

Modern Optics
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Lecture - 03
Maxwell's equations and electromagnetic waves (Contd.)

In this section we will discuss the Plane Electromagnetic Waves in Free Space.

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Contents:

- ✓ Plane EM waves in Free Space
- ✓ Relations and directions of $\vec{E}$, $\vec{B}$ and $\vec{k}$ vectors
- ✓ Polarisation of the EM waves

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The contents are as follows: first we will consider plane electromagnetic waves equation in free space which is derived from the basic set of Maxwell's equations. Then we will look at the electric field, the magnetic field and the propagation vector, their directions, and how they are oriented in the space and their relative orientation also. Followed by this we will discuss the polarisation properties of electromagnetic waves in the free space.

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EM wave in vacuum

$$\nabla^2 \psi = \frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2}$$

where ψ represents $\vec{E}$ or $\vec{B}$ or any of their components

$$\nabla^2 \vec{E} = \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} \quad \nabla^2 \vec{B} = \frac{1}{c^2} \frac{\partial^2 \vec{B}}{\partial t^2}$$

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So, we first consider the basic wave equation which is in this form that del square of psi equal to 1 upon c square del square psi del t square. Here, the psi represents the electric field or the magnetic field or any of their components. So, if we consider the electric field or their component or any of their components then we can write this equation in this form. Whereas, if we considered the magnetic field or their components we will consider this equation, basically both of them are of the same form that is this equation.

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plane wave solution

$$\nabla^2 \psi = \frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2}$$

where ψ represents $\vec{E}$ or $\vec{B}$ or any of their components

$$\psi(x, y, z, t) = \psi_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$
$$= \psi_0 e^{i[(k_x x + k_y y + k_z z) - \omega t]}$$

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Then we look for the solution for this equation, as we know we have already seen that the plane wave solution is given by this equation in the Cartesian coordinate system that the total psi solution is equal to psi 0 and e to the power of k dot r minus omega t. Which, if you decompose this k dot r into the Cartesian coordinates then you can write that psi 0 equal to e to the power of i k x x plus k y y plus k z z minus omega t.

Your omega t it is the frequency of the electromagnetic waves, k x k y and k z are the components of the propagation constants in the three orthogonal Cartesian coordinate directions. So, this plane wave solution for the electromagnetic waves in vacuum is the same where, we will use this psi for a particular electric field or a magnetic field component.

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$\vec{E}$ and $\vec{B}$ fields and their components

$\vec{E} = \hat{i}E_x + \hat{j}E_y + \hat{k}E_z$	$\vec{B} = \hat{i}B_x + \hat{j}B_y + \hat{k}B_z$
$E_x = E_{x0} e^{i[(k_x x + k_y y + k_z z) - \omega t]}$	$B_x = B_{x0} e^{i[(k_x x + k_y y + k_z z) - \omega t]}$
$E_y = E_{y0} e^{i[(k_x x + k_y y + k_z z) - \omega t]}$	$B_y = B_{y0} e^{i[(k_x x + k_y y + k_z z) - \omega t]}$
$E_z = E_{z0} e^{i[(k_x x + k_y y + k_z z) - \omega t]}$	$B_z = B_{z0} e^{i[(k_x x + k_y y + k_z z) - \omega t]}$

