

Introduction to Quantum Mechanics
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Lecture – 06
Ehrenfest's theorems for the expectation values of x and p

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Ehrenfest's theorems

$$\langle x \rangle = \int |\psi(x,t)|^2 x dx$$
$$\langle p \rangle = \int \psi^*(x,t) \frac{\hbar}{i} \frac{\partial \psi}{\partial t} dx$$
$$\langle p \rangle = m \frac{d\langle x \rangle}{dt} = m v$$
$$\langle p \rangle = m \frac{d\langle x \rangle}{dt} \quad \text{--- (I)}$$
$$\left\langle -\frac{\partial V}{\partial x} \right\rangle = \frac{d\langle p \rangle}{dt} \quad \text{--- (II)}$$

In this lecture, I am going to talk about Ehrenfest's theorems, and they kind of tell you how the expectation values of p and x change with time, and these theorems tell you how expectation value of p and x change with time and they gave it some meaning which we have been sort of anticipating. So, the recall that expectation value x is define as integration $\psi^* x \psi dx$. And we said earlier that this is like the average value of x. Similarly, average value of p is defined as integration $\psi^* \frac{\hbar}{i} \frac{\partial \psi}{\partial x} dx$ and we said this is like the average value of momentum for a particle n wave function psi state. What Ehrenfest theorem tell you is the more meaning to this and show that p can be return as $m \frac{dx}{dt}$. So, this is like $m v$.

So, the expectation value of p that is average momentum that we have been talking about is $m \frac{d\langle x \rangle}{dt}$ of the average value of x, this is first theorem. Second one tells you about the rate of change of p and that tells if I take the expectation value of minus $\frac{\partial V}{\partial x}$ that is the derivative of potential v to the minus sign in front which is the force this is equal to $\frac{d\langle p \rangle}{dt}$ of average value of p. So, these again tell us at these expectation values as defined

earlier when we talking about equivalence of Heisenberg and showed in this picture, do have a meaning of being the average position and average momentum of a particle in state sign.

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$$\langle x \rangle = \int \psi^* x \psi dx$$

$$\frac{d}{dt} \langle x \rangle = \int \frac{\partial \psi^*}{\partial t} x \psi dx + \int \psi^* x \frac{\partial \psi}{\partial t} dx$$

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V(x)\psi$$

$$-i\hbar \frac{\partial \psi^*}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi^*}{\partial x^2} + V(x)\psi^*$$

$$\frac{d}{dt} \langle x \rangle = \left(\frac{1}{-i\hbar} \right) \left[\int -\frac{\hbar^2}{2m} \left(\frac{\partial^2 \psi^*}{\partial x^2} \right) \cdot x \cdot \psi + \int \psi^* V(x) x \psi \right]$$

$$+ \left[\int \frac{\hbar^2}{2m} \left(\frac{\partial^2 \psi}{\partial x^2} \right) x \cdot \psi^* - \int \psi^* V(x) x \psi \right]$$

$$\frac{d}{dt} \langle x \rangle = \frac{i}{\hbar} \int \frac{\hbar^2}{2m} \left(\frac{\partial^2 \psi}{\partial x^2} x \psi^* - \frac{\partial^2 \psi^*}{\partial x^2} x \psi \right) dx$$

So, let us now proof these. x average is nothing but integration $\psi^* x \psi dx$ and d by dt of x is therefore, going to be integration $d \psi^* x \psi dx$ plus integration $\psi^* x d \psi$ by $dt dx$. Notice that if you talking about the change of expectation value of x and p in time, the wave function necessarily has to be time dependent in a stationary state these things do not change in time as we showed earlier because the probability density is independent of time.

Now, I know that $i \hbar$ cross $d \psi$ over dt is nothing but minus \hbar cross square over $2 m$ $d^2 \psi$ over dx square plus $Vx \psi$. And minus $i \hbar$ cross $d \psi^*$ over dt is equal to minus \hbar cross square over $2 m$ $d^2 \psi^*$ over dx square plus $Vx \psi^*$. I can substitute these value in the equation about to get d by dt of x is equal to 1 over $i \hbar$ cross with the minus sign integration of minus \hbar cross square over $2 m$ $d^2 \psi^*$ over dx square times x times ψ . I put this in the bracket plus integration $\psi^* V x x \psi$ that is from the first term. So, this correspond to this term.

For the second term, I am going to have minus integration \hbar crosses square over $2 m$ $d^2 \psi$ over dx square I am write in $d \psi$ by dt times x times ψ^* minus integration $\psi^* V x \psi dx$ I have taken $i \hbar$ cross out. So, let us see now this is perfectly fine, this

is except that the sign is going to be plus, yes. So, this is fine. Now, I can cancel the second term. And what I am left with is $\frac{d}{dt} \langle x \rangle$ is equal to $\frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(\frac{\partial^2 \psi}{\partial x^2} \cdot x \psi^* - \frac{\partial^2 \psi^*}{\partial x^2} \cdot x \psi \right) dx$ minus I have written has $\frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(\frac{\partial \psi}{\partial x} \cdot x \psi^* - \frac{\partial \psi^*}{\partial x} \cdot x \psi \right) dx$ over $\int \left(\frac{\partial \psi}{\partial x} \cdot \psi^* - \frac{\partial \psi^*}{\partial x} \cdot \psi \right) dx$ that is what I am left with.

