

**Theoretical Mechanics**  
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**Module No # 02**  
**Lecture No # 08**  
**Conservation Laws**

Okay in this section what we are going to do is look at the conservation laws as the conservation laws which we obtain from Newtonian mechanics but of course in case of Lagrangian mechanics you can do little more.

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The slide contains handwritten text and equations on a light blue background. At the top, the title "Conservation Laws" is written in a light blue font. Below the title, the text "2n 2<sup>nd</sup> ordered DE." is written in blue ink. In the center, the equation  $\frac{\partial L}{\partial q_i} - \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) = 0$  is written in blue ink. Below the equation, the text "q<sub>i</sub>(t)" is written in blue ink. At the bottom, the text "First integrals" is written in blue ink, followed by the equation  $f(q, \dot{q}, t) = \text{constant}$  in blue ink.

Now what we have here in Lagrangian equation are  $2n$  differential second ordered differential equations and they given by  $\frac{\partial L}{\partial q_i} - \frac{d}{dt} = 0$ . So these are the  $2n$  differential equations now when you find the complete solution to these systems of differential equations you actually find each of these coordinate function are coordinate as a function of time  $t$  given some initial conditions. And it is in terms of  $t$  and  $2n$  arbitrary constants or constants which are given in terms of either initial conditions or the boundary conditions.

Now what will happen is this most of the time we will not be able find the complete solutions but even if we do not find the complete solutions it is still alright because if we can do at least one integral that helps us lot and in fact in the number of cases what we find is that we can actually

do one integral which is also called as first integrals and these first integrals typically can be put in a neat functional form like this.

So some function of  $y$  sorry  $q$ ,  $\dot{q}$  and  $t = \text{constant}$  so what will happen is in this equation there is only  $\dot{q}$  so you have first order differential equations in  $t$ . So once we get this is also important because this does give us the information about the nature of the actual system. Now the conservation laws actually are built at this level itself.

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Cyclic Co-ordinates : If  $L$  ind of  $q_j$  for some  $j$

$$\Rightarrow \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_j} \right) = 0 \Rightarrow \frac{\partial L}{\partial \dot{q}_j} = \text{Constant}$$

Ex  $L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) \Rightarrow \frac{\partial L}{\partial \dot{x}} = \text{Constant}$   
 $\Rightarrow m\dot{x} = \text{Constant}$

So to take an example consider the case of cyclic co-ordinates now in cyclic coordinates if Lagrangian is independent of one of the  $q$  variable say independent of  $q_j$  for some  $j$  okay for some specific  $j$ . If it is independent of  $q_j$  then  $\frac{\partial L}{\partial q_j}$  becomes 0 and you immediately get rule so  $\frac{d}{dt}$  of  $\frac{\partial L}{\partial \dot{q}_j} = 0$  which means  $\frac{\partial L}{\partial \dot{q}_j}$  actually is conserved quantity.

It remains constant throughout the motion of and if you look at this quantity here if  $q$ 's are (0) (04:23) Cartesian coordinates then you immediately identify this quantity to be the linear momentum. So for example for case of a free particle in two dimensions. So lagrangian is half  $m \dot{x}^2 + \dot{y}^2$  and this Lagrangian is independent of both  $x$  and  $y$  if I just look at  $x$  then this immediately tell us because Lagrangian is independent of  $x$  that  $\frac{\partial L}{\partial \dot{x}}$  must be constant and that in fact is nothing but  $m\dot{x} = \text{constant}$  and  $m\dot{x}$  remember is nothing but the  $x$  component of the linear momentum.

So in this problem of course also  $y$  is absent from Lagrangian my double sorry my dot that is linear momentum in  $y$  direction also we will be conserved in fact the net linear gets conserved.

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Cyclic Co-ordinates : If  $L$  ind of  $q_j$  for some  $j$

$$\Rightarrow \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_j} \right) = 0 \Rightarrow \frac{\partial L}{\partial \dot{q}_j} = \text{Constant} \leftarrow$$

Ex  $L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) \Rightarrow \frac{\partial L}{\partial \dot{x}} = \text{Constant}$   
 $\Rightarrow m\dot{x} = \text{constant}$

Ex  $L = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2)$

$\theta$  is not present  $\Rightarrow$

$$\frac{\partial L}{\partial \dot{\theta}} = m r^2 \dot{\theta} = \text{Constant}$$

Ang Mom.