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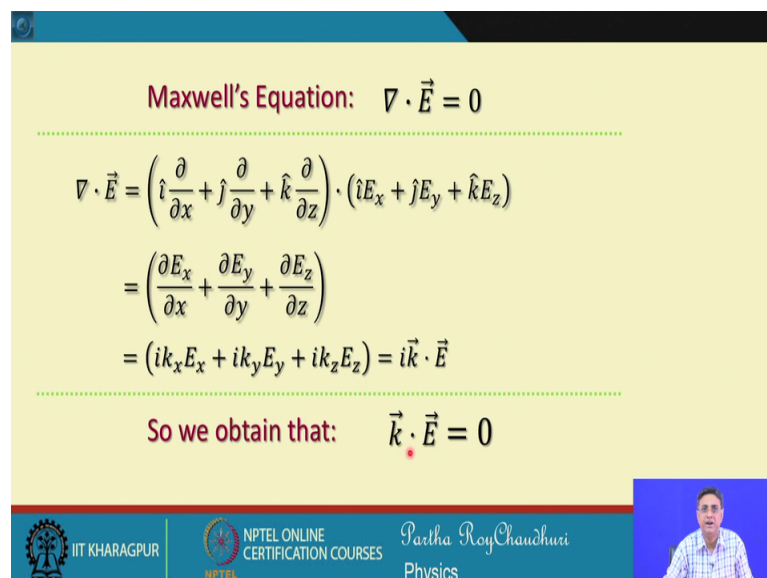
So, if I considered that the electric field in the Cartesian space it will be represented by this quantity that is iE x plus jE y and kE z. So, these are the three cartesian components of the electric field E x, E y and E z. And similarly, for the magnetic field we write that the components of the magnetic field along x is equal to B x, along y it is B y and along z this B z.

Now, each of the Cartesian components that is E x, E y and E z will follow the plane wave solution, as we have seen the plane wave equation will be satisfied by each of the Cartesian components of the electric field as well as the magnetic field. So, the electric field component E x will be equal to some amplitude of the electric field E x 0 into e to

the power of this space factor. And similarly, for E_y we have an amplitude of the electric field for the y component which is E_{y0} and the space factor and similarly for E_z .

So, this is a complete set of the field components representing the total electric field, which is given by this equation. In the same way for the magnetic field we have this magnetic field components B_x , B_y and B_z . So, each of them is related to the amplitude with the space factor by this equation and for B_y by this equation and so on. So, we have seen explicitly the electric field components in terms of its amplitude and the space factors.

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Maxwell's Equation: $\nabla \cdot \vec{E} = 0$

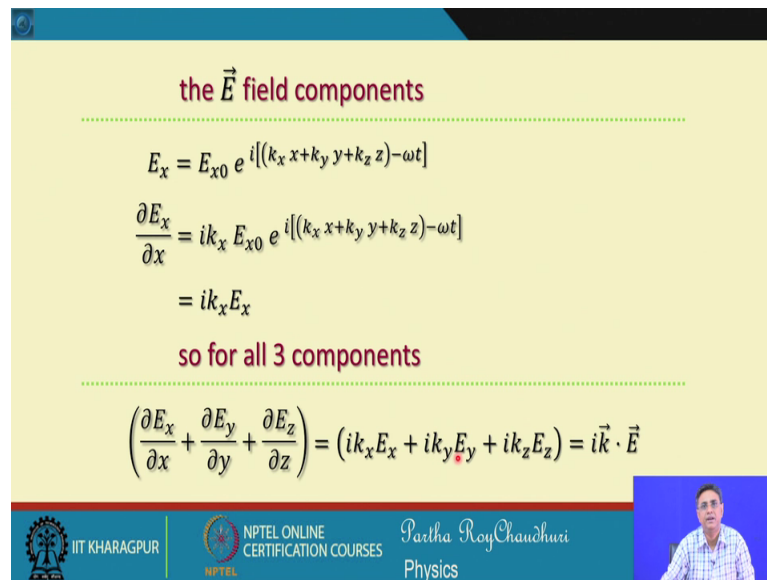
$$\begin{aligned}\nabla \cdot \vec{E} &= \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) \cdot (\hat{i} E_x + \hat{j} E_y + \hat{k} E_z) \\ &= \left(\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \right) \\ &= (ik_x E_x + ik_y E_y + ik_z E_z) = i \vec{k} \cdot \vec{E}\end{aligned}$$

So we obtain that: $\vec{k} \cdot \vec{E} = 0$

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Now, our objective is to look at the orientation, the relative orientation of the electric field and the magnetic field and the direction of propagation of the electromagnetic waves; when you considered the wave is propagating in the free space. So, for that first will bring in the first Maxwell's equation for free space that is $\nabla \cdot \vec{E} = 0$. Having seen that, we can see that this $\nabla \cdot \vec{E}$ is equal to this ∇ operator and this is the electric field. So, $\nabla \cdot \vec{E}$ can be represented by this $\nabla E_x + \nabla E_y + \nabla E_z$.

