

Introduction to Solid State Physics
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Lecture - 37
Reciprocal Lattice Vectors Part-II

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$\vec{R} = n_1 \vec{a}_1 + n_2 \vec{a}_2 + n_3 \vec{a}_3$ ← 3D lattice

$\rho(\vec{r}) = \rho(\vec{r} + \vec{R})$

$\vec{r} \rightarrow \vec{r} + \vec{R}$

$\therefore \rho(\vec{r} + \vec{R}) = \sum_{\vec{G}} \rho_{\vec{G}} e^{i \vec{G} \cdot (\vec{r} + \vec{R})}$

$= \underbrace{\sum_{\vec{G}} \rho_{\vec{G}} e^{i \vec{G} \cdot \vec{r}}}_{\rho(\vec{r})} e^{i \vec{G} \cdot \vec{R}}$

T.I $\Rightarrow \rho(\vec{r}) = \rho(\vec{r} + \vec{R}) \therefore e^{i \vec{G} \cdot \vec{R}} = 1$

Welcome back to the discussion. And in the last lecture, we had seen that you can take the Fourier transform of a lattice of points whose density at any point is given by this expression. In this is in three dimension, we have done it for one dimension, but this is the general case of a 3D lattice of points. And the requirement rho of r is the number density of points at any point r in the lattice.

So, you have a lattice of points a periodic structure, here I am showing it in two dimension ok. And for any point r in this lattice, the number density at that point is given by rho of r. And the requirement of translation in variance which is an a very important aspect of a Bravais lattice translational invariance requires that if I go from this point r and I translate by an amount which is given by my lattice vector r bar, where this is r plus capital R, then the density at r plus capital R namely at this point should be exactly identical to the density at this point and that is written out here.

And this then satisfies the condition of a Bravais lattice, where this vector R is n 1 a 1 plus n 2 a 2 plus n 3 a 3. This is your translation vector R of the lattice; n 1, n 2, n 3 are

the indices of the components along a_1, a_2, a_3 , which are the fundamental translation vectors of the lattice. And this translation invariance, so then rho of r you can do a Fourier expansion, rho of r you can write it as summation over some vector G rho G are the Fourier components e raised to $i \bar{G} \cdot \bar{r}$ where I am taking the Fourier expansion in terms of this vector G. And if I do translation invariance, then the requirement for translation invariance gives me this condition that e raised to $\bar{G} \cdot \bar{R}$, R is your lattice translation vector e raised to $i \bar{G} \cdot \bar{R}$ should be equal to 1.

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$$e^{i \bar{G} \cdot \bar{R}} = 1 ; \quad \bar{R} = n_1 \bar{a}_1 + n_2 \bar{a}_2 + n_3 \bar{a}_3$$

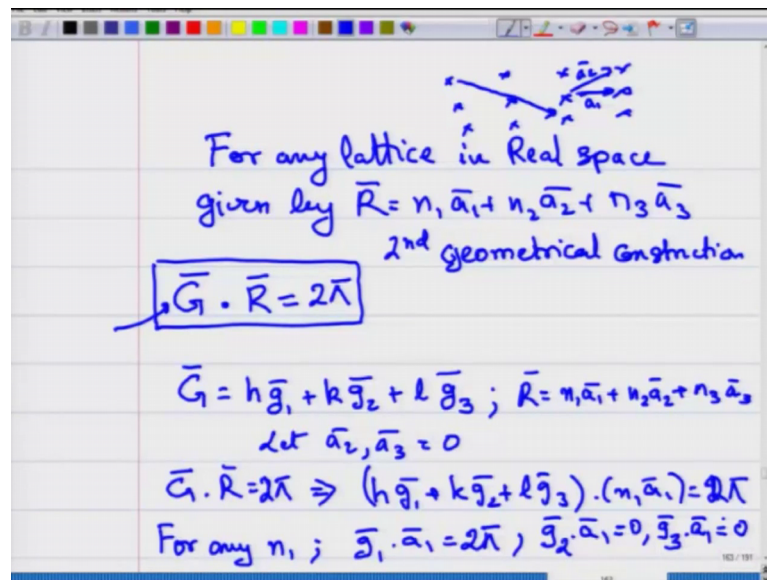
$$\Rightarrow \boxed{\bar{G} \cdot \bar{R} = 2\pi m}$$

