

Introduction to Quantum Mechanics
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Lecture – 06

Heisenberg's formulation of quantum mechanics: the quantum condition

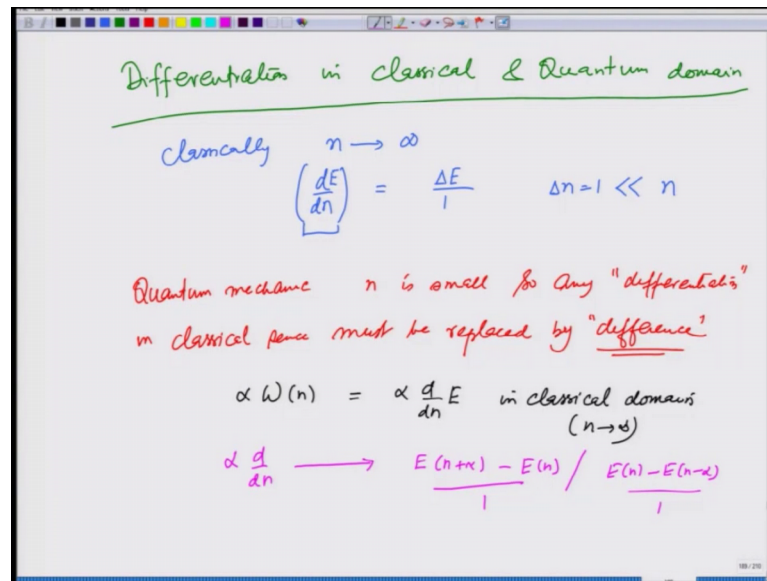
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$x(t) = \{ C_{n, n-\alpha}, \omega_{n, n-\alpha} \}$
 $(C_{n, n-\alpha} e^{i \omega_{n, n-\alpha} t})$
 $y(t) = \{ D_{n, n-\alpha}, \omega_{n, n-\alpha} \}$
 $x(t) y(t) \neq y(t) x(t)$
 not necessarily equal | $x(t) \cdot x(t) =$
 Obtain the quantum condition in terms
 $\{ C, \omega \}$

In the previous lecture, we introduced the quantum mechanical formulation of kinematic quantities by Heisenberg, we said if there is a quantity $x(t)$ it is represented by a collection of numbers C_n , let us say $n - \alpha$ and the corresponding frequency $\omega_{n, n - \alpha}$ and the time dependence of $C_{n, n - \alpha}$ is given as $e^{i \omega_{n, n - \alpha} t}$, then this collection of numbers gives this, there could be another quantity $y(t)$ which can be represented by collection of $D_{n, n - \alpha}$ and $\omega_{n, n - \alpha}$ and there I said that $x(t) y(t)$ is not equal to $y(t) x(t)$. So, what I should say is not necessarily equal to.

For example if I have $x(t)$ and multiplied by $x(t)$ that would; obviously, be the same if we change it here, but the way we calculate the representation has become, now matrix multiplication and what we want to do in this lecture is obtain the quantum condition in terms of these numbers C and ω that is our goal.

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So, again we will be going back to the correspondence principle, but before that let us now look for the time being differentiation in classical and quantum domain.

We have been doing it, but now let us explicitly say it classically, the quantum number n is very large and therefore, we can afford to do differentiation and which is the same as ΔE suppose i , this dE by dn by one because Δn which is one is much much much much smaller than n in the classicology in quantum mechanics this dE by dn would not be possible because numbers are small. So, in quantum mechanics n is small. So, any differentiation in classical sense must be replaced by difference that is how we distinguish between the 2.

So, for example, recall that $\omega(n) = \alpha \frac{dE}{dn}$ was equal to $\alpha \frac{dE}{dn}$ in classical domain; that means, n tending to infinity in quantum domain $\alpha \frac{dE}{dn}$ will go over to $\frac{E(n+\alpha) - E(n)}{1}$ that is the difference or $\frac{E(n) - E(n-\alpha)}{1}$ divided by 1 because the Δn is 1, this is the difference when we make a transition from classical to quantum. It has to be differences and $\alpha \frac{dE}{dn}$ is going to be this difference where the 2 quantum numbers differ by α and we are going to make use of it that is why I wrote it.

