much.

(Refer Slide Time: 03:13)

Evolution equations of E_1 and E_2

$$\frac{\partial E_2}{\partial z} = i \frac{\omega d}{cn_2} E_1^2 e^{-i\Delta k z}$$

$$\frac{\partial E_1}{\partial z} = i \frac{\omega d}{cn_1} E_1^* E_2 e^{i\Delta k z}$$

$E_2 \leftarrow E_1$

$E^{(o)}$

 $E^{(2o)}$

Conversion efficiency is very low so we can consider E_1 to be constant.

Dr. Samudra Roy
Department of Physi

Because, the power is coming from E 2 is coming from E one. So, E 1 basically generate gives the power to E 2 and that is why in the source term E 1 is sitting. So, very little amount of power can be transformed from E 1 to E 2. So, we can consider that E 1 is constant. Under such condition we can able to find out what is the E 2 value; that means, directly integrating this equation. When we directly integrate this equation we get something. So, let us find out what we are getting.

(Refer Slide Time: 03:47)

Phase Matching condition

$$\Delta k = 0 \Rightarrow k_2 - 2k_1 = 0$$

$\frac{\partial E_2}{\partial z} = i \frac{\omega d}{cn_2} E_1^2 e^{-i\Delta k z}$

 $\frac{\partial E_1}{\partial z} = i \frac{\omega d}{cn_1} E_1^* E_2 e^{i\Delta k z}$

$\eta(z) = \frac{P_2(z)}{P_1} = \frac{2}{\epsilon_0 n_2 c^3 A} \left(\frac{d\omega}{n_1}\right)^2 P_1 \frac{\sin^2(\Delta k z / 2)}{(\Delta k / 2)^2}$

$\Delta k = 0$

$\eta(z)|_{\Delta k \rightarrow 0} = \frac{P_2(z)|_{\Delta k \rightarrow 0}}{P_1} = \frac{2}{\epsilon_0 n_2 c^3 A} \left(\frac{d\omega}{n_1}\right)^2 P_1 z^2$

Dr. Samudra Roy
Department of Physi

Yeah, so, after integrating we get the value of E^2 which is proportional to z^2 that we have shown in the last class and then we figure what is the value of the efficiency. So, this is basically efficiency defined by η which is the ratio of P_2 by P_1 . I just mentioned in the previous slide which have some value related to this is a quantity this is a constant quantity, which we should not bother about that, but this is the quantity which is very important because the efficiency is entirely depend on that.

Now, what happened if I want to improve this efficiency? We know that if the value of efficiency itself is very small, but still it is a possibility that we can improve this value. So, in order to improve the value of this efficiency η it is better that we can we plot this as a function of Δk . So, Δk is what Δk is the phase mismatch. So, if I plot this entire quantity for a fixed z , if I plot this entire quantity efficiency as a function of Δk then we will get this kind of function which is nothing, but the sinc function. So, it is order of say $\sin^2 \Delta k$ divided by Δk^2 . So, it is something like this. Here, in state of Δk we are plotting $\Delta k z$ by 2 and Δk by 2.

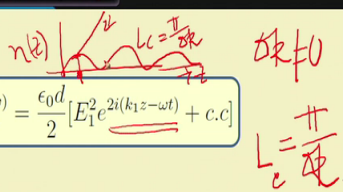
So, now if we plot this things we can see very clearly that we have a maxima here. So, this point the function reaches to a maxima and the other points you can see that there is the efficiency is very less value. Say in order to have the maxima, what is the condition if I see here Δk is 0 is the condition for which we are getting maxima. So, when we get the maxima at Δk equal to 0. So, this is the condition and this is the value of the efficiency and this efficiency we can see is now proportional to z^2 , that we have mentioned, ok. This portion we understood because this is the old thing that we are going to do.

