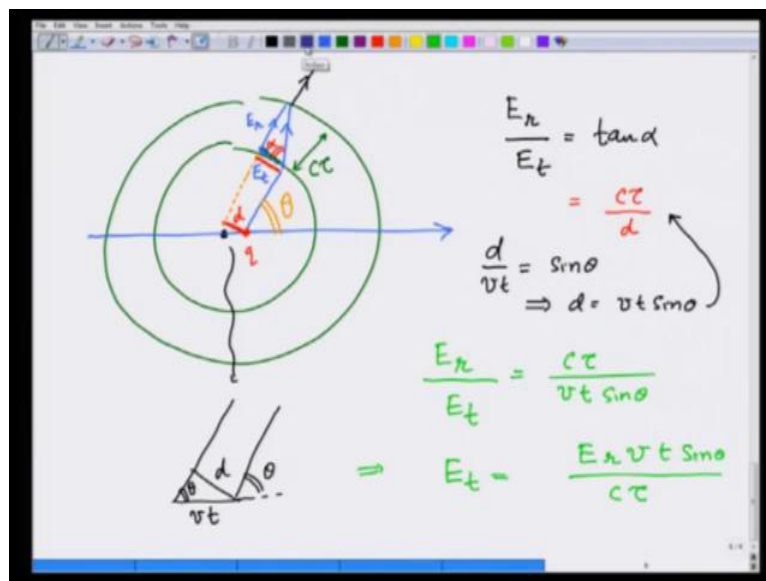


**Introduction to Electromagnetism**  
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**Lecture - 65**  
**Radiation from an accelerating charge – II**

So far we have established that one a charged accelerates the region or corresponding to acceleration, propagates out with speed  $c$ , and it has an electric field; that is, has a tangential component, all the components which perpendicular to direction of propagation of that region. In this lecture we are going to show that this field indeed has a  $1$  over  $r$  dependence and carries out power; that is  $1$  over  $r$  square, and that means, this is indeed radiation field.

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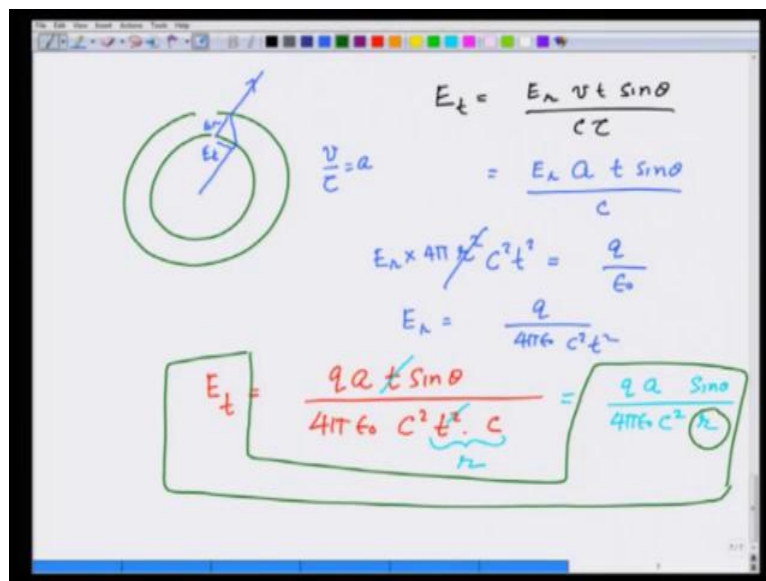


So, let me make a neat picture now concentrating only on one line of force. So, here is my inner circle, here is my outer circle, here is my charge later, here is my charge initially. This is the field line initially, this is the field line later, and I will connect them. And this charge is moving this way; charge  $q$ , and this has developed a radial component, and there is of course, component which is tangential component and the radial component. This is  $e$  plus calls it tangential, and the other forces  $e$  radial. Let this angle where I am observing the field, and the direction of acceleration be  $\theta$ . Let me

again connect the lines, this other line is like this. I have  $e$  and this distance  $c \tau$ , and I need this distance also.