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Quantum condition $\oint p dx = n h$

Heisenberg: do not keep this in integral form

$\oint p dx = n h + \text{Constant}$

$\frac{d}{dn} \oint p dx = h$

$\left[\oint p dx \right]_n - \left[\oint p dx \right]_{n-1} = h$

$\underline{n \rightarrow \infty} \quad \oint p dx = m \oint v^2 dt = m \int_0^T v^2 dt$

Now, recall that the quantum condition that gave us results before Heisenberg paper was this condition $\oint p dx$ equals $n h$ first thing Heisenberg said; do not keep this in integral form why because this $\oint p dx$ could be $n h$ plus some constant and we do not know this constant. So, it should be replaced by a differential form and therefore, what he said is replaced by taking derivative with respect to n on both sides $\oint p dx$ equals h and once you integrate it that constant may come out which would basically in quantum mechanical sense this would mean that integral $\oint p dx$ for n minus integral $\oint p dx$ for n minus 1 is equal to h , this is what he first replaces the quantum condition by and now we are going to derive this in terms of C and ω . So, let us do that now.

So, when we write $\oint p dx$ in the n tending to infinity regime you see again, we are going to make use of correspondence principle, I will derive some results from n equals infinity and then immediately translate them to quantum mechanical language and that is the use of correspondence principle that Heisenberg made, it is going to be equal to mass times v square $d t$ over a period which I can also write as mass integral 0 to t v square $d t$ that is the action.

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$$\begin{aligned}
 n \rightarrow \infty \quad x_n(t) &= \sum_{\alpha} C_{\alpha} e^{i\alpha \omega(n) t} \\
 v(t) &= \sum_{\alpha} (i\alpha C_{\alpha}) e^{i\alpha \omega(n) t} \\
 v^2(t) &= \sum_{\alpha} \sum_{\beta} (i\alpha C_{\alpha} \omega(n)) (i\beta C_{\beta} \omega(n)) e^{i(\alpha+\beta) \omega(n) t} \\
 \int_0^T v^2(t) dt &= \sum_{\alpha} \sum_{\beta} [-\alpha\beta C_{\alpha} C_{\beta} \omega^2(n)] \int_0^T e^{i(\alpha+\beta) \omega(n) t} dt \\
 T &= \frac{2\pi}{\omega(n)} \\
 \Rightarrow \int_0^T e^{i(\alpha+\beta) \omega(n) t} dt &= T \quad \text{for } \alpha = -\beta \\
 &= 0 \quad \text{for all other } \alpha \neq \beta
 \end{aligned}$$

So, let us evaluate this is still in the n tending to infinity regime x t is summation, let us try it $\alpha C_{\alpha} e^{i\alpha \omega(n) t}$ which depends on this n t or the n th level. So, v this n th level t is going to be equal to summation $\alpha i\alpha C_{\alpha} e^{i\alpha \omega(n) t}$ and there is a $\omega(n)$ also $\alpha \omega(n)$ which depends on n t that is v α or v t .

So, v^2 t is going to be equal to be equal to summation α summation β $i\alpha C_{\alpha} \omega(n)$ times $i\beta C_{\beta} \omega(n)$ $e^{i(\alpha+\beta) \omega(n) t}$ and when I do integration 0 to t v^2 t , dt is going to be summation α summation β , I can write this as minus $\alpha\beta C_{\alpha} C_{\beta} \omega^2(n)$ for the n th level $e^{i(\alpha+\beta) \omega(n) t} dt$ 0 to t .

And let me remind you this t is nothing but 2π over $\omega(n)$, the period for the n th orbit the for the lowest frequency all the other are the harmonics. And therefore, this implies that integral 0 to t $e^{i(\alpha+\beta) \omega(n) t} dt$ is going to be equal to t for α equals minus β and 0 for all other α and β because in that case i raise to $i(\alpha+\beta) \omega(n) t$ add capital T would be same at 0 and t .

