

**Introduction to Quantum Mechanics**  
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**Lecture – 06**  
**Solution of the stationary-state Schrodinger equation for a simple harmonic oscillator**

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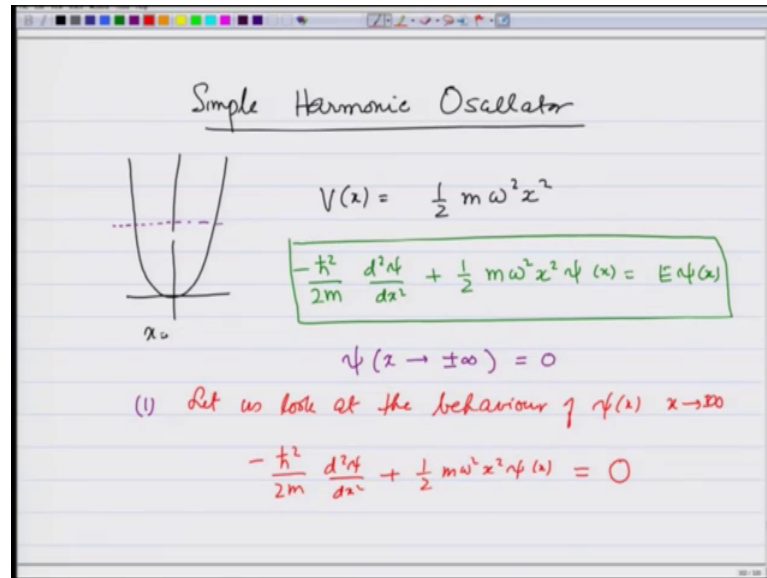
The image shows a digital whiteboard with handwritten text and equations. At the top, it says "Wave equation for material waves". Below that, it says "Schrodinger Equation" and shows the 3D equation: 
$$-\frac{\hbar^2}{2m} \nabla^2 \psi_n(x,y,z) + V(x,y,z) \psi_n(x,y,z) = E_n \psi_n(x,y,z)$$
 The 1D version of the equation is boxed in red: 
$$-\frac{\hbar^2}{2m} \frac{d^2 \psi_n}{dx^2} + V(x) \psi_n(x) = E_n \psi_n(x)$$
 At the bottom, it says "Equation is to be solved with boundary conditions And gives Eigenvalues & Eigenfunctions".

In the previous lecture, we introduced the wave equation for material waves and this the Schrodinger equation which is minus  $\hbar$  cross square over  $2m$ , I am writing the 3 dimensional origin  $\psi(x, y, z)$  plus  $V(x, y, z) \psi(x, y, z)$  equals  $E \psi(x, y, z)$  and we saw that with the boundary conditions, it gave us Eigen function, so I am going to now put,  $E_n$  as a subscript. So, that it indicates that there is a Eigen function  $\psi_n$  corresponding to energy  $E_n$  in one d. This equation was minus  $\hbar$  cross square over  $2m$  d<sup>2</sup>  $\psi$  over d x square plus  $V(x) \psi(x)$  equals  $E \psi(x)$  and again I am going to put that  $n$  to denote that there is an Eigen function.

It satisfies the boundary conditions for those Eigen values  $E_n$ . So, this is equation you solve it for one dimensional particle in a box problem and found the energy Eigen values to be the same as those given by old quantum theory the equation is to be solved to be solved with boundary conditions and gives Eigen energies and Eigen functions. So, this is a completely new way of looking at the quantum problem I am solving a wave

equation and through the Eigen values and Eigen energy is I am getting my answers and we yet to interpret psi, but let us solve in this lecture this equation for another one dimensional problem that we already know the answer of.

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And that is a simple harmonic oscillator for which a particle moves in a potential let us see the potential is centered around  $x$  equals 0.

