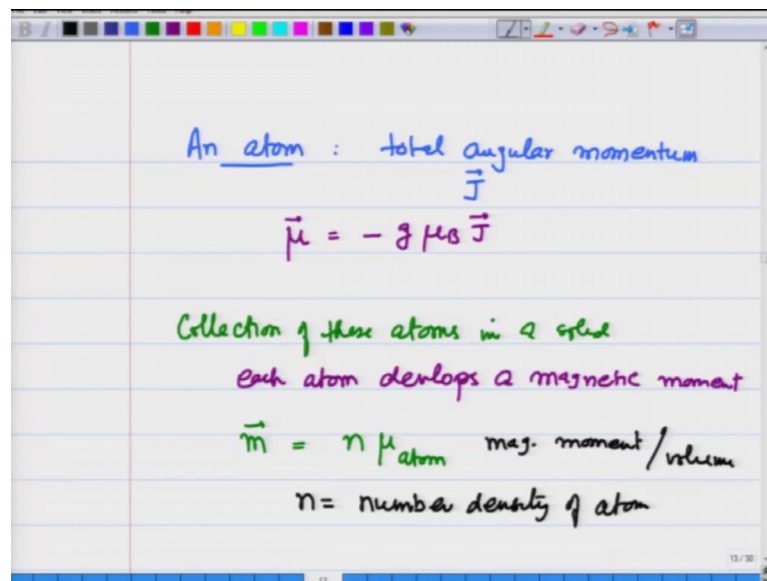


Introduction to Solid State Physics
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Lecture – 76
Paramagnetism in solids - II
Temperature dependence of paramagnetic susceptibility and Curie's Law

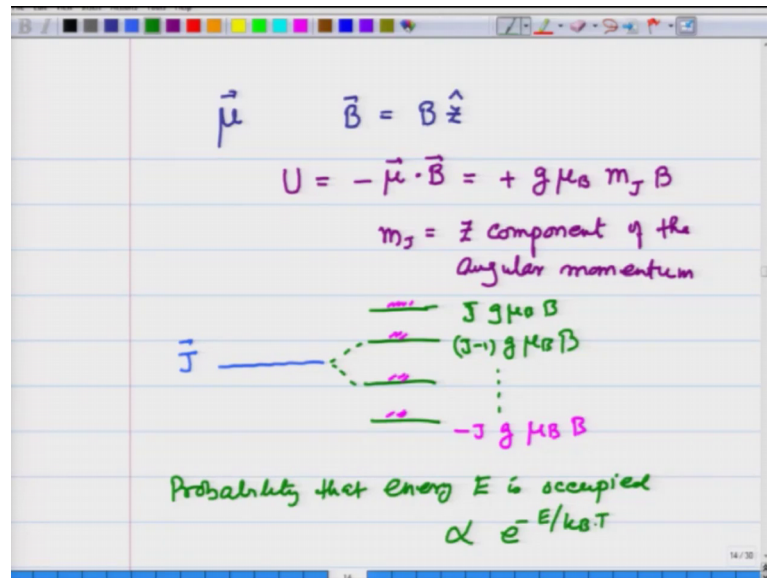
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In the previous lecture, I derived the magnetic moment and the Lande g factor for an atom. And what we concluded there is that an atom has a total angular momentum J , it will then have the corresponding magnetic moment μ which is equal to minus $g \mu_B$ times J .

Now, we have this collection of these atoms in a solid. So, we have a collection of these atoms in a solid. Then if I apply a magnetic field, each atom develops a magnetic moment and assuming that these atoms or these magnetic moments do not interact with each other. So, each atom is giving its own magnetic moment. The total magnetic moment m of the system is going to be equal to n times μ_{atom} , where n is the number density of atoms. So, this m should be called the magnetic moment per unit volume or the magnetization. So, what I want to now calculate is how this μ_{atom} is given as a function of temperature.

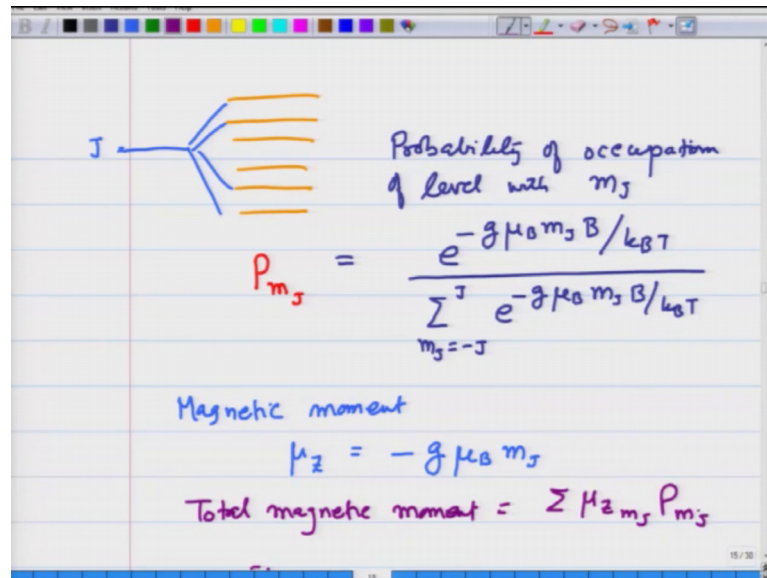
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Now, when this system the magnetic moment is put in a magnetic field let us say field B is in the z direction so, am going to write this as $B z$. I choose z direction because that is easy to calculate things with. Then the energy of each atom is going to be minus μ dot B and this is going to be equal to plus $g \mu_B m_J$ times B , where m_J is the z component of the angular momentum. So, I am going to have this atom which had initially a level J , with J angular momentum is going to split, is going to have all these different levels.

