

Introduction to Solid State Physics
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Lecture – 71
Calculating carrier density in semiconductors – II

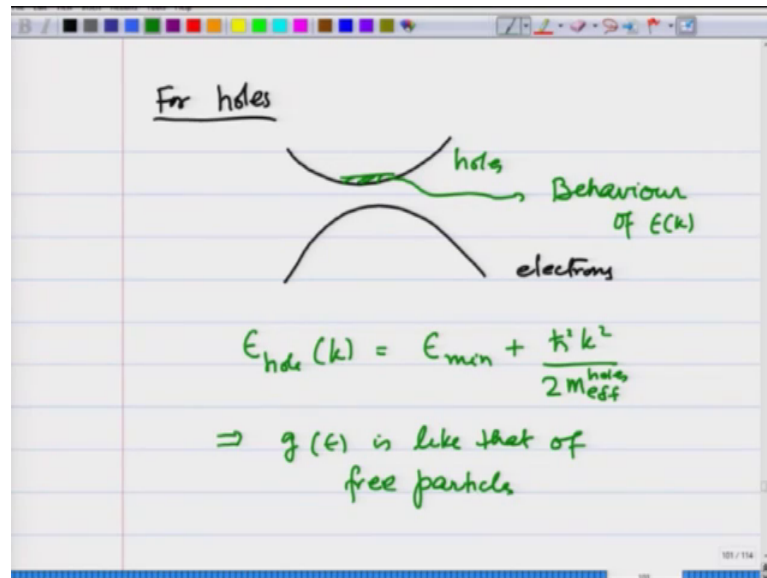
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$$f_{\text{hole}} = \frac{1}{e^{(E_{\text{hole}} + E_F)/k_B T} + 1}$$
$$n = \int g(\epsilon) f_e(\epsilon) d\epsilon$$
$$n_{\text{holes}} = \int g_{\text{holes}}(\epsilon) f_h(\epsilon) d\epsilon$$

g for electrons \equiv that for free electrons
g for holes \equiv g for free particles } m_{eff}

In the previous lecture I derived f_{hole} which is equal to 1 over e raised to ϵ_{hole} plus ϵ_F over $k_B T$ plus 1 and we had also seen that the number of electrons n is going to be equal to $g(\epsilon)$ this is all about electrons f let me write this e as a function of E dE and number of holes is going to be integration $g_{\text{holes}}(\epsilon) f_h(\epsilon) dE$. I had also argued that, g for electrons is identical to that for free electrons. And now I will show that g for holes is also equivalent to g for free particles. The only difference in these two from free particles is going to be that now I have $m_{\text{effective}}$ instead of m_{free} .

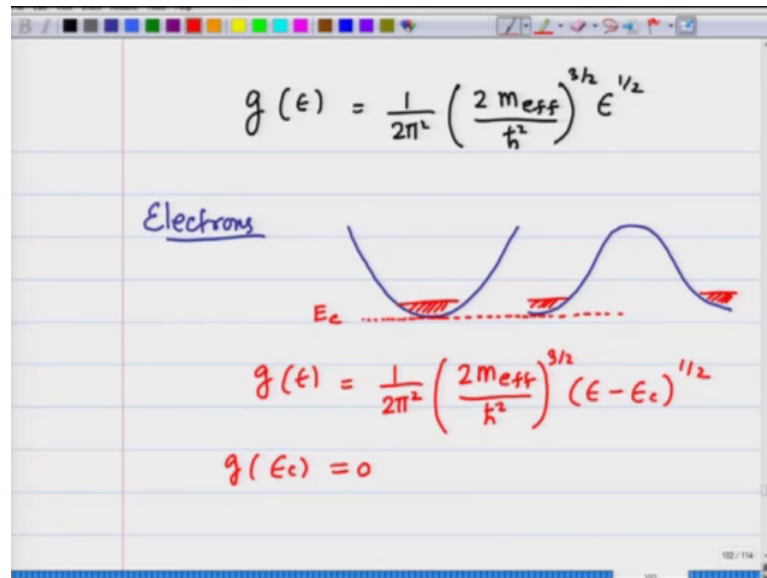
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Now for the holes if I have an energy band like this for electrons. We have seen that for holes the e becomes minus e so the band goes like this. This is for the holes and again, near the bottom the behaviour of ϵ vs k it is like that of free particles. So, ϵ hole as a function of k is like this ϵ_{min} which we can easily take to be 0 plus $\hbar^2 k^2$ over $2m_{\text{eff}}$ of holes.