Construct something called as the
Reciprocal lattice \rightarrow associated with
a crystal

'K' \rightarrow I_B

And this suggests that $\bar{G} \cdot \bar{R}$ is $2\pi m$.

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So, for any lattice in real space given by \bar{R} is equal to $n_1 \bar{a}_1 + n_2 \bar{a}_2 + n_3 \bar{a}_3$. For any lattice which is in real space your actual three-dimensional space that we are living in. For any lattice in real space where you can generate the lattice using this translation vector \bar{R} corresponding to any such lattice you can generate another lattice which is given by this vector \bar{G} , where $\bar{G} \cdot \bar{R}$ should be equal to 2π . And this vector \bar{G} is corresponding to taking the Fourier transform of the real lattice. So, corresponding to any lattice in real space which is generated by using this condition namely you have set of points in real space.

So, given a crystalline solid you first make a Bravais lattice using this translation vector, where you can go from any point on the lattice to any other point using this sort of a translation vector. You will have to define your \bar{a}_1 and \bar{a}_2 appropriately whatever is your choice that you want to do, and then using this you can generate your lattice points, but you can make another construction. And that second geometrical construction is with these vectors \bar{G} , where these vectors \bar{G} satisfy the condition that $\bar{G} \cdot \bar{R}$ is equal to 2π ok.

Now, we write this vector in terms of we have three components of this vector. We write it has h times \bar{g}_1 plus k times \bar{g}_2 plus l times \bar{g}_3 ; and \bar{R} you already know is equal to $n_1 \bar{a}_1 + n_2 \bar{a}_2 + n_3 \bar{a}_3$.

Now, let us say let a_2 and a_3 be equal to 0 ok, then what does this condition mean $\bar{G} \cdot \bar{R}$ is equal to 1 implies h times g_1 plus k times g_2 plus l times g_3 dot $n_1 a_1$ is equal to 1, you just satisfy put in this a_2 equal to a_3 and put it back here this is your condition. For any arbitrary value of n_1 , this condition can be satisfied by considering g_1 dot a_1 should be equal to 2π and so the condition that $\bar{G} \cdot \bar{R}$ is equal to 2π , g_2 dot a_1 should be equal to 0 and g_3 dot a_1 is equal to 0.

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$$\bar{G} \cdot \bar{R} = 2\pi \quad e^{i\bar{G} \cdot \bar{R}} = 1$$

$$\rightarrow \bar{g}_1 \cdot \bar{a}_1 = 2\pi, \bar{g}_2 \cdot \bar{a}_1 = 0, \bar{g}_3 \cdot \bar{a}_1 = 0$$

$$\parallel \text{If } \bar{a}_2 \neq 0, \bar{a}_1, \bar{a}_3 = 0$$

$$\rightarrow \bar{g}_2 \cdot \bar{a}_2 = 2\pi, \bar{g}_1 \cdot \bar{a}_2 = 0, \bar{g}_3 \cdot \bar{a}_2 = 0$$

$$\rightarrow \bar{g}_3 \cdot \bar{a}_3 = 2\pi, \bar{g}_1 \cdot \bar{a}_3 = 0, \bar{g}_2 \cdot \bar{a}_3 = 0$$

$$\boxed{\bar{g}_i \cdot \bar{a}_j = 2\pi \delta_{ij}}; \delta_{ij} = 0 \quad i \neq j$$

$$i=1, j=1 \quad \bar{g}_1 \cdot \bar{a}_1 = 2\pi, i=1, j=2 \Rightarrow \bar{g}_1 \cdot \bar{a}_2 = 0, \dots$$

