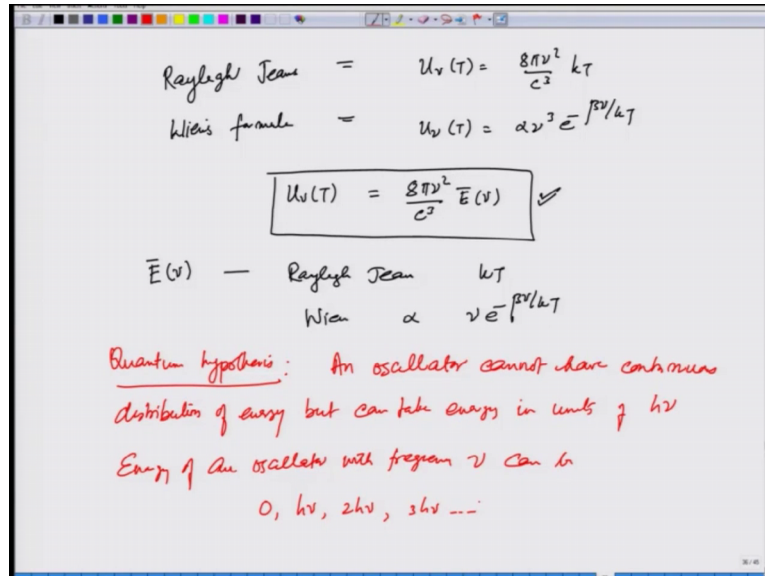


Introduction to Quantum Mechanics
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Lecture - 07

Black Body Radiation VII - Quantum hypothesis and Planck's distribution formula

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We have seen that there are 2 formula; Rayleigh Jean's and Wien's formula that give you radiation density; this gives you $u_{\nu} T$ equals $\frac{8\pi\nu^2}{c^3} kT$ and this gives you $u_{\nu} T$ equals $\sum \alpha \nu^3 e^{-\beta\nu/kT}$ and (Refer Time: 00:47) is somewhere in between because the curve does not fit both ends one limit or the other. And to resolve this and also $u_{\nu} T$ should be of the form $\frac{8\pi\nu^2}{c^3} E_{\nu}$. This is the correct formula.

This formula is derived in an exact manner. So, what we need is E_{ν} which according to Rayleigh Jeans is kT and according to Wien is proportional to $\nu e^{-\beta\nu/kT}$ and truth is somewhere in between. So, this is resolved by quantum hypothesis, it says that an oscillator cannot have continuous distribution of energy, but can take energy in units of $h\nu$; h is some constant. So, energy of an oscillator with frequency ν can be $0, h\nu, 2h\nu, 3h\nu$ and so on, but cannot be continuous classically we had continuous distribution, but now it cannot be continuous.

Let us see how this helps in getting the right formula or the plan's formula for correct formula for the distribution of spectral density of the radiation.

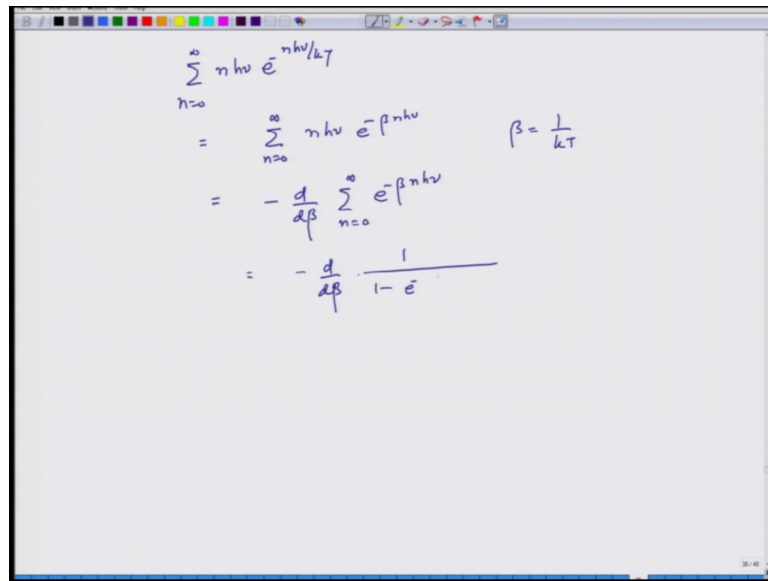
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Suppose an oscillator has energy $n h \nu$
 At temperature T , the probability of it having energy $n h \nu$
 $\propto e^{-n h \nu / k T}$
 Probability = $\frac{e^{-n h \nu / k T}}{\sum_{n=0}^{\infty} e^{-n h \nu / k T}}$
 $= \frac{e^{-n h \nu / k T}}{(1 - e^{-h \nu / k T})}$
 The average energy $\bar{E}(\nu) = \sum_{n=0}^{\infty} n h \nu \times \frac{e^{-n h \nu / k T}}{(1 - e^{-h \nu / k T})}$
 $= \frac{1}{1 - e^{-h \nu / k T}} \sum_{n=1}^{\infty} n h \nu e^{-n h \nu / k T}$

