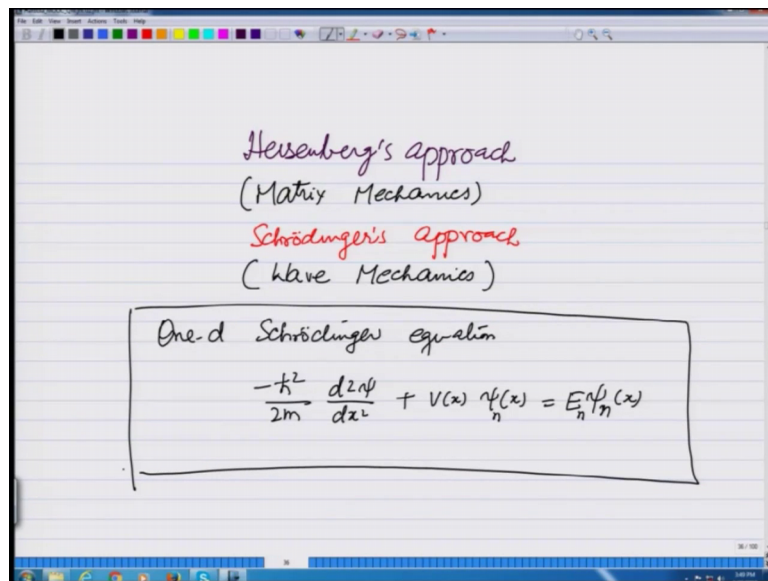


Introduction to Quantum Mechanics
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Lecture - 01

Solution of the Schrodinger equation for a particle in one and 2 delta function potentials

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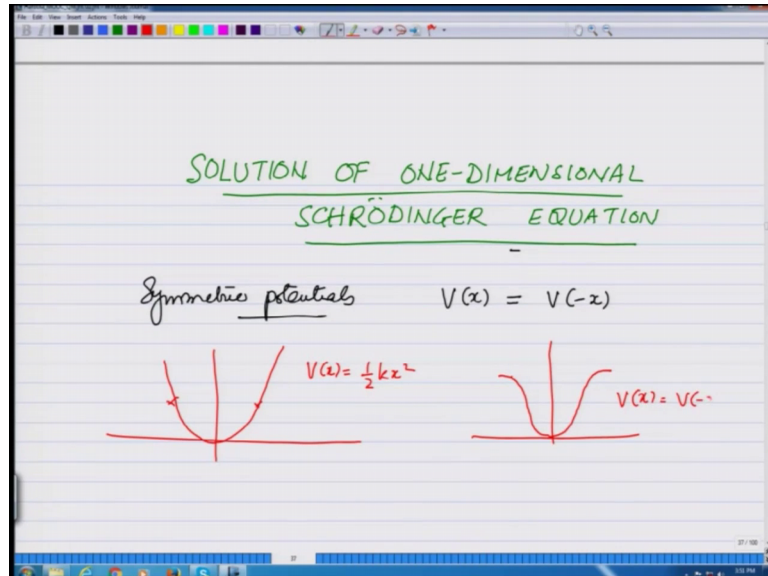


What we have done so far is the version which is known as the Heisenberg's approach which is essentially what we called matrix mechanics and as one of the students in the forum also pointed out it does become difficult after certain simple problems and the other version will not about is the Schrodinger's approach which is nothing but what I will call wave mechanics and its considers particles has having associated with sensor for the wave function the one d Schrodinger equation. And we have solved it in 2 cases is written as minus \hbar cross square over $2m$ d $^2 \psi$ over dx^2 plus $V \times \psi = E \psi$ and I am going to put in n here I as subscript to show that these are I can functions which satisfied boundary condition.

So, the satisfaction of boundary condition gives you the Eigen values energy E_n and the corresponding Eigen functions. So, this is the approaches that we are going to now follow in the coming lectures and apply Schrodinger equations solve problems some standard and some non standard wave, we are going to use numerical approach to solve

the problem we are going to start with very simple problems and slowly increase the difficulty level.

