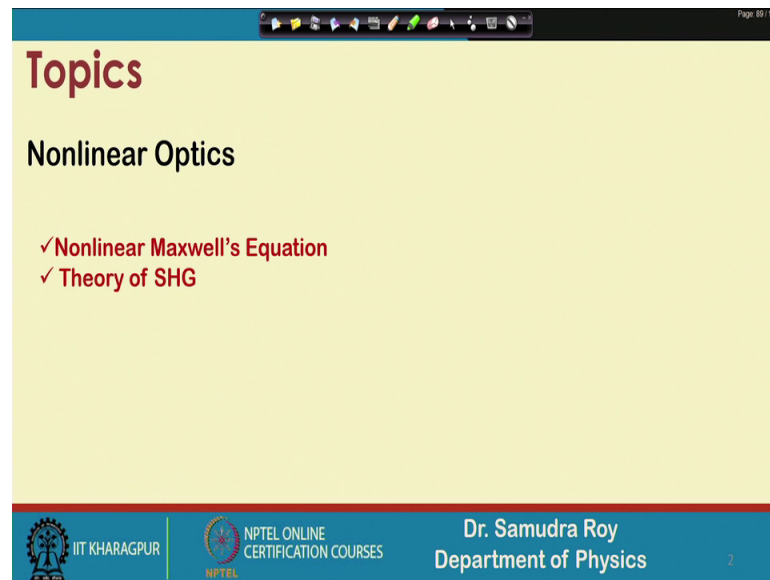


Introduction to Non-Linear Optics and its Applications
Prof. Samudra Roy
Department of Physics
Indian Institute of Technology, Kharagpur

Lecture – 16
Nonlinear Maxwell's Equation

So, welcome back students to the next class of Introduction to Non-Linear Optics and its Application course. So, in the previous class, we learned about the sum and different frequency generation and also we know how the second harmonics are generated. Today is the lecture number 16 and let us see what we have in today's lecture.

(Refer Slide Time: 00:38)



Topics

Nonlinear Optics

- ✓ Nonlinear Maxwell's Equation
- ✓ Theory of SHG

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy
Department of Physics

7

So, today we have 2 very important concept; one is non-linear Maxwell's equation and we will after having after having the knowledge of non-linear Maxwell's equation, we going to learn the theory of second harmonic generation. So far, roughly we understand how the second harmonics are generated. So, when we launch an electric field, because of the frequency mixing, different frequency components are generated. And if we launch one electric field with frequency ω , then we find that we can generate frequency 2ω which is called the second harmonic generation. But the detail theory is still missing because we do not know exactly what is happening inside the medium and how the second harmonics are evolving, how the second harmonic fields are evolving are still quite hazy.

So, in order to understand the detailing of the second harmonic generation or the basic physics or what is going on inside the material, it is important to know, what is the meaning of non-linear Maxwell's equation. Well, the Maxwell's equation is the governing equation the main governing equation of the electric field. So, when the electric field is inside the medium, what happened that it will interact with the medium and if the interaction is non-linear in nature, inside the Maxwell's equation, we have some kind of non-linear components. And this non-linear components is nothing, but the non-linear polarization that we have been learning for last few classes.

So, let us try to find out how to how to calculate this non-linear Maxwell's equation or how to derive this non-linear Maxwell's equation.

(Refer Slide Time: 02:24)

Nonlinear Maxwell's Equation

Handwritten red annotations on the left side of the slide:

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}$$

$$\vec{D} = \epsilon \vec{E} + \vec{P}_{NL}$$

Equations shown on the slide:

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}$$

$$\vec{P} = \vec{P}_L + \vec{P}_{NL}$$

$$\vec{P}_L = \epsilon_0 \chi^{(1)} \vec{E}$$

$$\vec{D} = \epsilon_0 \vec{E} + \epsilon_0 \chi^{(1)} \vec{E} + \vec{P}_{NL}$$

$$\vec{D} = \epsilon_0 (1 + \chi^{(1)}) \vec{E} + \vec{P}_{NL}$$

$$\vec{D} = \epsilon \vec{E} + \vec{P}_{NL}$$

Page: 88 / 148

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physics

The derivation is quite straightforward and it is something like the way we generate Maxwell's wave equations exactly we will generate in the similar way the non non-linear Maxwell's equation or non-linear Maxwell's wave equation. So, let us start with our fundamental expressions that is D, which is epsilon 0 plus P displacement vector is epsilon 0 E plus P. Now it is not necessary that D and E should be in the same direction because of this polarization can be a function of electric field. But if I launch the electric field that not ensured that the polarization will be in the same direction if the medium specially if the medium is an isotropic in nature, then there is a I mean if I launch an electric field one direction, then the polarization will be another direction.