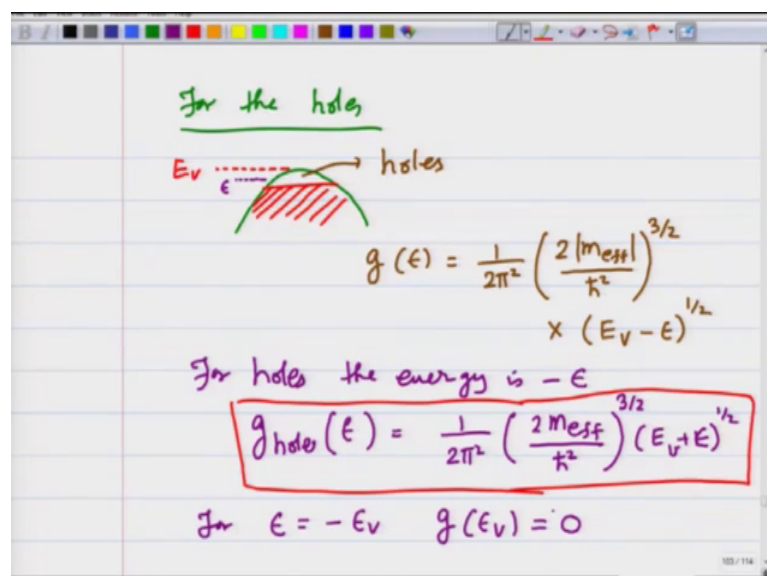
So, the energy or the dependence of the energy on k is pretty much like free particles and this implies $g(E)$ is like that of free particles. So, now to calculate the number of electrons and number of holes what we do is fix our integration limit and let us do that now. Before that let me write the formula, g as a function of E is equal to $\frac{1}{2\pi^2}$ and let me write this m_{eff} over \hbar^2 raised to $3/2$ ϵ raised to $1/2$.

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Now, when we are doing electrons and holes we have to put the proper integration limits. So, let us see for electrons what we are going to do? Electrons if there is a band like this or even like this the electrons are filled near the minimum and if this minimum is written as E_c or E_c Conduction then it is easy to see that I am going to have $g(E)$ is equal to $\frac{1}{2\pi^2} \left(\frac{2m_{eff}}{\hbar^2} \right)^{3/2} (E - E_c)^{1/2}$. So, $g(E_c) = 0$ and it rises, to some energy which is not very high as long as the electrons are up to the level where the energy dependence is k^2 this is what the $g(E)$ is going to be. How about for the holes?

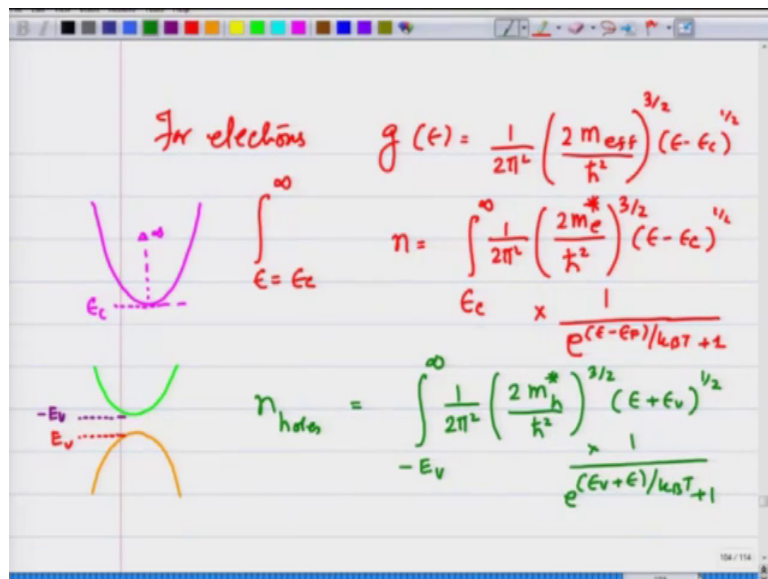
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Again let me make a band like this, the electrons are filled up to certain point and this empty region these are the holes. So, the distribution of holes is exactly the same as that of electrons. So, I am going to have g of E is equal to 1 over 2π square 2 mod m effective. Let me write it mod m effective because electron effective mass here is negative over h cross square raised to 3 by 2 times E_V minus E raised to 1 half where E_V is the top of the valence band, this is E_V . So, I want E_V minus E to be a positive number and this is E , E is right here.

