

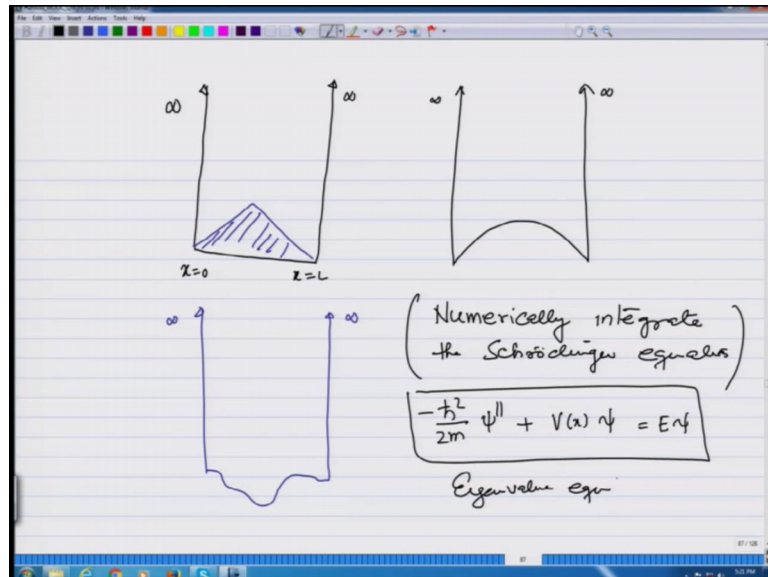
Introduction to Quantum Mechanics
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Lecture - 03

Numerical solution of a one dimensional Schrodinger equation for bound states- I

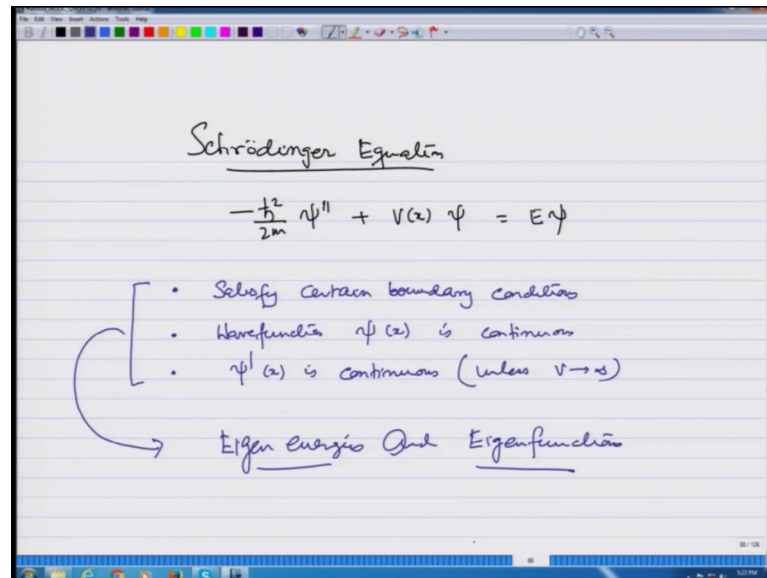
In the previous 2 lectures, we have considered 2 cases; one of a particle moving in a delta function potential and particle moving in a finite size box for which the Schrodinger equation could be solved analytically; however, all potentials are not; I mean able to see this kind of treatment.