So, we will come to this point later, but here more importantly we need to introduce the non-linear polarization. So, now, the polarization can be divided into 2 part; one is the linear one and another is the non-linear part. The linear part is introduced in this way and it is proportional to the electric field as usual. And if that is the case, then the D can be represented as $\epsilon_0 E$ plus $\epsilon_0 \chi^{(1)} E$ which is a linear term and then P non-linear term. We just left the P non-linear term as P^{nl} vector, ok. After having the value of D , now it is time to write in compact way, and if I do, then it should be $\epsilon_0 (1 + \chi^{(1)}) E$ and we know that $\epsilon_0 (1 + \chi^{(1)})$ is nothing but ϵ_0 and P non-linear will be there.

So, if I look from this equation to this equation, what essentially is changing? D was $\epsilon_0 E$ plus P this P is a P total. Now I combine this P linear term here and in this P linear term, we have the information of the refractive index that is the material information, that is in we basically inserting that material information into the system through refractive index and that is why, this ϵ_0 become ϵ and since a non-linear part is associated with this.

So, this non-linear part will be left alone and it will be simply that. So, from here to here, we make such certain modification and the modification is essentially we just separate out the non-linear polarization term and the rest part is same. If non-linear polarization is 0, then D equal to $\epsilon_0 E$ is the standard form that we usually find when we calculate the Maxwell's equations and the corresponding derivations ok.

(Refer Slide Time: 05:56)

Nonlinear Maxwell's Equation

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}$$

$$\vec{P} = \vec{P}_L + \vec{P}_{NL}$$

$$\vec{P}_L = \epsilon_0 \chi^{(1)} \vec{E}$$

$$\vec{D} = \epsilon_0 \vec{E} + \epsilon_0 \chi^{(1)} \vec{E} + \vec{P}_{NL}$$

$$\vec{D} = \epsilon_0 (1 + \chi^{(1)}) \vec{E} + \vec{P}_{NL}$$

$$\vec{D} = \epsilon \vec{E} + \vec{P}_{NL}$$

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physics

Ok.

(Refer Slide Time: 06:03)

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}; \quad \vec{\nabla} \times \vec{H} = \frac{\partial \vec{D}}{\partial t}$$

$\vec{B} = \mu_0 \vec{H}$

$$\vec{\nabla} \times \vec{\nabla} \times \vec{E} = -\frac{\partial(\vec{\nabla} \times \vec{B})}{\partial t} = -\mu_0 \frac{\partial(\vec{\nabla} \times \vec{H})}{\partial t} = -\mu_0 \frac{\partial^2 \vec{D}}{\partial t^2}$$

$$\vec{\nabla}(\vec{\nabla} \cdot \vec{E}) - \nabla^2 \vec{E} = -\mu_0 \frac{\partial^2 \vec{D}}{\partial t^2} = -\mu_0 \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} - \mu_0 \frac{\partial^2 \vec{P}_{NL}}{\partial t^2}$$

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\nabla^2 \vec{E} - \mu_0 \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} = \mu_0 \frac{\partial^2 \vec{P}_{NL}}{\partial t^2}$$

$\frac{\partial^2 \vec{D}}{\partial t^2} = \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$

$\vec{D} = \epsilon \vec{E} + \vec{P}_{NL}$

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physics

After having the knowledge of D with non-linear polarization as shown here, in the this is the expression that we figure out. Now we start deriving the Maxwell's equation with this condition with this condition. So, now, again we will going to use 2 Karl equation. So, Karl of E is minus del B del t. Karl of H is minus del D del t. So, now, if I take Karl both the side from this equation, so, Karl of B will be this and if I take Karl of H both the Karl of H both the cases. So, I use of Karl of ah if I from this first equation, if I do Karl

of B, it will be curl of B and this curl of B I use this one and then it will be curl of H with μ_0 .

Because, B is $\mu_0 H$ and then curl of H I will go to use this term. So, it will be the second derivative with respect to second order derivative with respect to time and D. So, this curl of curl of E can be represented in a standard form we know it is a grad of it is a divergence grad of divergence of E minus of grad square E and then the left hand side right hand side is minus epsilon 0 D 2 D d t square. This term is now represented into 2 part; one is this and another is this, because my D is epsilon 0 E plus P non-linear.

So, when I write double derivative of D essentially in the right hand side I need to write this is not this is epsilon, ok. So, let me clear it and do it once again.