So, suppose an oscillator has energy $n h \nu$ and at temperature T , the probability of it having energy $n h \nu$ is proportional to $e^{-n h \nu / k T}$; that is the probability proportional to. So, the probability will be equal to $e^{-n h \nu / k T}$ and I sum over all this n and $e^{-n h \nu / k T}$ and it vary from 0 to infinity, energy could be all the way up to infinity n starting from 0. So, this is the probability because this is proportional to $e^{-n h \nu / k T}$, I normalized it by this and you can see that the net the probability of having any value is going to be one; some energy, right.

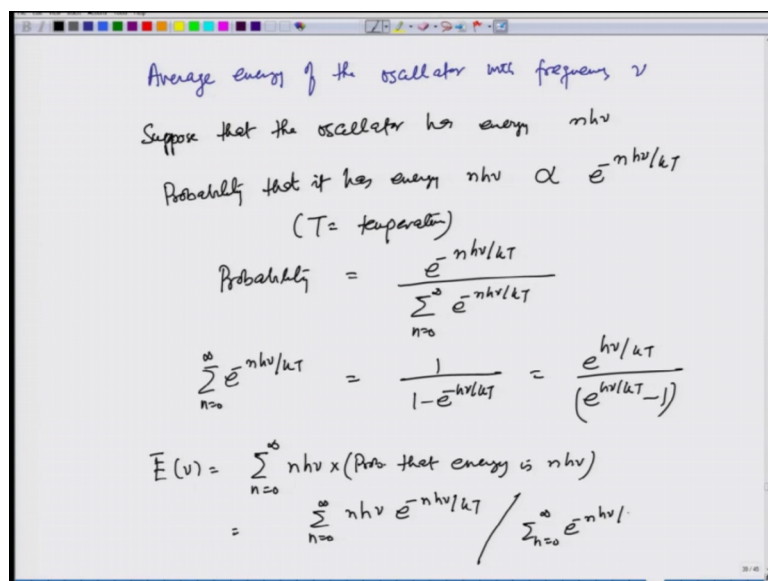
So, now this is the probability which I can write as $e^{-n h \nu / k T}$ divided by I can add this as the geometric series which is going to be $1 - e^{-h \nu / k T}$. So, the average energy $\bar{E}(\nu)$; $\bar{E}(\nu)$ is going to be energy $n h \nu$ times the probability of having that energy which is $e^{-n h \nu / k T}$ over $1 - e^{-h \nu / k T}$ summed over n , n equal 0 to infinity. So, I am going to write this as $1 / (1 - e^{-h \nu / k T}) \sum_{n=1}^{\infty} n h \nu e^{-n h \nu / k T}$ and have to evaluate this let us do that.

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$$\begin{aligned} & \sum_{n=0}^{\infty} n h \nu e^{-n h \nu / k T} \\ &= \sum_{n=0}^{\infty} n h \nu e^{-\beta n h \nu} \quad \beta = \frac{1}{k T} \\ &= -\frac{d}{d\beta} \sum_{n=0}^{\infty} e^{-\beta n h \nu} \\ &= -\frac{d}{d\beta} \frac{1}{1 - e^{-\beta h \nu}} \end{aligned}$$

So, to calculate summation $n h \nu e^{-n h \nu / k T}$ n equal to 0 to infinity, I am going to write this as write $k T$ equals β . So, I am going to write this as summation n equals 0 to infinity $n h \nu e^{-\beta n h \nu}$ where β is $1 / k T$ and this can be written as minus d by $d \beta$ of summation n equals 0 to infinity $e^{-\beta n h \nu}$, but this sum we just now did. So, this is minus d by $d \beta$ of $1 / (1 - e^{-\beta h \nu})$ even this energies we want to calculate the average energy of the oscillator with frequency ν .