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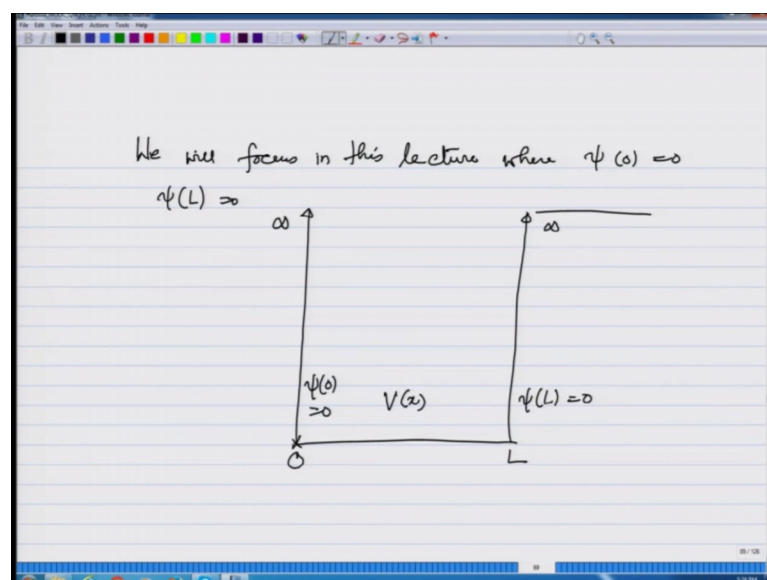
For example, even if I take a potential where the potential becomes infinity at the boundaries let us say x equal 0 and x equals 1, but change the potential in between suppose this like this or in the potential which is like this or in general a potential which is some arbitrary shape in between how do I solve the Schrodinger equation and for that we have to numerically integrate the Schrodinger equation, how do I numerically integrate at differential equation, I will write now focus principally on solving the Schrodinger equation which is a second order differential equation ψ double prime plus $V \times \psi$ equals $e \psi$ and this is an Eigen value equation; that means, I have to satisfy certain boundary conditions then only the solution is acceptable.

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So, let me just give you that for the Schrodinger equation minus h cross square over 2 m psi double prime plus V x psi equals E psi I am supposed to number 1 satisfy certain boundary conditions the wave function psi x is continuous and third psi prime x is continuous unless we goes to infinity, I have to satisfy all these 3 conditions and then this together gives me the Eigen energies and eigenfunctions.

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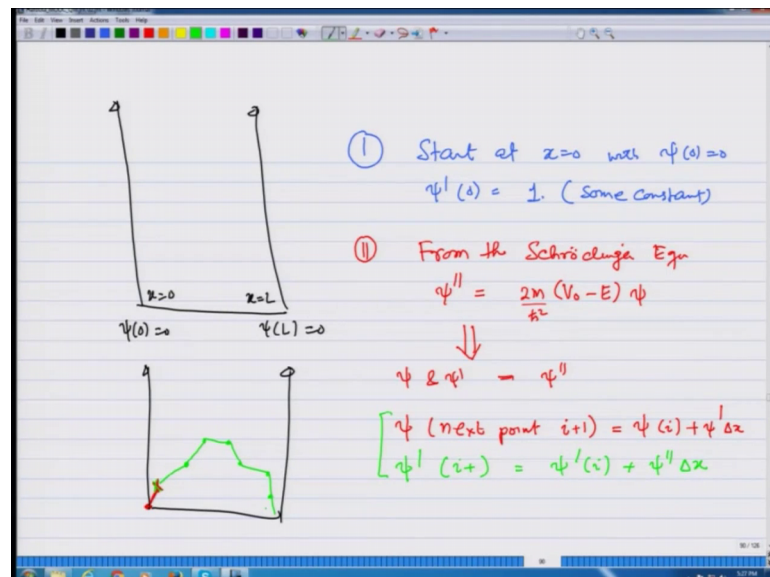


So, what we need to do is solve this equation satisfying the boundary conditions and making sure that psi n psi prime both are continuous to start learning about numerical

techniques we will focus in this lecture on a simpler problem where $\psi(0)$ is 0 and $\psi(L)$ is 0 that will also make us understand the problem better.

So, this is the problem with infinite (Refer Time: 04:24) high V goes to infinity at 2 points and remains infinity outside between 0 and L , V is some shape $V(x)$. So, ψ is 0 at this edge and ψ is 0 at the other edge and we have to build the solution in between. So, in numerical solution of differential equations we actually construct the solution how do we do that? We start integrating the ψ' and ψ'' to update our size and all that.