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the $\vec{E}$ field components

$$E_x = E_{x0} e^{i[(k_x x + k_y y + k_z z) - \omega t]}$$

$$\frac{\partial E_x}{\partial x} = i k_x E_{x0} e^{i[(k_x x + k_y y + k_z z) - \omega t]}$$

$$= i k_x E_x$$

so for all 3 components

$$\left(\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \right) = (i k_x E_x + i k_y E_y + i k_z E_z) = i \vec{k} \cdot \vec{E}$$

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Now, since we know that the electric field E_x satisfies this equation, the solution of electric field E_x is equal to $E_{x0} e^{i(k_x x + k_y y + k_z z) - \omega t}$. Where ω is the frequency of the electromagnetic waves, k_x , k_y and k_z as I have mentioned these are the three Cartesian components of the propagation vectors in the free space. So, $\nabla \cdot E_x$ when you operate on E_x when you operate on E_x then it will detect this $i k_x$ and will leave the E_{x0} into this part, that is E_x as it is.

So, by doing so, we get $\nabla \cdot E_x$ is equal to $i k_x E_x$ and this is for only the x component $\nabla \cdot E_x$. If we considered the other 3 component, other 2 components that is $\nabla \cdot E_y$ and $\nabla \cdot E_z$ then we will end up with similar expressions for $\nabla \cdot E_y$ and $\nabla \cdot E_z$ which you can put together in this form. So, we get that $\nabla \cdot E_x + \nabla \cdot E_y + \nabla \cdot E_z$ is equal to this quantity $i k_x E_x + i k_y E_y + i k_z E_z$, which is which can be written as $i \vec{k} \cdot \vec{E}$.

Now, from this $\nabla \cdot \vec{E}$ we have seen that $i \vec{k} \cdot \vec{E}$ is equal to $i \vec{k} \cdot \vec{E}$, but $\nabla \cdot \vec{E}$ itself is equal to 0. Therefore, you readily obtain that $\vec{k} \cdot \vec{E}$ is equal to 0. This tells you that $\vec{k}$ and $\vec{E}$, because their dot product is 0. So, they must be perpendicular to each other. So, we conclude by saying that $\nabla \cdot \vec{E}$ is equal to $\vec{k} \cdot \vec{E}$ is equal to 0, which means that the wave vector $\vec{k}$ is perpendicular to the electric field $\vec{E}$ when the electromagnetic wave is propagating in the free space.

Now, if we consider instead of $\nabla \cdot \vec{E} = 0$ the first equation, the second equation that is $\nabla \cdot \vec{B} = 0$ and we go back with this equation $\nabla \cdot \vec{B} = 0$, you can again write this equation $\nabla \cdot \vec{B}$ will be equal to $i \frac{\partial}{\partial x} B_x + j \frac{\partial}{\partial y} B_y + k \frac{\partial}{\partial z} B_z$. As a result we will obtain this $\nabla \cdot \vec{B}$ will be equal to $\frac{\partial}{\partial x} B_x + \frac{\partial}{\partial y} B_y + \frac{\partial}{\partial z} B_z$.

Now, in the same way if we calculate this $\frac{\partial}{\partial x} B_x$ we will obtain $i k_x B_x$ and so on. Therefore, this $\frac{\partial}{\partial x} B_x + \frac{\partial}{\partial y} B_y + \frac{\partial}{\partial z} B_z$ will be equal to $i k_x B_x + i k_y B_y + i k_z B_z$, which will be equal to $i \vec{k} \cdot \vec{B}$ in the same way.