Now, you can see that if m_J is positive, energy is positive. So, with m_J positive I am going to have the level go up so, let us say this is $J g \mu_B B$, next one is going to be J minus 1 $g \mu_B B$ and so on and the lowest level is going to be minus $J g \mu_B B$. These are the energy levels into which this will split and at a temperature T , the atoms are going to be distributed among these levels. However, the number of atoms in each level is going to be different. That is because the probability that energy E is occupied by an atom is going to be proportional to e raised to minus E over k Boltzmann times T .

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The image shows a digital whiteboard with handwritten notes. On the left, a blue line labeled 'J' branches into five horizontal orange lines, representing energy levels. To the right of these lines, the text reads 'Probability of occupation of level with m_J '. Below this, the probability formula is written as $P_{m_J} = \frac{e^{-g\mu_B m_J B / k_B T}}{\sum_{m_J=-J}^J e^{-g\mu_B m_J B / k_B T}}$. Further down, the text 'Magnetic moment' is written, followed by the formula $\mu_z = -g\mu_B m_J$. At the bottom, the formula for total magnetic moment is given as $\text{Total magnetic moment} = \sum \mu_z m_J P_{m_J}$.

Therefore, if you put it all together, here is this atom with level given by J and this has split into different levels right let me show these levels by different colors. Each level has a different probability of occupation. So, probability of occupation of level with z component m_J is going to be equal to e raised to minus $g \mu_B m_J B$ over $k_B T$ divided by, I have to normalize the probability sum over m_J varying from minus J to J e raised to minus $g \mu_B m_J B$ over $k_B T$. And let me label this as probability that level m_J is occupied.

Now, the magnetic moment μ in the z direction for this level is going to be minus $g \mu_B m_J$. So, if I multiply each of this magnetic moment by the corresponding probability and add it up I am going to get the total magnetic moment. So, total magnetic moment is going to be summation of μ_z for m_J so, times the probability that this m_J level is occupied.

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The image shows a digital whiteboard with the following content:

$$\mu_z(T) = \frac{\sum_{m_J=-J}^J -g \mu_B m_J e^{-g \mu_B m_J B / k_B T}}{\sum_{m_J=-J}^J e^{-g \mu_B m_J B / k_B T}}$$

At low temperatures

$$\mu_B B / k_B T \ll 1$$

$$\mu_z(T) = \frac{\sum_{m_J=-J}^J -g \mu_B m_J \left(1 - \frac{g \mu_B m_J B}{k_B T}\right)}{\sum_{m_J=-J}^J \left(1 - \frac{g \mu_B m_J B}{k_B T}\right)}$$

So, let us do that. So, μ_z is going to be minus $g \mu_B m_J e^{-g \mu_B m_J B / k_B T}$ summed over m_J equals minus J to J divided by summation m_J equals minus J to J $e^{-g \mu_B m_J B / k_B T}$. This is the net magnetic moment and if you like, I can call this as a function of temperature and as a function of applied field B . Now, what we will find is that at low temperatures and we are going to specify what these low temperatures are later in the next lecture when I do certain examples.

Let me take this $\mu_B B / k_B T$ to be much less than 1. Then I can expand this exponential term and write μ_z as a function of temperature as equal to summation m_J where m_J varying from minus J to J minus $g \mu_B m_J$ times $1 - \frac{g \mu_B m_J B}{k_B T}$ divided by summation m_J equals minus J to J $1 - \frac{g \mu_B m_J B}{k_B T}$. Under this approximation this is what μ_z is going to be. Let me write this again on the next page.

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The image shows a whiteboard with the following handwritten equations:

$$\mu_z(T) = \frac{\sum_{m_J=-J}^J -g \mu_B m_J \left(1 - \frac{g \mu_B m_J B}{k_B T}\right)}{\sum_{m_J=-J}^J \left(1 - \frac{g \mu_B m_J B}{k_B T}\right)}$$

$$\sum_{m_J=-J}^J m_J = 0$$

$$\mu_z(T) = \frac{\sum_{m_J=-J}^J \frac{g^2 \mu_B^2 B}{k_B T} m_J^2}{(2J+1)}$$

$$\sum_{J=1}^J m_J^2 = \frac{J(J+1)(2J+1)}{6} \quad \text{---} \times 2$$

So, I have μ_z as a function of temperature is equal to summation m_J equals minus J to J minus $g \mu_B m_J \left(1 - \frac{g \mu_B m_J B}{k_B T}\right)$ divided by summation m_J equals minus J to J $1 - \frac{g \mu_B m_J B}{k_B T}$. Let us look at the first term in the numerator. Summation m_J equals minus J to J is equal to 0 because they are equal number of negative and positive terms. And therefore, I can write $\mu_z(T)$ equals summation m_J equals minus J to J $\frac{g^2 \mu_B^2 B}{k_B T} m_J^2$ divided by the second term in the denominator will be 0 and I get $2J + 1$.

