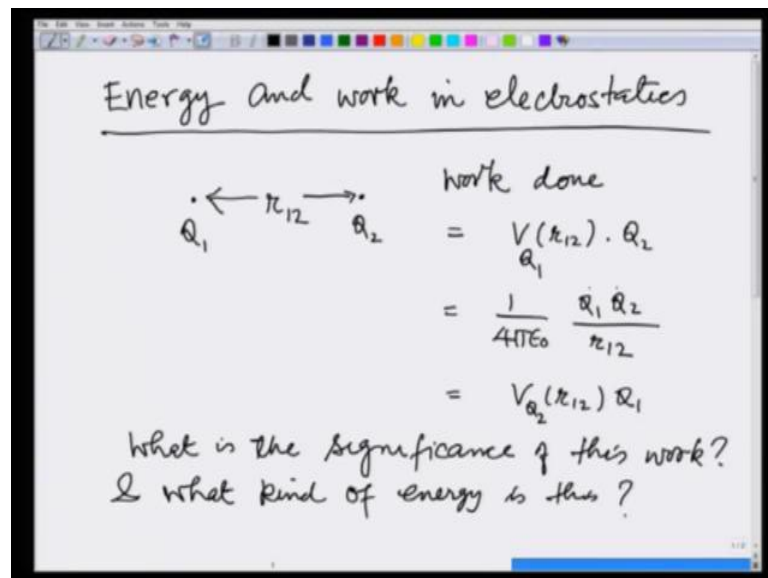


Introduction to Electromagnetism
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Lecture - 23
Energy of a Charge Distribution – I

We have looked at how charges or charge distributions give rise to electric field and potential, and how to calculate them. Now on, for 2 or 3 lectures, we are going to focus on what happens to the energy to system, when we assemble a set of charges or assemble a distribution of charge.

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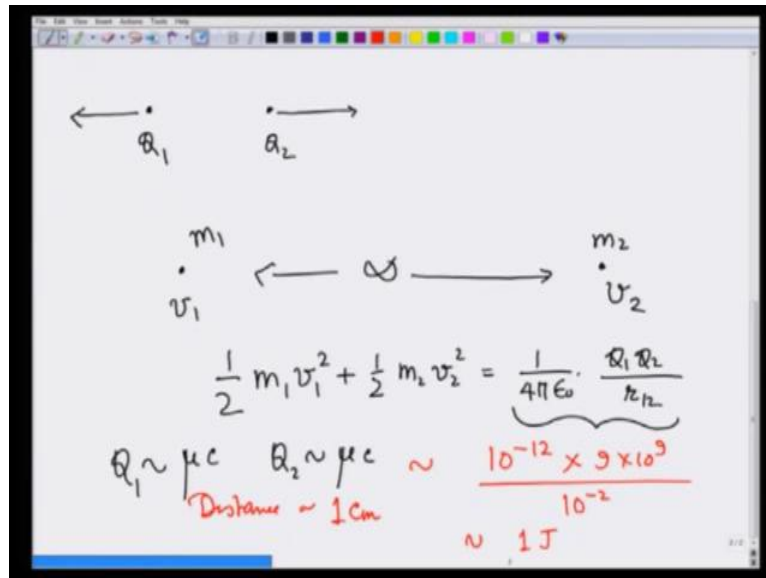


So, what we are going to work on is the energy and work in electrostatics. Let us start with the simplest example, if I have a point charge Q_1 let us say, and I bring another point charge Q_2 over distance, let us say r_{12} from it. Then by the definition of the potential the work done in doing so will be equal to the potential at r_{12} due to charge Q_1 times Q_2 and this is equal to $\frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{r_{12}}$, you see this is symmetry in Q_1 and Q_2 .

So, I could also have a sort of this work as the work of bringing charge Q_1 over distance r_{12} from Q_2 or the work in bringing Q_1 in the potential of Q_2 at a distance r_{12} from Q_2 . What we are interested in now is, what happens when I bring in assembly of charges many, many charges. And finally, we go to a charge distribution. So, that it is described

by charge density rho. Now, first what is the significance of this work and what kind of energy is this? Well, since this is energy due to position of charges, it is potential energy and this significance is that this is the amount of work that we have done in an assembly in this charges.