(Refer Slide Time: 08:03)

The slide contains the following mathematical content:

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}; \quad \vec{\nabla} \times \vec{H} = \frac{\partial \vec{D}}{\partial t}$$

$$\vec{\nabla} \times \vec{\nabla} \times \vec{E} = -\frac{\partial(\vec{\nabla} \times \vec{B})}{\partial t} = -\mu_0 \frac{\partial(\vec{\nabla} \times \vec{H})}{\partial t} = -\mu_0 \frac{\partial^2 \vec{D}}{\partial t^2}$$

$$\vec{\nabla}(\vec{\nabla} \cdot \vec{E}) - \nabla^2 \vec{E} = -\mu_0 \frac{\partial^2 \vec{D}}{\partial t^2} = -\mu_0 \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} - \mu_0 \frac{\partial^2 P_{NL}}{\partial t^2}$$

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\vec{D} = \epsilon \vec{E} + \vec{P}_{NL}$$

$$\nabla^2 \vec{E} - \mu_0 \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} = \mu_0 \frac{\partial^2 P_{NL}}{\partial t^2}$$

Handwritten red annotations show the expansion of the double derivative of D:

$$\frac{\partial^2 \vec{D}}{\partial t^2} = \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} + \frac{\partial^2 P_{NL}}{\partial t^2}$$

So, I am making the double derivative of D d is this. So, it will right hand side it is simply the double derivative of E plus the double derivative of P n l. So, I making if I use this it will be simply this quantity. So, now, grad dot E is 0. So, this term will cancel out. So, eventually we have these is equal to the right hand side. So, these is the equation that is given in inside the border. We have one very important aspect we find in this equation that in normally this term is not there.

(Refer Slide Time: 08:52)

Slide 82/152

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}; \quad \vec{\nabla} \times \vec{H} = \frac{\partial \vec{D}}{\partial t}$$

$$\vec{\nabla} \times \vec{\nabla} \times \vec{E} = -\frac{\partial(\vec{\nabla} \times \vec{B})}{\partial t} = -\mu_0 \frac{\partial(\vec{\nabla} \times \vec{H})}{\partial t} = -\mu_0 \frac{\partial^2 \vec{D}}{\partial t^2}$$

$$\vec{\nabla}(\vec{\nabla} \cdot \vec{E}) - \nabla^2 \vec{E} = -\mu_0 \frac{\partial^2 \vec{D}}{\partial t^2} = -\mu_0 \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} - \mu_0 \frac{\partial^2 \vec{P}_{NL}}{\partial t^2}$$

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\vec{D} = \epsilon \vec{E} + \vec{P}_{NL}$$

$$\nabla^2 \vec{E} - \mu_0 \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} = \mu_0 \frac{\partial^2 \vec{P}_{NL}}{\partial t^2}$$

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physics

When we consider everything is linear or the polarization is a linear polarization, we have this is equal to 0 and then we have some kind of solution because this is a plane wave the plane wave solution is there. So, we have already done in our basic ah basic linear optics course. So, now, in state of having 0 term what we have is a source term.

(Refer Slide Time: 09:23)

Slide 83/152

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}; \quad \vec{\nabla} \times \vec{H} = \frac{\partial \vec{D}}{\partial t}$$

$$\vec{\nabla} \times \vec{\nabla} \times \vec{E} = -\frac{\partial(\vec{\nabla} \times \vec{B})}{\partial t} = -\mu_0 \frac{\partial(\vec{\nabla} \times \vec{H})}{\partial t} = -\mu_0 \frac{\partial^2 \vec{D}}{\partial t^2}$$

$$\vec{\nabla}(\vec{\nabla} \cdot \vec{E}) - \nabla^2 \vec{E} = -\mu_0 \frac{\partial^2 \vec{D}}{\partial t^2} = -\mu_0 \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} - \mu_0 \frac{\partial^2 \vec{P}_{NL}}{\partial t^2}$$

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\vec{D} = \epsilon \vec{E} + \vec{P}_{NL}$$

$$\nabla^2 \vec{E} - \mu_0 \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} = \mu_0 \frac{\partial^2 \vec{P}_{NL}}{\partial t^2}$$

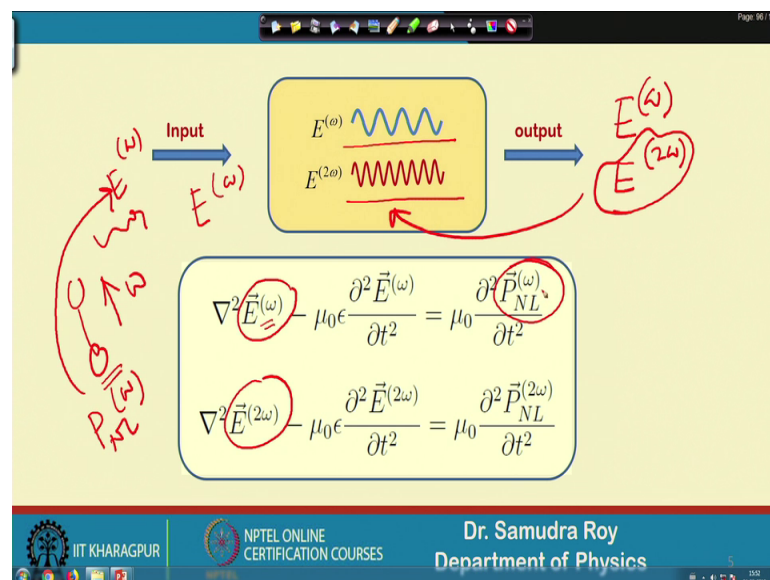
IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physics