So, the condition that $\bar{G} \cdot \bar{R}$ is equal to 2π gives us our first condition that g_1 dot a_1 should be equal to 2π , g_2 dot a_1 should be equal to 0, g_3 dot a_1 should be equal to 0. Similarly, take a_2 not equal to 0, a_1 and a_3 equal to 0. And you will get the condition that g_2 dot a_2 should be equal to 2π , g_2 dot a_1 should be equal to 0, and g_3 dot a_2 and g_1 dot a_2 should be equal to 0 and g_3 dot a_2 should be equal to 0.

So, this says the above conditions g_1 dot a_1 equal to 2π , and g_2 and g_3 dot a_1 is equal to 0 suggest that g_2 and g_3 are perpendicular to a_1 . So, this directly gives us way to construct g_1 , g_2 and g_3 . Similarly, g_2 is having some component along a_2 , but g_1 and g_3 are perpendicular to a_2 ok. And you can write down the same thing g_3 dot a_3 is equal to 2π , and g_1 dot a_3 should be equal to 0, and g_2 dot a_3 should be equal to 0.

So, the condition that $\bar{G} \cdot \bar{R}$ which is coming from the condition that $\bar{G} \cdot \bar{R}$ is equal to 1 this gives us this condition gives us a way to construct these vectors g_1 , g_2 and g_3 . In general I can rewrite the all these conditions as g_i dot a_j is $2\pi \delta_{ij}$

where δ_{ij} is equal to 0 for $i \neq j$. So, you put i equal to 1, and j equal to 1, then $\mathbf{g}_1 \cdot \mathbf{a}_1$ is equal to 2π ; i equal to 1 and j equal to 2 implies $\mathbf{g}_1 \cdot \mathbf{a}_2$ is equal to 0 and so on. So, this is a condition that one uses to find out these vectors $\mathbf{g}_1, \mathbf{g}_2$ and \mathbf{g}_3 .

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Real space lattice $\bar{\mathbf{R}} = n_1 \bar{\mathbf{a}}_1 + n_2 \bar{\mathbf{a}}_2 + n_3 \bar{\mathbf{a}}_3$

$\bar{\mathbf{G}}, \bar{\mathbf{G}} \cdot \bar{\mathbf{R}} = 2\pi \Rightarrow$

$\bar{\mathbf{G}} = h \bar{\mathbf{g}}_1 + k \bar{\mathbf{g}}_2 + l \bar{\mathbf{g}}_3 \Rightarrow \bar{\mathbf{g}}_i \cdot \bar{\mathbf{a}}_j = 2\pi \delta_{ij}$

$$\bar{\mathbf{g}}_1 = 2\pi \frac{\bar{\mathbf{a}}_2 \times \bar{\mathbf{a}}_3}{\bar{\mathbf{a}}_1 \cdot (\bar{\mathbf{a}}_2 \times \bar{\mathbf{a}}_3)}$$

$\bar{\mathbf{g}}_1 \cdot \bar{\mathbf{a}}_1 = 2\pi \frac{\bar{\mathbf{a}}_1 \cdot (\bar{\mathbf{a}}_2 \times \bar{\mathbf{a}}_3)}{\bar{\mathbf{a}}_1 \cdot (\bar{\mathbf{a}}_2 \times \bar{\mathbf{a}}_3)} = 2\pi$

So, corresponding to any real space lattice, real space lattice which is generated by the translation vector $n_1 \mathbf{a}_1 + n_2 \mathbf{a}_2 + n_3 \mathbf{a}_3$, one can find another lattice where $\bar{\mathbf{G}} \cdot \bar{\mathbf{R}}$ is equal to 2π . And this gives the condition and this $\bar{\mathbf{G}}$ is written as h times $\bar{\mathbf{g}}_1$ plus k times $\bar{\mathbf{g}}_2$ plus [vocalize-noise] using the above condition if you know $\mathbf{a}_1, \mathbf{a}_2, \mathbf{a}_3$, one can construct $\mathbf{g}_1, \mathbf{g}_2, \mathbf{g}_3$ by using the conditions $\mathbf{g}_i \cdot \mathbf{a}_j$ is equal to $2\pi \delta_{ij}$.