So, the potential that is given for the particle of mass  $m$  is one half  $m$  omega square  $x$  square and therefore, the wave equation or the Schrodinger's equation for this becomes minus  $\hbar$  cross square over  $2m$   $d^2\psi$  over  $dx^2$  plus one half  $m$  omega square  $x$  square  $\psi$  equals  $E\psi$  and this equation is to be solved with the boundary condition and boundary condition here is going to be because now suppose I take a particle of energy  $E$ ; it can go anywhere and you will see actually wave can also penetrate through the region of finite potentials. So,  $\psi$  as  $x$  goes to plus or minus infinity should be 0 that is a boundary condition we are going to solve this equation with because far away from the origin is hardly any chance that I will find the particle and therefore, the wave function must vanish there.

Let us see; how do we go about solving equation. So, first let us look at the behavior of  $\psi$  for extending to very large value plus or minus infinity in that case, I am going to have minus  $\hbar$  cross square over  $2m$   $d^2\psi$  over  $dx^2$  plus one half  $m$  omega square  $x$  square  $\psi$  as  $x$  goes to plus or minus infinity the term  $m$  omega square  $x$  square

becomes very very large and  $\frac{d^2\psi}{dx^2}$  can also become very large. So,  $E$  is going to be insignificant. So, I can approximately write this as 0 that will give me the asymptotic solutions it is not the true solution just the solution when  $x$  goes to plus or minus infinity now let us solve this.

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The image shows a whiteboard with the following handwritten mathematical steps:

$$x \rightarrow \pm\infty \quad -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \frac{1}{2} m\omega^2 x^2 \psi(x) = 0$$

$$\frac{d^2\psi}{dx^2} - \frac{m^2\omega^2}{\hbar^2} x^2 \psi(x) = 0$$

$$\psi(x) = A e^{-\frac{m\omega}{2\hbar} x^2}$$

$$\frac{d\psi}{dx} = A \cdot -\frac{m\omega}{\hbar} x e^{-\frac{m\omega}{2\hbar} x^2}$$

$$\frac{d^2\psi}{dx^2} = A \left(\frac{m\omega}{\hbar}\right)^2 x^2 e^{-\frac{m\omega}{2\hbar} x^2} - A \frac{m\omega}{\hbar} e^{-\frac{m\omega}{2\hbar} x^2}$$

$$\lim_{x \rightarrow \infty} \psi(x) \sim A \left(\frac{m\omega}{\hbar}\right)^2 x^2 e^{-\frac{m\omega}{2\hbar} x^2}$$

$$\psi(x \rightarrow \pm\infty) \sim A e^{-\frac{m\omega}{2\hbar} x^2}$$

So, I have for  $x$  tending to plus or minus infinity  $-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \frac{1}{2} m\omega^2 x^2 \psi(x) = 0$  or I can write  $\frac{d^2\psi}{dx^2} - \frac{m^2\omega^2}{\hbar^2} x^2 \psi(x) = 0$ . Since we are only interested in the solution which is true in this limit, I am going to ignore any terms that are smaller than  $x^2$ . So, I can in that limit write  $\psi(x)$  as some constant  $e$  raised to minus  $\frac{m\omega}{\hbar}$  cross  $x^2$ . In fact,  $-\frac{m\omega}{2\hbar} x^2$ .

Let us see; how does this work out. So, if I take  $\frac{d\psi}{dx}$  by  $\frac{d}{dx}$  this is going to be equal to  $A$  times minus  $\frac{m\omega}{\hbar}$  cross  $x$   $e$  raised to minus  $\frac{m\omega}{2\hbar}$  cross  $x^2$  and if I now take the second derivative  $\frac{d^2\psi}{dx^2}$ . They will come out to be  $A$  times  $\frac{m\omega}{\hbar}$  cross  $x^2$   $e$  raised to minus  $\frac{m\omega}{2\hbar}$  cross  $x^2$  minus  $A$   $\frac{m\omega}{\hbar}$  cross  $e$  raised to minus  $\frac{m\omega}{2\hbar}$  cross  $x^2$ .