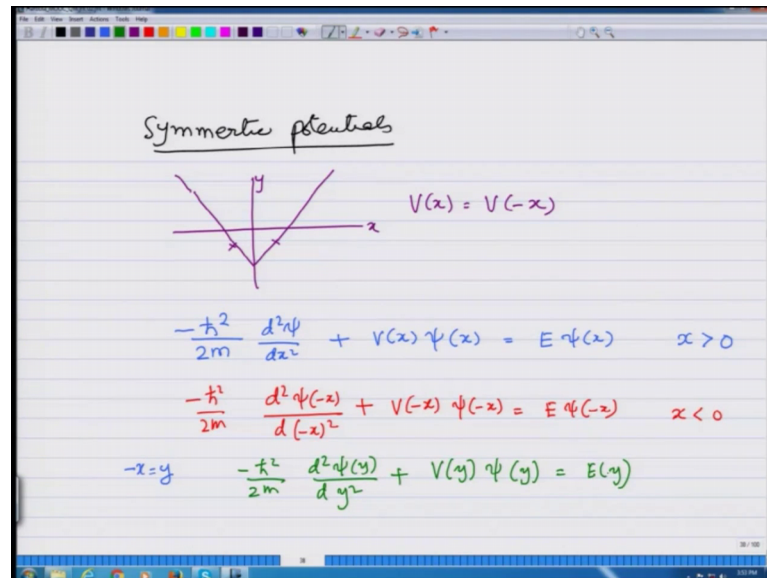
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So, from now onward what we are concentrating on is the solution of the time one dimensional Schrodinger equation and since I am starting the simple problems, I will start with something called the symmetric potentials and what I mean by that is that if I go from positive x to negative x ; the potential does not change.

So, $V(x)$ is $V(-x)$ for example, if I take a potential which is simple harmonic that is $V(x) = \frac{1}{2} kx^2$ its same for the negative axis and the positive axis. So, this is a symmetric potential any potential that has exactly the same value on 2 sides of the y axis is a symmetric potential and does not confine to potential being only positive.

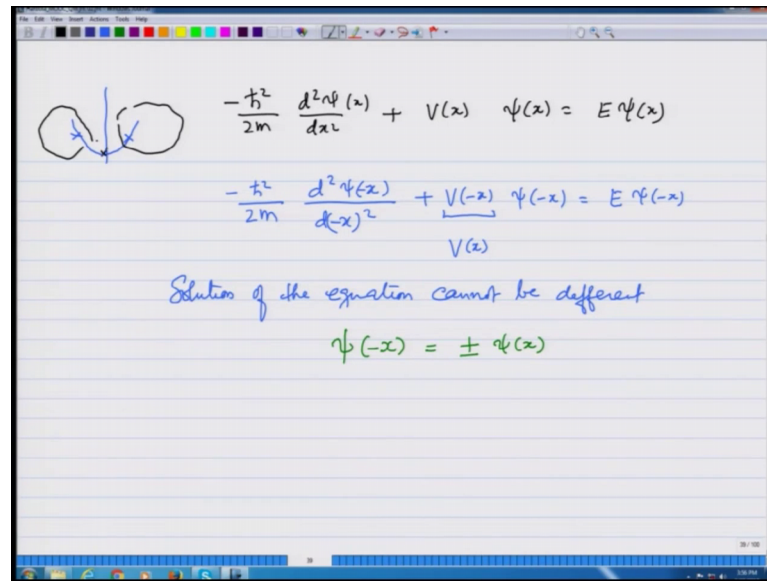
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So, let me just write some more examples. So, I am talking about symmetric potentials, for example, I could write this potential which is again the same on 2 of sides the y axis, this is also $V(x) = V(-x)$, this is also symmetric potential and why I am focusing on these write now is that they have a special advantage they provide special advantage while solving the Schrodinger equation let me look at the Schrodinger equation for such potentials. So, the Schrodinger equation is this plus $V(x)\psi(x) = E\psi(x)$, let me write this only for $x > 0$, the equation for $x < 0$ is therefore, going to be from this equation minus $\frac{\hbar^2}{2m} \frac{d^2\psi(-x)}{d(-x)^2} + V(-x)\psi(-x) = E\psi(-x)$, all I have done is change x to $-x$.

Now keep in mind that $-x$ is same as y substitute $-x = y$ then the equations becomes $-\frac{\hbar^2}{2m} \frac{d^2\psi(y)}{dy^2} + V(y)\psi(y) = E\psi(y)$. So, this is basically going back and forth important thing, now is going to be that when I take this equation.

