

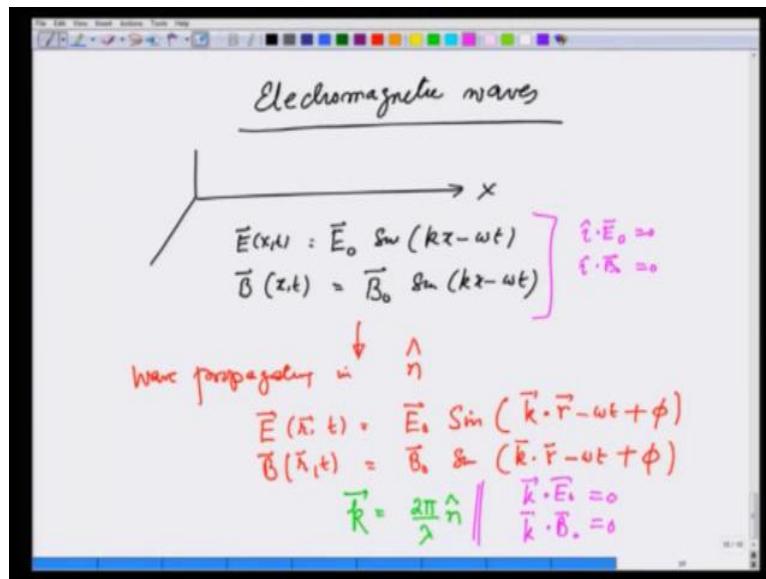
Introduction to Electromagnetism
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Module - 07

Lecture - 60

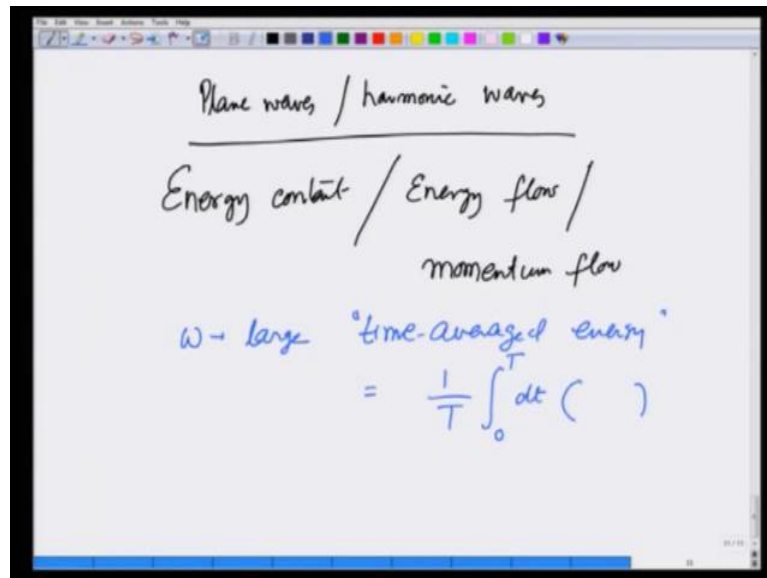
Energy and intensity and momentum carried by electromagnetic waves

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In the previous two lectures we talked about electromagnetic waves that come naturally out of Maxwell's equations. In particular we focused on wave propagating in the x direction and wrote this as $E(x,t) = E_0 \sin(kx - \omega t)$ and $B(x,t) = B_0 \sin(kx - \omega t)$. If you generalize this I can in general write for wave propagating in the direction of unit vector \hat{n} I can write $E(\vec{r}, t) = E_0 \sin(\vec{k} \cdot \vec{r} - \omega t + \phi)$ or $B(\vec{r}, t) = B_0 \sin(\vec{k} \cdot \vec{r} - \omega t + \phi)$. So, I can write this as $k \cdot r - \omega t$ and B as a function of r and p as equal to $B_0 \sin(k \cdot r - \omega t)$. And if I want two more general I will add a phase ϕ that can make a cosine or whatever where the vector \vec{k} is equal to $2\pi/\lambda$ in the direction of the propagation.

(Refer Slide Time: 02:28)



In the previous lecture we saw the waves for propagating I direction $\mathbf{k} \cdot \mathbf{e} = 0$ and $\mathbf{k} \cdot \mathbf{b} = 0$. In the same manner in general I am going to have $\mathbf{k} \cdot \mathbf{e} = 0$. That means the electric field is perpendicular to the direction of propagation and similarly $\mathbf{k} \cdot \mathbf{b} = 0$ and $\mathbf{e} \times \mathbf{b}$ is going to be in the direction of propagation. With this background what you want to focus on is again waves from this time which are harmonic and plane waves. So, plane waves, harmonic, mono chromatic or moving with frequency ω waves and see what is the energy content of these waves, how much energy do they carry and how much momentum do they contain. Remember, these are the quantity we have had already calculated for electro static field.