(Refer Slide Time: 06:13)

Phase relationship

$$E^{(2\omega)} = \frac{1}{2}[E_2 e^{i(k_2 z - 2\omega t)} + c.c.]$$

$$P_{NL}^{(2\omega)} = \frac{\epsilon_0 d}{2}[E_1^2 e^{2i(k_1 z - \omega t)} + c.c.]$$

$n(z)$ 

$L_c = \frac{\pi}{\Delta k}$

$$\phi_{P_{NL}^{(2\omega)}} = k_2 z - 2\omega t$$



$$\phi_{E^{(2\omega)}} = 2(k_1 z - \omega t)$$

$$\Delta\phi = \phi_{P_{NL}^{(2\omega)}} - \phi_{E^{(2\omega)}} = (k_2 - 2k_1)z = \Delta k z$$


↓

$$\Delta\phi|_{z=L_c} = (k_2 - 2k_1)L_c = \Delta k L_c = \Delta k \frac{\pi}{\Delta k} = \pi$$

Hence $P_{NL}^{(2\omega)}$ and $E^{(2\omega)}$ are out of phase at $z = L_c$. So after L_c , ω will be generated from 2ω and efficiency will drop.

Dr. Samudra Roy
Department of Physi




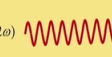
Now, the question is what is the phase relationship between so, in the previous curve we see that electric field is launched here and second harmonics is generating here. So, second harmonic basically is generating because P non-linear 2ω is generating.

(Refer Slide Time: 06:34)

Evolution equations of E_1 and E_2

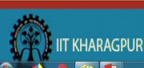
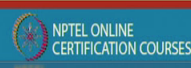
$$\frac{\partial E_2}{\partial z} = i \frac{\omega d}{cn_2} E_1^2 e^{-i\Delta k z}$$

$$\frac{\partial E_1}{\partial z} = i \frac{\omega d}{cn_1} E_1^* E_2 e^{i\Delta k z}$$


$E^{(\omega)}$ 
 $E^{(2\omega)}$ 

$P_{NL}^{(2\omega)} \rightarrow E^{(2\omega)}$

Conversion efficiency is very low so we can consider E_1 to be constant.

Dr. Samudra Roy
Department of Physi



So, this is the reason why E_2 is generating. So, P non-linear 2ω is a quantity which is important because this is the source term for second harmonic. So, now, the next thing is to understand what is the phase relationship between E_2 , which is generated and the corresponding source term that is constantly feeding the energy to E_2 ; that means, P

non-linear is a non-linear polarization term that is vibrating with a frequency 2ω and because of that we are getting the frequency the field with frequency 2ω , but one thing you should note that this is also a travelling wave because of this term.

So, both are travelling wave and we find in the next previous class that if the velocity of these two travelling wave was same, then we will get the efficiency maximum. So, the conversion of energy become maximum and the reason is that. So, though these two waves are propagating together and the velocity phase velocity is matching and if their phase velocity is matching that eventually leads to the phase matching condition $\tilde{k} = 0$.

So, here we will try to find out what is the corresponding phase. So, the phase of this quantity is written as phase $P_{\text{non-linear } 2\omega}$, this is a phase of non-linear polarization term which is nothing, but k_2 this is written wrongly. So, this phase should be this phase. So, this quantity should be $E_{2\omega}$. So, this is the phase of 2ω frequency. So, 2ω frequency basically give rise to this phase $k_2 z$ and the phase of this non-linear polarization the suffix is now reversed here somehow. So, please note that then the phase is $2k_1 z - \omega t$.

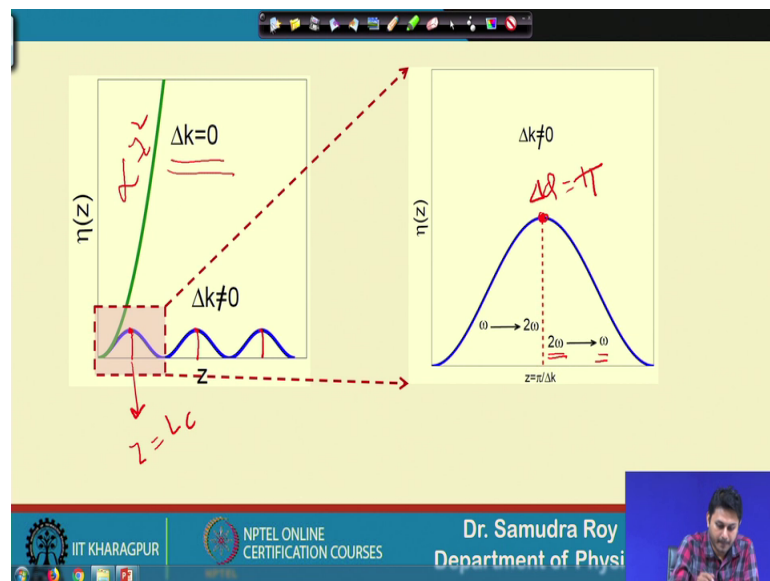
So, now if I want to find out what is the difference between this two phase we will get an interesting result. So, difference between these two phase is nothing, but this minus is this. So, if you do the difference then we will have $k_2 - 2k_1$ which is Δk into z . So, Δk into z is the phase difference between this two.