This is basically is source term. This P non-linear term basically is source term of the electric field that mean, whatever the electric field will be there inside the medium in this differential equation we have a source term. And this source term is essentially coming

with a fact that mine P non-linear is there. So, P non-linear is basically feeding some kind of energy to generate electric field and that is the same this picture we have already shown with a schematic diagram, that when the dipoles are vibrating in a non-linear fashion, you will start generating some kind of electric field.

And this electric field will have a different frequency component exactly the similar concept we have through this equation, where we find after deriving the Maxwell's equation considering the non-linear polarization term. We have something like this a source term.

(Refer Slide Time: 10:26)



So, the next thing is what happened when we have the electric field do the system having frequency components. So, we have already derived the non-linear Maxwell's wave equations. So, this is the picture. We will launch an electric field and we are we will get we put some kind of input and we will get output outside. But in the system inside the system what happened that we have 2 frequency components.

(Refer Slide Time: 10:54)

The propagating electric fields

$$E^{(\omega)} = \frac{1}{2} [\tilde{E}_1 e^{i(k_1 z - \omega t)} + c.c.]$$

$$E^{(\omega)} = \frac{1}{2} [\tilde{E}_1 e^{i\phi_1} e^{i(k_1 z - \omega t)} + c.c.]$$

$$\tilde{E}_1 = \tilde{E}_1 e^{i\phi_1} \text{ (Complex amplitude)}$$

$$E^{(\omega)} = \frac{1}{2} [E_1 e^{i(k_1 z - \omega t)} + c.c.]$$

$$E^{(2\omega)} = \frac{1}{2} [E_2 e^{i(k_2 z - 2\omega t)} + c.c.]$$

Diagram showing two sine waves: $E^{(\omega)}$ (blue) and $E^{(2\omega)}$ (red).

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physics

So if I launch, if I launch electric field E of ω , if I consider only the second harmonic generation, then inside this system is second harmonic field will generate. So, output; obviously, we will have E of ω and E of 2ω so; that means, E of 2ω will going to be here somewhere. So, in this figure, we find that inside the material which is non-linear in nature. We have the fundamental wave and also the second harmonic both are present.

So, for both the cases the non-linear Maxwell's equation of the non-linear wave equation is valid. So, for 2 waves we have 2 equations. Basically, the equation is same; the electric field component E of ω and electric field component 2ω are different. But the equation is same for electric field component E of ω what happened the source term is P non-linear ω as I mentioned. So, when we launch the electric field, what happened the dipole? If this is the dipole, it will going to vibrate at frequency ω and start generating E field of ω .

So, that is why, the source term is nothing, but the P non-linear which is vibrating at the frequency ω . So, that is the source of the electric field which is generating and the frequency of the electric field as same as the frequency at which the polarization is vibrating. Here, the polarization is vibrating with the frequency ω non-linear polarization and as a result, we are getting some kind of electric field which is also vibrating with a frequency ω exactly in the similar way.

(Refer Slide Time: 12:51)

The slide illustrates the generation of a second-harmonic electric field. It shows an input consisting of two electric fields, $E^{(\omega)}$ and $E^{(2\omega)}$, entering a system. The output is an electric field $E^{(2\omega)}$. Below this, two wave equations are presented:

$$\nabla^2 \vec{E}^{(\omega)} - \mu_0 \epsilon \frac{\partial^2 \vec{E}^{(\omega)}}{\partial t^2} = \mu_0 \frac{\partial^2 \vec{P}_{NL}^{(\omega)}}{\partial t^2}$$

$$\nabla^2 \vec{E}^{(2\omega)} - \mu_0 \epsilon \frac{\partial^2 \vec{E}^{(2\omega)}}{\partial t^2} = \mu_0 \frac{\partial^2 \vec{P}_{NL}^{(2\omega)}}{\partial t^2}$$

Handwritten red annotations include a dipole diagram with a circle and an upward arrow labeled (2ω) , and labels P_{NL} and $E^{(2\omega)}$ with arrows pointing to the corresponding terms in the equations.