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$\nabla \cdot \vec{E}(\vec{r}, t) = \vec{k} \cdot \vec{E} = 0$

wave vector $\vec{k}$ is perpendicular to $\vec{E}$

In the same way

$\nabla \cdot \vec{B}(\vec{r}, t) = \vec{k} \cdot \vec{B} = 0$

wave vector $\vec{k}$ is perpendicular to $\vec{B}$

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Therefore, we obtain in the same way the $\nabla \cdot \vec{B}$ since, $\nabla \cdot \vec{B}$ is equal 0 $\vec{k} \cdot \vec{B}$ is also equal to 0; which will lead to the conclusion that the wave vector $\vec{k}$ is perpendicular to $\vec{B}$. So, we have been able to achieve two relations that is $\vec{k}$ is perpendicular to $\vec{E}$ and $\vec{k}$ is also perpendicular to $\vec{B}$.

This means that $\vec{k}$ and $\vec{B}$ are perpendicular, $\vec{E}$ and $\vec{k}$ their also perpendicular. But, that does not ensure that $\vec{E}$ and $\vec{B}$ whether they are perpendicular or not. So, there is a more rigorous way of looking at the relation between $\vec{k}$, $\vec{B}$ and $\vec{E}$; so to do that we consider Maxwell's third equation the curl equation.

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Maxwell's Equation: $\nabla \times \vec{E}(\vec{r}, t) = -\frac{\partial \vec{B}(\vec{r}, t)}{\partial t}$

LHS:

$$\nabla \times \vec{E} = \begin{pmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & E_z \end{pmatrix} \cdot \frac{\partial E_z}{\partial y} = ik_y E_{z0} e^{i[(k_x x + k_y y + k_z z) - \omega t]} = ik_y E_z$$

$$(\nabla \times \vec{E})_x = \frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = ik_y E_z - ik_z E_y = i(\vec{k} \times \vec{E})_x$$

taking all 3 components $\nabla \times \vec{E} = i(\vec{k} \times \vec{E})$



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Now, the curl of E if we have to represent this we write in this way that $\hat{i} \nabla \times \hat{j} \nabla \times \hat{k} \nabla \times \hat{i} \nabla \times \hat{j} \nabla \times \hat{k} \nabla \times \hat{i} \nabla \times \hat{j} \nabla \times \hat{k}$. So, this is the curl operator operating on E which if we take the x component of this curl then it will be $\nabla \times \nabla \times \hat{i} \nabla \times \hat{j} \nabla \times \hat{k} \nabla \times \hat{i} \nabla \times \hat{j} \nabla \times \hat{k}$ of E y. So, this quantity is the x component of the curl which will give you $ik_x \nabla \times \hat{j} \nabla \times \hat{k} \nabla \times \hat{i} \nabla \times \hat{j} \nabla \times \hat{k}$ minus $ik_z \nabla \times \hat{j} \nabla \times \hat{k} \nabla \times \hat{i} \nabla \times \hat{j} \nabla \times \hat{k}$ because, $\nabla \times \hat{i} \nabla \times \hat{j} \nabla \times \hat{k}$ will detach only the y component that is k_y will be out. So, $ik_y \nabla \times \hat{j} \nabla \times \hat{k} \nabla \times \hat{i} \nabla \times \hat{j} \nabla \times \hat{k}$ and this factor will remain as it is which is equal to $ik_y \nabla \times \hat{j} \nabla \times \hat{k}$.

So, $\nabla \times \vec{E}$ the x component of this quantity will give you $ik \nabla \times \vec{E}$ the x component of that and if we consider all the other 2 components, namely the $\nabla \times \vec{E}$ y component and $\nabla \times \vec{E}$ z component we will end up with two more expressions; $ik \nabla \times \vec{E}$ y component, $ik \nabla \times \vec{E}$ z component. And then, if we put together all the 3 components we can write that $\nabla \times \vec{E}$ is equal to $ik \nabla \times \vec{E}$.