So, now if we try to find out what is the phase difference at the coherence length; we know what is the coherence length, the coherence length is the length at which we have the efficiency maxima when we plot this things under $\Delta k = 0$ condition and then L_c if I put it here in place of z we will have Δk multiplied by L_c .

Now, we know from the previous lecture that L_c was π divided by Δk . So, at this point at this point we will have the efficiency maximum when $\Delta k \neq 0$, this curve you should remember. So, it was something like this and then one efficiency is going. So, it is efficiency as a function of z . So, this was proportional to z^2 , but this is a sinc function and we are plotting as a function of z here. So, this point here we have L_c which was π divided by Δk , that was the figure we had in the previous class.

So, now, if I want to find out exactly this point what is the phase difference between this two we find that the phase difference is exactly π . So, π phase difference is there at this point where the efficiency is maximum Δk equal to 0. Now, after that what happened, if I go beyond that then we find that the efficiency is now dropping down. So, that means, P NL and E are out of phase exactly at the point z equal to L_c . So, after L_c what happened ω will not going to generate 2ω rather ω will be generated from 2ω . So, that is why the efficiency will go down.

(Refer Slide Time: 11:35)



So, let us go to the next slide to understand the fact this is the figure and this figure is basically gives you the entire scenario what is going on. So, here the efficiency versus z curve and this is basically the condition this curve is proportional to z square and it is happening when Δk is equal to 0 that we have already figured out. But, this is the more realistic case when Δk not equal to 0, the efficiency will be not that much, but we still have some value where efficiency reaches to a maxima; for example, these are the points.

Now, these points here when efficiency is maxima it is called the coherence length. So, at this particular coherence length what happened, the efficiency reaches to a maxima. So, from this figure if I now expand this portion only then we find that x this goes to a maxima at this particular point and then what happened it is go down.

So, here the exactly phase mismatch $\Delta\phi$ was π . So, the phase of the propagating non-linear polarization term and the phase of the electric field having frequency 2ω component is now entirely out of phase and as a result what happened? The efficiency goes down because now, 2ω is generating from ω is generating from 2ω ; exactly the opposite phenomena is happening because their phase is now reverse, that is why the efficiency is going down, ok.

(Refer Slide Time: 13:18)

Quantum Mechanical Picture

Two photons of frequency ω merge to form a single photon of frequency 2ω .

Energy conservation $\hbar\omega + \hbar\omega = 2\hbar\omega = \hbar(2\omega)$

Momentum conservation $\hbar k_1 + \hbar k_1 = \hbar k_2$

$2k_1 = k_2$
(phase matching condition)

Vector Diagram
(Collinear phase matching)

The slide features a diagram on the left showing two energy levels with two upward arrows labeled $\hbar\omega$ merging into a single downward arrow labeled $\hbar(2\omega)$. On the right, a vector diagram shows two horizontal arrows labeled $\hbar k_1$ pointing right, and a single horizontal arrow labeled $\hbar k_2$ pointing right below them. The slide footer includes the IIT Kharagpur logo, NPTEL Online Certification Courses text, and Dr. Samudra Roy, Department of Physics.

After that what we have, the quantum mechanical picture of the entire system which is important. So, we know that in second harmonic generation ω frequency from ω frequency we are getting 2ω frequency. So, that means, from energy conservation we can have that two photon, this is a photon picture, two photon can merge to generate one photon of frequency ω . So, 2 photon of frequency ω and ω where the energy is $\hbar\omega$ and $\hbar\omega$, if I add this two we will have $2\hbar\omega$. So, that means, two photons are merging together to generate a photon of 2ω frequency.