We can say that, if this is my dipole, suppose and it is now vibrating with a frequency 2ω and when it is vibrating with a 2ω , my non-linear polarization can be represented 2ω . This is the way it is vibrating and as a result, it is generating a frequent generating electric field with the same frequency 2ω . So, that means, the generation of 2ω is essentially supported by the non-linear polarization which is having which is having a component of 2ω .

Now, if I look to this differential equation, in this differential equation in the source term, both the cases we have P non-linear ω and P non-linear 2ω , 2 source term are there. This source term basically responsible for generating electric field with ω . This source term is important for generating the electric field 2ω , but the equation is same for both the cases; only the components are different. So, in the next step, that is a very important to find out what should be this value and what should be this value so that, we can solve at the end of the day we can solve our Maxwell's equation ok.

The electric field, we know very clearly what should be the electric field because; we have been using the form of this electric field for long. So, electric field with frequency component ω we just use as $E = E_0 e^{i(kz - \omega t)}$ plus complex conjugate of that with a half term, so that, it this is real this is the normal way to represent the electric field and that is the frequency component of E ω , but here one

thing you should remember that this E_1 and this is a phase, but some additional phase should also be there.

So, here we write E_1 tilde which is a real quantity and then E to the power $i k_1 z - \omega t$ that is the component related to the plane wave and on top of that we have some additional phase which is ϕ_1 . So, this ϕ_1 can be represented as E to the power $i \phi_1$ multiplied by E_1 and write only the $k_1 z - \omega t$ in a separate way. So, if I do then I can write E_1 is E_1 tilde E to the power of $i \phi_1$.

So, now, this is complex quantity. So, this is amplitude, but normally the concept is amplitude is a always real quantity, but here we absorb the phase additional phase term to make it more general and find that the amplitude is now complex in nature. So, that in the future, we may have the term like E_1^* this is the term which basically gives you the idea that E_1 which is essentially the amplitude may have some phase associated with that. So, that is why it is a complex.

(Refer Slide Time: 16:28)

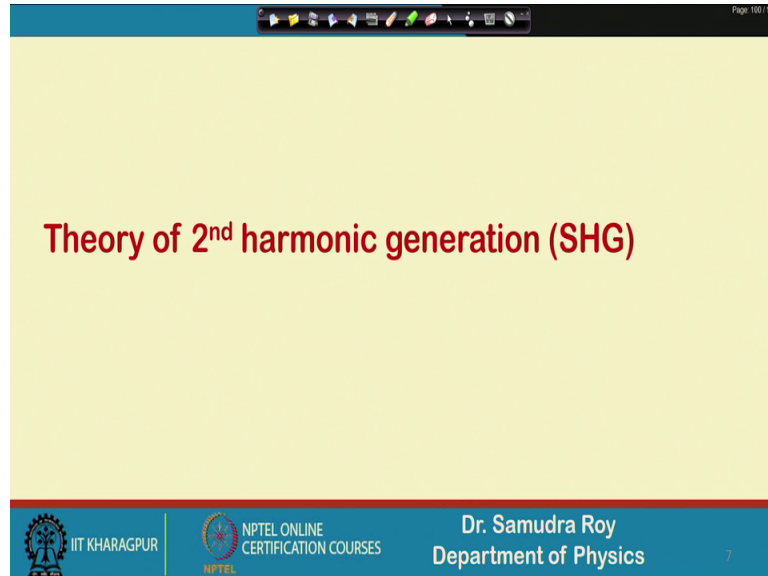
The slide is titled "The propagating electric fields". It features a man in a pink shirt in the foreground. The background slide contains the following content:

- Two wave diagrams: $E^{(\omega)}$ (blue wavy line) and $E^{(2\omega)}$ (red wavy line).
- Equations on the left: $E^{(\omega)} = \frac{1}{2} [\tilde{E}_1 e^{i(k_1 z - \omega t)}]$ and $E^{(2\omega)} = \frac{1}{2} [\tilde{E}_2 e^{i(k_2 z - 2\omega t)}]$.
- Equations on the right, enclosed in a red box: $E^{(\omega)} = \frac{1}{2} [E_1 e^{i(k_1 z - \omega t)} + c.c.]$ and $E^{(2\omega)} = \frac{1}{2} [E_2 e^{i(k_2 z - 2\omega t)} + c.c.]$.
- Bottom right text: "Dr. Samudra Roy, Department of Physics".