Now for holes, the energy is minus that of the electron. So, g holes for energy of hole equal to ϵ I am going to have 1 over 2π square 2 . Now I can write m effective without the mod sign because hole effective mass becomes positive automatically h cross square raised to 3 by 2 E_V minus minus I will have plus E over raised to 1 half. So, this is the density of states in terms of the hole energy, E_V is a given number. Notice that for ϵ hole equals minus E_V g E_V is equal to 0 . So, this is all consistent. How about the limits?

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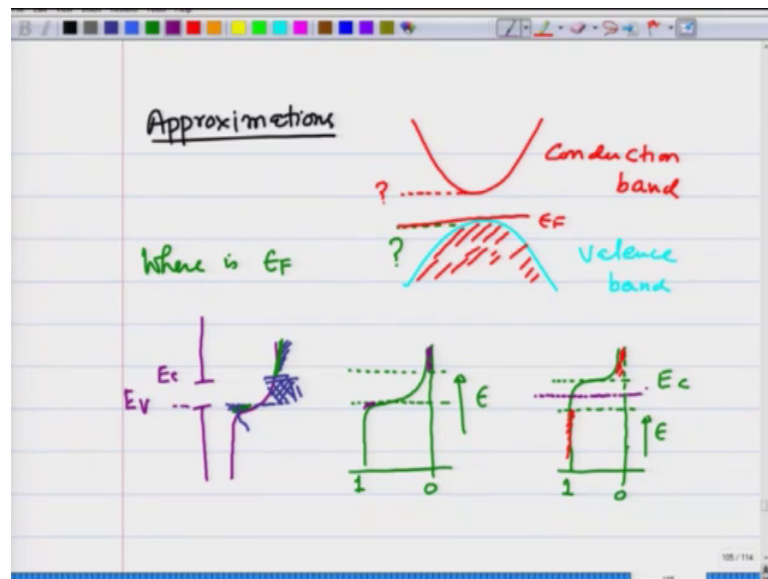


Now, for electrons g E was equal to 1 over 2π square 2 m effective in the conduction band over h cross square raised to 3 by 2 E minus E_C raised to 1 half. So, integration limit is going to be from E equals E_C to infinity. So, this is a band and we are starting E from E_C and all the way up to infinity. So, the number of electrons n is therefore going to be integration E_C to infinity 1 over 2π square 2 m_e by that we will understand that

this is electron mass but effective mass. To indicate it is effective let me also put a star on top. $\frac{3}{2} E_C - E_F$ raised to $1/2$ times the probability which is going to be $1 / e^{(E - E_F) / k_B T + 1}$. This is the number of electrons.

Similarly, number of holes n_{holes} is going to be let me again determine the limits. Here was the electronic band and here is a band for the holes. If this was the top of the valence band was E_V the bottom of the hole band is going to be $-E_V$. So, I am going to have number of holes is equal to $\int_{-E_V}^{\infty} \frac{1}{2\pi^2} \frac{m_{hole}^*}{h^2} \sqrt{E + E_V} \frac{1}{e^{(E - E_F) / k_B T + 1}}$. This is the integration that we got to vary out. So, we have now determined the formula that I should be applying to calculate the number of holes and number of electrons.

