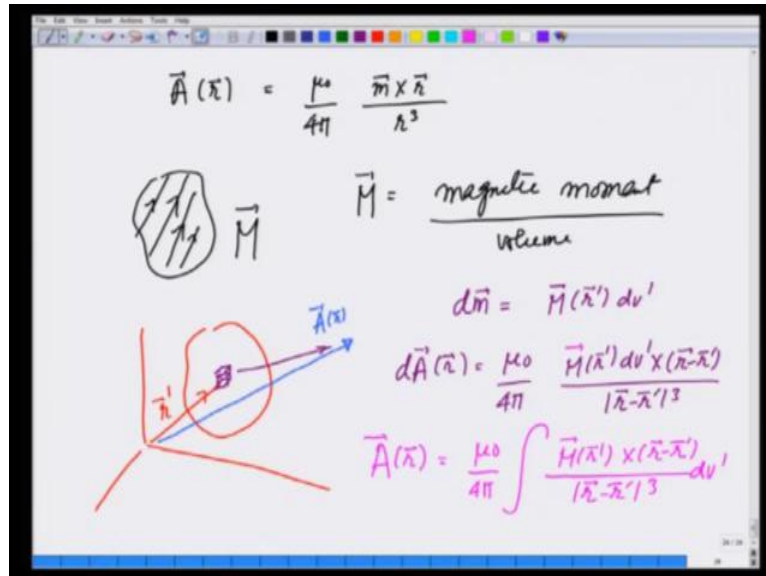


Introduction to Electromagnetism
Prof. Manoj K. Harbola
Department of Physics
Indian Institute of Technology, Kanpur

Lecture - 43
Magnetic Materials – II: Bound Current Densities

(Refer Slide Time: 00:12)



We have seen in the previous lecture that a point dipole gives me a vector potential, which is μ_0 over 4π \vec{m} cross \vec{r} over r^3 . Using this now we are going to show, if I have a material, magnetic material which has a permanent magnetization \vec{M} , something similar to the permanent polarization in the electrostatic case. So, that \vec{M} actually represents magnetic moment per unit volume.

Then this is equivalent to bound currents and surface bound currents. Again, similar to where a polarization in electrostatic case was equivalent to bound charges both the bulk bound charge and surface bound charge. Again to do this, we take help of vector potential, so when I have this system, let us make this is at \vec{r}' and suppose I am calculating vector potential at \vec{r} . Then I will take a small volume here near \vec{r}' and this will have a magnetic moment $d\vec{m}$, this is a small volume, which is equal to \vec{M} at \vec{r}' dV' .

And therefore, the vector potential $d\vec{A}$ differential vector potential at \vec{r} , due to this is going to be μ_0 over 4π $\vec{M}(\vec{r}') dV' \times (\vec{r} - \vec{r}')$, where $\vec{r} - \vec{r}'$

prime is this vector from that point to the point of interest, divided by r minus r prime cubed. And therefore, the vector potential A at r is nothing but mu 0 over 4 pi integral M r prime, this is a vector cross r minus r prime divided by r minus r prime cubed d v prime.

(Refer Slide Time: 02:47)

$$\vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{M}(\vec{r}') \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} dv'$$

$$= \frac{\mu_0}{4\pi} \int \vec{M}(\vec{r}') \times \vec{\nabla}' \frac{1}{|\vec{r} - \vec{r}'|} dv'$$

↓ $\vec{\nabla} \times (f\vec{A}) = \vec{\nabla} f \times \vec{A} + f \vec{\nabla} \times \vec{A}$

$$\vec{\nabla}' \times \left(\frac{1}{|\vec{r} - \vec{r}'|} \vec{M}(\vec{r}') \right) = \vec{\nabla}' \frac{1}{|\vec{r} - \vec{r}'|} \times \vec{M}(\vec{r}') + \frac{1}{|\vec{r} - \vec{r}'|} \vec{\nabla}' \times \vec{M}(\vec{r}')$$

$$\vec{\nabla}' \times \left(\frac{1}{|\vec{r} - \vec{r}'|} \vec{M}(\vec{r}') \right) = \vec{\nabla}' \frac{1}{|\vec{r} - \vec{r}'|} \times \vec{M}(\vec{r}') + \frac{1}{|\vec{r} - \vec{r}'|} \vec{\nabla}' \times \vec{M}(\vec{r}')$$

Now, we are going to do some vector manipulations here and show that this is equivalent M r prime cross gradient with respect to prime of 1 over r minus r prime d v prime. Now, I am going to use a vector identity, which is curl of f, where f is a scalar function times A is equal to gradient of f cross A plus f curl of A. Therefore, curl of 1 over r minus r prime which is like f M r prime and I should be taking curl with respect to the prime variable.

So, notice prime here should be equal to gradient of with respect to prime variable 1 over r minus r prime cross M r prime plus 1 over r minus r prime, curl of with respect to prime variable r prime. Let me write it again in different color, so that you see it clearly. Curl with respect to prime variable of 1 over r minus r prime M r prime is equal to gradient with respect to prime variable of r minus r prime cross M r prime plus 1 over r minus r prime curl of M r prime.