That is why you can easily check this by doing this differentiation and which in the limit of extending to infinity, I can write approximately as  $A \frac{m\omega}{\hbar} x^2 e^{-\frac{m\omega}{2\hbar} x^2}$  and you can see that this

cancels the other term in the equation therefore, it satisfies the equation. So, could this be the solution by itself, but right now what we have found is that  $\psi(x)$  as  $x$  tends to plus or minus infinity goes as  $A e^{-\frac{m\omega}{2\hbar} x^2}$ .

Let us now substitute this in the original equation the true Schrodinger equation and see what the answer comes out to be. So, if I put this in the original Schrodinger equation.

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$$-\frac{\hbar^2}{2m} \psi'' + \frac{1}{2} m \omega^2 x^2 \psi \quad \psi'' = \frac{d^2 \psi}{dx^2}$$

$$-\frac{\hbar^2}{2m} \left[ A \left( \frac{m\omega}{\hbar} \right)^2 x^2 \psi(x) - A \frac{m\omega}{\hbar} \psi(x) \right]$$

$$+ \frac{1}{2} m \omega^2 x^2 \psi(x)$$

$$\psi(x) = A e^{-\frac{m\omega}{2\hbar} x^2}$$

$$-\frac{1}{2} m \omega^2 x^2 \psi(x) + \frac{\hbar\omega}{2} \psi(x) + \frac{1}{2} m \omega^2 x^2 \psi(x)$$

$$\psi(x) = A e^{-\frac{m\omega}{2\hbar} x^2} = \frac{\hbar\omega}{2} \psi(x) \rightarrow 0 \text{ as } x \rightarrow \pm\infty$$

Which is minus  $\hbar^2$  over  $2m$  let me write  $\psi''$  for  $\frac{d^2 \psi}{dx^2}$  plus  $\frac{1}{2} m \omega^2 x^2 \psi$  where  $\psi''$  is nothing, but  $\frac{d^2 \psi}{dx^2}$ , then I find that this is going to be minus  $\frac{\hbar^2}{2m}$  and this we have found to be equal to their constant  $A \frac{m\omega}{\hbar} x^2 \psi$  minus  $A \frac{m\omega}{\hbar} \psi$  plus  $\frac{1}{2} m \omega^2 x^2 \psi$  times  $A \psi$ .

In fact, in some writing side even this  $A$  should not be there. So, let me; I can just cut this out, I can write that  $\psi(x)$  is  $A e^{-\frac{m\omega}{2\hbar} x^2}$ . So, I get minus  $\frac{1}{2} m \omega^2 x^2 \psi$  plus  $\frac{\hbar\omega}{2} \psi$  plus  $\frac{1}{2} m \omega^2 x^2 \psi$  and I cancel these 2 terms. So, I get  $\frac{\hbar\omega}{2} \psi$  you notice therefore, that this wave function  $\psi(x)$  equals some constant  $A e^{-\frac{m\omega}{2\hbar} x^2}$  goes to 0 as  $x$  tends to plus or minus infinity and satisfies the equation that when I put it in this equation gives you a constant time  $\psi$ .

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$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \frac{1}{2} m\omega^2 x^2 \psi(x) = E\psi(x)$$

$$\psi(x) = A e^{-\frac{m\omega}{2\hbar} x^2}$$

$$\text{LHS} = \frac{\hbar\omega}{2} A e^{-\frac{m\omega}{2\hbar} x^2}$$

$$\psi(x) = A e^{-\frac{m\omega}{2\hbar} x^2} \text{ is an eigenfunction with the energy eigenvalue}$$