Now, how can we write, is there a simple way to write this down? There is of course, a simple way to write this down $\bar{\mathbf{g}}_1$ can be written as $2\pi \mathbf{a}_2 \times \mathbf{a}_3$ divided by $\mathbf{a}_1 \cdot \mathbf{a}_2 \times \mathbf{a}_3$, one expression which actually satisfies this condition is writing $\bar{\mathbf{g}}_1$ like this. And one can easily check $\bar{\mathbf{g}}_1 \cdot \mathbf{a}_1$ is $2\pi \mathbf{a}_1 \cdot \mathbf{a}_2 \times \mathbf{a}_3$ divided by $\mathbf{a}_1 \cdot \mathbf{a}_2 \times \mathbf{a}_3$, this is nothing else but 2π .

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Handwritten notes on a whiteboard:

$$\bar{g}_1 \cdot \bar{a}_2 = 2\pi \frac{\bar{a}_2 \cdot (\bar{a}_2 \times \bar{a}_3)}{\bar{a}_1 \cdot (\bar{a}_2 \times \bar{a}_3)} = 0$$

$$\bar{g}_1 \cdot \bar{a}_3 = 0$$

$$\bar{g}_i \cdot \bar{a}_j = 2\pi \delta_{ij}$$

$$\bar{g}_1 = 2\pi \frac{\bar{a}_2 \times \bar{a}_3}{\bar{a}_1 \cdot (\bar{a}_2 \times \bar{a}_3)}$$

$$\bar{g}_2 = 2\pi \frac{\bar{a}_3 \times \bar{a}_1}{\bar{a}_2 \cdot (\bar{a}_3 \times \bar{a}_1)}$$

$$\bar{g}_3 = 2\pi \frac{\bar{a}_1 \times \bar{a}_2}{\bar{a}_3 \cdot (\bar{a}_1 \times \bar{a}_2)}$$

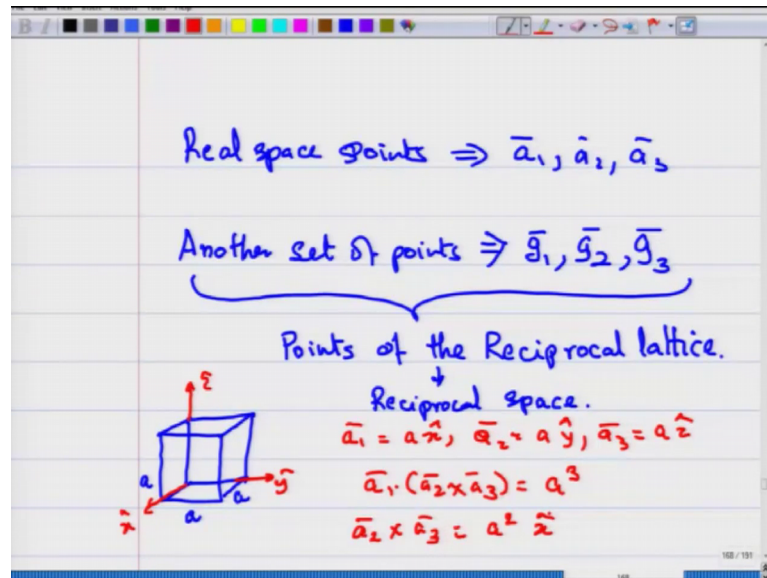
$\bar{a}_1 \cdot (\bar{a}_2 \times \bar{a}_3)$
 = volume of cell $(\bar{a}_1, \bar{a}_2, \bar{a}_3)$
 $\bar{a}_2 \cdot (\bar{a}_3 \times \bar{a}_1) = \bar{a}_3 \cdot (\bar{a}_1 \times \bar{a}_2)$

