

Tapestry of Field theory: Classical & Quantum, Equilibrium & Nonequilibrium Perspectives

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Week - 02

Lecture – 13

Hamiltonian formalism, okay. So, we need some more energy for this one. So, the way we write Hamiltonian for particles, we can write Hamiltonian for fields. Okay. So, momentum density, we have momentum for particles. So, what is this dL by dx dot, partial L partial x dot.

So, instead of that we have partial L , not partial, okay, it is a functional derivative. L is a function, so we get partial L partial δL delta π dot. Now, i is a component, it could be x component, y component, g component, I get partial L partial π dot. Now, L is integral, $L d^4x$.

So, when I go inside, then it is going to be partial. So, this is total L , d cube not d , sorry, d cube. So, total L I have to do the functional derivative because it is function of any ϕ . But when I go inside, this Lagrangian density is function of ϕ , then I do partial derivative. Okay.

Now, that is momentum. I defined, this is a definition, Hamiltonian density is π_i , this is π_i , π_i dot minus L is very similar to what we do for particles, $\pi_i x_i$ dot minus L . Okay. And we define total Hamiltonian is integral over space. Okay.

So, now, equation of motion. So, like we have x dot is derivative of H , you know, with π dot. So, we have ϕ_i dot equal to partial H , no, there is no derivative. This is old slide actually. ΔH delta π_i and we have π_i dot is derivative ΔH delta ϕ_i .

Okay. So, I think I have an example. So, we have this example. For this I am going to write down the Hamiltonian. Now, this Lagrangian, Lagrangian is function of ϕ , ϕ dot, but Hamiltonian is function of ϕ and π .

Okay. So, this dot is not there. So, ΔL delta ϕ . Now, this dot is there, ΔL ϕ_i dot, like ΔL delta x_i dot is there, no, for, and that is, will act on derivative, which is ϕ_i dot, this one, square. You take the derivative, I will get ϕ dot. So, π is nothing but

phi dot, which is like momentum is \dot{x} , very similar.

The formula is in fact motivated from particles. So, Hamiltonian will be $\pi \dot{\phi} - L$. Now, π is nothing but $\dot{\phi}$. So, this guy will give us, okay. So, what will I get? $\dot{\phi}$ into a $\dot{\phi}$, which is $\dot{\phi}^2$ minus this one, minus half $\dot{\phi}^2$ square, that is this one.

Okay. Then this minus minus becomes plus, so plus $\text{grad } \phi^2$ plus $V \phi$. Okay. Now, $\dot{\phi}$ is nothing but π . So, this is my Hamiltonian density. It should be function of π and ϕ and $\text{grad } \phi$, but not $\dot{\phi}$.

Okay. Don't write $\dot{\phi}$ for Hamiltonian. The way for particles, we don't write \dot{x} for Hamiltonian, we write only momentum. Okay. So, equation of motion. So, $\dot{\phi}$ is $\delta H / \delta \pi$, this is not there, $\delta H / \delta \pi = 1$ for this is 1D.

So, this is going to give us $\dot{\phi} = \delta H / \delta \pi = \pi$. Okay. Now, this is not dot here. So, $\pi = \delta H / \delta \dot{\phi}$ with a minus sign, this is a minus sign here.

Okay. Now so, we get $\delta H / \delta \pi$. So, this is a functional derivative. Okay. So, functional derivative, I again apply the logic of functional derivative is function of the derivative ϕ as well as $\text{grad } \phi$ will come.

Okay. Both the terms will come. So, I get $\partial H / \partial \phi - \text{div}(\partial H / \partial \text{grad } \phi)$. Okay. This is, this will come and this one is the Hamiltonian. So, function of derivative ϕ will give us, will come from here $V \phi$.

So, that is $V \phi$ and $\partial H / \partial \text{grad } \phi$ will give us $\text{grad } \phi$ here and divergence. So, divergence of a grad is Laplacian. Okay, that. So, this is $\dot{\phi}$, $\dot{\phi}$ and $\dot{\phi}$ is, well $\pi = \dot{\phi}$. So, this is nothing but $\ddot{\phi}$ or this, it is square.

Okay. Right. So, that I am just using this, this relation. So, this equation. Okay. So, field with a potential, equation, Euler Lagrange equation is this. And there are potential, there are in fact, we will encounter fields with potential.

You might have seen the free energy, you know, for, for phase transition is $\phi^2 - \phi^4$. So, these are the potential. Okay. And so, this is the example. I think I will do the other example later, because the thing is already is close to saturation.

So, this one is a quite important one. So, you can write down Lagrange for Schrodinger

equation, non-relativistic, non-relativistic Schrodinger equation and I will do it in the next class. Thank you.