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$$\begin{aligned} \frac{d}{dt} \langle x \rangle &= \frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(\frac{\partial^2 \psi}{\partial x^2} \cdot x \psi^* - \frac{\partial^2 \psi^*}{\partial x^2} \cdot x \psi \right) dx \\ &= \frac{i\hbar}{2m} \int \frac{\partial}{\partial x} \left(\frac{\partial \psi}{\partial x} \cdot x \psi^* - \frac{\partial \psi^*}{\partial x} \cdot x \psi \right) dx \\ &\quad + \frac{i\hbar}{2m} \int \left\{ -\frac{\partial \psi}{\partial x} \cdot \frac{\partial \psi^*}{\partial x} \cdot x - \frac{\partial \psi}{\partial x} \cdot \psi^* \right. \\ &\quad \left. + \frac{\partial \psi^*}{\partial x} \cdot \frac{\partial \psi}{\partial x} \cdot x + \frac{\partial \psi^*}{\partial x} \cdot \psi \right\} dx \\ \frac{d}{dt} \langle x \rangle &= \frac{i\hbar}{2m} \int \left(-\psi^* \frac{\partial \psi}{\partial x} + \frac{\partial \psi^*}{\partial x} \psi \right) dx \\ \int \frac{\partial \psi^*}{\partial x} \psi dx &= \int \frac{\partial}{\partial x} (\psi^* \psi) - \int \psi^* \frac{\partial \psi}{\partial x} dx \end{aligned}$$

So, what I am getting is $\frac{d}{dt} \langle x \rangle$ is equal to $\frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(\frac{\partial^2 \psi}{\partial x^2} \cdot x \psi^* - \frac{\partial^2 \psi^*}{\partial x^2} \cdot x \psi \right) dx$ minus I have written has $\frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(\frac{\partial \psi}{\partial x} \cdot x \psi^* - \frac{\partial \psi^*}{\partial x} \cdot x \psi \right) dx$ over $\int \left(\frac{\partial \psi}{\partial x} \cdot \psi^* - \frac{\partial \psi^*}{\partial x} \cdot \psi \right) dx$ that I can write as $\frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(\frac{\partial \psi}{\partial x} \cdot \psi^* - \frac{\partial \psi^*}{\partial x} \cdot \psi \right) dx$ that the first term. Now, I have to make some correction because there are extra terms going to appear so $\frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(-\psi^* \frac{\partial \psi}{\partial x} + \frac{\partial \psi^*}{\partial x} \psi \right) dx$ minus $\frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(-\frac{\partial \psi}{\partial x} \cdot \frac{\partial \psi^*}{\partial x} \cdot x - \frac{\partial \psi}{\partial x} \cdot \psi^* + \frac{\partial \psi^*}{\partial x} \cdot \frac{\partial \psi}{\partial x} \cdot x + \frac{\partial \psi^*}{\partial x} \cdot \psi \right) dx$. Now, the first term can be integrated directly and if ψ vanish at infinity boundary condition gives you that this is going to be 0.

In the second term, these two terms are going to be cancel. So, what I am left with is $\frac{d}{dt} \langle x \rangle$ is equal to $\frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(\psi^* \frac{\partial \psi}{\partial x} - \frac{\partial \psi^*}{\partial x} \psi \right) dx$ with the minus sign plus $\frac{i}{\hbar} \frac{\hbar^2}{2m} \int \left(\frac{\partial \psi}{\partial x} \cdot \psi^* - \frac{\partial \psi^*}{\partial x} \cdot \psi \right) dx$ integrated over dx . The second term $\int \left(\frac{\partial \psi}{\partial x} \cdot \psi^* - \frac{\partial \psi^*}{\partial x} \cdot \psi \right) dx$ can be return as integration of $\frac{\partial}{\partial x} (\psi^* \psi)$ minus integration of $\psi^* \frac{\partial \psi}{\partial x}$ over dx . Again the first integral vanishes because ψ vanish that large distances.

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The image shows a digital whiteboard with a toolbar at the top. The derivation is written in blue ink on a lined background. It starts with the equation $\frac{d}{dt} \langle x \rangle = \frac{i\hbar}{2m} \int -\psi^* \frac{\partial \psi}{\partial x} dx \times 2$. The '2' in the denominator and the '2' at the end of the integral are crossed out. This is followed by $= \frac{1}{m} \int (\psi^* \frac{\hbar}{i} \frac{\partial \psi}{\partial x}) dx$ and $= \frac{1}{m} \langle p \rangle$. A blue box encloses the final result: $\frac{d}{dt} \langle x \rangle = \frac{\langle p \rangle}{m}$. Below this, the text 'Second theorem:' is written, followed by the equation $\frac{d}{dt} \langle p \rangle = \langle -\frac{\partial V}{\partial x} \rangle$.