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Maxwell's Equation: $\nabla \times \vec{E}(\vec{r}, t) = -\frac{\partial \vec{B}(\vec{r}, t)}{\partial t}$

RHS :

$$\frac{\partial \vec{B}}{\partial t} = \frac{\partial}{\partial t} (iB_x + jB_y + kB_z) \quad \frac{\partial B_x}{\partial t} = i\omega B_{x0} e^{i[(k_x x + k_y y + k_z z) - \omega t]}$$

$$= -i\omega (iB_x + jB_y + kB_z) \quad = i\omega B_x$$

$$= -i\omega \vec{B}$$

therefore, this term $\frac{\partial \vec{B}}{\partial t} = -i\omega \vec{B}$

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So, this is again a very interesting relation now, we will have to evaluate the right hand side that is $\nabla \times \vec{B}$. So, because you know that $\vec{B}$ vector can be represented with the Cartesian forms like $iB_x + jB_y + kB_z$. And if you take the time derivative of the Cartesian component, then we will see that $\nabla \times \vec{B}$ we will take out this ω and i factor outside. And this equation, expression for the magnetic field x component will remain as it is. As a result we can write $\nabla \times \vec{B}$ is equal to $i\omega \vec{B}$.

So, this is only for the x component of the magnetic field when you take the time derivative. And similarly, if you do the same thing for the other components that is $\nabla \times \vec{B}$ of B_y $\nabla \times \vec{B}$ of B_z then we can write this equation all put together in this form which is equal to $i\omega \vec{B}$. And this is a very interesting way of looking at things that the right hand side just boils down to $i\omega \vec{B}$. So, $\nabla \times \vec{B}$ equal to $i\omega \vec{B}$. Now, we have ready expression for $\nabla \times \vec{E}$ as well as $\nabla \times \vec{B}$.

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$$\nabla \times \vec{E}(\vec{r}, t) = -\frac{\partial \vec{B}(\vec{r}, t)}{\partial t}$$

$$\nabla \times \vec{E} = i(\vec{k} \times \vec{E}) \quad \frac{\partial \vec{B}}{\partial t} = -i\omega \vec{B} \Rightarrow \vec{k} \times \vec{E} = \omega \vec{B}$$

$\vec{B}$ is perpendicular to $\vec{E}$
 $\vec{B}$ is perpendicular to $\vec{k}$

$$\hat{k} \times \vec{E} = \frac{\omega}{k} \vec{B} = c \vec{B}$$

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If we know put together so, del cross E has given us i k cross E and del cross B has given as del B del t equal to minus i omega B. So, we can use these two expressions to arrive at the conclusion that k cross E k cross E must be equal to omega B. So, that is a very beautiful relation which says that the B vector is perpendicular to E and B vector is perpendicular to k. That means B vector must be perpendicular to the plane containing the vector k and the vector E. So, B is perpendicular to E and B is perpendicular to k.

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In the same way

$$\nabla \times \vec{B}(\vec{r}, t) = \mu_0 \epsilon_0 \frac{\partial \vec{E}(\vec{r}, t)}{\partial t}$$

$$\vec{B} \times \vec{k} = \frac{\omega}{c^2} \vec{E}$$

$$\nabla \times \vec{B} = i(\vec{k} \times \vec{B}) \quad \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t} = -\frac{i\omega}{c^2} \vec{E} \Rightarrow \vec{k} \times \vec{B} = -\frac{\omega}{c^2} \vec{E}$$