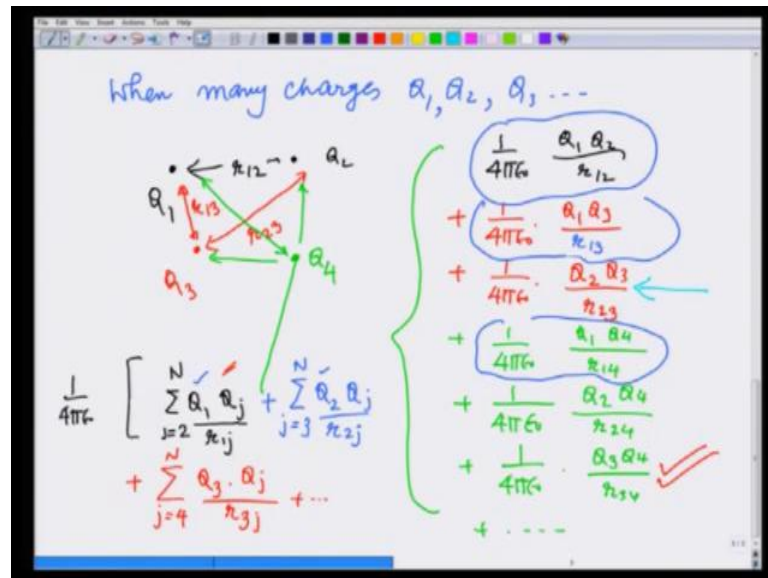
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And therefore, if I leave Q 1 and Q 2 by themselves, if they are same charges they repel from each other and therefore, it start rushing away. Finally; reaching infinite distance let us say this is infinite with speed v 2 and v 1. So, that their energy have let us say the mass of first charge is m 1 and mass of the second charge is m 2, then the total kinetic energy m 1 v 1 square plus 1 half m 2 v 2 square will be equal to this initial energy that we said is in the charge distribution and therefore, this will be equal to 1 over 4 pi Epsilon 0 Q 1 Q 2 over r 1 2.

What kind of magnitude of energy are we talking about? Let us say Q 1 is of the order of micro coulomb, Q 2 is of the order of micro coulomb, then this potential energy we are talking about and let me will be of the order of 10 raise to minus 12 times 9 times 10 raise to 9 divided by the distance is say 1 centimeter between them. So, this will be 10 raise to minus 2 of the order of 1 2, this is a kind of energy we are talking about. If this is micro coulombs, then the energy is of the order of 1 joule if they at distance of 1 centimeter. Imagine, what would happen if they were 1 coulomb, then the energy would be 10 raise to 9. We can already see that 1 coulomb is a lot and lot of charges.

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Next, what happens when many charges Q_1, Q_2, Q_3 and so on are there, let us assemble them one by one. So, I have charge Q_1 let us bring Q_2 to a distance r_{12} and I am doing, so I have done work which is $\frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{r_{12}}$. Let me now bring in charge Q_3 . Let us say this r distance, r_{13} from charge Q_1 and r_{23} from charge Q_2 , in bringing this charge I am doing work against a force provided by Q_1 and force due to Q_2 .

And therefore, the energy is going to be sum of the potential energy due to these two charges. So, I am going to get the additional terms which are $\frac{1}{4\pi\epsilon_0}$, potential energy of Q_3 due to charge Q_1 , and notice that this is all using principle of super position, we just adding it up. So, this is going to be $\frac{Q_1 Q_3}{r_{13}}$ plus $\frac{1}{4\pi\epsilon_0} \frac{Q_2 Q_3}{r_{23}}$.

Next let us bring in charge Q_4 , this will now be coming in the potential of these three charges. So, now, I will add three contributions, so this is going to be $\frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_4}{r_{14}}$ plus $\frac{1}{4\pi\epsilon_0} \frac{Q_2 Q_4}{r_{24}}$ plus $\frac{1}{4\pi\epsilon_0} \frac{Q_3 Q_4}{r_{34}}$, and then I can keep adding these terms. If you look at the pattern and all these, you will notice that I can write this as, if I have charge Q_1 if I bring next $n-1$ charges, then I have Q_j summation j equals 2 through N . Of course, there is $\frac{1}{4\pi\epsilon_0}$ outside.

So, these are the terms let me point them out here and presence of one I have this term, I have this term, so Q_2, Q_3, Q_4 and so on. Next, in bringing these charges I also done work against potential due to Q_2 . So, it will be $Q_2 Q_j$ I have already accounted for interaction between 1 and 2, so j will be equal to 3 to N . So, this r_{1j} , this is next term is divided by r_{2j} , next I will be counting the potential energy due to bringing all these charges in presence of charge 3, so I will have q_3 times q_j .

However, notice that I have already accounted for $Q_1 Q_3$ in this term and $Q_2 Q_3$ in the second term, and therefore I have j equals 4 through N over r_{3j} and so on. The last term the red one accounts for this term, this term was already accounted for the second term and so on.