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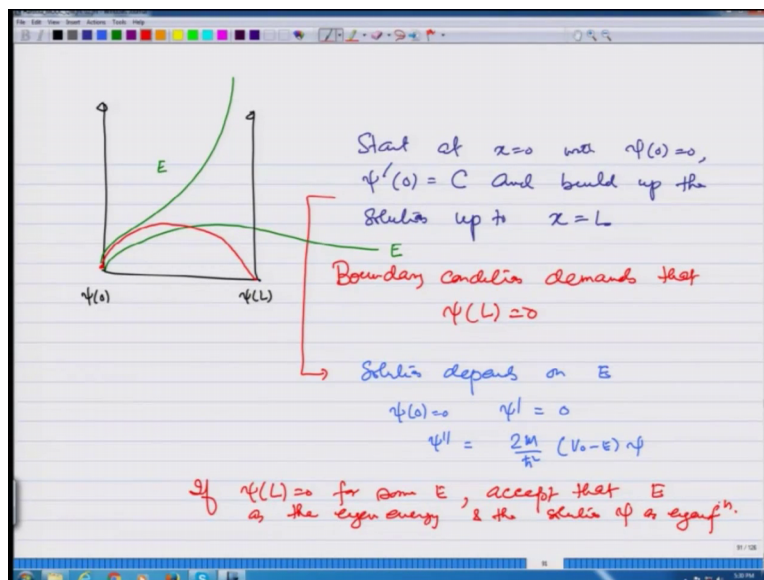
So, one technique could be in this case in which case the $\psi(0)$ on 2 sides $\psi(0)$ and $\psi(L)$ both are 0 this is $x=0$ $x=L$. So, one way of constructing could be start at x equal to 0 with $\psi(0)=0$ and $\psi'(0)$ equals some finite number let us say 1 or some constant.

I will keep showing what I am doing here. So, in this box, I have let me show ψ by red $\psi(0)$ and some slope what we can do with this is find the value at the next point. So, number 2 from the Schrodinger equation ψ'' is equal to $V_0 - E$ $2m$ over \hbar^2 times ψ . So, I know this implies that at any point I know ψ and ψ' , it gives me ψ'' also. So, from this point, on the left hand side started $x=0$, I have update it and so I can write ψ at the next point, let me call it $i+1$ th point is equal to ψ at point i plus $\psi'(i) \Delta x$ and I can also update ψ'

prime at $i + 1$ th point as ψ prime at i th point plus ψ double prime Δx and how do I know ψ double prime of I know ψ x some point I know ψ double prime also.

Once I do that. So, I found everything on this cross from here I again build my solution to the next point and from that to the next point, from that to the next point, from that to the next point and so on, I can keep building the solution because I am keeping terms only up to the first order, I have to make sure that this Δx is really small in the second order can be neglected now once we start with this process.

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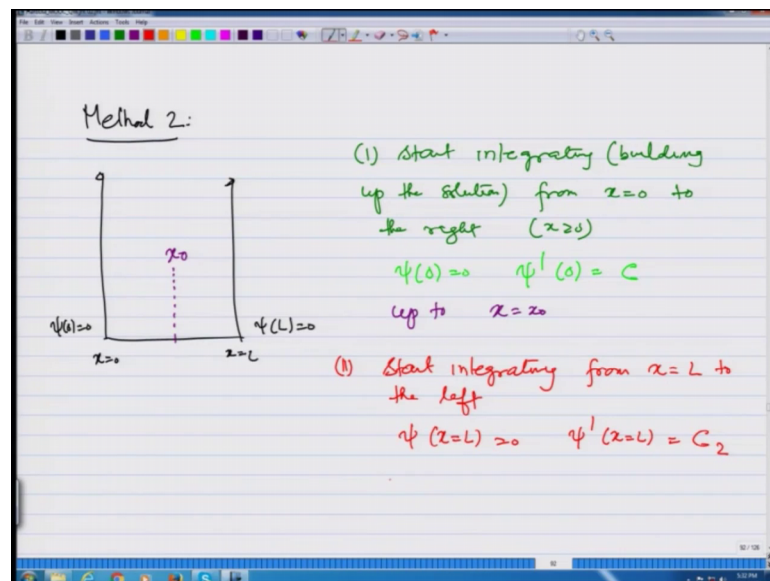


So, in this case, I am taking the case where the ψ vanishes on both sides ψ equals 0 ψ equals L . So, start at x equal to 0 with $\psi(0)$ equals 0 ψ prime 0 equals some constant. So, let us write it C , I wrote one earlier and build up the solution up to x equals L . Now boundary condition demands that $\psi(L)$ be equal to 0, remember now when I am starting from ψ equal to ψ at x equals 0, I am building up a solution and the solution depends on E ; what is that happened because I had a $\psi(0)$ equal to 0 ψ prime equal to 0, but I have ψ double prime at any point which is given as $2m$ over h cross square V_0 minus E ψ .