And you can show similarly that $\bar{g}_1 \cdot \bar{a}_2$ will be 2π a 2 dot a 2 cross a 3 is equal to a 1 dot a 3 this is 0, because a 2 is along the plane a 2 is perpendicular to a 2 cross a 3. So, this dot product has to be 0. And similarly $\bar{g}_1 \cdot \bar{a}_3$ is equal to 0, a $\bar{g}_1 \cdot \bar{a}_3$ is equal to 0. So, $\bar{g}_i \cdot \bar{a}_j$ which is equal to $2\pi \delta_{ij}$ can be generated using the \bar{g}_i which actually satisfy this condition is given as 2π a 2 cross a 3 divided by a 1 dot a 2 cross a 3, \bar{g}_2 which is 2π a 3 cross a 1 divided by a 2 dot a 3 cross a 1, and \bar{g}_3 is 2π a 1 cross a 2 divided by a 3 dot a 1 cross a 2.

One can show that if you use $\bar{g}_1, \bar{g}_2, \bar{g}_3$ of this form they will satisfy this condition. So, from $\bar{a}_1, \bar{a}_2, \bar{a}_3$, you can generate another set of vectors $\bar{g}_1, \bar{g}_2, \bar{g}_3$ which can be expressed in terms of $\bar{a}_1, \bar{a}_2, \bar{a}_3$, but they satisfy a very fundamental condition that $\bar{g}_i \cdot \bar{a}_j$ is equal to $2\pi \delta_{ij}$, this is a very necessary condition.

Now, a few words about this denominator $\bar{a}_1 \cdot \bar{a}_2 \times \bar{a}_3$ is nothing else but the volume of the cell which is defined by these vectors $\bar{a}_1, \bar{a}_2, \bar{a}_3$. And $\bar{a}_1 \cdot \bar{a}_2 \times \bar{a}_3$ is equal to $\bar{a}_2 \cdot \bar{a}_1 \times \bar{a}_3$ which is also equal to $\bar{a}_3 \cdot \bar{a}_1 \times \bar{a}_2$. So, at the denominator, all of this is nothing else but the volume of the real space, this is the volume of the cell in real space which is defined by your fundamental translation vectors \bar{a}_1, \bar{a}_2 and \bar{a}_3 .

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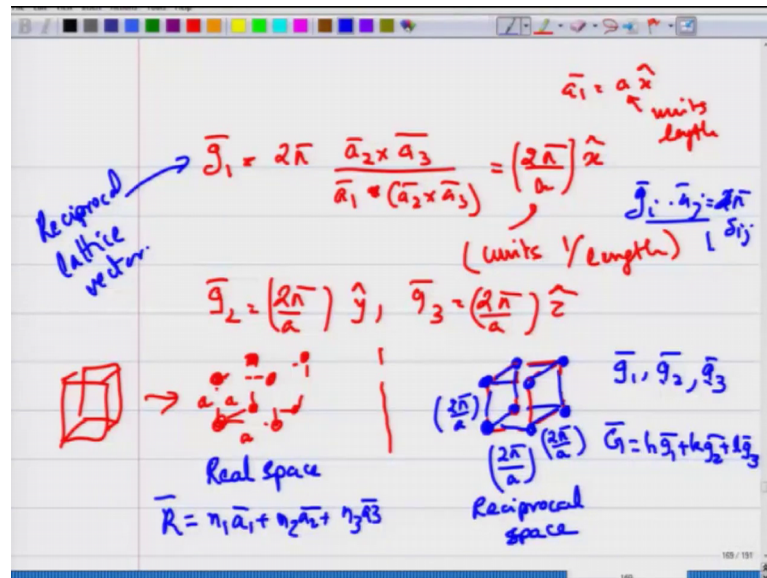


The points which are generated so I have said that you have a set of real space points which are generated with a 1, a 2 and a 3 and you have another you can generate another set of points which are g 1, g 2 and g 3. The points which are generated using this are called the points of the reciprocal lattice. This is a very important concept that corresponding to every crystal, you have a real space lattice. And corresponding to every crystal you can also do another geometric construction where you can define the reciprocal space lattice g 1, g 2 and g 3.