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Now, some approximations; so, when I have these conduction band and the valence band, where is epsilon F or the Fermi energy? Is it at the top here question mark, is it at the bottom here question mark or is it somewhere else? So, let us see, at 0 temperature all these states the valence band is completely filled. And therefore, E_F can be at the top of the conduction band there is a possibility. On the other hand as the temperature is increased you will see, now let me make this slightly differently. Here is the gap and here is the valence band. This is E_V this is E_C on top. The Fermi Fermi-Dirac distribution would change like this.

Roughly if E_F was on top of the valence band all the electrons from here would go out and go into the conduction band. There is no electrons in between region all the electrons will be here. Whatever the number of electrons that are getting out of here should be going into the conduction band, but you notice that the probability that has reduced out here is larger than the probability of occupation in the conduction band, because if the Fermi level is here let me make it again more clearly here is the gap the Fermi Fermi-Dirac distribution would go like this.

Where this is 0, 1 and this is E . So, the number of electron that I have gone out of here will be larger than the number of electrons that have reached in the conduction band. On the other hand look at the other possibility. If the Fermi level was at E_C then, the Fermi Fermi-Dirac distribution as temperature rises would be like this. Where this is 1, this is 0 and this is E . In this case hardly any electrons are going out of the valence band, but many electrons are here in the conduction band in both the cases, case in the middle and on the right there is a disbalance. Where as I know in an intrinsic semiconductor the number of electrons that go out from the valence band should go only to the conduction band and so, they should be equal.

So, I am going to have E_F somewhere in the middle. When you are looking at these slides please go over this argument yourself and convince yourself that it will be somewhere in the middle. If it is on top of the valence band or at the bottom of the conduction band there is there will be this is going to be a disbalance in the number of electrons that leave the valence band and those number of electrons that reach the conduction band.

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The image shows a digital whiteboard with handwritten mathematical expressions. At the top, there are two exponential terms: $e^{(E_c - E_F)/k_B T} \sim \frac{E_{gap}}{2}$ and $e^{(E_v + E_F)/k_B T} \sim \frac{E_{gap}}{2}$. Below these, the bandgap energies for Silicon and Germanium are listed: $E_{gap}(Si) \sim 1.2 \text{ eV}$ and $E_{gap}(Ge) \sim 0.8 \text{ eV}$. A red note states $\frac{E_{gap}}{2k_B T} \rightarrow \text{large number}$. Finally, two inequalities are written: $e^{(E_c - E_F)/k_B T} \gg 1$ for electrons and $e^{(E_v + E_F)/k_B T} \gg 1$ for holes.

Now, if it is somewhere in the middle let us look at these numbers. Now, e raised to E minus ϵ_F over $k_B T$ is going to be of the order of E_{gap} divided by 2. Similarly for the holes, E plus ϵ_F over $k_B T$ is going to be of the order of E_{gap} over 2. I should be little careful and write E_c here and E_v here. And E_{gap} for silicon is about 1.2 electron volts. E_{gap} germanium is about 0.8 eV.

So, you can see that E_{gap} divided by $2k_B T$ is going to be a large number. If that is the case, then e raised to $\epsilon_c - \epsilon_F$ over $k_B T$ is going to be much larger than 1 for electrons, and e raised to $\epsilon_v + \epsilon_F$ over $k_B T$ is also going to be much larger than 1 for holes and therefore I can neglect that one. And I can write f_c as e raised to $\epsilon_c - \epsilon_F$ over $k_B T$ and f_v is equal to e raised to $\epsilon_v + \epsilon_F$ over $k_B T$. It is understood that here E refers to ϵ_c here, which is minus ϵ_F missing electron.

(Refer Slide Time: 16:16)

Handwritten notes on a whiteboard:

$$f_e = e^{-(E-E_F)/k_B T}$$

$$f_h = e^{-(E+E_F)/k_B T} \quad \begin{array}{l} E = E_{\text{hole}} \\ = -E_{\text{massy electron}} \end{array}$$

$$n = \frac{1}{2\pi^2} \left(\frac{2m_e}{\hbar^2} \right)^{3/2} \int_{E_c}^{\infty} (E - E_c)^{1/2} e^{-(E-E_c)/k_B T} dE$$

$$\int_{E_c}^{\infty} (E - E_c)^{1/2} e^{-(E-E_c)/k_B T} dE = \int_0^{\infty} \sqrt{z} e^{-\beta z} dz$$

$$\beta = \frac{1}{k_B T}, \quad z = (E - E_c) \quad \frac{1}{\beta^{3/2}} \Gamma\left(\frac{3}{2}\right)$$