So, simple thing here, how to find out what is my E_{ω} and $E_{2\omega}$. So, in the slide we can find that E_{ω} is nothing but $E_1 E$ to the power of $k_1 z - \omega t$ $E_{2\omega}$ E to the power of $k_2 z - 2\omega t$ $E_2 E$ to the power $k_2 z - 2\omega t$ because this is 2ω . So, that is why, 2ω frequency is there k_1 is a propagation constant for the frequency component ω or the fundamental wave and k_2 is the frequency component of second harmonic. So, these E_{ω} and $E_{2\omega}$ are the

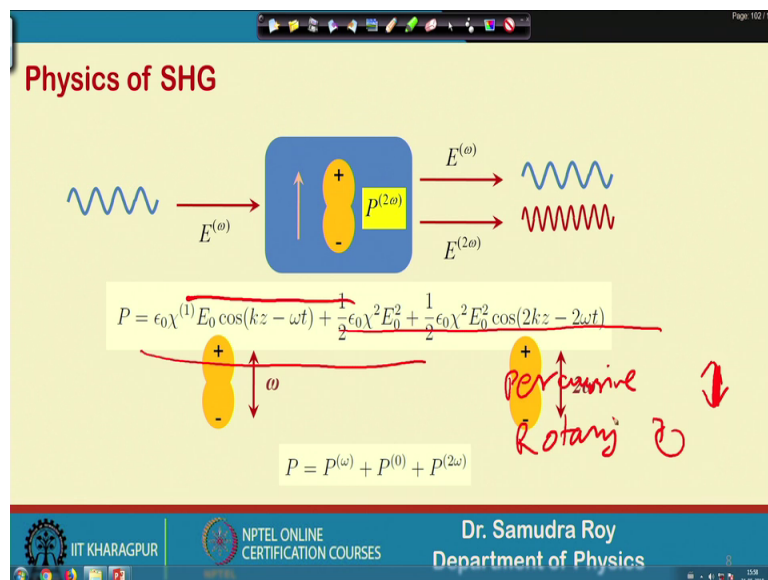
electric field and this electric field has this kind of forms remember that E_1 and E_2 are now complex amplitude in general ok.

(Refer Slide Time: 17:12)



Now, we going to learn as I mentioned the theory of second harmonic generation.

(Refer Slide Time: 17:18)



So, if I go back to the picture, as I mention that this picture, we have already shown in our previous classes that if I launch an electric field of $E \omega$, the polarization will now going to vary by these are the polarization terms.

Several time I mention that for this polarization term, we have a frequency component P omega and P 2 omega that is important here I am not now I am not considering about the optical rectification term. So, this is now vibrating with the frequency omega. This is now vibrating with the frequency 2 omega. This is a source term of 2 omega electric field and this is a source term of omega electric field ok.

(Refer Slide Time: 18:10)

Theory of SHG

$E^{(\omega)}$ (blue wavy line)
 $E^{(2\omega)}$ (red wavy line)

$E = E^{(\omega)} + E^{(2\omega)}$

Handwritten notes: $\sqrt{E^{(\omega)}}$, $\sqrt{E^{(2\omega)}}$, $P_{NL}^{(\omega)}$, $P_{NL}^{(2\omega)}$

$P_{NL} = \epsilon_0 \chi^{(2)} E^2$
 $\chi^{(2)} = 2d$

$P_{NL} = 2\epsilon_0 d E^2 = 2\epsilon_0 d [E^{(\omega)} + E^{(2\omega)}]^2$

$P_{NL} = \frac{\epsilon_0 d}{2} [E_1 e^{i(k_1 z - \omega t)} + E_2 e^{i(k_2 z - 2\omega t)} + c.c.]^2$

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physics

So, I know what is my electric field, right. So, electric field I have already derived. Now, the next quantity that we need to derive it is P non-linear of 2. So, I need to. So, E omega I know E 2 omega, I know next thing is to find out P non-linear omega. What will be the form of that or P non-linear of 2 omega? Because, in order to use the Maxwell's equation or solve Maxwell's equation, this four quantity is important;, this term and this term we know what is the what is the functional form of that, but this and this we need to figure out.

So, if I want to find out the second harmonic. So, P non-linear is nothing, but epsilon 0 chi 2 E square. Now, here we you going to use this chi 2 as 2 D this is some historical reason people are using this D mac matrix rather using chi 2. So, we will also follow the same thing and going to use in different books you find that this D they represent the chi 2 in terms of D.

But there is a related relation between chi 2 and D and that relation is chi 2 is equal to 2 multiplied by D. So, now, if I use this chi 2, then my P non-linear is 2 epsilon D and E

square. Now, this E square term is containing 2 frequency component omega and 2 omega. So, omega 1 plus omega 2 and square of that so, P non-linear is now having E 1 E to the power of i k z minus omega t E 2 E to the power of i k z minus omega t and plus complex conjugate of that and square of that and half of that. Because, if I go back to our previous slide, let us go back to our previous slide.