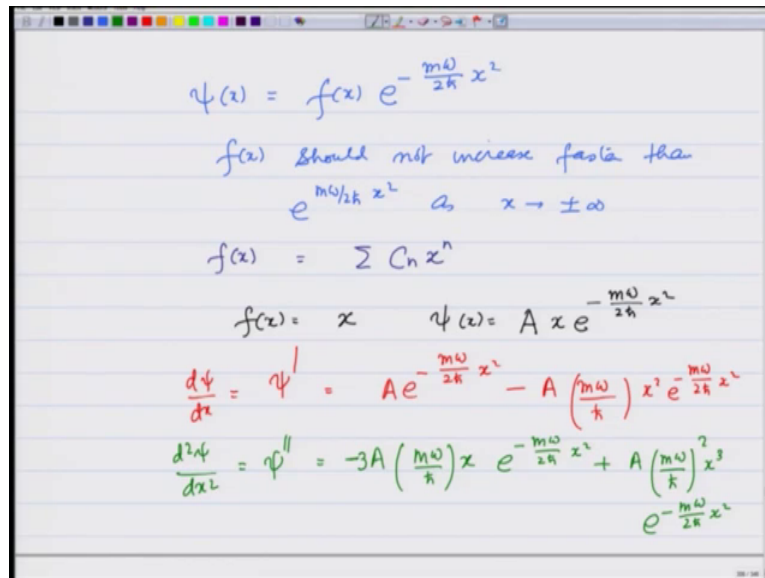
$$E = \frac{\hbar\omega}{2}$$

The image also contains a graph of the wavefunction  $\psi(x)$  versus  $x$ . The graph shows a parabolic potential well (black curve) and a Gaussian wavefunction (red curve) centered at  $x=0$ . The vertical axis is labeled  $\psi(x)$  and the horizontal axis is labeled  $x$ .

So, what I have shown you is that in the Schrodinger equation  $-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \frac{1}{2} m\omega^2 x^2 \psi = E\psi$  the Schrodinger equation for a particle in a simple harmonic oscillator potential, if I take  $\psi(x)$  to be  $A e^{-\frac{m\omega}{2\hbar} x^2}$  and substitute this in the equation. So, left hand side gives me  $\frac{\hbar\omega}{2} A e^{-\frac{m\omega}{2\hbar} x^2}$ ; that means, this itself  $\psi(x)$  equals  $A e^{-\frac{m\omega}{2\hbar} x^2}$  is an Eigen function because it satisfy a boundary conditions with the Eigen value  $E = \frac{\hbar\omega}{2}$ .

Recall that the lowest energy in the Heisenberg's formulation of quantum mechanics was also  $\frac{\hbar\omega}{2}$ . So, this represents the lowest energy Eigen state lowest energy Eigen function and how does this function look if you plot it, it is maximum in the middle and then goes down as  $e^{-x^2}$  this is how it looks. So, this is  $\psi(x)$  I plotted against  $x$  you can see that it has minimum number or largest possible wavelength and therefore, it is really the lowest energy system I mean this is I am just telling you know very qualitative way what about other Eigen states.

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$$\psi(x) = f(x) e^{-\frac{m\omega}{2\hbar} x^2}$$

$f(x)$  should not increase faster than  $e^{\frac{m\omega}{2\hbar} x^2}$  as  $x \rightarrow \pm\infty$

$$f(x) = \sum C_n x^n$$

$$f(x) = x \quad \psi(x) = A x e^{-\frac{m\omega}{2\hbar} x^2}$$

$$\frac{d\psi}{dx} = \psi' = A e^{-\frac{m\omega}{2\hbar} x^2} - A \left(\frac{m\omega}{\hbar}\right) x^2 e^{-\frac{m\omega}{2\hbar} x^2}$$

$$\frac{d^2\psi}{dx^2} = \psi'' = -3A \left(\frac{m\omega}{\hbar}\right) x e^{-\frac{m\omega}{2\hbar} x^2} + A \left(\frac{m\omega}{\hbar}\right)^2 x^3 e^{-\frac{m\omega}{2\hbar} x^2}$$