Once I have done this the rest of the job is very easy. So, number of electrons is going to be equal to $\frac{1}{2\pi^2} \left(\frac{2m_e}{\hbar^2} \right)^{3/2} \int_{E_c}^{\infty} (E - E_c)^{1/2} e^{-(E-E_c)/k_B T} dE$. The first term is a constant.

It remains out and the second term $\int_{E_c}^{\infty} (E - E_c)^{1/2} e^{-(E-E_c)/k_B T} dE$ can be written as this is raised to 1/2 can be written as integration 0 to infinity square root of $Z e^{-\beta Z} dZ$ where β is $1/k_B T$ and Z equals $\epsilon - \epsilon_c$. This integral is very straightforward it is $\frac{1}{\beta^{3/2}} \Gamma\left(\frac{3}{2}\right)$. This is this function.

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$$\begin{aligned}
 n &= \frac{1}{2\pi^2} \left(\frac{2m_e^*}{\hbar^2} \right)^{3/2} \int_0^\infty \sqrt{z} e^{-\beta z} dz \\
 &= \frac{1}{2\pi^2} \left(\frac{2m_e^*}{\hbar^2} \right)^{3/2} (k_B T)^{3/2} \Gamma\left(\frac{3}{2}\right) \frac{\sqrt{\pi}}{2} \\
 &= 2 \left(\frac{m_e^* k_B T}{2\pi \hbar^2} \right)^{3/2} e^{-(E_c - E_F)/k_B T} \\
 \text{Mistake : } & \frac{1}{2\pi^2} \left(\frac{2m_e^*}{\hbar^2} \right)^{3/2} \int_{E_c}^\infty (E - E_c)^{1/2} e^{-\frac{(E - E_F)}{k_B T}} dE \\
 &= \frac{1}{2\pi^2} \left(\frac{2m_e^*}{\hbar^2} \right)^{3/2} e^{-\frac{(E_c - E_F)}{k_B T}} \int_{E_c}^\infty (E - E_c) e^{-\frac{(E - E_c)}{k_B T}} dE
 \end{aligned}$$

So, let me rewrite this again. So, number of electrons is equal to 1 over 2 pi square 2 m e star over h cross square raised to 3 by 2 integration 0 to infinity square root of Z e raised to minus beta z d Z and this is equal to 1 over 2 pi square 2 m e star over h cross square 3 by 2. 1 over beta raised to 3 by 2 which is k B T raised to 3 by 2 gamma 3 by 2 and gamma 3 by 2 is nothing but I cut this and instead write a square root of pi by 2. And therefore this whole thing can be written as 2 m e star k B T over 2 pi h cross square raised to 3 by 2 and there is a factor outside minus E C minus E F over k B T.

So, what is the mistake I did? If you are careful you would have seen it the integral was 1 over 2 pi square 2 m e star over h cross square raised to 3 by 2 integration E C to infinity E minus E C raised to 1 half e raised to minus E minus E F over k B T d E. And this can be written as 1 over 2 pi square 2 m e star over h cross square raised to 3 by 2 e raised to minus E C minus E F over k B T integration E C to infinity E minus E C e raised to minus E minus E C over k B T d E. Rest of the calculation goes through because of this mistake I had not gotten this term into the picture which I have now taken care off.