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Average energy of the oscillator with frequency ν
Suppose that the oscillator has energy $n h \nu$
Probability that it has energy $n h \nu \propto e^{-n h \nu / k T}$
($T = \text{temperature}$)

$$\text{Probability} = \frac{e^{-n h \nu / k T}}{\sum_{n=0}^{\infty} e^{-n h \nu / k T}}$$
$$\sum_{n=0}^{\infty} e^{-n h \nu / k T} = \frac{1}{1 - e^{-h \nu / k T}} = \frac{e^{h \nu / k T}}{(e^{h \nu / k T} - 1)}$$
$$\begin{aligned} \bar{E}(\nu) &= \sum_{n=0}^{\infty} n h \nu \times (\text{Prob that energy is } n h \nu) \\ &= \frac{\sum_{n=0}^{\infty} n h \nu e^{-n h \nu / k T}}{\sum_{n=0}^{\infty} e^{-n h \nu / k T}} \end{aligned}$$

So, that let us suppose that the oscillator has energy $n h \nu$ then according to Boltzmann, the probability that it has energy $n h \nu$, this proportional to e raised to minus $n h \nu$ over $k T$ where T is the temperature.

And therefore, probability by itself is going to be e raised to minus $n h \nu$ over $k T$ divided by summation n equals 0 to infinity e raised to minus $n h \nu$ over $k T$. Now we have normalized the whole thing, right. Now summation n equal 0 to infinity e raised to minus $n h \nu$ over $k T$ is nothing but 1 over 1 minus e raised to minus $h \nu$ over $k T$ which is equal to e raised to $h \nu$ over $k T$ over e raised to $h \nu$ over $k T$ minus 1 .

So, this we have calculated. Therefore, average energy $e \nu$ is going to be equal to energy $n h \nu$ times probability that energy is $n h \nu$ summed over n equal 0 to infinity. So, I calculate multiply by the probability of it having energy $n h \nu$ and then sum over. So, this comes out to be equal to summation n equal 0 to infinity $n h \nu e$ raised to minus $n h \nu$ over $k T$ divided by that whole thing summation n equal 0 to infinity e raised to minus $n h \nu$ over $k T$.

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The image shows a whiteboard with handwritten mathematical derivations. The top part shows the calculation of the average energy $\sum n h \nu e^{-n h \nu / k T}$ by differentiating the partition function $\sum_{n=0}^{\infty} e^{-n \beta h \nu}$ with respect to β . The bottom part shows the final expression for the average energy $E(\nu)$ in terms of $h \nu$ and $k T$.

$$\begin{aligned} \sum n h \nu e^{-n h \nu / k T} &= \sum_{n=0}^{\infty} n h \nu e^{-n \beta h \nu} & \beta &= \frac{1}{k T} \\ &= -\frac{d}{d\beta} \sum e^{-n h \nu \beta} \\ &= -\frac{d}{d\beta} \frac{1}{1 - e^{-h \nu \beta}} \\ &= \frac{1}{(1 - e^{-h \nu \beta})^2} \times e^{-h \nu \beta} \times h \nu \\ &= \frac{h \nu e^{h \nu / k T}}{(1 - e^{-h \nu / k T})^2} \end{aligned}$$

$$\begin{aligned} E(\nu) &= \frac{h \nu e^{-h \nu / k T}}{(1 - e^{-h \nu / k T})^2} \times \frac{(1 - e^{-h \nu / k T})}{(1 - e^{-h \nu / k T})} \\ &= \frac{h \nu e^{-h \nu / k T}}{1 - e^{-h \nu / k T}} = \left(\frac{h \nu}{e^{h \nu / k T} - 1} \right) \end{aligned}$$

So, let us now calculate this, to do this I do a trick summation $n h \nu e$ raised to minus $n h \nu$ over $k T$, I write as summation n equal 0 to infinity $n h \nu e$ raised to minus $n \beta h \nu$ where β is 1 over $k T$ and this can be written as minus d by $d \beta$ of summation e raised to minus $n h \nu \beta$, but this we have calculated. This comes out to be minus d by

d beta of 1 over 1 minus e raised to minus h nu beta which is then equal to 1 over 1 minus e raised to minus h nu beta square times e raised to minus h nu beta times h nu.

So, this is equal to h nu e raised to minus h nu over k T divided by 1 minus e raised to minus h nu over k T whole square. Therefore, e average nu is equal to h nu e raised to minus h nu over k T over 1 minus e raised to minus h nu over k T square times whatever we had calculated earlier as summation over e raised to minus n h nu over k T and that is 1 over 1 minus e raised to minus h nu over k T which comes out to be h nu e raised to minus h nu over k T divided by 1 minus e raised to minus h nu over k T which is equal to h nu over e raised to h nu over k T minus 1.