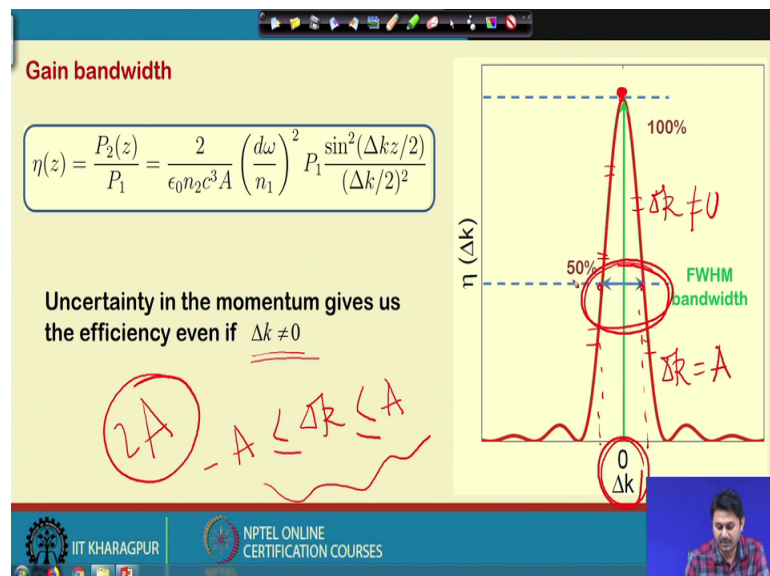
So, this is the quantum mechanical picture, where we can see that one photon and another photon is merged to generate another 2ω photons. So, this is the photon with frequency 2ω . In the similar way if I try to find out the momentum conservation then we can say that the momentum of two photon which is $\hbar k_1 + \hbar k_1$, where k_1 is the propagation conserve propagation constant of the

photon having frequency ω and then the momentum conservation suggest that if this two are added then we are getting another momentum $\hbar k_2$, which is the momentum of the wave that is having frequency component 2ω .

Now, if I simplify this or if I just cut this \hbar altogether \hbar cross then whatever the expression we have is nothing, but the phase matching equation. So, the phase matching condition is nothing, but the conservation of the momentum. So, when the phase matching condition is there we can say the momentum is conserved and when the momentum is conserved or under that condition we have greater efficiency.

So, in vector diagram we can write this things in this way. So, one photon with $\hbar k_1$ momentum and another photon which $\hbar k_2$ momentum, they add up to generate another photon which having the propagation constant k_2 . This is the collinear kind of phase matching because all this vectors are in similar direction. But, in the future class we will find that there are non collinear phase match is still possible where more than two waves are associated to generate a third wave. So, three wave mixing process are there. Then we will have a different condition of this phase matching in terms of this vector picture or this vector diagram well.

(Refer Slide Time: 16:10)



After having the knowledge of all this things, now it is time to understand what is the meaning of this gain bandwidth. So, what so far we find that efficiency reaches to a maxima when Δk equal to 0, that is the thing we have already find.

So, this is the point where efficiency reaches to a maxima under the condition that Δk equal to 0. So, Δk equal to 0 or in absolute phase matching condition we have a maxima. Now, the question is if I plot this things as a function of Δk we can see that even if Δk is not equal to 0, for example, this point, this point, this point, this point or this point, where Δk is not equal to 0 in all the cases Δk is not equal to 0, we still have some amount of efficiency.

So, Δk not equal to 0, still give some amount of efficiency. Obviously, the efficiency is going down very sharply both the sides of Δk around Δk equal to 0, but the thing is that Δk equal to 0 is not the only condition to have the efficiency. Efficiency still is there when Δk not equal to 0 and the reason behind that the un uncertainty in the momentum basically gives the efficiency even Δk not equal to 0, because there is quantum mechanics there is a uncertainty.