So, depending on what; E is ψ prime is given and I use all these 3 to build up my solution. So, $\psi(L)$ is not going to be 0 until I really have that Eigen value. So, what could happen? For example, if I build up the solution one solution could be like this, I am showing in green this is not acceptable. So, this is some e second solution go to do like

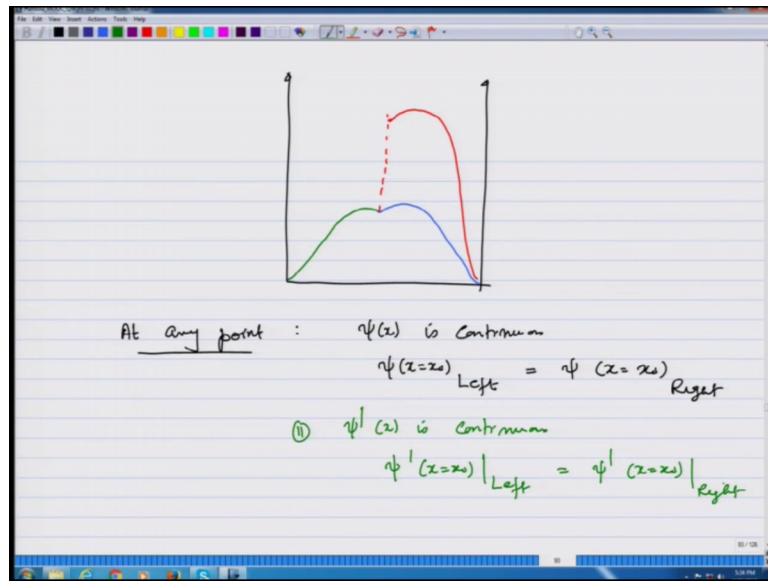
this some other E unless I have exactly write E that at least goes to 0 on the other side and that is acceptable. So, if $\psi(L)$ is equal to 0 for some E accept that E as the Eigen energy and the solution ψ as Eigen function and then you can keep increasing E and find more and more and more states. Since the particle is bound in this case between 0 and L ; that means, V goes to infinity outside this boundary there going to be infinite number of states in this case and so you can keep scanning over E and find more and more states.

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The other method could be again using the property is the solution is the Schrodinger equation and this case again I have $\psi(0)$ equals 0 $\psi(L)$ equals 0 x equals 0 x equals L number 1 start integrating or what I say building up the solution from x equals 0 to the right, where write means x greater than 0. So, you start with $\psi(0)$ equal 0 $\psi'(0)$ equals sum number and then you know ψ'' from ψ and you start integrating out up to certain point x , let us call it x_0 up to x equals x_0 second start integrating again building up solution from x equals L to the left. So, what you will do is ψ at x equals L is 0 take ψ' at x equals L to be some other constant C to ψ'' I know.

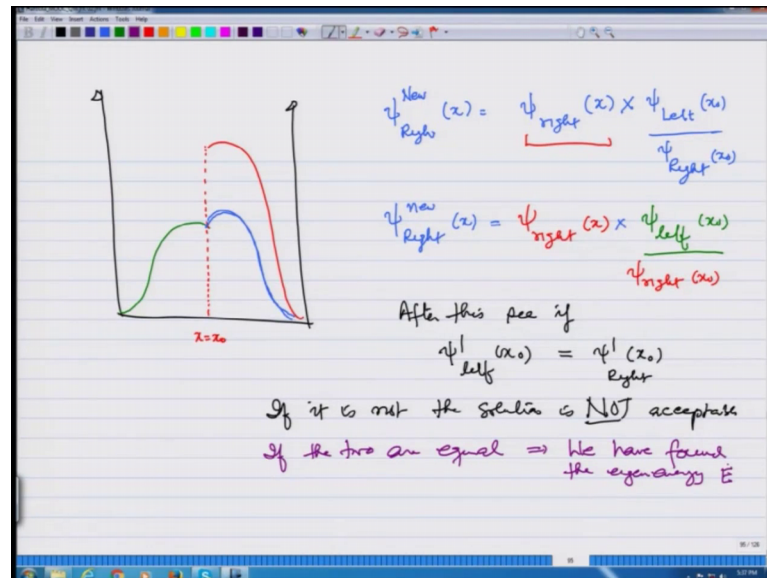
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So, you start again integrating to the left and up to x equals x_0 . So, what I have we done what we have done is in this potential, we will started a building up a solution from here, let us say this is what comes out from this side and we started building up the solution from the right side and let us say this comes out like this it need not match the other situation could be that let me show it in red this fellow could be like this what we should do is at x equals x_0 the solution should match what does that mean at any point what do we have for the acceptable solution $\psi(x)$ is continuous; that means, $\psi(x)$ equals x_0 from left should be equal to $\psi(x)$ equals x_0 from right and number 2 $\psi'(x)$ is continuous. So, you should have $\psi'(x)$ equals x_0 from the left is equal to $\psi'(x)$ equals x_0 coming from the right.