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The image shows a digital whiteboard with handwritten mathematical equations and a diagram. On the left, a potential well is sketched with a vertical axis and two regions. To the right, the Schrödinger equation is written for a symmetric potential $V(x) = V(-x)$. The equations are:

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi(x)}{dx^2} + V(x) \psi(x) = E \psi(x)$$
$$-\frac{\hbar^2}{2m} \frac{d^2 \psi(-x)}{d(-x)^2} + \underbrace{V(-x)}_{V(x)} \psi(-x) = E \psi(-x)$$

Below the equations, it is noted that the solutions cannot be different, leading to the parity property:

$$\psi(-x) = \pm \psi(x)$$

Minus \hbar^2 crosses square over $2m$ $d^2 \psi(x)$ plus $V(x) \psi(x) = E \psi(x)$ and change to the other side of the y axis. So, that is some potential is there an instead of writing here I want to write the equation for the other side this is going to be minus \hbar^2 crosses square by over $2m$ I can just change signs $\psi(x)$ minus sign here over d minus x square plus V minus x ψ minus x equals $E \psi$ minus x this same as $V(x)$.

So, since V is the same on the both sides the solution of the equation cannot be different write $V(x)$ and $V(-x)$ are one and the same, think there will be exactly the same number as the function of. So, as the function of minus x , we gives exactly same number as the function of x . So, solutions cannot be different at this differ by constant.

So, $\psi(-x)$ is going to be plus or minus $\psi(x)$, I choose that constant to be one even if I take to some constant c I can always take it that to be one multiplied by constant does not really change this solution, important thing is that if $V(-x)$ and $V(x)$ other same then the way solution evolves as the function of x on the right hand sides. So, let me show this in black the way this solution would evolve from a boundary condition, here 0 to the right hand side would evolve exactly in the same way on the left hand side and therefore, the 2 solutions are going to be changed by only by a factor of plus or minus 1.

So, this is an important property of the wave function for symmetric potentials which I am going to make use of.

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Example: Particle in a δ function potential

$$V(x) = -V_0 \delta(x)$$

There is a bound state for this potential
 $E < 0$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} - V_0 \delta(x) \psi(x) = E \psi(x)$$

If $x \neq 0$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E \psi(x) = -|E| \psi(x) \quad (E < 0)$$

So, in solving Schrodinger equation let me take the first example of particle in a delta function potential and what I mean by that is that the potential $V(x)$ is equal to minus V_0 delta of x , how does this potential look? This potential is going to look like this, if this is the x axis, this is the y axis, it is going to be very deep at x equal to 0, instead going to minus infinity and 0 everywhere else this could be a representation of a potential which is very deep and has very small width classically whether this can afford a bound state or not we do not know, but quantum mechanically I am going to show you that this affords some state.

So, there is a bound state for this potential; what do we mean by that bound state; that means, total energy e is going to be less than 0 the particle is going to be bound near the delta function.

Let us see how it is that comes out of bound. So, the Schrodinger equation for this is minus \hbar^2 cross square over $2m$ $d^2\psi$ over dx^2 minus V_0 delta x $\psi(x)$ equals $e \psi(x)$ and what that means, is that if x is not equal to 0, then I have minus \hbar^2 cross square over $2m$ $d^2\psi$ over dx^2 is equal to $e \psi(x)$ and this means minus modulus of $e \psi(x)$ because e is less than 0 and what will happen at x equal to 0, I show you in a minute.

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$x \neq 0 \quad -\frac{\hbar^2}{2m} \psi''(x) = -|E| \psi$
 $\psi'' - \frac{2m|E|}{\hbar^2} \psi = 0$
 $\psi(x) = e^{\sqrt{\frac{2m|E|}{\hbar^2}} x} \quad \text{or} \quad e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$

Since $\psi(x)$ represents a bound state
 The wavefunction $\psi(x) \rightarrow 0$ as $x \rightarrow \pm \infty$