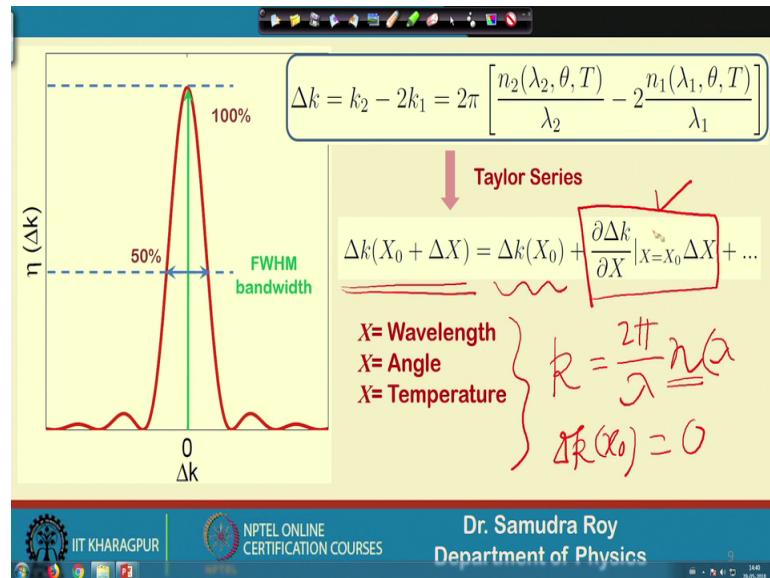
So, Δk equal to 0 condition is not the prime condition to have a second harmonic you can still have second harmonic when Δk not equal to 0. So, that means, there is a band width of this efficiency curve. So, we have a certain amount of the range of Δk . So, if I say Δk is less than equal to some value A and minus A so, in between this two values which is $2A$. So, in between this two values if I say this is this quantity where we have the 50 percent efficiency, I have Δk equal to say some value called A .

So, now the question is how do we find out this bandwidth and why this bandwidth is generated as I mentioned quantum mechanically if Δk not equal to 0, then still we have some kind of efficiency, but classically also we find out the functional form in such that even Δk not equal to 0, give us some kind of efficiency. So, now, we measure this bandwidth, now width can be measure in a different way. So, we consider full with half maxima.

So, what is a full with half maxima? So, we have a bell kind of curve, this is a bell shaped some sort of bell shaped curve and here if this is the maxima then the half of the maxima is around this point. So, half of the maxima we have, now try to find out what is a width of that point that particular point. So, that is why it is called full width at half maxima. This is the maxima that means, the 50 percent efficiency and try to find out what is the width of this thing. So, we define fully the half maxima in such a way that we have some kind of width of Δk .

So, delta k is a width here in this point we have some amount of width for which we are getting the efficiency and this efficiency is around 50 percent of the main peak efficiency.

(Refer Slide Time: 20:11)



Well, which of the parameters on which this things is depend on? That is a important quantity, that is a important thing to understand. So, let us first see what is delta k? Delta k is k 2 minus k 1; k 2 is the propagation constant. If I write this propagation constant in terms of refractive index then it should be simply. So, k is 2 pi by lambda into refractive index, in free space n is 1. So, we have k equal to 2 pi by lambda, but if the space has some kind of refractive index if the medium has some kind of refractive index then we can write k equal to 2 pi by lambda into n.

Now, this refractive index n may be function of many things. First of all it is a function of lambda no doubt about that. So, whatever the lambda I am launching, so that means, I have to launch very precise lambda. So, lambda 2 and lambda 1 if I say these are the fixed quantity still n 2 can be varied with temperature and the angle at which I am launching. So, theta here is a angle and temperature which is a very popular parameter to change the refractive in the specially in the crystal.

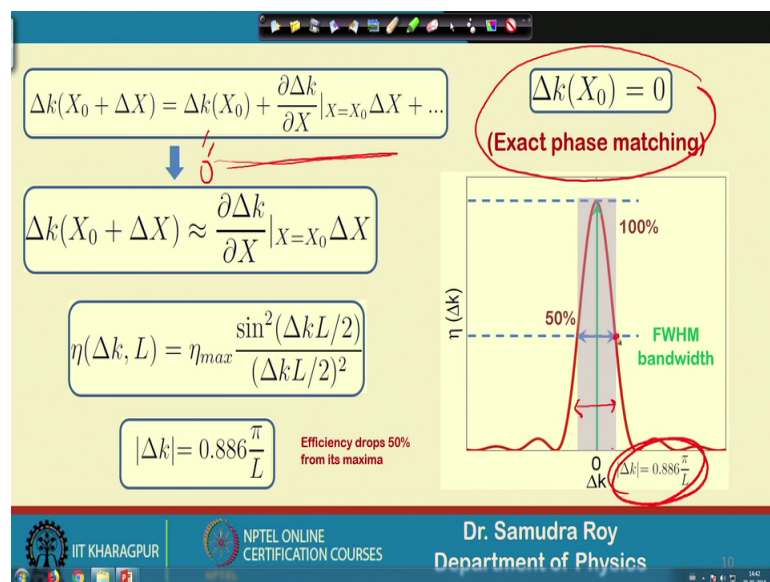
So, these are the parameters which are depending I mean which basically because on I mean this parameter basically n 2 is depending on. So, now, if I want to change want to find out what is the variation based on this parameter we can say delta k will vary when

this one of this parameter will vary. So, this x can be any parameter, may be it is a wavelength, may be it is an angle or temperature.