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$$\int_0^T v^2(t) dt = \sum_{\alpha} (\alpha^2 \omega^2(n) C_{\alpha} C_{-\alpha} T)$$

$$= 2\pi \sum_{\alpha} \alpha^2 \omega(n) C_{\alpha} C_{-\alpha}$$

$$\boxed{n \rightarrow \infty \quad \oint p dx = 2\pi m \sum_{\alpha} \alpha^2 \omega(n) C_{\alpha} C_{-\alpha}}$$

↓ Translate this into quantum result

$$\frac{d}{dn} \oint p dx = h$$

$$2\pi m \frac{d}{dn} \sum_{\alpha} \alpha^2 \omega(n) C_{\alpha} C_{-\alpha} = h$$

So, that is the period and only when alpha equals minus beta or beta equals minus alpha then this quantity becomes one and therefore, the answer becomes capital T and therefore, I can write integral P square t d t 0 to t is equal to summation alpha beta is now replaced by minus alpha. So, this is going to be alpha square omega square n C alpha C minus alpha t that is all and 1 omega and t multiplication gives me 2 pi. And therefore, I can write this as 2 pi summation alpha square omega n C alpha C minus alpha that is the classical result.

So, in the limit of n tending to infinity integration p dx comes out to be there was an m also there. So, I get 2 pi m summation alpha; alpha square omega n C alpha C minus alpha, right. Now I need to that is the classical result or intend to infinite result. Now I need to translate this into quantum result and not only that I need to write d by d n of integral p dx is equal to h. So, let us write that if I do that d by D n d by d n of summation alpha 2 pi m, I can write outside alpha square omega n C alpha C minus alpha is equal to h. Now I am going to make use of the correspondence principle.

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The whiteboard contains the following handwritten equations and annotations:

$$2\pi m \frac{d}{dn} \sum_{\alpha} \alpha^2 \omega(n) C_{\alpha} C_{-\alpha} = h$$

Annotations below the first equation:

- A purple arrow points from $\alpha^2 \omega(n)$ to $\alpha(\omega(n)) = \omega_{n, n-\alpha}$.
- Green arrows point from C_{α} and $C_{-\alpha}$ to $C_{n, n-\alpha}$ and $C_{n-\alpha, n}$ respectively.
- Below $C_{n-\alpha, n}$, it is noted that $C_{n-\alpha, n} = C_{n, n-\alpha}^*$.

$$\rightarrow 2\pi m \frac{d}{dn} \sum_{\alpha} \alpha \left[\omega_{n, n-\alpha} |C_{n, n-\alpha}|^2 \right] = h$$

$$\frac{\alpha d}{dn} \quad 2\pi m \sum_{\alpha} \alpha \frac{d}{dn} \left[\omega_{n, n-\alpha} |C_{n, n-\alpha}|^2 \right] = h$$

$$2\pi m \sum_{\alpha} \left[\omega_{n+\alpha, n} |C_{n+\alpha, n}|^2 - \omega_{n, n-\alpha} |C_{n, n-\alpha}|^2 \right] = h$$

So, I have $2\pi m \frac{d}{dn}$ of summation α ; $\alpha^2 \omega(n) C_{\alpha} C_{-\alpha}$ is equal to h when I do quantum mechanics, what does this mean? Remember, C_{α} and all these things are corresponding to the n th level.

So, C_{α} ; let me write each number separately, C_{α} would be $C_{n, n-\alpha}$ $C_{-\alpha}$ would be nothing but $C_{n-\alpha, n}$ which is same as $C_{n, n-\alpha}^*$ and $n-\alpha$, we have done this in the previous lecture and from here, I am going to take $\alpha \omega(n)$ as $\omega_{n, n-\alpha}$. So, when I translate, I will be very careful in this I will write $2\pi m \frac{d}{dn}$ of summation α ; α remains in the brackets, I am going to have $\omega_{n, n-\alpha} \text{ modulus of } C_{n, n-\alpha} \text{ square}$ is equal to h .