So, let us see now I have  $e$  tangential or  $e$  radial divided by  $e$  tangential is equal to tangent of  $\alpha$ , where  $\alpha$  is this angle;  $e$  radial divided by  $e$  tangential is tangent of  $\alpha$ . This tangent of  $\alpha$  is also the same, as the distance  $c \tau$  divided by the distance here, which I am showing by red right next to it, which is the same distance is this. Let us call a  $d$ ;  $c \tau$  over  $d$ . Now, I will calculate  $d$ . Let me take this picture to the bottom and show that, this is  $d$ , this angle is  $\theta$ , this distance  $d$  at the lower distance. If this is  $\theta$ , this is also  $\theta$ . So, I have  $d$  over  $v t$  equals  $\sin$  of  $\theta$ , which implies that  $d$  is equal to  $v t \sin \theta$ . Substituting this here I get  $e$  radial divided by  $e$  tangential is equal to  $c \tau$  over  $v t \sin \theta$ , which give me  $e$  tangential is equal to  $e$  radial  $v t \sin \theta$  divided by  $c \tau$ . So, now we got a working formula with which we now start working.

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So, we have got  $e$  tangential  $e$  equals  $e$  radial  $v t \sin \theta$  over  $c \tau$ . Where again I make this picture inner circle, outer circle, the field line is very much like this. This is  $e$  tangential, this is  $e$  radial. Now,  $v$  over  $\tau$  is the acceleration, and therefore, I can write this formula as  $e$  radial acceleration times  $t \sin \theta$  over  $c$ . By Gauss's I have  $e$  radial times  $4 \pi r^2$ ;  $r$  is nothing, but  $c \tau$   $c$  square  $t$  square, so I can cut this and write this as  $c$  square  $t$  square is equal to total charge  $q$  divide by  $\epsilon_0$ ; and therefore,  $e$  radial is  $q$  over  $4 \pi \epsilon_0 c^2 t^2$ . And this immediately gives me  $e$  tangential is equal to  $q a t \sin \theta$

divided by  $4\pi\epsilon_0 c^2 r$ . Again I am going to cancel one of the  $t$ 's, and write  $ct$  as  $r$  again, because  $r$  is the distance up to which, this tangential component of electric field has moved, and I can write this as;  $qa \sin\theta$  divided by  $4\pi\epsilon_0 c^2 r$ . Notice, we accomplished  $4\pi$  for looking for. We have found that the tangential electric field is proportional to  $1$  over  $r$ , how far this tangential component has moved, and rest of the constants.

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The image shows a whiteboard with a diagram and mathematical derivations. The diagram on the left depicts a charge moving with acceleration  $a$  (indicated by a red arrow) and the resulting electric field lines. The derivations on the right are as follows:

$$E_t = \frac{qa \sin\theta}{4\pi\epsilon_0 c^2 r}$$

$$B_t = \frac{E_t}{c} = \frac{qa \sin\theta}{4\pi\epsilon_0 c^3 r}$$

$$|\vec{S}| = \epsilon_0 E^2 c = \frac{1}{\mu_0} \vec{E} \times \vec{B}$$

$$= \frac{1}{\mu_0} \frac{q^2 a^2 \sin^2\theta}{16\pi^2 \epsilon_0^2 c^3 r^2}$$

$$= \frac{q^2 a^2}{16\pi^2 \epsilon_0 c^3} \left( \frac{\sin^2\theta}{r^2} \right)$$

So, what we have gotten in this case is, that  $e$  tangential is equal to  $qa \sin\theta$  over  $4\pi\epsilon_0 c^2 r$ . If I look at this charge which was given acceleration in this direction, it develops tangential field which is zero here,  $\sin\theta$  is maximum here at  $\pi/2$ . It is very large here, opposite in direction to the acceleration. Smaller here, even smaller here, becomes zero at  $\theta$  is equal to zero. The other way on the other side, very large and elimination again zero.