We will also focus on that ω is going to be large. So, we are talking about something like the ten raise fifteen, ten raise ten, ten raise to fourteen. That means, in one second this so many oscillation we are talking about light waves or microwaves. In that case, it does not make sense of energy variation but will be talking about time averaged energy and what we mean by this will be taking one over the time period and integrate whatever quantity we get from 0 to T over time. Whatever quantity energy density or the energy flow the momentum flow. So at such high frequency what we see is for example if I am looking at this light I am not seeing it varying I am seeing a constant light because it varies ten raise to fifteen times per second and therefore I see only the average light my eyes somehow averages it out I will

(Refer Slide Time: 04:12)

$$\vec{E} = \vec{E}_0 \sin(\vec{k} \cdot \vec{r} - \omega t) \quad \vec{B} = \vec{B}_0 \sin(\vec{k} \cdot \vec{r} - \omega t)$$

Energy content

$$\frac{1}{2} \epsilon_0 \vec{E}^2$$

$$= \frac{1}{2} \epsilon_0 E_0^2 \sin^2(\vec{k} \cdot \vec{r} - \omega t)$$

$$\frac{1}{2} \epsilon_0 E_0^2 \frac{1}{T} \int_0^T \sin^2(\vec{k} \cdot \vec{r} - \omega t) dt$$

$$= \frac{1}{4} \epsilon_0 E_0^2$$

$$\frac{1}{2\mu_0} B_0^2$$

$$= \frac{1}{2\mu_0} B_0^2 \sin^2(\vec{k} \cdot \vec{r} - \omega t)$$

Time averaged

$$\frac{1}{4\mu_0} B_0^2$$

$$= \frac{1}{4\mu_0} \frac{E_0^2}{c^2} = \frac{1}{4} \epsilon_0 E_0^2$$

Energy density in an EM wave = $\frac{1}{2} \epsilon_0 E_0^2$

So, we have e equal to $e_0 \sin k \cdot r - \omega t$ and b is equal to $b_0 \sin k \cdot r - \omega t$ if you want you can add π also doesn't really make a difference. Let's calculate the energy content recall that energy in the electric field is one half $\epsilon_0 e$ square and therefore in this case it is going to become one half $\epsilon_0 e_0$ square that is the amplitude square sine square $k \cdot r - \omega t$. Recall I said that this is varying in time very fast so I will be talking about average energy. Average energy is going to be one half $\epsilon_0 e_0$ square 0 to t one over t sine square $k \cdot r - \omega t$ at a given point r is fix minus ωt and this you know from your ordinary mathematics is one half.

So, this becomes one fourth of $\epsilon_0 t_0$ square. Let us now calculate what the energy in the magnetic field is magnetic is. Magnetic field you recall that it carries energy one over two $\mu_0 b_0$ square sine square $k \cdot r - \omega t$ which when time average will give me another factor of two. So, this becomes one over four $\mu_0 b_0$ square recall from previous lecture that b_0 is nothing but e_0 over c . So, this becomes one over four $\mu_0 e_0$ square but one over c square is $\mu_0 \epsilon_0$. So, this becomes one fourth $\epsilon_0 e_0$ square therefore the energy contained in the electric field and magnetic field component of this electromagnetic wave is exactly the same which is the one fourth $\epsilon_0 e_0$ square and you add the two and you get the energy density in an electromagnetic $e m$ wave to be one half $\epsilon_0 e_0$ square. So, when this wave exist for example this light is coming to me then in between average energy per unit volume is one half $\epsilon_0 e_0$ square.

(Refer Slide Time: 07:15)

Energy flow

Intensity = Poynting vector

Poynting vector = $\frac{1}{\mu_0} \vec{E} \times \vec{B}$

$$= \frac{1}{\mu_0} \hat{n} |\vec{E}_0| |\vec{B}_0| \sin^2(\vec{k} \cdot \vec{r} - \omega t)$$

$$= \frac{1}{\mu_0} \hat{n} \frac{|\vec{E}_0|^2}{c} \sin^2(\vec{k} \cdot \vec{r} - \omega t)$$

$$= \frac{1}{2} \frac{|\vec{E}_0|^2}{\mu_0} \sqrt{\frac{\epsilon_0}{\mu_0}} \hat{n}$$

$$= \frac{1}{2} \sqrt{\frac{\epsilon_0}{\mu_0}} E_0^2 \hat{n}$$

Energy flow: $c u$

$$= \frac{1}{\sqrt{\epsilon_0 \mu_0}} \frac{1}{2} \epsilon_0 E_0^2$$

$$= \frac{1}{2} \sqrt{\frac{\epsilon_0}{\mu_0}} E_0^2$$

So, this wave let us again take it as it is propagating in x direction if I take a small box here of unit volume for small volume dv then we content this is going to be one half $\epsilon_0 E_0^2 dv$ how about the energy flow which is related to intensity, which is related to pointing vector. The pointing vector is given as one over μ_0 $\vec{E} \times \vec{B}$ and we all ready seen that $\vec{E} \times \vec{B}$ is in the direction of propagation. So, I can write this as one over μ_0 direction \hat{n} E_0 and modulus of $B_0 \sin^2(\vec{k} \cdot \vec{r} - \omega t)$ we recall that is again E_0 by c . So, I can write this as one over μ_0 in the direction of propagation E_0^2 over $c \sin^2(\vec{k} \cdot \vec{r} - \omega t)$ when i take the time average $\sin^2(\vec{k} \cdot \vec{r} - \omega t)$ is going to give me a factor of two one half and therefore this becomes one half E_0^2 over μ_0 and c is nothing but square root of $\epsilon_0 \mu_0$ one. So, this become one half square root of ϵ_0 over μ_0 E_0^2 in the direction of propagation.