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The image shows a whiteboard with the following handwritten text and equations:

N charges

$$W = \frac{1}{4\pi\epsilon_0} \sum_{i=1}^N \sum_{j>i}^N \frac{Q_i Q_j}{r_{ij}}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{1}{2} \sum_{i=1}^N \sum_{\substack{j=1 \\ i \neq j}}^N \frac{Q_i Q_j}{r_{ij}}$$

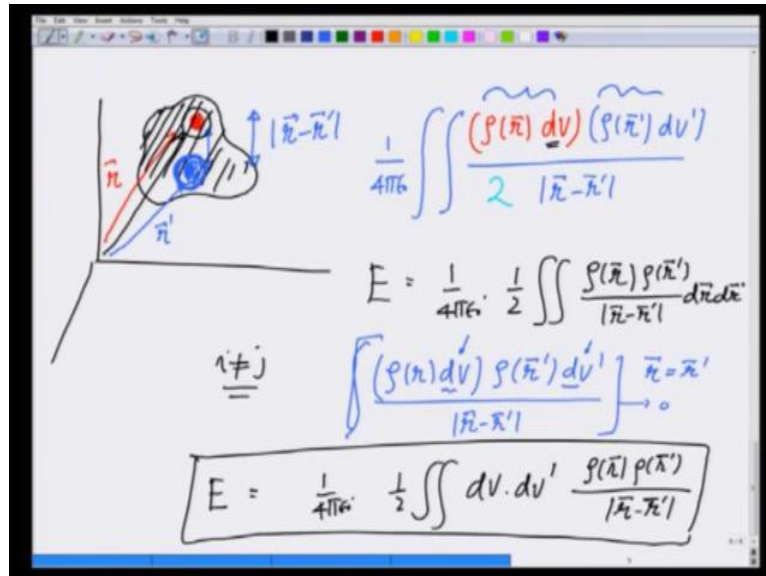
$$E = W = \frac{1}{4\pi\epsilon_0} \frac{1}{2} \sum_{i=1}^N \sum_{\substack{j=1 \\ i \neq j}}^N \frac{Q_i Q_j}{r_{ij}}$$

So, in general when I have an assembly of charges, when I have N charges I can write the potential energy is equal to 1 over $4\pi\epsilon_0$ summation i equal to 1 through N and presence of each other charge I am doing to work from j , which is greater than i up to N , $Q_i Q_j$ divided by r_{ij} . I would like you to expand this and see that this sets well you got derived earlier. I can also write this as 1 over $4\pi\epsilon_0$ $1/2$ of i equals 1 through N j equals 1 through N ; that means, I am now allowing j to vary not just from greater than i to N , but all the j 's.

So, I will be double counting each $Q_i Q_j$, so I divide this by a factor of half here. So, you keeping i not equal to j $Q_i Q_j$ over r_{ij} . So, that is the expression for the energy of

an assembly of charges. So, let us write this again I should be careful, let me write not this not as V, but as W that is work done. So, energy of this assembly of charges is equal to the work done bringing all these charges together, which is going to be $\frac{1}{4\pi\epsilon_0} \frac{1}{2} \sum_{i=1}^N \sum_{j=1, j \neq i}^N \frac{q_i q_j}{r_{ij}}$, this is the energy that is there in assembly of two charges.

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Let us now look at what happens when I go to a distribution in what; that means, is if I have a charge distribution. I can think of this charge distribution as if there is a charge somewhere in this volume shown by red they where charge somewhere in this volume shown by blue, and there is an interaction between let us write the position of charge and red to be let me show this by red r and the one by blue by r prime, then the distance between them is going to be modulus of r minus r prime.

Then the interaction energy between these two is going to be density at r times volume small volume d v that is the charge there, then density at r prime times of volume at r prime divided by the distance between them. And then I integrate over v and v prime and to 1 over 4 pi Epsilon 0 when we do this integration; that means, I am taking this charge multiplying and adding all these charges, and therefore I should also be dividing by a factor of two here, just like we did in counting a charges Q i and Q j.

So, if final expression for the energy then comes out to be $\frac{1}{4\pi\epsilon_0} \frac{1}{2} \int \int \frac{\rho(r)\rho(r')}{|r-r'|} dv dv'$. You may ask what

happened to $i \neq j$ term, vol we are taking these volumes here all this red volume here or the blue volume here infinitesimal small. And therefore, when I multiply $\rho(\mathbf{r}) dV$ and $\rho(\mathbf{r}') dV'$ with $r = r'$ this becomes infinitesimally very, very, very small.

In the sense that this divided by $r - r'$ close to 0 and therefore, there is no problem of self energy like $Q_i, i = j$ term in the previous expression that we derive earlier here where we put $i \neq j$. So, when you have a continuous distribution, so that these volumes of microscopic or macroscopically very small. So, that the product goes to 0 then there is no problems this remain 0 and there is no problem of self interaction. So, finally, we write E of distribution of charge which is $\frac{1}{4\pi\epsilon_0} \iint \rho(\mathbf{r}) \rho(\mathbf{r}') \frac{dV dV'}{|\mathbf{r} - \mathbf{r}'|}$. In the next lecture, we are going to solve some examples of energy of some distribution of charge.