Now, that we can also generate by writing  $\psi(x)$  as  $\sum f(x) e^{-\frac{m\omega}{2\hbar} x^2}$  where  $f(x)$  should not increase faster than  $e^{\frac{m\omega}{2\hbar} x^2}$  as  $x$  tends to plus or minus infinity if it increases faster, then the wave function would not go to 0 as  $x$  goes to plus minus of infinity. So,  $f(x)$  should not increase faster. In fact, what I am going to do is a  $f(x)$  as a finite polynomial some constant  $C_n x^n$  and see what happens. So, let us now first try  $f(x)$  equals  $x$ . So, that the corresponding  $\psi(x)$  is going to be  $\sum$  constant  $A x e^{-\frac{m\omega}{2\hbar} x^2}$  then the corresponding  $\psi'$  which is  $d\psi/dx$  is equal to  $A e^{-\frac{m\omega}{2\hbar} x^2} - A \left(\frac{m\omega}{\hbar}\right) x^2 e^{-\frac{m\omega}{2\hbar} x^2}$ .

Let us take its double derivative. So,  $d^2\psi/dx^2$  which is  $\psi''$  is going to be equal to  $-3A \left(\frac{m\omega}{\hbar}\right) x e^{-\frac{m\omega}{2\hbar} x^2} + A \left(\frac{m\omega}{\hbar}\right)^2 x^3 e^{-\frac{m\omega}{2\hbar} x^2}$  even I differentiate the second term that also gives me an  $x$  term and therefore, this is going to become actually  $3A \left(\frac{m\omega}{\hbar}\right) x e^{-\frac{m\omega}{2\hbar} x^2} - A \left(\frac{m\omega}{\hbar}\right)^2 x^3 e^{-\frac{m\omega}{2\hbar} x^2}$  when I substitute all this in the Schrodinger equation what I am going to get is as follows.

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$$-\frac{\hbar^2}{2m} \psi'' + \frac{1}{2} m \omega^2 x^2 \psi = E \psi$$

$$-\frac{\hbar^2}{2m} \left[ -3A \left( \frac{m\omega}{\hbar} \right) x e^{-\frac{m\omega}{2\hbar} x^2} + A \left( \frac{m\omega}{\hbar} \right)^2 x^3 e^{-\frac{m\omega}{2\hbar} x^2} \right]$$

$$+ \frac{1}{2} m \omega^2 x^2 A x e^{-\frac{m\omega}{2\hbar} x^2}$$

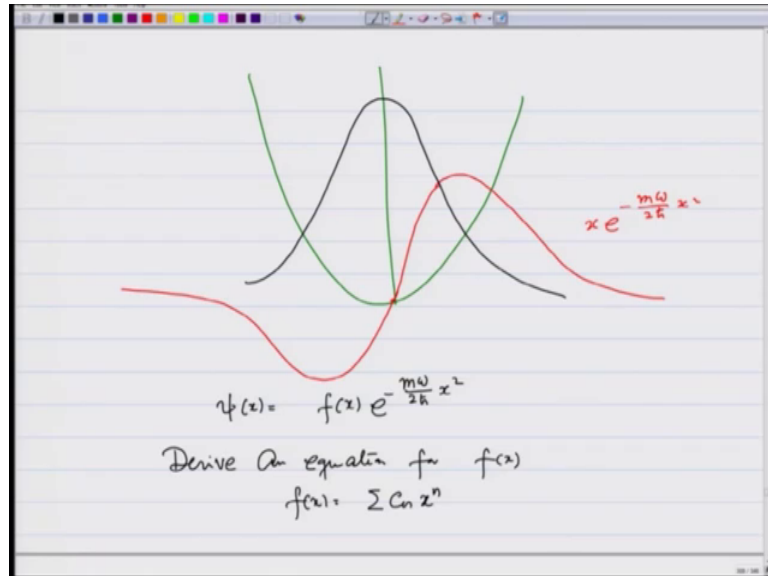
$$= + \left( \frac{3\hbar\omega}{2} \right) A x e^{-\frac{m\omega}{2\hbar} x^2} = \left( \frac{3\hbar\omega}{2} \right) \psi(x)$$

An eigenvalue                      wavefunction

So, the Schrodinger equation is minus h cross square over 2 m psi double prime plus one half m omega square x square psi equals e psi and I am going to substitute this. So, I get minus h cross square over 2 m inside I am going to get minus 3 A m omega over h cross times x e raised to minus m omega over 2 h cross x square plus A m omega over h cross square x cube e raised to minus m omega over 2 h cross x square plus 1 half m omega square x square times A x e raised to minus m omega over 2 h cross x square if you look at it carefully you will notice that this term cancels with this term.