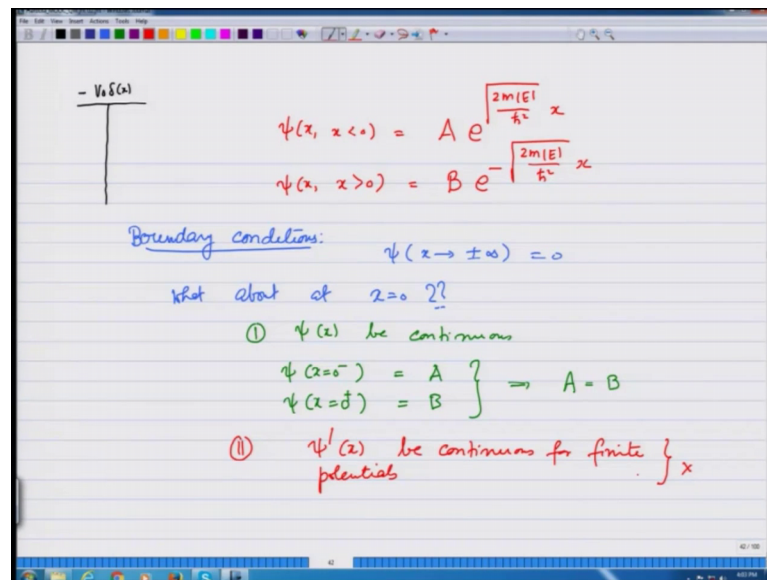
$\Rightarrow \psi(x, x < 0) = A e^{\sqrt{\frac{2m|E|}{\hbar^2}} x}$
 $\psi(x, x > 0) = B e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$

So, for x not equal to 0, I have minus \hbar^2 over $2m$ ψ'' , let me write it as $\psi'' - \frac{2m|E|}{\hbar^2} \psi = 0$ and therefore, $\psi'' - \frac{2m|E|}{\hbar^2} \psi = 0$. This is a positive quantity and therefore, you can immediately write the solutions. So, $\psi(x)$ is going to be $e^{\sqrt{\frac{2m|E|}{\hbar^2}} x}$ or $e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$ that is the solution since this is the bound state right let me make the potential again.

So, that it is you know mine; what we are solving for the potential is right here, very deep minus infinity this is $V_0 \delta(x)$ and 0 everywhere else since $\psi(x)$ represents bound state; that means, the wave function $\psi(x)$ must go to 0 as x goes to plus or minus infinity; that means, there is a almost 0 probability of finding the particle far away from the potential.

And this immediately implies that $\psi(x)$ for $x < 0$ is going to be $e^{\sqrt{\frac{2m|E|}{\hbar^2}} x}$ with some constant A , in front and $\psi(x)$, $x > 0$ is going to be some constant $B e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$, I have deliberately chosen these constants to be different because I do not want to make a point.

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So, in this case that we are considering a delta potential $V = V_0 \delta(x)$, these solutions are $\psi(x, x < 0) = A e^{\sqrt{2m|E|/\hbar^2} x}$ and $\psi(x, x > 0) = B e^{-\sqrt{2m|E|/\hbar^2} x}$. What about boundary conditions $\psi(x \rightarrow \pm\infty) = 0$ is automatically satisfied.

Now, we have written the solution for $x > 0$ and $x = 0$. What about at $x = 0$ is the equation, recall what we demanded from the solution of the Schrodinger equation 1 $\psi(x)$ be continuous and this means at $x = 0$, the value should be the same from both sides.

Now, if you look at the solutions $\psi(x = 0)$, if you come from the left hand side that is $x < 0$, this is A and $\psi(x = 0)$ is B , if you come from the right hand side. So, I will denote that; I will write the upper 1 as 0^- and the 2 should be the same and this implies $A = B$. 2, we had demanded that $\psi'(x)$ that is the derivative of ψ with respect to x be continuous for finite potentials.

Now, in this case a particular case that we are dealing with the potential is not finite at $x = 0$ it goes to minus infinity therefore, this condition has to be replaced by some other condition and we will do that now.

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2nd Condition (ψ' be continuous)

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} - V_0 \delta(x) \psi(x) = E \psi(x)$$

$$\int_{0^-}^{0^+} \left[-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} - V_0 \delta(x) \psi(x) - E \psi(x) \right] dx$$

$$\int_{0^-}^{0^+} \delta(x) dx = 1$$

$$\int_{0^-}^{0^+} f(x) \delta(x) dx = f(0)$$

$$-\frac{\hbar^2}{2m} \left. \frac{d\psi}{dx} \right|_{0^-}^{0^+} - V_0 \psi(0) = E \int_{0^-}^{0^+} \psi(x) dx$$

$$-\frac{\hbar^2}{2m} [\psi'(0^+) - \psi'(0^-)] = V_0 \psi(0)$$

So, the second condition which is psi prime be continuous, it is going to replace by something else and let us see how that is replaced we have the Schrodinger equation $\frac{d^2 \psi}{dx^2} - V_0 \delta(x) \psi(x) = E \psi(x)$.