So, now Δk will go to vary with a small change of Δx , then what happened? If I now make a Taylor series we call this is a small amount Δx is a small amount if I make a Taylor series we have $k(X_0)$; X_0 is exact value for which the phase matching is there. So, $\Delta k(X_0)$ we can safely say this is a 0, because this is the absolute phase matching condition. All the condition is valid here so, it is an absolute phase matching condition or x is 0 is a value for which the phase matching is there.

But, if we have some kind of deviation from x_0 then we have a next higher order term which is this $\Delta k(X) \approx \frac{\partial \Delta k}{\partial X} \bigg|_{X=X_0} \Delta X$; that means, this is the term which basically the measurement of the amount of deviation of the Δk . So, how to find this term so, we need to know how the Δk is changing as a function of X ? X may be wavelength angle or temperature as I say. So, how this things is changing and when this things are changing this rate and if I multiplied with ΔX , we can find the total amount of change Δk because of the change of ΔX .

(Refer Slide Time: 23:32)



So, next we will go to find that what is basically the value. So, this is the amount of value we have for Δk expression the value of Δk as a function of X_0 plus ΔX by expanding this as a Taylor series. As I mentioned for exact phase matching condition

delta k X 0 has to be 0, see if I put this term equal to 0, then the rest term is something like this.

Now, we know there is a bandwidth of the efficiency. So, I can figure out what is my delta k in this shaded region. If I calculate we find that my delta k is this value. At this value it is not equal to 0, mind it this delta k is not equal to 0, but this delta k is 0.886 for sinc function this is the value you can, where the efficiency drops 50 percent from its maxima. So, if I change delta k from 0 to some value this value this efficiency reach to that point, in the negative side also it will reach to this point. So, that means, my delta k here is some value where my width is something.

So, in order to find the width I need to calculate, what is this delta k. So, delta k is this value. So, what we will do? We now have the delta k in our hand. So, delta k is the amount that is changing. So, I can put this delta k here. So, let us go to the next slide may be that will be easier for understand.

(Refer Slide Time: 25:15)

$$\Delta k(X_0 + \Delta X) \approx \left. \frac{\partial \Delta k}{\partial X} \right|_{x=X_0} \Delta X$$

$$|\Delta k| = 0.886 \frac{\pi}{L}$$

$$|\Delta X| = 0.886 \frac{\pi}{L \left(\left| \frac{\partial \Delta k}{\partial X} \right|_{x=X_0} \right)}$$

$$|\Delta X|_{FWHM} = 0.886 \frac{2\pi}{L \left(\left| \frac{\partial \Delta k}{\partial X} \right|_{x=X_0} \right)}$$

Handwritten notes:
 $x_0 + \Delta X$
 $L = \text{Given}$
 $\frac{\partial \Delta k}{\partial X} \Big|_{x=x_0}$

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physics

Well, this is the slide we have. So, this is the quantity I want to find out this is basically the tolerance that up to which delta X, how far the parameter X can be change. So, X 0 is the exact condition for phase matching. Now, we change this parameter to delta X amount. Now, the question is how far we can extend this delta X. So, if I now calculate from this. So, delta X has to be delta k is this amount. If I put delta k equal to 0.006 pi by L then 0.006 pi by L divided by this quantity basically gives you my delta X.

So, ΔX tolerance or ΔX 4 with half maxima is this amount. Now, if L value is given Δk , the rate of change of Δk with respect to the given parameter X at $X = X_0$ point is given if these two terms are given, then we can readily find out what is the tolerance limit of ΔX . ΔX is the value which can change the amount and now, we can find exactly if these two values are given that what should be the tolerance of ΔX parameter which may be angle, which may be temperature and which may be wavelength.

So, with this note so, let me conclude this class. So, in the next class we will extend this concept and try to find out more on second harmonic generation and how to increase the efficiency and all these issues. So, with that thank you very much and see you in the next class.