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$$n = 2 \left(\frac{m_e^* k_B T}{2\pi \hbar^2} \right)^{3/2} e^{-(E_C - E_F)/k_B T}$$

Number of holes

$$n_h = \frac{1}{2\pi^2} \left(\frac{2m_h^*}{\hbar^2} \right)^{3/2} \int_{-E_V}^{\infty} (E + E_V)^{1/2} e^{-\frac{(E + E_F)}{k_B T}} dE$$

$$= \frac{1}{2\pi^2} \left(\frac{2m_h^*}{\hbar^2} \right)^{3/2} e^{-\frac{E_V}{k_B T}} \int_{0}^{\infty} (E + E_V)^{1/2} e^{-\frac{E + E_V}{k_B T}} dE$$

$$= \frac{1}{2\pi^2} \left(\frac{2m_h^*}{\hbar^2} \right)^{3/2} e^{-\frac{E_V - E_F}{k_B T}} \cdot (k_B T)^{3/2} \frac{\sqrt{\pi}}{2}$$

So, number of electrons in the conduction band is $2 \frac{m_e^* k_B T}{2\pi \hbar^2}$ cross square raised to 3 by 2 $e^{-\frac{E_C - E_F}{k_B T}}$. How about the number of holes? The number of holes n_h is going to be $\frac{1}{2\pi^2} \frac{2m_h^*}{\hbar^2}$ cross square raised to 3 by 2, where m_h^* tells you the effective mass of the hole integration minus E_V to infinity $(E + E_V)^{1/2} e^{-\frac{E + E_F}{k_B T}} dE$. Which now I am going to write as equal to $\frac{1}{2\pi^2} \frac{2m_h^*}{\hbar^2}$ cross square raised to 3 by 2 $e^{-\frac{E_V - E_F}{k_B T}}$ integration minus E_V to infinity $(E + E_V)^{1/2} e^{-\frac{E + E_V}{k_B T}} dE$.

This is exactly the same integral as we encountered in the case of electrons and therefore I can write this as, equal to $\frac{1}{2\pi^2} \frac{2m_h^*}{\hbar^2}$ cross square $e^{-\frac{E_V - E_F}{k_B T}}$ times $(k_B T)^{3/2} \frac{\sqrt{\pi}}{2}$ that is the answer that we get for this. So, the number of holes also comes out to be exactly similar to that of number of electrons. And I get n_h is equal to $\frac{m_h^* k_B T}{2\pi \hbar^2}$ cross square raised to 3 by 2 $e^{-\frac{E_V - E_F}{k_B T}}$. This is the number of holes.

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$$n_h = 2 \left(\frac{m_h^* k_B T}{2\pi \hbar^2} \right)^{3/2} e^{(E_V - E_F)/k_B T}$$

I have used E & ϵ interchangeably

Product of $(n \times n_h)$

$$= 4 \left(m_e^* m_h^* \right)^{3/2} \left(\frac{k_B T}{2\pi \hbar^2} \right)^3 e^{-(E_C - E_F)/k_B T} \times e^{(E_V - E_F)/k_B T}$$
$$= 4 \left(m_e^* m_h^* \right)^{3/2} \left(\frac{k_B T}{2\pi \hbar^2} \right)^3 e^{-(E_C - E_V)/k_B T}$$

Now, what a caution I have used E and epsilon interchangeably. So, that I hope I was going along you understood. Now the product of n the number of electrons and number of holes is therefore going to be equal to $m_e^* m_h^*$ raised to 3 by $2 k_B T$ over $2\pi \hbar$ cross square raised to 3, e raised to minus E_C minus E_F over $k_B T$ times e raised to E_V minus E_F over $k_B T$. And this comes out to be $m_e^* m_h^*$ raised to 3 by $2 k_B T$ over $2\pi \hbar$ cross square raised to 3 e raised to minus E_C minus E_V raised to divided by $k_B T$. So, let us write it again because I want to spend some time on this formula.

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Numa of electrons \times numba of holes

$$n n_h = 4 \left(m_e^* m_h^* \right)^{3/2} \left(\frac{k_B T}{2\pi \hbar^2} \right)^3 e^{-(E_C - E_V)/k_B T}$$

NOTICE: (i) There is no E_F in the formula
(ii) The formula depends only on the energy gap & T.