That, means if there is an electromagnetic, plane electromagnetic wave going in some direction in the same direction per unit area per unit time the energy that flows is one half square root of ϵ_0 and μ_0 E_0^2 square. Let us also see from the energy density point of view the energy flow if you recall form one or previous lecture is going to be nothing but c time energy density u which is going to be one over square root of $\epsilon_0 \mu_0$ times one half $\epsilon_0 E_0^2$ which is same as one half square root of ϵ_0 over μ_0 E_0^2 square same answer.

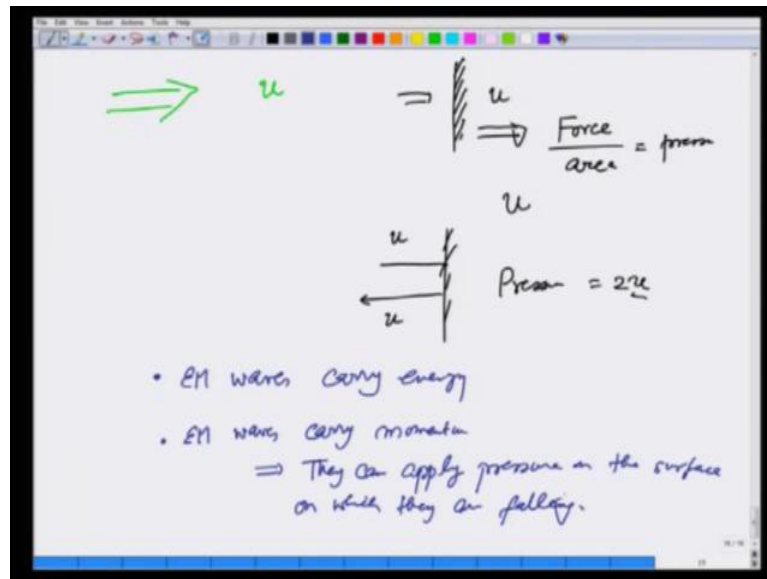
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The image shows a handwritten derivation on a whiteboard. It starts with the word "Momentum:" followed by "momentum density =". The main equation is $\frac{|\vec{S}|}{c^2} = \frac{MC^2 T^{-1}}{L^2 \cdot T L^{-1}} = \frac{MLT^{-1}}{L^3}$. A circled term $\frac{|\vec{S}|}{c^2}$ is equated to $\frac{cu}{c^2} = \frac{u}{c}$. Below this, the expression $\frac{1}{2} \sqrt{\frac{\epsilon_0}{\mu_0}} \frac{E_0^2}{c^2}$ is shown. The final result is "Momentum flow = c. momentum density = $\frac{1}{2} \sqrt{\frac{\epsilon_0}{\mu_0}} \frac{E_0^2}{c} = u$ ".

So, that is energy that is carried by the electromagnetic wave or that is the intensity that we received. This is also related to the momentum carried by electromagnetic wave. Recall that the momentum density is nothing but s over c square let us check that modulus of s is nothing but energy $m l$ square t minus two per unit area per unit time which then becomes equal to momentum density $m l t$ minus one over l cube. So, if I divide by C Square here C Square is nothing but l square t minus two. This cancels from doing c square one of the l I cancel and make it l by get $m l t$ minus one over l cube. So, this is the momentum density and therefore if there is a plane electromagnetic wave going the momentum density is given as one half square root of ϵ_0 over μ_0 e_0 square over c square. Given this momentum density what is the movement flow.

The momentum flow is nothing but c times of momentum density which is going to be equal to one half ϵ_0 over μ_0 e_0 square over c . We, write them in terms of u so s is nothing but c times u , u is the energy density in the electromagnetic field over c square which is μ_0 over c and the momentum flow is nothing but u itself. If, there is momentum flow that means there is an electromagnetic wave going this way you carry the momentum which is equal to u per unit area per unit time.

(Refer Slide Time: 12:20)



What does that momentum do? Suppose if I put a screen here which absorbs all this radiation falling if all the radiation falling is that absorbed that means all the momentum u per unit area per unit time which is coming gets absorbed in this. So, this thing should get momentum in this direction which is momentum per unit time per unit area areas nothing but force per unit area which is equal to pressure. So, an electromagnetic wave falling on a surface in which it gets absorbed applies the pressure equal to energy density.

If, the wave is reflected perfectly then the momentum coming in per unit area, per unit time would be u . Momentum going out with u again under pressure would be equal to two. So, I conclude this lecture by restating again that electromagnetic waves carry energy which we see in the form of light or heat. Which comes from the fact that electromagnetic waves carry momentum and this implies they can apply pressure on the surface on which they are formed.