Since the discontinuity is at $x = 0$, let me integrate this equation with respect to x from 0^- to 0^+ what does that mean? Pictorially if this is the potential $-V_0 \delta(x)$ I am integrating from slightly left from a to slightly right I am integrating in this interval shown by green here; that means, its include the delta function and keep in mind that integration of $\delta(x) dx$ from 0^- to 0^+ is 1 any function $f(x) \delta(x) dx$ integrated over from 0^- to 0^+ is equal to $f(0)$.

So, therefore, what do I get from the equation $-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} - V_0 \delta(x) \psi(x) = E \psi(x)$ when $\frac{d^2 \psi}{dx^2}$ is integrated I get $\frac{d\psi}{dx}$ and this is from 0^- to 0^+ minus $V_0 \psi(0)$ is equal to $E \int_{0^-}^{0^+} \psi(x) dx$ which I can write as $0^+ - 0^-$ psi at 0 as the interval is made smaller and smaller. The right hand side is going to go to 0. So, the only conclusion I can draw now is that $-\frac{\hbar^2}{2m} [\psi'(0^+) - \psi'(0^-)] = V_0 \psi(0)$. So, this is what the delta function does if the delta function was not there psi prime would have been continuous, but because the delta function there is a discontinuity in psi prime and that is what leads to know the energy.

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$$-\frac{\hbar^2}{2m} [\psi'(0^+) - \psi'(0^-)] = V_0 \psi(0)$$

$$\psi(x > 0) = A e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$$

$$\psi(0^+) = A \sqrt{\frac{2m|E|}{\hbar^2}} \iff \psi'(x > 0) = -A \sqrt{\frac{2m|E|}{\hbar^2}} e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$$

$$\psi(x < 0) = A e^{+\sqrt{\frac{2m|E|}{\hbar^2}} x}$$

$$\psi(0^-) = A \sqrt{\frac{2m|E|}{\hbar^2}} \iff \psi'(x < 0) = +\sqrt{\frac{2m|E|}{\hbar^2}} A e^{+\sqrt{\frac{2m|E|}{\hbar^2}} x}$$

$$-\frac{\hbar^2}{2m} \times -2 \sqrt{\frac{2m|E|}{\hbar^2}} A = V_0 A$$

$$\frac{2m|E|}{\hbar^2} = \left(\frac{mV_0}{\hbar^2}\right)^2 = \frac{m^2 V_0^2}{\hbar^4}$$

So, let us write this now. So, the boundary condition at 0 is now minus h cross square over 2 m psi prime at 0 plus minus psi prime at 0 minus is equal to V 0 psi 0. Now psi at x greater than 0 is nothing but e raise to minus square root of 2 m mod e over h cross square x and therefore, there is a factor a also psi prime for x greater than 0 is equal to minus a times square root of 2 m mod a over h cross square e raise to minus square root of 2 m mod a over h cross square x and psi x less than 0 is a e raise to plus square root of 2 m mod a over h cross square x and psi prime for x less than 0 is going to be equal to plus square root of 2 m mod e over h cross square a e raise to square root of 2 m mod a over h cross square x.

So, this immediately implies that psi prime 0 plus is equal to minus a square root of 2 m mod a over h cross square lower side immediately implies that psi prime 0 minus is equal to a square root of 2 m mod a over h cross square and psi 0 of course, A. So, if I substitute this in the equation above blue equation I get minus h cross square over 2 m times minus 2 square root of 2 m mod a over h cross square is equal to V 0 times A, this A also floating round here and we cancel this a and I get 2 m mod e over h cross square is equal to m V 0 over h cross square is 2 also cancels. So, this is equal to m square V 0 square over h cross raised to 4.