Another example would be if I take the same situation but write the Lagrangian in terms of polar coordinate then now you see this particular Lagrangian is independent of theta but it is not independent of  $r$  okay. So what gets conserved here because this theta is absent theta is not present this immediately tells us that  $\frac{\partial L}{\partial \dot{\theta}}$  and this in this case will turn out to be  $m r^2 \dot{\theta}$  must be constant.

And this we immediate recognize as angler momentum. So in both these cases if a particular coordinate is missing from the Lagrangian or Lagrangian is independent of the coordinate the corresponding quantity  $\frac{\partial L}{\partial \dot{q}_j}$  remains constant and by analogy with these two examples if I am taking about Cartesian coordinate then the corresponding quantity  $\frac{\partial L}{\partial \dot{x}}$  little component of linear momentum if it was theta then it is a angler momentum.

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Conjugate momentum or Generalized momentum

$$p_j = \frac{\partial L}{\partial \dot{q}_j}$$

So by analogy with these quantity  $\frac{\partial L}{\partial \dot{q}}$  is in fact named as conjugate momentum or generalize momentum. And remember this is always conjugate to a given coordinate so I will write it as  $p_j$  defined as  $\frac{\partial L}{\partial \dot{q}_j}$  and the conservation theorem now says that if  $L$  is independent of  $q_j$  then the corresponding conjugate momentum  $p_j$  is conserved. And also we have remember that  $p_j$  is not necessarily a momentum it does not necessarily dimension of linear momentum.

What we have shown here is that in one case it turns out to be linear momentum in an another case it turns out to be angler momentum in more complicated cases which I will do in the next week. For example if you have charge particles in electromagnetic field even then we can write down the Lagrangian for this system and in that case the conserved quantity for a particular coordinate is in fact is not just the mechanical momentum of the mast or massive charges but it actually is the combination of the mechanical momentum + the momentum carried by the electromagnetic field themselves.

See the total of these two is actually conserved so what we are talking about here is not necessarily the momentum of the system or even the momentum itself. But some quantity which we call as conjugate momentum or generalize momentum that is the quantity that is conserved in this case okay. Now in a same way I can also derive from this point on I can also actually derive the normally momentum conservation law which we get form the Newtonian mechanics.

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$N$  particles  $r_i \quad i=1, \dots, N$  ,  $q=(q_1, \dots, q_n)$   
 $q_j$  for some  $j$  :  
 for each  $i$   $\vec{r}_i(q_1, \dots, q_j + \epsilon, \dots, q_n, t) = \vec{r}_i(q_1, \dots, q_j, q_n, t) + \vec{n} \epsilon$  for small  $\epsilon$   
 $\dot{\vec{r}}_i(q_1, \dots, q_j + \epsilon, \dots, q_n, t) = \dot{\vec{r}}_i(q_1, q_2, \dots, q_j, \dots, t)$   
 $T = \sum \frac{1}{2} m_i \dot{\vec{r}}_i^2$  ind of  $q_j$   $\frac{\partial T}{\partial q_j} = 0$   
 $\frac{\partial T}{\partial q_j} = \sum_i m_i \dot{\vec{r}}_i \cdot \frac{\partial \dot{\vec{r}}_i}{\partial \dot{q}_j} = \sum m_i \dot{\vec{r}}_i \cdot \frac{\partial \vec{r}_i}{\partial q_j}$   
 $= \sum m_i \dot{\vec{r}}_i \cdot \vec{n} \epsilon$   
 $= \vec{P}_{tot} \cdot \vec{n} \epsilon$

Say take for example we have a mechanical system so there are  $n$  particles and the coordinate or position vectors of these  $n$  are given by  $n$  okay. And we are using generalize coordinate in this case first I will for simplicity use only unconstraint system so I have generalize coordinates which we denote by  $q$  but includes  $q$  small  $n$  these many coordinates okay. And what I will do is this see it is very common for example one of the examples even we were discussing the generalize coordinates it is very common to choose one of the coordinates to be center of mass coordinates.