$\vec{B}$ is perpendicular to $\vec{E}$
 $\vec{k}$ is perpendicular to $\vec{E}$

$$\vec{B} \times \hat{k} = \frac{\omega}{kc^2} \vec{E} = \frac{1}{c} \vec{E}$$

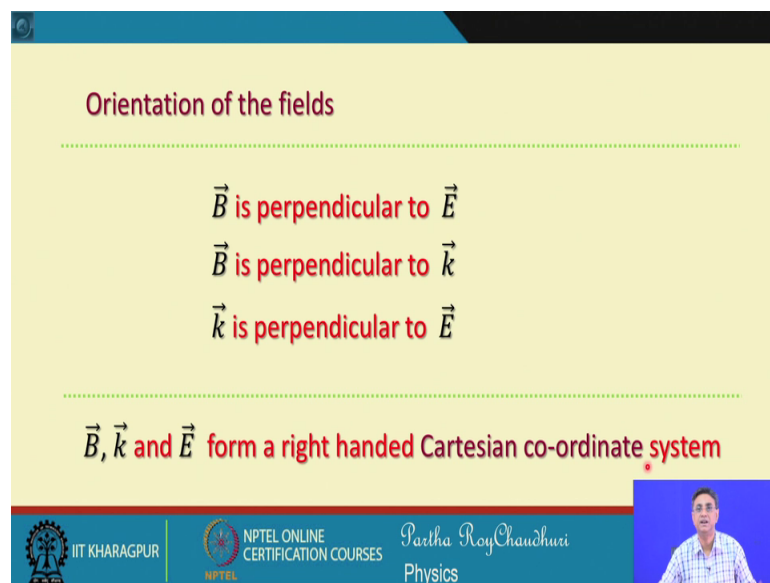
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Now, in the same way if we proceed for the fourth equation that is $\nabla \times \mathbf{B}$ is equal to $\mu_0 \epsilon_0 \nabla \times \frac{\partial \mathbf{E}}{\partial t}$. If we evaluate this equation and in the same way, if we evaluate this left hand part of the calculation that is $\nabla \times \mathbf{E}$ $\nabla \times \mathbf{B}$ is equal to $i k \times \mathbf{B}$. And for the right hand side if I do this derivative time derivative of the electric field, I arrive at $\frac{1}{c^2} \nabla \times \frac{\partial \mathbf{E}}{\partial t}$ because, $\mu_0 \epsilon_0$ is $1/c^2$. So, this quantity will become $i \omega$ by c^2 and the vector $\mathbf{E}$. So, I still get the equation for $\nabla \times \mathbf{E}$ in this form $\nabla \times \mathbf{B}$ and $\nabla \times \mathbf{E}$ in this form.

Now, if I use this left hand side and right hand side together I may say that $k \times \mathbf{B}$ must be equal to $i \omega/c^2 \mathbf{E}$. So, look at this equation this tells you that k is perpendicular to $\mathbf{E}$ and $\mathbf{B}$ is perpendicular to $\mathbf{E}$; that means the electric field $\mathbf{E}$ is perpendicular to the plane containing the vectors k and $\mathbf{B}$. So, this is the complementary part which we have seen earlier for the electric field.

But all three of them they are very nicely related can see that if I remove, if I want to get rid of this minus sign we can reverse this position that is $\mathbf{B} \times k$ will be equal to $\omega/c^2 \mathbf{E}$. So, $\mathbf{B}$, k and $\mathbf{E}$ they form this equation from here you can conclude that $\mathbf{B}$ is perpendicular to $\mathbf{E}$ and k is perpendicular to $\mathbf{E}$ also.

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Orientation of the fields

$\vec{B}$ is perpendicular to $\vec{E}$
 $\vec{B}$ is perpendicular to $\vec{k}$
 $\vec{k}$ is perpendicular to $\vec{E}$

$\vec{B}$, $\vec{k}$ and $\vec{E}$ form a right handed Cartesian co-ordinate system

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Now, we will look at the consequence of these two findings $\mathbf{B}$ is perpendicular to $\mathbf{E}$, $\mathbf{B}$ is perpendicular to k and also k is perpendicular to $\mathbf{E}$. This means that $\mathbf{B}$, k , $\mathbf{E}$ all of them

are perpendicular to each other, all of them are or at right angles to each other and they form a right handed Cartesian coordinate system. We will see that how they are they are connected to each other in terms of the direction, in terms of the coordinate access.