So, what you could do is match the 2 solutions was by multiplying by a factor and then compare the ψ' values.

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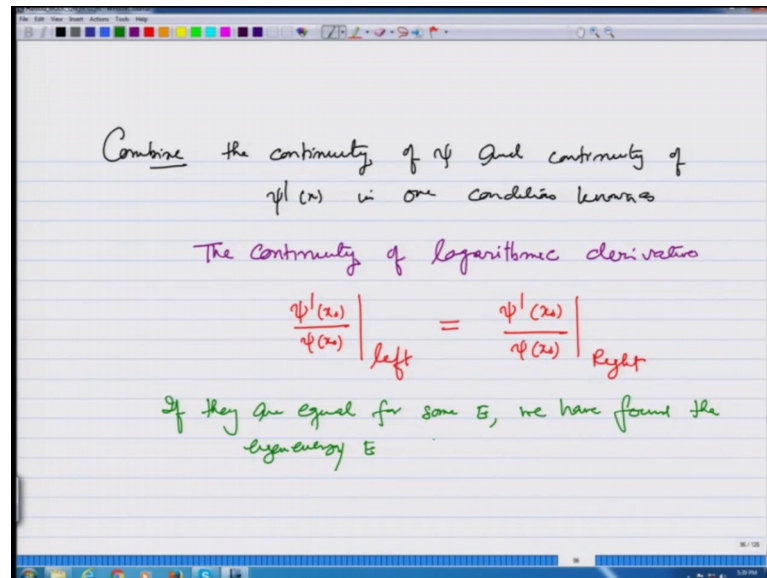


So, again let me make this picture and you have a solution coming from the left it could be like this solution coming from the right which could be like this. Now let us see if it at this point where we are matching which is x equals x_0 . Since we started with different slopes first the wave function did not match. So, what we need to do is multiply the right hand side to the factor that brings it down makes the wave function match at this point. So, I can write ψ_{Right} at x , I could write as ψ_{Right} at x times ψ_{Left} at x_0 divided by ψ_{Right} at x_0 , this will make the match. So, this will make it make the curve to be this blue one right. So, let me write this new.

So, from red we have done this one let me write this one as red. So, I have made $\psi_{\text{Right}}^{\text{New}}$ x equals ψ_{Right} x times ψ_{Left} x_0 divided by ψ_{Right} x_0 and then after that this is step one and after this, this C if ψ_{Left}' x_0 is same as ψ_{Right}' x_0 coming from right if it is not the solution is not acceptable on the other hand if the 2 are equal then we have found the Eigen energy E . So, for some energy they are equal and that is Eigen energy they will not be equal for all energies after all this rescaling has been done.

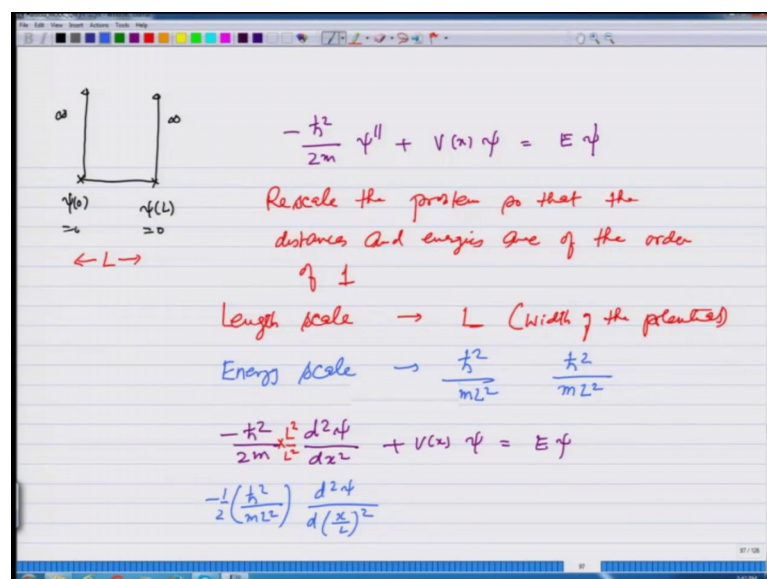
So, this the way to find Eigen value and the eigenfunction in 2 different ways one by making sure that the wave function is smooth all over; that means, ψ and ψ' are the same no matter which side I integrate the wave function form and the other I build us smooth wave function starting from one boundary and make sure that the other boundary condition is also satisfied both are acceptable ways.