Now, how about the pointing vector or  $b$  tangential;  $b$  tangential I have already seen as  $e$  tangential  $c$ , and therefore, it is going to be  $qa \sin\theta$  over this is  $c^2$ ;  $4\pi\epsilon_0 c^2 r$ . The pointing vector  $s$ , which is going to be radially out, is going to be equal to one half  $\epsilon_0 e^2$  times  $c$ , which is the same as  $1$  over  $\mu_0 e$  cross  $b$ . So, this is equal to one half of  $\epsilon_0 c$  times  $e^2$ , which is  $q^2 a^2 \sin^2\theta$  divided by  $16\pi^2 \epsilon_0^2 c^3 r^2$ . So, the pointing vector, this is magnitude and direction is readily outward is, I will cancel one of the  $\epsilon_0$  zeros. I

will cancel one of the c's becomes c cubed. So, this is equal to, this half is also not there right now, because I am not taking the average. Half you remember from before, because of the harmonic nature, this half is not there, is equal to  $\frac{2 q^2 a^2}{16 \pi \epsilon_0 c^3} \frac{\sin^2 \theta}{r^2}$ .

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The image shows a whiteboard with the following handwritten equations:

$$S = \frac{q^2 a^2}{16 \pi \epsilon_0 c^3} \frac{\sin^2 \theta}{r^2}$$

Below this, there is a diagram of a point charge  $q$  with a double-headed arrow indicating oscillation. To the right, the acceleration  $a$  is given as:

$$a = \ddot{x} = (A \sin \omega t)$$

$$= -\omega^2 A \sin \omega t$$

Then, the magnitude of the radiation field  $|\vec{S}|$  is calculated as:

$$|\vec{S}| = \frac{q^2 \omega^4 A^2}{16 \pi \epsilon_0 c^3} \frac{\sin^2 \omega t \sin^2 \theta}{r^2}$$

Finally, the average power is found by integrating over one cycle  $T$ :

$$\text{Average Power} = \frac{1}{T} \int_0^T = \frac{q^2 \omega^4 A^2}{32 \pi \epsilon_0 c^3} \left( \frac{\sin^2 \theta}{r^2} \right)$$

So, you see the power radiated is also going as  $1/r^2$ , which is very satisfying, because that's how power from a point source grows. Let us see now what happens if I take this point charge, and start moving back and forth  $q$ , which is making an oscillating motion, harmonic oscillating motion. So, we already saw that  $S$  is equal to  $\frac{q^2 a^2}{16 \pi \epsilon_0 c^3} \frac{\sin^2 \theta}{r^2}$ . In this case,  $a$  is equal to  $\ddot{x}$ , where  $x$  is nothing, but some amplitude  $\sin \omega t$ . So, which is going to be  $-\omega^2 A \sin \omega t$ .

And therefore, the power radiated by an oscillating charge would be equal to  $\frac{q^2 \omega^4 A^2}{32 \pi \epsilon_0 c^3} \frac{\sin^2 \theta}{r^2}$ . This is  $\sin \omega t$  not  $\sin^2 \omega t$ , over  $16 \pi \epsilon_0 c^3} \frac{\sin^2 \theta}{r^2}$ ; the average power. By average power mean, we average it over cycle so that we only observe that the frequency is very large, give me a factor of half, because of this. This was a mistake I was making earlier when I wrote the factor of half, so this is give me  $\frac{q^2 \omega^4 A^2}{32 \pi \epsilon_0 c^3} \frac{\sin^2 \theta}{r^2}$ . This is the average power that is given out by oscillating charge. Oscillating charge, since this moves away from the centre also

makes a dipole which is equal to  $q \cdot x$ . So, I can even think of this as an oscillating dipole, which is giving out radiation. Radiation is maximum in the direction perpendicular to the dipole, and it is zero in the direction of the dipole. How about the total power radiated.