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so we see

$$\vec{k} \times \vec{E} = \omega \vec{B}$$

$$\vec{B} \times \vec{k} = \frac{\omega}{c^2} \vec{E}$$

take $\vec{k}$ along z and $\vec{E}$ along x so $\vec{B}$ along y and $\vec{k}$ along z

$k_z \times E_x = \omega B_y$
➡
 $B_y \times k_z = \frac{\omega}{c^2} E_x$

then $\vec{B}$ happens to be along y then $\vec{E}$ happens to be along x

Consistent ! $\{\vec{E}, \vec{B}, \vec{k}\} \rightarrow \{x, y, z\}$ right handed triad

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So, you see $\vec{k} \times \vec{E} = \omega \vec{B}$, $\vec{k} \times \vec{B} = \frac{\omega}{c^2} \vec{E}$. If we consider this example that, let us suppose that I take the electromagnetic wave to be propagating along the z direction. Then I can write $\vec{k}$ equal to $k_z \hat{z}$ and if I consider the electric field is oriented along the x direction, then $k_z \times E_x$ should give me ωB_y . So, that should yield B_y ; that means, if I take $\vec{k}$ along z and $\vec{E}$ along x I should expect that $\vec{B}$ should be oriented along y, $\vec{z} \times \vec{x} = \vec{y}$.

And if I assume that $\vec{B}$ happens to be along y axis and then $\vec{k}$ is along z axis which I have already considered, then $\vec{B} \times \vec{k} = B_y \times k_z = \frac{\omega}{c^2} E_x$; that means, $\vec{y} \times \vec{z} = \vec{x}$. So, they again consistent then $\vec{E}$ happens to be along x. So that means, this is consistent, because I have started with the assumption that $\vec{k}$ is along z, $\vec{E}$ is along x, the result was that $\vec{B}$ is along y,

Now, I have considered that $\vec{B}$ is along y and $\vec{k}$ is along z so, it yields that $\vec{E}$ is along x. So, they are consistent and we can write in a compact form that is $\vec{E}, \vec{B}, \vec{k}$ will represent that x y z the axis of the Cartesian coordinate system and it forms a right handed triad of vectors. So, this is how we could see that the electric field, the magnetic field and the direction of propagation they are normal to each other. And because, the

direction of propagation is such that the electric field and magnetic field they are always perpendicular and there are also perpendicular to each other than the wave is a transverse wave.

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so we see

$\vec{k} \times \vec{E} = \omega \vec{B}$

$\vec{B} \times \vec{k} = \frac{\omega}{c^2} \vec{E}$

take $\vec{k}$ along z and $\vec{E}$ along x so $\vec{B}$ along y and $\vec{k}$ along z

$k_z \times E_x = \omega B_y$
➡
 $B_y \times k_z = \frac{\omega}{c^2} E_x$

then $\vec{B}$ happens to be along y then $\vec{E}$ happens to be along x

Consistent ! $\{\vec{E}, \vec{B}, \vec{k}\} \rightarrow \{x, y, z\}$ right handed triad

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So, this is how we can also look at the nature of the electromagnetic waves, the transverse nature of the electromagnetic waves in the free space. So, pictorially if we represent this is your x axis, this is your y axis and this is your z axis. So, along the x axis if you consider the electric fields are oriented then along the y axis the magnetic fields will be oriented. And as a matter of fact that orthogonal triad of vectors will give you that the propagation direction will be along the z axis. So E, B, k forms a right handed Cartesian coordinate system.

Now, we discussed this fact for the case of electromagnetic waves in free space and you know not only for free space, but for most dielectrics this permeability μ is equal to μ_0 . Those this is a nonmagnetic property and mostly μ you can approximate as μ_0 . Therefore, this B becomes equal to $\mu_0 H$, where H is the magnetic field, B was the magnetic induction vector. In that case as well for the rest part of all formalisms we can replace that B by H.

So, if you orient the electric field along the x direction and H along y direction then again the propagation can vector will be along z direction. So, in that case E, H and k they will form the triad of vectors. So, in free space we can see that E, H, k they form the

triad of vectors and also E , B and k that forms the triad of vectors. Because E and B , H and k they are related in the same direction.