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The image shows a handwritten derivation on a whiteboard. At the top, the equation $\frac{2m|E|}{\hbar^2} = \frac{m^2 V_0^2}{\hbar^4 k^2}$ is written. Below it, the equation $|E| = \frac{m V_0^2}{2\hbar^2}$ is boxed. Underneath the box, the text "OR" is written. Then, the equation $E = -|E| = -\frac{m V_0^2}{2\hbar^2}$ is written. At the bottom, a graph shows a wave function $\psi(x)$ that is zero for $x < 0$ and $x > a$, and has a positive peak between $x = 0$ and $x = a$. The energy $E = -\frac{m V_0^2}{2\hbar^2}$ is written to the right of the graph.

So, what I have is from the boundary condition $2 m \text{ mod } A$ over h cross square is equal to m square V_0 square over h cross raise to 4, we can again cancels in terms this gives me h cross square $1 m$ cancels and therefore, I get mod of E equals $m V_0$ square over $2 h$ cross square or energy E which is equal to minus mod E is equal to minus $m V_0$ square $2 h$ cross square; how does the wave function look the wave function if I plot it and on the x axis is exponentially decaying on the right hand side exponentially decaying on the left hand side $d \psi$ by $d x$ is discontinuous at x equal 0 and the energy e equals minus $m V_0$ square $2 h$ square that is the solution for a particle moving in a delta function potential you can see it is bound and it has energy just slightly negative that is why it is bound and the wave function is exponentially decaying on the 2 sides.

There are no other bound states in this potential. So, this is the answer for this if you want to normalize the wave function right. So, you have to see now is calculated the wave function ψx .

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$$\psi(x) = \begin{cases} A e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x} & x > 0 \\ A e^{\sqrt{\frac{2m|E|}{\hbar^2}} x} & x < 0 \end{cases}$$

A is yet not fixed

$$\int_{-\infty}^{\infty} |\psi(x)|^2 dx = A^2 () = 1$$

This will normalize the wavefn
 And give the total probability
 of finding particle anywhere in
 the interval $-\infty < x < \infty = 1$

Let me write it in 1 bracket e; e raise to minus square root 2 m mod E, we have calculated over h cross square x mod x greater than 0 A e raise to square root of 2 m mod E over h cross square x or x less than 0 A is yet not fixed if you want to fix it, what we will do is we will take psi x square integrated from minus infinity to infinity which will give me sum A square times sum constant and fix it equal to 1 and this will normalize the wave function and give the total probability of finding particle anywhere in the interval x minus infinity to infinity to be equal to 1; then I have a normalized wave function.

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Example 2: Two δ -functions

$$V(x) = -V_0 \delta(x+a) - V_0 \delta(x-a)$$

$$V(-x) = V(+x)$$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} - V_0 \delta(x+a) \psi - V_0 \delta(x-a) \psi = E \psi$$

For $(x \neq a, x \neq -a)$
$$-\frac{\hbar^2}{2m} \psi'' = E \psi = -|E| \psi$$

So, this is a very simple problem to start with the wave out generalized this would be 2. So, let me take example 2. So, the way to generalize this week 2 delta functions. So, in this case, I have the x axis, this is x equal 0 and 2 make the problem simple, what I will do is, I will put the delta function symmetrically about x equals 0 and both are the same $V_0 \delta(x - a)$ and this will minus $V_0 \delta(x + a)$. So, again the potential is symmetric $V(x) = V(-x)$ and therefore, the solution is going to be either symmetric or anti symmetric in this case it. So, happens that we get 2 bound states and let me show you how. So, the Schrodinger equation is $-\frac{\hbar^2}{2m} \psi'' - V_0 \delta(x - a) \psi - V_0 \delta(x + a) \psi = E \psi$.

For $x \neq a$ or $x \neq -a$ together equation is $-\frac{\hbar^2}{2m} \psi'' = E \psi$ which is $-\frac{\hbar^2}{2m} \psi'' = -|E| \psi$.