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$$S = \frac{q^2 \omega^4 A^2}{32\pi \epsilon_0 c^3} \frac{\sin^2 \theta}{r^2} = \frac{\text{Power}}{\text{Area}}$$

$$\text{Total power} = \int \frac{q^2 \omega^4 A^2}{32\pi \epsilon_0 c^3} \frac{\sin^2 \theta}{r^2} \cdot d(\Omega) \cdot r^2$$

$$= \frac{q^2 \omega^4 A^2}{32\pi \epsilon_0 c^3} \times 2\pi \int_{-1}^1 (1 - \cos^2 \theta) d(\cos \theta)$$

$$= \frac{q^2 \omega^4 A^2}{16 \epsilon_0 c^3} \times \frac{4}{3}$$

$$\text{Total Power} = \frac{q^2 \omega^4 A^2}{12\pi \epsilon_0 c^3}$$

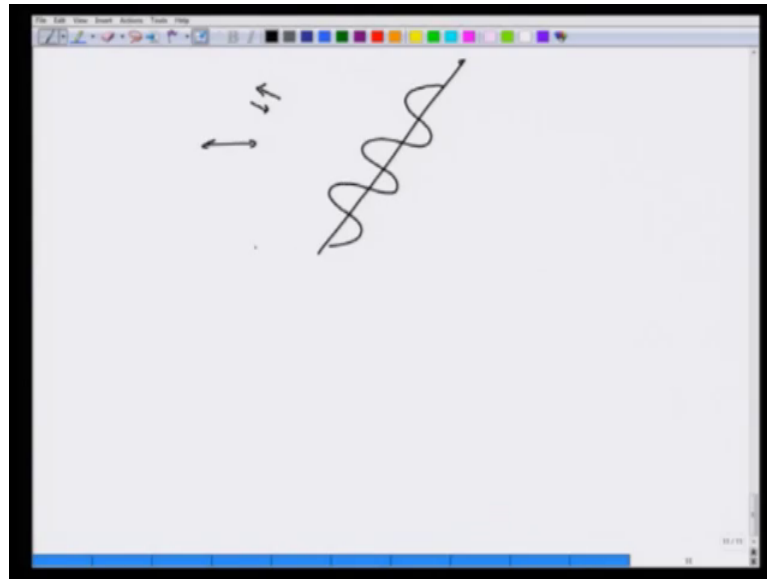
$$p = qA \sin \omega t$$

The total power radiated. So, this  $S$  is  $q^2 \omega^4 A^2$  over  $32\pi \epsilon_0 c^3$  times  $\sin^2 \theta$  over  $r^2$ , which is nothing but power per unit area. So, total power is going to be equal to  $q^2 \omega^4 A^2$  over  $32\pi \epsilon_0 c^3$  times  $\sin^2 \theta$  over  $r^2$  times  $d(\Omega)$  times  $r^2$ . Remember this is the area in direction  $r$ . Now,  $r^2$  cancel, there is no  $\phi$  dependence so that  $2\pi$  will come out, so this is integrated over. This becomes  $q^2 \omega^4 A^2$  over  $32\pi \epsilon_0 c^3$  times  $2\pi$  times integration  $\sin^2 \theta$ , I can write as  $1 - \cos^2 \theta$  from  $\cos \theta$  going from  $-1$  to  $1$ ,  $2\pi$  over  $32\pi$  gives me  $1/16$ .

So, this is going to be  $q^2 \omega^4 A^2$  over  $16 \epsilon_0 c^3$  times  $1 - \cos^2 \theta$  from  $\cos \theta$  gives me  $4/3$ . So, this is equal to  $q^2 \omega^4 A^2$  over  $12 \epsilon_0 c^3$ . Yes, I am missing a factor of  $\pi$ , if you go back to the formula this was  $\pi^2$  here, and therefore, I have  $\pi^2$  here, and  $\pi^2$  here,  $\pi^2$  here,  $\pi^2$  here. So, finally,  $\pi^2$  here, and therefore, finally, I have a  $12\pi$ . This is total power radiated by a charge oscillating back and forth or a dipole of dipole moment  $p = q \cdot a$ . If it is oscillating back and forth you have a

time dependence  $\sin \omega t$  that this dipole gives out; this is known as lambda formula for radiation. So, what I have shown in this lecture is very qualitative, drawing field lines, and things like those, that an accelerating charge develops a field, which is perpendicular to direction of propagation; it goes as  $1$  over  $r$  and therefore, is the radiation field.

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Now, imagine what will happen when this charge is moving back and forth. If it is moving back and forth, and therefore sometimes the field will be in one direction, sometimes is in the other direction, as the way propagates the field goes up and down like this, and this is the harmonic way of going out.