(Refer Slide Time: 04:55)

The whiteboard shows the following derivation:

$$\vec{\nabla}' \left(\frac{1}{|\vec{r}-\vec{r}'|} \right) \times \vec{M}(\vec{r}') = \vec{\nabla}' \times \left(\frac{1}{|\vec{r}-\vec{r}'|} \vec{M}(\vec{r}') \right) - \frac{1}{|\vec{r}-\vec{r}'|} \vec{\nabla}' \times \vec{M}(\vec{r}')$$

$$\int_{dv'} \vec{M}(\vec{r}') \times \vec{\nabla}' \left(\frac{1}{|\vec{r}-\vec{r}'|} \right) = \int_{dv'} \underbrace{-\vec{\nabla}' \times \left(\frac{\vec{M}(\vec{r}')}{|\vec{r}-\vec{r}'|} \right)}_{\text{from above}} + \int_{dv'} \frac{\vec{\nabla}' \times \vec{M}(\vec{r}')}{|\vec{r}-\vec{r}'|}$$

$$\int (\vec{\nabla} \times \vec{v}) dV = \int (\vec{n} \times \vec{v}) dS$$

Red arrows indicate the flow of information: one arrow points from the first term of the second equation to the first term of the third equation, and another points from the second term of the second equation to the second term of the third equation. A red circle with arrows is drawn around the first term of the third equation.

Let us invert this sense put 1 over r minus r prime with respect to prime cross M r prime is equal to curl of 1 over r minus r prime M r prime minus 1 over r minus r prime curl prime M r prime. And therefore, M r prime cross del prime of 1 over r minus r prime, I have just switched m n gradient. So, it will become a minus sign becomes minus curl of M r prime over r minus r prime plus curl of M r prime over r minus r prime.

And therefore, the integral of this over d v prime and this integral d v prime, integral d v prime, this again I am going to use an identity here, which says that integral of curl of a vector quantity over a volume is equal to n cross that vector quantity over the surface. So, what this means is that, if I have a closed surface enclosing a volume, this volume integral is equal to n cross v, where n is coming out of the surface over this surface, we can use this now here in the first integral.

(Refer Slide Time: 07:18)

$$\begin{aligned}
 & \int \vec{M}(\vec{r}') \times \vec{\nabla}' \frac{1}{|\vec{r} - \vec{r}'|} dV' \\
 &= - \int \vec{\nabla}' \times \left(\frac{\vec{M}(\vec{r}')}{|\vec{r} - \vec{r}'|} \right) dV' \\
 & \quad + \int \frac{\vec{\nabla}' \times \vec{M}(\vec{r}')}{|\vec{r} - \vec{r}'|} dV' \\
 &= \int - \frac{\hat{n}' \times \vec{M}(\vec{r}')}{|\vec{r} - \vec{r}'|} ds' + \int \frac{\vec{\nabla}' \times \vec{M}(\vec{r}')}{|\vec{r} - \vec{r}'|} dV'
 \end{aligned}$$

And therefore, I write $\vec{M}(\vec{r}') \times \vec{\nabla}' \frac{1}{|\vec{r} - \vec{r}'|}$ integral dV' is equal to which is minus gradient prime cross $\vec{M}(\vec{r}') / |\vec{r} - \vec{r}'|$ dV' plus integral curl of $\vec{M}(\vec{r}') / |\vec{r} - \vec{r}'|$ dV' . As equal to minus $\hat{n}' \times \vec{M}(\vec{r}') / |\vec{r} - \vec{r}'|$ ds' integral plus integral prime over $\vec{M}(\vec{r}') / |\vec{r} - \vec{r}'|$ dV' . So, what it means is, if I have a volume, where is there some magnetic moment \vec{M} , then the surface $\hat{n}' \times \vec{M}$ with the minus sign gives me the surface current this term. And this is the like the bulk current, why, because I also know that \vec{A} can be written $\vec{A} = \vec{A}_s + \vec{A}_b$.

(Refer Slide Time: 08:59)