$$\frac{d}{dt} \langle x \rangle = \frac{i\hbar}{2m} \int -\psi^* \frac{\partial \psi}{\partial x} dx \times 2$$
$$= \frac{1}{m} \int (\psi^* \frac{\hbar}{i} \frac{\partial \psi}{\partial x}) dx$$
$$= \frac{1}{m} \langle p \rangle$$

$$\frac{d}{dt} \langle x \rangle = \frac{\langle p \rangle}{m}$$

Second theorem: $\frac{d}{dt} \langle p \rangle = \langle -\frac{\partial V}{\partial x} \rangle$

So, what I am left with is $\frac{d}{dt}$ of x is equal to $i\hbar$ cross over $2m$ integration of minus $\psi^* \frac{\partial \psi}{\partial x} dx$ times 2. This two cancels, and I can write this as $\frac{1}{m}$ integration $\psi^* \frac{\hbar}{i} \frac{\partial \psi}{\partial x} dx$, and this is nothing but $\frac{1}{m}$ expectation value of p . So, what I have shown is the $\frac{d}{dt}$ of x is expectation value of p divided by m . This is a first Ehrenfest theorem that again tells us that the expectation value of p or expectation value of x can be thought of as the average values of the momentum and position of the particle. The second theorem we have to show that the $\frac{d}{dt}$ of expectation value of p is equal to nothing but minus $\frac{\partial V}{\partial x}$ expectation value which is the expectation value of the force.

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$$\frac{d}{dt} \langle p \rangle = \frac{d}{dt} \int \psi^* \frac{\hbar}{i} \frac{\partial \psi}{\partial x} dx$$

$$= \int \left(\frac{\partial \psi^*}{\partial t} \right) \cdot \frac{\hbar}{i} \frac{\partial \psi}{\partial x} dx + \int \psi^* \frac{\hbar}{i} \frac{\partial}{\partial x} \left(\frac{\partial \psi}{\partial t} \right) dx$$

Write $\frac{\partial \psi^*}{\partial t}$ and $\frac{\partial \psi}{\partial t}$ using the TDSE

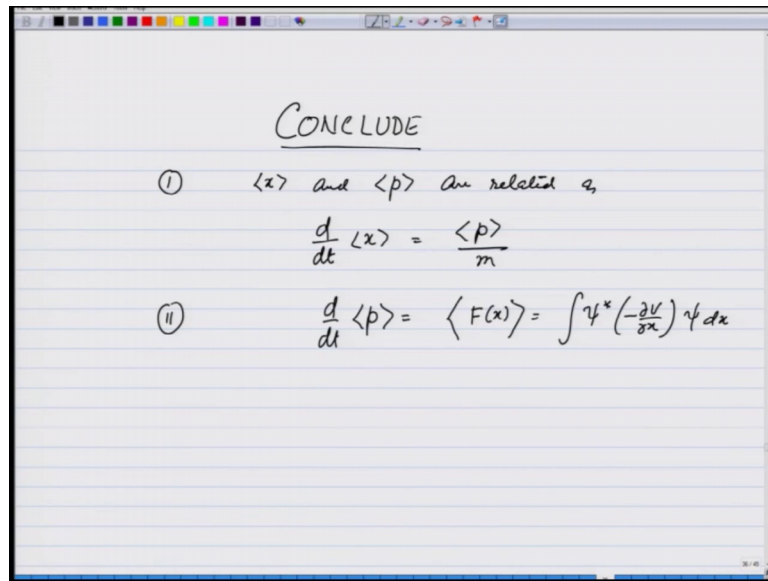
$$\frac{d}{dt} \langle p \rangle = \left\langle -\frac{\partial V}{\partial x} \right\rangle = \int \psi^* \left(-\frac{\partial V}{\partial x} \right) \psi dx$$

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To prove Ehrenfest second theorem we again do what we did earlier for proving the first theorem that is write d by dt of p as d by dt of this expectation value which is ψ^* \hbar $\frac{\partial \psi}{\partial x}$ dx . And this can be expanded has equal to integration $\frac{\partial \psi^*}{\partial t}$ \hbar $\frac{\partial \psi}{\partial x}$ dx plus integration ψ^* \hbar $\frac{\partial}{\partial x}$ $\left(\frac{\partial \psi}{\partial t} \right)$ dx of d ψ over dt dx . I interchange the partial derivative is with a respect to x and t that we can do.

The next steps going to be write d ψ^* dt and d ψ over dt using the time dependent Schrodinger equation as we did earlier. And we do that and substitute this in the equation an d by dt p equals after the manipulations and partial integration and all that, you get this equal to minus dv by dx expectation value which is nothing but integration of ψ^* $-\frac{dv}{dx}$ ψ dx which is the second theorem of Ehrenfest which tells you that average value of p or the expectation value of the p changes according to Newton's laws. I have not done the derivation here I am leaving that has an exercise, because it is a very simple derivation although little longish. So, I leave it for you.

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So, what we have shown in this lecture is conclude one at x and p are related as d by dt of x is expectation value p by m . And two - that d by dt of expectation value of p is nothing but the expectation value of the force which is same as ψ star minus dv by dx ψ dx . So, second tells you that these quantities are the average p and average x for particle n state of ψ .