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Using quantum hypothesis & Boltzmann probability

$$\bar{E}(v) = \frac{h\nu}{e^{h\nu/kT} - 1}$$

$$U_\nu(T) = \frac{8\pi\nu^2}{c^3} \cdot \frac{h\nu}{e^{h\nu/kT} - 1} = \frac{8\pi h \nu^3}{c^3} \cdot \frac{1}{e^{h\nu/kT} - 1}$$

Notice

- $U_\nu(T) = \nu^3 f(\nu/T)$
- $\int_{\nu/T \rightarrow \infty}$ (small λ) $e^{h\nu/kT} - 1 \sim e^{h\nu/kT}$
 $U_\nu(T) = \frac{8\pi h \nu^3}{c^3} e^{-h\nu/kT}$ (Wien's formula)
- $\int_{\nu/T \rightarrow 0}$ $e^{h\nu/kT} - 1 \sim 1 + h\nu/kT$
 $U_\nu(T) = \frac{8\pi \nu^2}{c^3} kT$ (Rayleigh-Jeans)

So, what we have found is that using quantum hypothesis and Boltzmann probability that average energy of the oscillator is h nu over e raised to h nu over k T minus 1 and therefore, u nu T is going to be 8 pi nu square over C cubed h nu over e raised to h nu over k T minus 1, notice which is 8 pi h nu cubed over C cubed 1 over e raised to h nu over k T minus 1, notice u nu T is of the form nu cubed f nu over T number 1. Number 2 if nu over T goes to infinity that is very large or small lambda then e raised to h nu over k T minus 1 will go to e raised to h nu over k T and u nu T becomes 8 pi h nu cubed over C cubed e raised to minus h nu over k T which is Wien's formula.

And if ν/T goes to 0 which is very small, then e raised to $h\nu/kT$ goes to 1 plus $h\nu/kT$ and u_ν/T goes to $8\pi h\nu^3/c^3 kT$ which is Rayleigh Jeans.

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The formula

$$u_\nu(T) = \frac{8\pi h\nu^3}{c^3} \frac{1}{e^{h\nu/kT} - 1}$$

fits the experimental results

It confirms Quantum Hypothesis

$$\int u_\nu(T) d\nu = \int_0^\infty \frac{8\pi h\nu^3}{c^3} \frac{1}{e^{h\nu/kT} - 1} d\nu$$

$$\nu/T = x$$

$$= \int_0^\infty \frac{8\pi h T^3 x^3}{c^3} \frac{1}{e^{hx/k} - 1} T dx$$

$$= \frac{8\pi h}{c^3} T^4 \int_0^\infty \frac{x^3}{e^{hx/k} - 1} dx \propto T^4$$

So, this formula derived by quantum hypothesis goes correctly to Wien's formula and Rayleigh Jean's formula in the right limits much more important, the formula u_ν/T equals $8\pi h\nu^3/c^3 kT$ fits the experiments. So, this is the formula for the spectral density and it confirms quantum hypothesis that is that the oscillator with frequency ν can take energies only in units of $h\nu$.

How about Stefan Boltzmann's law? So, let us integrate $u_\nu/T d\nu$, it comes out to be $8\pi h\nu^3/c^3 kT$ going from 0 to infinity, substitute ν/T to be equal to some x then this formula becomes $8\pi h T^3 x^3/c^3 kT dx$ which is nothing, but $8\pi h/c^3 k T^4 \int_0^\infty \frac{x^3}{e^{hx/k} - 1} dx$ which is proportional T^4 and you can determine the coefficient for the Stefan Boltzmann coefficient for this.

So, to conclude this lecture what we have shown you is that the Rayleigh Jean's formula and the Wien's formula are 2 limits of the correct formula which is derived from

quantum hypothesis. So, from now on we assume that the energy levels of an oscillator are quantized.

What other evidences for quantization could be there that we will see in the next lecture that you can extend this hypothesis further to include radiation that radiation should also be quantized and also a mechanical oscillators should be quantized. So, far we have said only electromagnetic oscillators, those oscillators that are giving out radiation are quantized, we will see that mechanical oscillators will also be quantized and they explain the behavior of specific heat of solids at low temperatures. So, that the evidence for quantization becomes stronger and stronger and then will take off from there.