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Now, we sometimes combine in fact not some time all the time combine the continuity of ψ and continuity of $\psi' x$ in one condition known as a continuity of logarithmic derivative. So, in this case what I do is I take $\psi' x_0$ over ψx_0 left and compare this with $\psi' x_0$ over ψx_0 coming from right and check if they are equal if they are equal for some e we have found the Eigen energy E , if not we again go to some other energy and check.

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So, we can scan over the energy and see what happens. So, these are going to be the 2 methods which we use in calculating the eigenfunctions and Eigen energies of a given problem a right now as I said earlier we have occurring on problems where the potential could be anything between x equals 0 and x equals L , but it is going to be infinity outside that region. So, that wave function is always going to be 0 at the 2 boundaries now the Schrodinger equation is minus \hbar^2 cross square over $2m$ psi double prime plus $V(x)$ psi equals E psi. Now these numbers \hbar cross and m are very small we are dealing with microscopic world. So, is good idea to rescale the problem? So, that the distances and energies are of the order of one; that means, we changing units rather than of the order of being 10 days to minus 36; 10 days to minus 19 and so on that will be very difficult to program in a computer.

So, what we do for example, in this case in the case of the problems that we are dealing with right now is that length scale we will take to be L that is the width of the potential and the energy scale naturally becomes \hbar^2 cross square over mL^2 square. So, with keep the lens scale to be \hbar^2 cross square over mL^2 square which is dimensionally correct. So, let us rewrite the Schrodinger equation for this. So, when we do that we have minus \hbar^2 cross square over $2m$ $d^2\psi$ over dx^2 plus $V(x)$ psi equals E psi let me multiply here by L here and divide by L^2 square. So, that I have the equation as minus \hbar^2 cross square over mL^2 square one half in front $d^2\psi$ over dy^2 plus $\bar{V}(y)$ psi equals \bar{E} psi divide.

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Divide by \hbar^2/mL^2 all over

$$-\frac{1}{2} \frac{d^2\psi}{d(x/L)^2} + \frac{V(x)}{(\hbar^2/mL^2)} \psi(x) = \frac{E}{(\hbar^2/mL^2)} \psi$$

$$y = \frac{x}{L}, \quad \bar{E} = \frac{E}{(\hbar^2/mL^2)}, \quad \bar{V} = \frac{V}{(\hbar^2/mL^2)}$$

$$\boxed{-\frac{1}{2} \frac{d^2\psi(y)}{dy^2} + \bar{V}(y) \psi(y) = \bar{E} \psi(y)}$$

$0 \leq y \leq 1$ for $0 \leq x \leq L$
 $L \sim \text{few } \text{\AA} \text{ (few } 10^{-10} \text{ m)}$

By $\frac{h^2 \psi''}{2m}$ over $\frac{h^2 \psi''}{2m}$ all over to get $-\frac{\hbar^2}{2m} \psi'' + V(x)\psi = E\psi$ plus $V(x)\psi$ is divide by $\frac{h^2 \psi''}{2m}$ is equal to $E \frac{2m}{\hbar^2}$ over $\frac{h^2 \psi''}{2m}$ over $m L^2 \psi''$ call $y = \frac{x}{L}$ and $E' = \frac{E}{\frac{\hbar^2}{2m L^2}}$ and $V' = \frac{V}{\frac{\hbar^2}{2m L^2}}$ to rewrite the equation as $-\frac{1}{2} \psi'' + V'(y)\psi = E'\psi$ and that is my equation in this dimensionless quantities y , E' and V' . And now I can take y between 0 and 1 for x varying between 0 and L .

Remember in quantum mechanics or in atomic quantities L is of the order of may be a few young storms which is few 10^{-10} meters. Now I can vary y between 0 and 1. So, I can take the increment in y to be of the order of point 1 and so on where as if I want to do the same thing in terms of x it would be much more difficult and numerically unstable.