Now, recall that I had done something with what $\frac{d}{dn}$ means in quantum mechanics and I am going to use that now. So, I am going to write this as $2\pi m \sum_{\alpha} \alpha \frac{d}{dn}$ of summation α , I can take $\frac{d}{dn}$ inside because that is operating only on the brackets square brackets. So, I am going to write this $\alpha \frac{d}{dn} \omega_{n, n-\alpha} \text{ modulus of } C_{n, n-\alpha} \text{ square}$ is equal to h and this is nothing but difference over α . So, I am going to write this as $2\pi m \sum_{\alpha} \left[\omega_{n+\alpha, n} |C_{n+\alpha, n}|^2 - \omega_{n, n-\alpha} |C_{n, n-\alpha}|^2 \right] = h$ α varies from minus infinity to plus infinity.

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Quantum condition

$$2\pi m \sum_{\alpha=-\infty}^{\infty} [\omega_{n+\alpha, n} |c_{n+\alpha, n}|^2 - \omega_{n, n-\alpha} |c_{n, n-\alpha}|^2] = h$$

$$= 2\pi m \sum_{\alpha=-\infty}^{\infty} \omega_{n+\alpha, n} |c_{n+\alpha, n}|^2 - 2\pi m \sum_{\alpha=-\infty}^{\infty} \omega_{n, n-\alpha} |c_{n, n-\alpha}|^2 = h$$

$$\beta = -\alpha \quad \sum_{\beta=-\infty}^{\infty} \omega_{n, n+\beta} |c_{n, n+\beta}|^2 = - \sum_{\beta=-\infty}^{\infty} \omega_{n+\beta, n} |c_{n+\beta, n}|^2$$

$\omega_{n, n+\beta} = -\omega_{n+\beta, n}$
 $|c_{n+\beta, n}| = |c_{n, n+\beta}|$

$4\pi m \sum_{\alpha=-\infty}^{\infty} \omega_{n+\alpha, n} |c_{n+\alpha, n}|^2 = h$

So, let me rewrite this to the quantum condition comes out to be $2\pi m$ summation α equals minus infinity to infinity $\omega_{n+\alpha, n}$ modulus of $C_{n+\alpha, n}$ square minus $\omega_{n, n-\alpha}$ modulus $C_{n, n-\alpha}$ mod square is equal to h , I can do even better and write this as $2\pi m$ summation α equals minus infinity to infinity $\omega_{n+\alpha, n} C_{n+\alpha, n}$ mod square minus $2\pi m$ summation α equals minus infinity to infinity $\omega_{n, n-\alpha}$ modulus $C_{n, n-\alpha}$ mod square equals h , I am going to change the second term this term.

I am going to replace β by minus α and write this as summation β equals infinity to minus infinity is the same thing, whichever order you run it, n $\omega_{n, n+\beta}$ modulus $C_{n, n+\beta}$ whole square which I can write as summation β equals minus infinity to infinity ω with a minus sign, I will change the $\omega_{n, n+\beta}$ modulus $C_{n, n+\beta}$ mod square I could do this because $\omega_{n, n+\beta}$ is minus $\omega_{n+\beta, n}$, n $C_{n, n+\beta}$ modulus is same as modulus $C_{n, n+\beta}$.

If I switch the 2, right, it does not really change whether transition is from n to $n+\beta$ or $n+\beta$ to n 2β . And therefore, I can write this quantum condition also as $2\pi m$ summation α equals minus infinity to infinity $\omega_{n+\alpha, n}$ modulus $C_{n+\alpha, n}$ plus α square and the second term comes from the other term. So, it becomes instead of 2 it becomes 4; $4\pi m$ equals h .

So, we have all the C_n s and ω s, we have the quantum condition transfer them and now we need to apply these conditions to calculate C and ω which we will do in the next lecture.