Use the formula to calculate n or n_h in an intrinsic semiconductor

$$\boxed{n_h = n}$$

So, we get number of electrons times number of holes n times n_h is equal to $m_e^* m_h^*$ raised to 3 by $2 k_B T$ over $2 \pi \hbar^2$ cross square raised to 3 $e^{\frac{-E_C - E_V}{k_B T}}$. Notice 1; there is no E_F in the formula number 2; the formula depends only on the energy gap and temperature T . Therefore this formula is going to be in general valid. No matter where the Fermi level is and this is very important. This is important because I can apply it even when impurities are there and Fermi level get shifted, we come to that later. But right now I can use this formula, use the formula to calculate n or n_h in an intrinsic semiconductor. Using the fact that n_h equals n in an intrinsic semiconductor and I also then use it to calculate the Fermi energy. So, let us do that now.

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The image shows a whiteboard with the following handwritten equations and text:

$$n = n_h \Rightarrow n^2 = 4(m_e^* m_h^*)^{3/2} \left(\frac{k_B T}{2\pi \hbar^2} \right)^3 e^{-\frac{\Delta E_{\text{gap}}}{k_B T}}$$

$$\Rightarrow n = n_h = 2(m_e^* m_h^*)^{3/4} \left(\frac{k_B T}{2\pi \hbar^2} \right)^{3/2} e^{-\frac{\Delta E_{\text{gap}}}{2k_B T}}$$

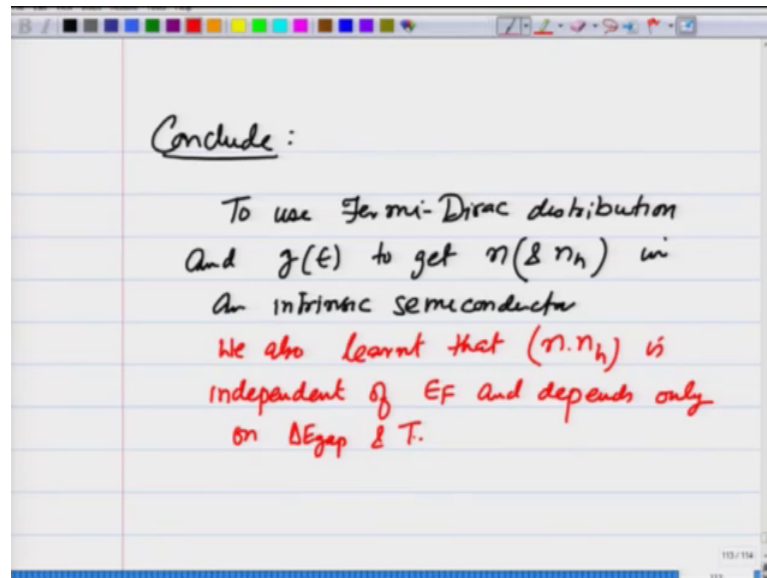
$$\Delta E_{\text{gap}} = (E_C - E_V)$$

The Fermi energy

$$n = 2 \left(\frac{m_e^* k_B T}{2\pi \hbar^2} \right)^{3/2} e^{-(E_C - E_F)/k_B T}$$

So, I have n equals n_h and therefore this implies n square is going to be equal to $m_e^* m_h^*$ raised to 3 by $2 k_B T$ over $2 \pi \hbar^2$ cross square cubed times $e^{\frac{-\Delta E_{\text{gap}}}{k_B T}}$. And therefore n equals n_h equals $m_e^* m_h^*$ raised to 3 by $4 k_B T$ over $2 \pi \hbar^2$ cross square raised to 3 by 2 $e^{\frac{-\Delta E_{\text{gap}}}{2k_B T}}$ where ΔE_{gap} is equal to $E_C - E_V$. So, this gives me the number of electrons or number of holes. How about the Fermi energy? Recall that n was equal to $2 m_e^* k_B T$ over $2 \pi \hbar^2$ cross square raised to 3 by 2 $e^{\frac{-E_C - E_F}{k_B T}}$ and it is similarly was given in terms of $E_F - E_V$. You can use these to calculate E_{gap} and I leave that as an exercise.

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So, what we have done in this lecture is use Fermi-Dirac distribution and g_E to get n and n_h is the same thing in an intrinsic semiconductor that is point number 1. Point number 2 we also learnt that n times n_h is independent of E_F and depends only on ΔE_{gap} and temperature T . In the next lecture we will use these to calculate the number of electrons and holes in a doped semiconductor.

Thank you.