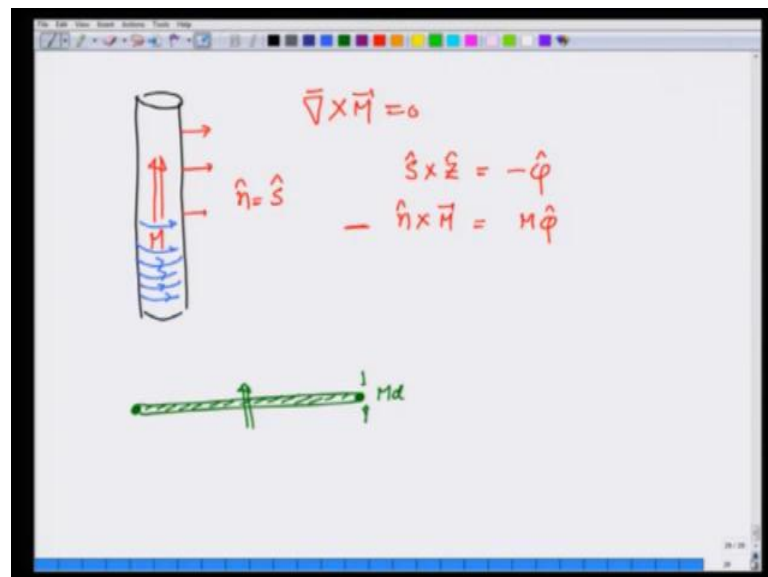
$$\begin{aligned}
 \vec{A}(\vec{r}) &= \int \frac{-\hat{n}' \times \vec{M}(\vec{r}')}{|\vec{r} - \vec{r}'|} ds' \\
 & \quad + \int \frac{\vec{\nabla}' \times \vec{M}(\vec{r}')}{|\vec{r} - \vec{r}'|} dV' \\
 &= \int \frac{\vec{K}(\vec{r}')}{|\vec{r} - \vec{r}'|} ds' + \int \frac{\vec{J}(\vec{r}')}{|\vec{r} - \vec{r}'|} dV'
 \end{aligned}$$

$\vec{K}_{\text{bound}} = -\hat{n}' \times \vec{M}$ $\vec{J}_{\text{bound}} = \vec{\nabla}' \times \vec{M}$

So, let us make it again here is a volume on which I have some M , then I can write A_r as equal to integral minus n prime cross M r prime over r minus r prime d s prime plus integral curl of M r prime over r minus r prime d v prime. And this should be equivalent to a surface current K at r prime over r minus r prime d s prime plus a bulk current J r prime over r minus r prime d v prime.

So, we get from this magnetic moment that it is equivalent to a surface current K , which is equal to minus n cross M and a bulk current J , which is curl of M . And since, these are not free currents, we call them bound currents, just like we had bound charges earlier.

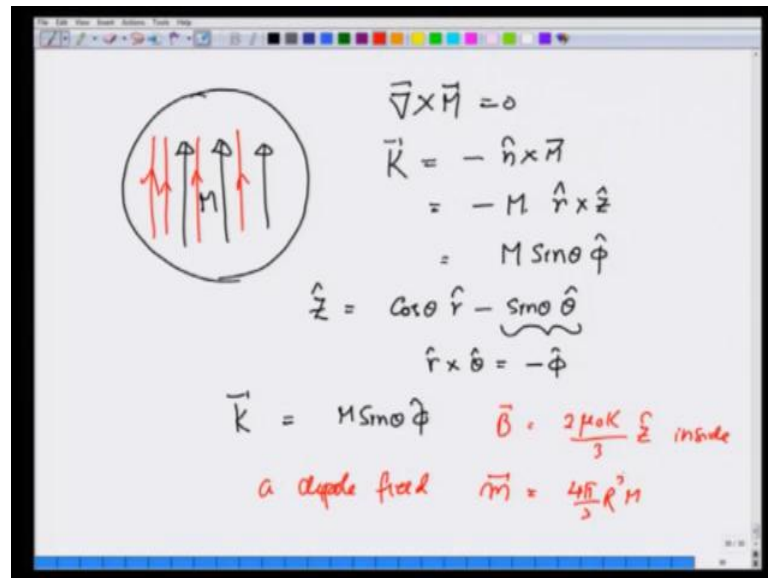
(Refer Slide Time: 10:19)



Let us take a couple of examples, if I have a long cylinder with magnetic moment in z direction, uniform magnetic moment in z direction, then curl of M is 0, n is coming out of the surface. So, n is in the same direction as S in the cylindrical coordinates and S cross z gives you minus ϕ , and therefore n cross M with a minus sign is nothing but $M \phi$.

So, this long cylindrical magnet having uniform magnetization is equivalent to having surface current, which is equal to M flowing in ϕ direction and no bulk current. So, this becomes like a solenoid, the field will be constant inside. So, it will be like a solenoid field. On the other hand, if I have a very flat thin wafer like or thin disc like thing with M in z direction, this will be like a ring of current, where the total current will be given by M times this width d , that will be the current going through.

(Refer Slide Time: 11:56)



This third example, I can take is that of a sphere with constant magnetization, then again curl of \vec{M} will be 0 and the surface current \vec{K} will be minus $\hat{n} \times \vec{M}$, which in this case will be minus $M \hat{r} \times \hat{z}$ unit vector. And you can check this is going to give me $M \sin \theta \hat{\phi}$; that is very easy to see, because \hat{z} is nothing but cosine of θ \hat{r} minus sin of θ $\hat{\theta}$ $\hat{r} \times \hat{\theta}$ gives me minus $\hat{\phi}$.

And therefore, \vec{K} becomes $M \sin \theta \hat{\phi}$, but this we have already solved, this we have solved in a previous lecture. And therefore, this is going to give me a constant magnetic field \vec{B} , which is $\frac{2\mu_0 K}{3} \hat{z}$ inside and a dipole field corresponding to \vec{m} equals $\frac{4\pi}{3} R^3 M$ outside.