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The image shows a whiteboard with handwritten notes. On the left, there is a diagram of a potential well $V(y)$ with a wavy bottom. To the right of the diagram is the Schrödinger equation: $-\frac{1}{2} \frac{d^2\psi}{dy^2} + V(y)\psi = E\psi$. Below this, it says "METHOD I: $\psi(0) = 0$, $\psi'(0) = 1$ ". Further down, it lists $\Delta x \approx 0.01, 0.005$ and $\psi(0) = 0, \psi'(0) = 1, \psi'' = 2(V-E)\psi$. There are two sets of equations in red ink, one for Δx and one for $2\Delta x$, showing how to update ψ and ψ' numerically. A small diagram of a horizontal line with a double-headed arrow labeled Δx is also present.

So, now we are able to integrate the equation. So, I have again I will take this box where in between there some potential let us call it we now y this is given in units of $\frac{h^2}{2m L^2}$. So, I am going to drop that bar in this write this as $-\frac{1}{2} \psi'' + V(y)\psi = E\psi$ method one is going to be number what we did we start with $\psi(0) = 0$ equals 0 take $\psi'(0) = 1$ let us take it to be one and then integrate. So, what we will do is we will divide this whole thing into

small ϵ of interval Δx . So, Δx could be of the order of point 0.1005 over so on, right some Δx very small number.

So, I will start from ψ equal to 0 and ψ' at 0 equals 1 ψ'' by the Schrodinger equation which is $2V - E \psi$, then I have ψ at Δx would be $0 + \psi'$ at 0 times $\Delta x + \psi''$ at 0 which is 0, right. Now Δx and now I can write ψ'' at Δx to be equal to $2V$ at $\Delta x - E \psi$ at Δx . This is one cycle completed from here I can go to the next step and write ψ at $2\Delta x$ would be equal to ψ at $\Delta x + \psi'$ at $\Delta x \Delta x + \psi''$ at $\Delta x \Delta x$ is equal to ψ' at $\Delta x + \psi''$ at Δx which we have already calculated at $\Delta x \Delta x$ and ψ'' at $2\Delta x$ is now updated to be equal to $2V$ at $2\Delta x - E \psi$ at $2\Delta x$.

This is another cycle we have completed this m should be way away and now I can build up the solution and so on.

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$$\psi_{i+1} = \psi_i + \psi'_i \Delta x$$

$$\psi'_{i+1} = \psi'_i + \psi''_i \Delta x$$

$$\psi''_{i+1} = (\psi_{i+1}) 2(V(x_{i+1}) - E)$$

\downarrow
 Check if $\psi(N) = 0$
 $|\psi(N)| < 10^{-3}, 10^{-4} \dots$
 You have found the eigenenergy E

So, in general what I am doing is I am writing ψ at $i + 1$ th point to be equal to ψ at i plus ψ' at i times Δx , I am updating ψ' at $i + 1$ as ψ' at i plus ψ'' at i times Δx and once I found ψ at $i + 1$ I can update ψ'' at $i + 1$ is equal to ψ at $i + 1$ times $2V$ at x at $i + 1$ minus E notice that we will depend on E and this way I can go all the way building up the solution at different points

all the way to the other side x equals L and then check if ψ let us say the last point is n is equal to 0 and numerically it is very difficult to make sure is 0.

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$\psi_i \quad (i=1 \dots N)$
 Normalize ψ

$$\int_0^L |\psi|^2 dx = \sum_{i=1}^N |\psi_i|^2 \Delta x = C_N$$

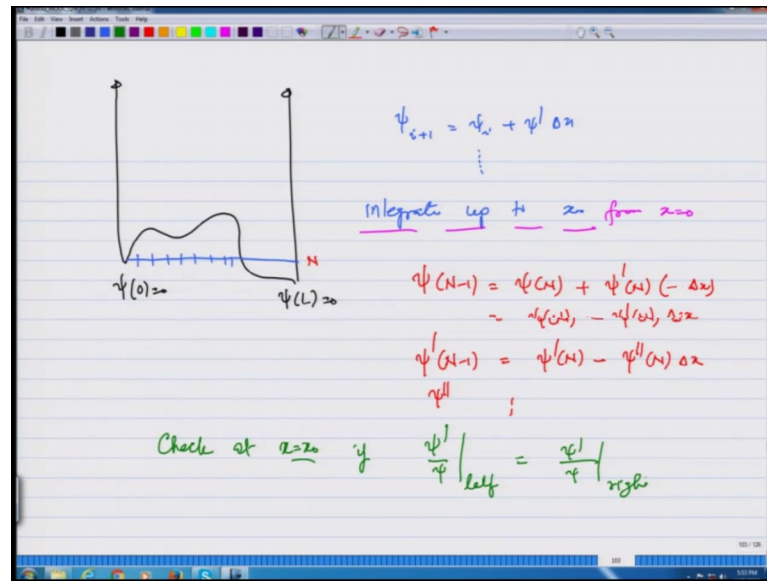
$$\psi_i^{new} = \frac{\psi_i}{\sqrt{C_N}} \quad \text{gives the normalized wavefunction}$$