(Refer Slide Time: 20:24)

The propagating electric fields

Handwritten formula: $P_{nl} = \epsilon_0 \frac{d}{dt} E^2 = 2\epsilon_0 \frac{d}{dt} \frac{1}{4} [E^{(\omega)} + E^{(\omega)^*}]^2$

Equations shown on the slide:

$$E^{(\omega)} = \frac{1}{2} [\tilde{E}_1 e^{i(k_1 z - \omega t + \phi_1)} + c.c.]$$

$$E^{(\omega)} = \frac{1}{2} [\tilde{E}_1 e^{i\phi_1} e^{i(k_1 z - \omega t)} + c.c.]$$

$$E_1 = \tilde{E}_1 e^{i\phi_1} \text{ (Complex amplitude)}$$

$$E^{(\omega)} = \frac{1}{2} [E_1 e^{i(k_1 z - \omega t)} + c.c.]$$

$$E^{(2\omega)} = \frac{1}{2} [E_2 e^{i(k_2 z - 2\omega t)} + c.c.]$$

Diagrams show waveforms for $E^{(\omega)}$ (a single wave) and $E^{(2\omega)}$ (a higher frequency wave).

Page: 102/101

IIT KHARAGPUR | NPTEL ONLINE CERTIFICATION COURSES | Dr. Samudra Roy, Department of Physics

So, in our previous slide, here you will find that E 1 and E 2 is this. Once we have E 1 and E 2 in this form, now we need to find out what is my P non-linear. So, P non-linear has different frequency component, but we are using this notation and then E square. So, when I make E square. So, E is now 2 of epsilon 0 D 1 by 4 and square means E omega plus E of 2 omega square and E omega is this.

So, that is why, this quantity plus this quantity plus complex conjugate and square of that exactly the same thing we have done in this that slide. So, I should go back to that slide where making the square term here. So, E 1 E to the power of k z minus omega t E 2 E to the power of k z minus omega 2 plus this square of that. So, P non-linear is this.

(Refer Slide Time: 21:39)

Nonlinear polarization for different frequency components



$$P_{NL} = \frac{\epsilon_0 d}{2} [E_1 e^{i(k_1 z - \omega t)} + E_2 e^{i(k_2 z - 2\omega t)} + c.c.]^2$$

Homework

$$P_{NL}^{(\omega)} = \frac{\epsilon_0 d}{2} [2E_1^* E_2 e^{i(k_2 - k_1)z - \omega t} + c.c.]$$

$$P_{NL}^{(2\omega)} = \frac{\epsilon_0 d}{2} [E_1^2 e^{2i(k_1 z - \omega t)} + c.c.]$$

$P^{(0)}$
 $+ P_{NL}^{(\omega)}$
 $+ P_{NL}^{(2\omega)}$



 Dr. Samudra Roy
 Department of Physics

So, now, what we will try to find out what is P non-linear omega and P non-linear 2 omega because, when we do this entire calculation it is for sure.



(Refer Slide Time: 21:51)

**Solution of Nonlinear Maxwell's Equation
(for fundamental wave)**

$$\nabla^2 \vec{E}^{(\omega)} - \mu_0 \epsilon \frac{\partial^2 \vec{E}^{(\omega)}}{\partial t^2} = \mu_0 \frac{\partial^2 \vec{P}_{NL}^{(\omega)}}{\partial t^2}$$

$$E^{(\omega)} = \frac{1}{2} [E_1 e^{i(k_1 z - \omega t)} + c.c.]$$

$$P_{NL}^{(\omega)} = \frac{\epsilon_0 d}{2} [2E_1^* E_2 e^{i(k_2 - k_1)z - \omega t} + c.c.]$$



 Dr. Samudra Roy
 Department of Physics

It is for sure that I will have I will have some component like P 0 P non-linear 0 plus P non-linear omega 1 sorry P non-linear omega because, there is no omega one plus P non-linear 2 omega this 3 components should be there. So, P non-linear omega is a frequency component that is generated with a frequency omega and P 2 omega is a frequency component of and P non-linear with frequency component 2 omega. So, now, if we do

that. So, I can also put this is some sort of homework that how one can generate P non-linear omega from this P non-linear omega with this term. But I will going to explain that, but I want you the student to make it by your own hand do that by your by your own hand.

So, let me let me try to understand that how this P omega components 2 omega components are there P omega and 2 omega. So, omega components will have. So, when I make a square of this term. So, what will happen? We have one term which is square of this; that means, 2 omega will be there one term we have a square of this things.