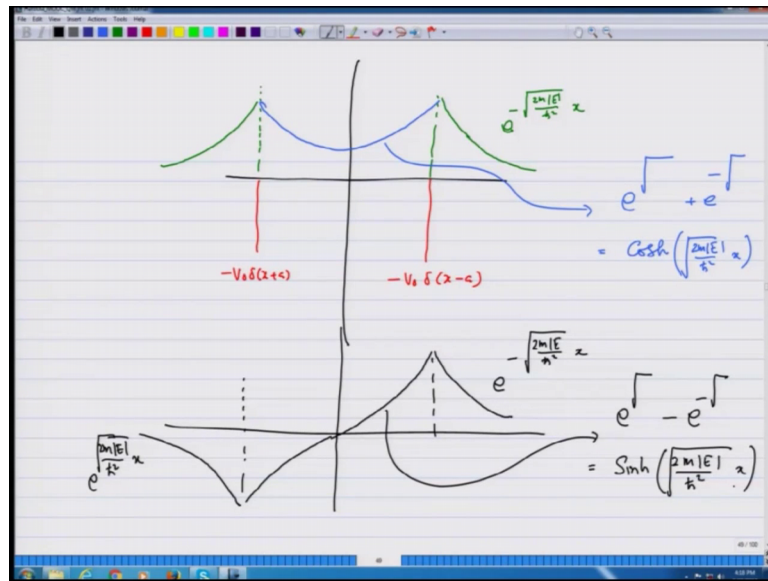
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The image shows a whiteboard with handwritten mathematical work. At the top, the Schrodinger equation is written as $-\frac{\hbar^2}{2m} \psi'' = -|E| \psi$. This is rearranged to $\psi'' - \frac{2m|E|}{\hbar^2} \psi = 0$. The general solution is given as $\psi(x) = e^{\sqrt{\frac{2m|E|}{\hbar^2}} x}$ or $e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$. Below this, it says "or a linear combination of the two". At the bottom, a potential energy diagram is drawn on a horizontal blue line. Two vertical red lines represent delta functions: one at $x = -a$ labeled $-V_0 \delta(x + a)$ and one at $x = a$ labeled $-V_0 \delta(x - a)$. A vertical blue line is drawn at $x = 0$.

So, the equation I am solving again is $-\frac{\hbar^2}{2m} \psi'' = -|E| \psi$ or $\psi'' - \frac{2m|E|}{\hbar^2} \psi = 0$. So, solutions again are $e^{\sqrt{\frac{2m|E|}{\hbar^2}} x}$ or $e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$ or a linear combination of the 2 now let me show you how you form this linear combination.

Even if you do an elaborate exercise you will get exactly the same answer. So, this is my x axis the potential is symmetric $V(x) = V(-x)$ minus $V_0 \delta(x - a)$ minus $V_0 \delta(x + a)$.

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Now, look not the wave function the upper side these are the potentials minus V naught delta x plus a minus V naught delta x minus a and on the right side the potential as to d k exponentially e raise to minus square root of $2 m$ mod E over h cross square x . So, on the x positive side, the solution has negative power of the x exponential on the left hand side. The solution again as to d k because the bound state in between I am going to have either symmetric or anti symmetric solutions the symmetric solution would be linear combination with two. And let me show on the lower side anti symmetric solution and that would be the solution decaying on one side and anti symmetric solution has to go to 0 necessarily through 0. It has to be 0 at the 0 and again explanation on the other side, this is e raise to minus square root of $2 m$ mod E over h cross square x . This is e raise to square root of $2 m$ mod E over h cross square x and in between it is an anti symmetric combination.

Symmetric combination is going to be a raise to the square root whole thing plus e raise to minus square root over thing which is nothing but the hyperbolic cosine function x and the lower case this is going to be an anti symmetric combination. So, is going to be e raised to square root of this whole thing minus e raise to minus square root of this whole thing and this is going to be sin hyperbolic of $2 m$ mod E over h cross square square root of h . So, I give you a qualitative argument based on symmetry how the wave function should look like as to what it linear combinations of these 2 solutions would be for the symmetric and anti symmetric wave function.