 METHOD II :

So, what you do is you say ψ n modulus is less than some very small number let us say 10^{-3} or 10^{-4} or so on if that is the case you accept it. So, the moment that happens you have found the Eigen energy E whichever E that happens for your found that energy and once you have found that energy you have also found ψ_i at different i equals 1 through n gives you the value of ψ_i at each point. So, now, you normalize ψ what you do you calculate integral ψ mod square dx 0 to L which will be nothing, but summation i mod ψ square Δx i equals 1 through n and this gives you sum number let us say C_N then ψ_i^{new} is equal to whatever ψ you calculated earlier divided by square root of C_N give the normalized wave function that is one way of finding the wave function in this case.

Method 2 which in this case is not will required, but I am giving it because it will require for cases that will discuss in the coming lectures. In method 2, you actually start building up the solution from one side you start building up the solution from the other side and you match a logarithmic derivative at some point x_0 once that matches that gives you the Eigen value.

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So, what you do is some potential in between we have psi L equal to 0 psi 0 equals 0. So, you start you again divide this whole thing into small-small-small intervals. So, you write psi i plus 1 equal psi i plus psi prime delta x and all those steps that we did earlier and integrate up to x 0 from x equal to 0 this done that is the last point is n, then you go psi n minus 1 equals psi n which is 0 plus psi prime n times minus delta x because you going left.

So, this is going to be psi n minus psi prime n whatever value you assume times delta x you go psi prime n minus one equals psi prime at n minus psi double prime n delta x and then to update psi double prime and so on. So, you build up the solution the other way and then check at x equals x 0 if psi prime over psi left same as psi prime over psi right how do we check that we can say that psi prime over psi left minus psi prime over psi, right, modulus is some less than 10 raise to minus 3 or 10 raise to minus 4 that you have to play around and see where you get the answer once that happens if it happens for certain E that is the Eigen value.

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$$\left| \frac{\psi}{\psi} \right|_{\text{left}} - \frac{\psi}{\psi} \Big|_{\text{right}} < 10^{-3}, 10^{-4} \dots$$

for E 's when this happens you have found the eigenenergy, and the corresponding ψ

$$\psi_{\text{right}}^{\text{new}}(x) = \psi_{\text{right}}(x) \cdot \frac{\psi_{\text{left}}(x_0)}{\psi_{\text{right}}(x_0)}$$

The graph shows a potential well on the left and a wave function that is zero at the boundaries and has a peak inside the well.

So, for E 's when this happens you have found the Eigen energy and the corresponding ψ not to make the ψ proper what you do is you come from left you come from right. So, what you do is ψ right at some x new has to be equal to ψ right at x times ψ left at x_0 to at website right at x_0 this will make sure that now the way function match at x equals x_0 and then you found the continuous wave function and in the normalized. So, what I have given you is in this very specific kind of problem where the wave function vanishes at the 2 boundaries how to find the wave function and the Eigen energies.

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CONCLUDE

We have learnt about integrating the Schrödinger equation for very specific problems

$$V(x) = \begin{cases} \infty & x \leq 0 \\ & x \geq L \end{cases}$$
$$= V(x) \quad 0 < x < L$$
$$\psi(0) = 0 = \psi(L)$$

So, to conclude we have learnt about integrating the Schrodinger equation for very specific problem where $V(x)$ is equal to infinity for $x \leq 0$ and $x \geq L$ and is equal to $V(x)$ for x in between.

So, $\psi(0)$ is 0 and so is $\psi(L)$ and we found a way of finding the solution, we will now in the next lecture consider cases where ψ it is not vanishes at these boundaries, but goes all the way to infinity out find solutions for those situation.