Now, use the series sum formula for summation m_J^2 equals 1 to J . Is going to be equal to $J(J+1)(2J+1)$ divided by 6 and I will substitute that in the expression above. Now this is going here multiplied by 2 because I have m_J going from minus J to minus 1 and m_J going from 1 to J and 0 anyway does not contribute.

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The image shows a digital whiteboard with the following handwritten equations and text:

$$\mu_z(T) = \frac{g^2 \mu_B^2}{k_B T} \cdot B \cdot \frac{2J(J+1)(2J+1)}{6 \times (2J+1)}$$

$$\mu_z(T) = \left\{ \frac{g^2 \mu_B^2 J(J+1)}{3k_B T} \right\} B$$

$$\mu = \chi_M B / \mu_0$$

$$\chi_M = \frac{n \mu_0 g^2 \mu_B^2 J(J+1)}{3k_B T} = \frac{C}{T}$$

$n =$ Number of atoms / volume / mole

So, I get $\mu_z(T)$ as equal to $g^2 \mu_B^2$ over $k_B T$ times B times $2J(J+1)$ divided by $6 \times (2J+1)$. And I cancel this term and this 2 goes with 6 and gives me 3 and therefore, I end up getting $\mu_z(T)$ as equal to $g^2 \mu_B^2 J(J+1)$ divided by $3k_B T$ times B . And if I compare this with the expression for susceptibility, μ is equal to $\chi_M B / \mu_0$. I get χ_M equals $n \mu_0 g^2 \mu_B^2 J(J+1)$ divided by $3k_B T$. Where n is the number of atoms per unit volume or per mole depending on how you want to write, what you want to calculate χ_M for. And this equals C/T where C is a constant.

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The image shows a digital whiteboard with the following handwritten text and equations:

$$\frac{\vec{m}}{\text{volume}} = \vec{M} = \frac{C}{T}$$

Curie's law

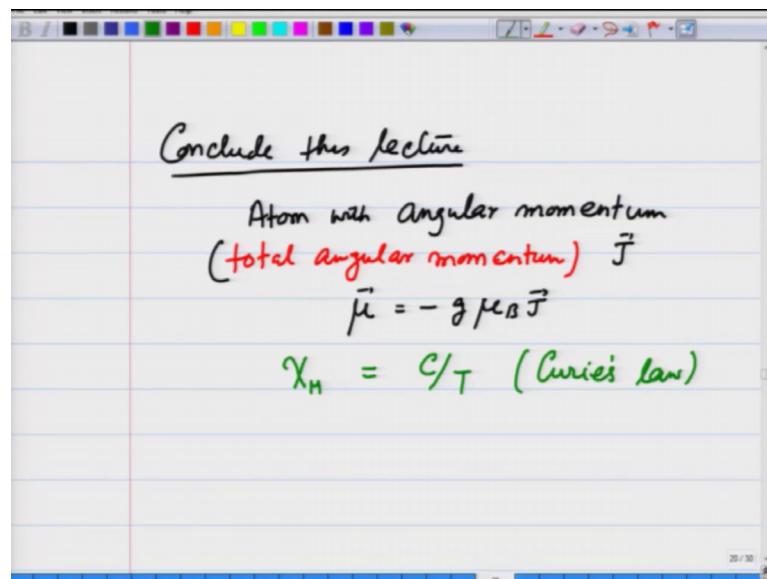
Paramagnetic susceptibility $\propto \frac{1}{T}$
 $= C/T$

Paramagnetic systems

So, what I found through this calculation is that, when I put this collection of atoms in a solid and the atomic magnetic moments are not interacting with each other; I get the magnetic moment m per unit volume, which is also known as the magnetization is equal to some function C divided by T , it is inversely proportional to T . This is known as Curie's law. What does that law say? It says that, paramagnetic susceptibility is proportional to 1 over T is goes as some constant divided by T . As temperature goes up the paramagnetic susceptibility goes down because the levels start filling roughly equally at very high temperature.

So, this is the introduction to paramagnetic systems. Systems where what we have done is, we have taken an atom see how the levels are split under Russell Saunders or L s coupling when a field is applied. And then populated them according to the Boltzmann's law and using that we found the paramagnetic susceptibility.

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So, to conclude this lecture, what we have found is that atom with angular momentum and when I say angular momentum, I mean the total angular momentum. So, let me specify total angular momentum J has magnetic moment μ equals minus $g \mu_B J$. And when I apply magnetic field these levels due to different z component of J split and populate differently and that leads to a paramagnetic susceptibility χ_m equals C over T , which is known as Curie's law. How these J s and L and s are calculated will be the subject of next lecture.