(Refer Slide Time: 23:42)

Nonlinear polarization for different frequency components

$$P_{NL} = \frac{\epsilon_0 d}{2} [E_1 e^{i(k_1 z - \omega t)} + E_2 e^{i(k_2 z - 2\omega t)} + c.c.]^2$$

Handwritten notes: $a = -\omega$, $b = 2\omega$, $\omega = \omega$, $a \rightarrow b = 2\omega$, $\omega = \omega$

$$P_{NL}^{(\omega)} = \frac{\epsilon_0 d}{2} [2E_1^* E_2 e^{i(k_2 - k_1)z - i\omega t} + c.c.]$$

$$P_{NL}^{(2\omega)} = \frac{\epsilon_0 d}{2} [E_1^2 e^{2i(k_1 z - \omega t)} + c.c.]$$

Handwritten expansion: $(a+b+a+b)^2 = a^2 + b^2 + 2ab + 2ba + 2ab + 2ba = a^2 + b^2 + 4ab$

So, we have 4 omega term will be there and complex conjugate terms are there so, that means, we have a term say this is a plus b plus a star plus b star complex conjugate means the star of that multiplication of a plus b plus a star plus b star it is something like this

So, when we do that it is nothing, but a plus b plus a star plus b star square of that when you do the square of all this things we will have a square term plus b square term plus a star square term plus b star square term plus 2 of a b this couple term plus b a star plus a star b star plus plus a b star this kind of terms you will get. So, now, every time you need to find out what is a frequency component of that. So, a square term has a frequency of 2 omega because, this is omega b square term will having a frequency component of 4

omega. Because, it is square of that a star square is a frequency component of minus of 2 omega this is minus of four omega.

But when the cross terms are there then, there is a possibility that we will have some terms which has a frequency component omega and how you get the frequency component of omega very easily you can find that if I make a star b which is which is this term then what is a frequency component a star has a frequency component of minus omega b as a frequency component of 2 omega if I add it a star b will be the frequency component 2 omega minus omega which is omega why add this multiplication the frequency component is inside the exponential terms.

So, it will be E to the power i omega and then it should be multiplied by E to the power of one is E to the power minus omega and multiplied by E to the power of 2 i of omega. So, when you multiplied that this frequency component will go into add up. So, when we add up then we will have a k term like k 2 minus k 1 and then the frequency of omega. So, you will find that we will have a frequency component of omega, but most importantly this term E 1 star and E 2.

In the similar way, if I try to find out what is the frequency component of 2 omega for P non-linear it is nothing, but the first one that is a square term because only when you make the square term only the first square; that means, a square will give you 2 omega. So, that is why the amplitude is epsilon E one square here in this case the amplitude is 2 a b. So, this 2 multiplication by a star b. So, that is why 2 multiplication of a star and b with the corresponding exponential terms.

So, this exponential term basically gives you the frequency components and when we have all the frequency components, I can determine what is my P non-linear 2 omega and what is my P non-linear omega. So, now, we have all the information in our hands.

So, we are in the position to solve the Maxwell's equation because the P non-linear term and the electric field term both are there the solution of the Maxwell's equation for fundamental waves how to find out how the E omega will going to evolve if I try to find out so, I will going to use.

(Refer Slide Time: 27:30)

**Solution of Nonlinear Maxwell's Equation
(for SH wave)**

$$P_{NL}^{(2\omega)} = \frac{\epsilon_0 d}{2} [E_1^2 e^{2i(k_1 z - \omega t)} + c.c.]$$

$$\nabla^2 \vec{E}^{(2\omega)} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}^{(2\omega)}}{\partial t^2} = \mu_0 \frac{\partial^2 \vec{P}_{NL}^{(2\omega)}}{\partial t^2}$$

$$E^{(2\omega)} = \frac{1}{2} [E_2 e^{i(k_2 z - 2\omega t)} + c.c.]$$

Dr. Samudra Roy
Department of Physics

We will go to use sorry we will go to this equation this another equation where E omega E omega is this and P non-linear omega is this that we just derived in the previous in the previous page also for 2ω also for 2ω we will have. This is how the second harmonics will go to evolve. We will have this equation this is a non-linear Maxwell's equation and P non-linear we just extract that P non-linear is nothing but, this where the frequency component P non-linear 2ω is the frequency component. And here, we find the frequency component is really 2ω the important thing is the amplitude here.

So, we will have the expression of P non-linear 2ω we have the already we have the expression of E . So, both the information is now in our hand. So, the next step is to solve the non-linear Maxwell's equations. So, now in the next class today we will go to I mean we understand how the non-linear Maxwell's equation can be useful to understand the second harmonic.

So, in the second harmonic, we find that it is important to find out the source term that is P non-linear 2ω and P non-linear omega and once we have the functional form of both P and E with different frequency component we can use that in the Maxwell's equation. And readily, we can find the evaluation equation of the second harmonic wave and the fundamental waves. So, in the future classes, we just solve these non-linear Maxwell's equation with certain approximation and try to find out how second harmonic

wave will going to evolve. So, with this note let me conclude here thank you very much for your attention. So, see you in the next class.

Thank you.