When you multiplied by minus h cross over 2 m and what you get is equal to plus 3 h cross omega by 2 times x e raised to minus m omega over 2 h cross x square which is nothing, but 3 h cross omega by 2 times psi x. So, again find the constant times the wave function and this is also therefore, an Eigen value and this is the corresponding wave function how does it look.

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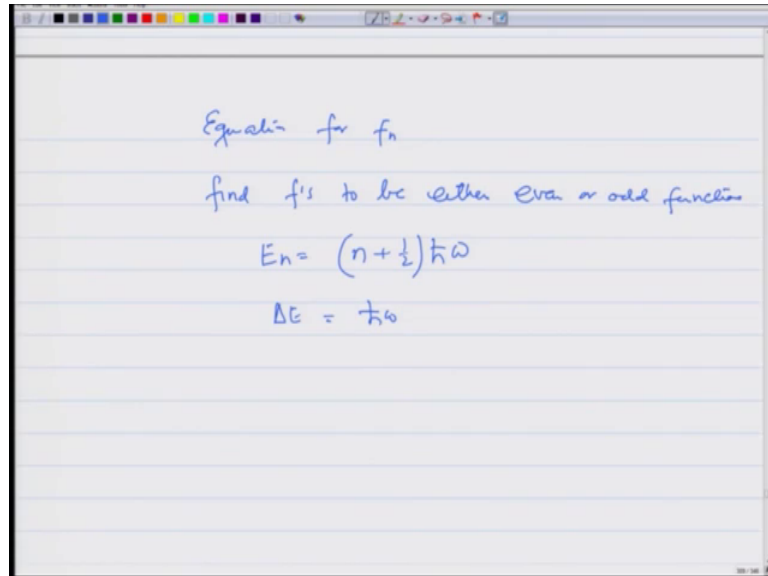


So, if I look at this potential this new wave function is 0 at the origin goes up and comes down on the other side it comes down and goes like this. This is  $x e^{-\frac{m\omega}{2\hbar}x^2}$  and compare to the ground state wave function which was looking like this you can see that this has more bigger more turning its lambda is shorter. So, its kinetic energy is higher.

So, it is higher energy in fact, you can show that as you go to higher and higher energies these are going to be wave function is more and more vehicles in general therefore, I can write  $\psi(x) = \sum f(x) e^{-\frac{m\omega}{2\hbar}x^2}$  and derive an equation for  $f(x)$  which I will give in the assignment and then write  $f(x)$  as  $\sum C_n x^n$  finite polynomial and solve for  $C_n$ s by matching the conditions in the 2 sides and when you do that what you find is  $f(x)$ s come out to be either even or odd.

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So, you derive an equation for  $f_n$  and find  $f_n$ 's to be either even or odd functions and energies come out to be  $n$  plus one half  $\hbar$  cross  $\omega$  which is exactly the same that we found in Heisenberg formulation of quantum mechanics.

So, this is another instance where the wave equation has given me the results already known  $\Delta E$  is  $\hbar$  cross  $\omega$ , but the answer that comes is through a very different approach by solving a wave equation with appropriate boundary conditions. The two approaches equivalent is Heisenberg's formulation the same as Schrodinger's formulation is just that is two different ways of writing this will explore next week for the time being we have now found a fundamental equation known as a Schrodinger equation which gives the same answers in the cases that we have solved as other theories. So, this is the fundamental equation of quantum mechanics.