So, corresponding to every real space point which is defined by linear combinations of a 1, a 2 and a 3, you can generate another set of geometrical points both are geometrical constructions for a crystal you can generate a real space lattice using linear combinations of a 1, a 2, a 3. There is another set of points which you can generate not in the real space ok. There is another set of points which you can generate with these vectors g 1, g 2 and g 3. Those are the points which are generated which are called as the points belonging to a reciprocal lattice and they are in a space called as the reciprocal space.

Simply put let us take a cubic lattice with lattice constant as a. And these are your x, y and z. For this cell, you can define a 1 as a x cap, a 2 as a y cap, a 3 as a z cap. a 1 dot a 2 cross a 3 will be the volume a cube, a 2 cross a 3 will be equal to a square y cap cross z cap which is x cap.

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So, g_1 which is 2π by a_2 cross a_3 divided by a_1 dot a_2 cross a_3 will be equal to 2π by a . So, g_1 is along the a_1 direction, if you recall a_1 is a \hat{x} cap, so g_1 is along the \hat{x} direction, but whereas in the real space a_1 is the reciprocal space vector has units of 1 over length. The real space has units of length, the reciprocal space has units of 1 over length. Similarly, g_2 will be 2π by a_2 and g_3 will be 2π by a_3 .

So, corresponding to a crystal which is cubic, you can in real space generate cubic set of points these are your cubic lattice points in real space whose spacings are a . And you can define your lattice vectors a_1, a_2, a_3 . And using those lattice vectors a_1, a_2, a_3 , you can generate the full space of lattice points, but you can also generate using g_1, g_2, g_3 , which is also a square lattice another set of points. Let me draw it in blue.

So, corresponding to this crystal with real space, you can also generate equivalently another set of points which is governed by g_1, g_2, g_3 as given here, which also happens to be a square lattice, but now the lattice constants here are 2π by a rather than units of length they are units of inverse of length. So, this is in real space, this is in reciprocal space, reciprocal because they are inverse of length ok.

So, corresponding to every real lattice, you can generate another lattice using your reciprocal lattice vectors. These vectors are called as reciprocal lattice vectors. And here you have a translation vector \vec{r} which is written as linear combination of $n_1 a_1 + n_2 a_2 + n_3 a_3$

$n_2 a_2 + n_3 a_3$, here your vector which is your translation vector in reciprocal space is h times g_1 plus k times g_2 plus l times g_3 .

So, you have for every lattice you can do two types of geometrical constructions, one is the real space which you are familiar with a_1, a_2, a_3 , linear combinations of a_1, a_2, a_3 . And corresponding to the same lattice in reciprocal space, you can generate another set of points with these vectors g_1, g_2, g_3 which you can calculate explicitly using the condition that $g_i \cdot a_j$ is equal to $2\pi \delta_{ij}$. This condition gave us these expressions which help us to calculate your reciprocal lattice points and the reciprocal lattice is inverse of length. And as you see for the cubic lattice it is $2\pi/a$; a is the lattice constant of your cubic lattice.

So, this is a very important concept of a reciprocal lattice corresponding to every real lattice you have also a reciprocal lattice. And it is very useful for understanding the properties of the crystal. By using the reciprocal lattice for a given lattice, you can understand various properties like diffraction of waves through the crystal which we will do next in our coming lecture.