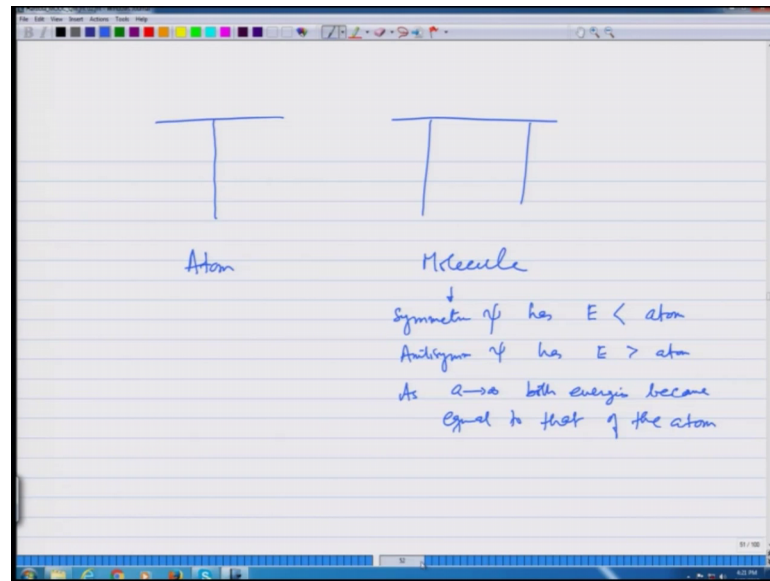
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The image shows a whiteboard with handwritten mathematical expressions. At the top, a set of curly braces contains two equations: $\psi(x) = \psi(-x)$ and $\psi(x) = -\psi(-x)$. Below this, the word "Mathematically:" is written. To its right, three equations are listed for different regions of x : $\psi(x) = A e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$ for $x > a$; $\psi(x) = B e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x} + C e^{+\sqrt{\frac{2m|E|}{\hbar^2}} x}$ for $-a < x < a$; and $\psi(x) = D e^{+\sqrt{\frac{2m|E|}{\hbar^2}} x}$ for $x < -a$. At the bottom, a sentence in purple ink reads: "To find the energy and eigenfunctions for a particle moving in a double δ -function potential".

So, in the first case $\psi(x)$ equals $\psi(-x)$ and the second case $\psi(x)$ equals minus $\psi(-x)$. This is a qualitative argument based on the symmetry if you do it mathematically, what you would say is $\psi(x)$ is equal to $A e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x}$ for $x > a$, $\psi(x)$ is equal to $B e^{-\sqrt{\frac{2m|E|}{\hbar^2}} x} + C e^{+\sqrt{\frac{2m|E|}{\hbar^2}} x}$ for $-a < x < a$ and $\psi(x)$ equals $D e^{+\sqrt{\frac{2m|E|}{\hbar^2}} x}$ for $x < -a$ and then match the boundary condition that the wave function is continuous at the boundary and so on and then you will find the solution.

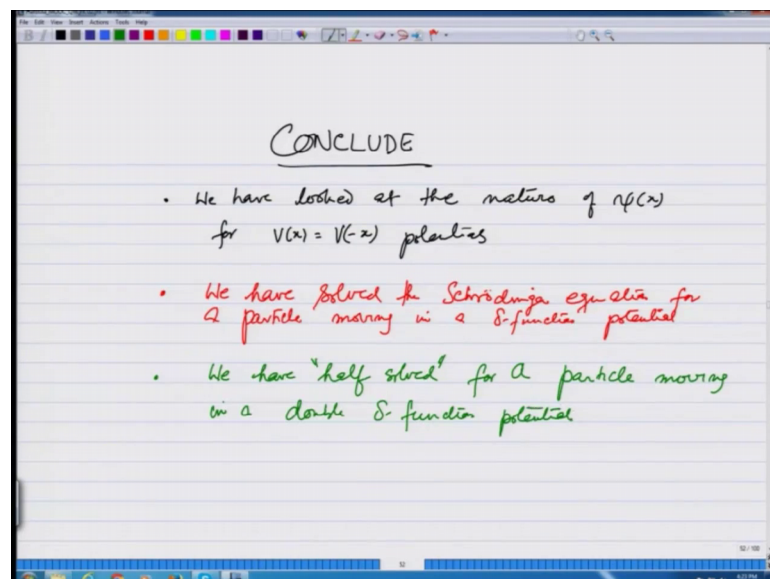
Now this I will be assigning you as an assignment problem to find the energy and eigenfunctions actually eigenfunctions I have already written for a particle moving in a double delta function potential.

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Significance of this problem is that a single delta function is like an atom and a double delta function is like a molecule and in this case, you would find that symmetric psi has energy lower than the atom. So, that is why molecules form and anti symmetric psi has energy greater than the atom and as a tends to infinity both energies become equal to that of the atom.

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So, let me conclude this lecture by stating that we have noted the nature of ψ for V equals V minus x potentials, we have solved the Schrodinger equation for a particle

moving in a delta function potential and third we have half solved. And I will give you western as the assignment for a particle moving in a double delta function potentially.