So let me assume that in this there is one variable  $q_j$  okay for some specific  $j$  and the role of this  $q_j$  is this when you take the system and if I keep all the  $q$ 's fixed and only change  $q_j$  then the entire system is translated by exactly the same amount. So each one of these  $r_i$ 's will translate by exactly the same amount. So what I mean is this so  $r_i$  at so for each  $i$   $r_i$  of  $q_1$   $q_j +$  some small epsilon and  $q_n$   $t$  this is equal to  $r_i$  at  $q_1$  to  $q_n$  including  $q_j$   $t +$  some  $n$  vector times epsilon okay.

So all position vector shift in the direction of  $n$  and by the same amount which is magnitude of  $n$  vector into that small constant epsilon okay. So this is of course for small epsilon okay what happens to the velocities? Velocity of course will remain constant because if it take  $r_i$  dot here at  $q_j +$  epsilon and all the way up to  $t$  this also will be equal to  $r_i$  dot at  $q_1, q_2, q_j$  and all the vector  $t$ . And because  $n$  times epsilon is constant for all of the position vectors its derivative would be 0.

So the all the velocities are unaffected by this translation of the entire system so what must happen is this? If the velocities are unaffected what happens to  $t$ ?  $t$  which is kinetic energy which is sum over half  $m_i r_i^2$  this is independent of  $q_j$  okay or in other words I will write this as  $\frac{\partial T}{\partial q_j} = 0$ . And look at what happens to  $\frac{\partial T}{\partial \dot{q}_j}$  so  $\frac{\partial T}{\partial \dot{q}_j}$  this is nothing but sum over  $i$   $m_i \dot{r}_i$  vector dotted with  $\frac{\partial r_i}{\partial \dot{q}_j}$ .

And this we remember is in fact equal to half  $m_i \dot{r}_i$  dot into  $\frac{\partial r_i}{\partial \dot{q}_j}$  remember while we derive the Lagrangian equations we had this identity proven. And this of course is equal to half sum over  $m_i \dot{r}_i$  dot but what is  $\frac{\partial r_i}{\partial \dot{q}_j}$  **ohh** look at this one here that is nothing but dotted with sorry yup the factor half of course will not be there sorry delete the factor of half okay this is nothing but  $\vec{n}$  vector times epsilon.

And what is this? This is nothing but total momentum which sum over  $m_i \dot{r}_i$  dot vector into epsilon that is it okay and what happens to the potential term. We have already seen that the kinetic energy term will be independent of  $q_j$  the potential energy term.

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$$\frac{\partial V}{\partial q_j} = \sum_i \vec{\nabla}_i V \cdot \frac{\partial \vec{r}_i}{\partial q_j} = \sum_i (-F_i) \cdot \vec{n} \epsilon$$

$$= -\vec{F}_{tot} \cdot \vec{n} \epsilon$$

Lagrange Eq for  $q_j$

$$\frac{d}{dt} (\vec{p}_{tot} \cdot \vec{n}) = \vec{F} \cdot \vec{n} = 0$$

$$\Rightarrow \vec{p}_{tot} = \text{constant.}$$

So  $\frac{\partial V}{\partial q_j}$  we will of course at least for this case assume that the potentials are independent of velocities okay in which case this is in fact nothing but sum over  $i$  and you have gradient of  $V$  with respect to  $\vec{r}_i$  into  $\frac{\partial r_i}{\partial q_j}$  and this is nothing but sum over  $i$  this is  $-\vec{f}_i$  into and again this is nothing but  $\vec{n}$  vector times epsilon. So the Lagrange's equation are actually in this case for  $q_j$  will be nothing but  $\frac{d}{dt}$  of total momentum dotted with  $\vec{n} = \sum_i \vec{f}_i$ .

I will write this as  $-\dot{L}$  total dotted with n times epsilon of course i can immediately remove the and what happens when you have  $\frac{\partial L}{\partial q_j}$  so if the Lagrangian is independent of  $q_j$  then we automatically know that this would be equal to 0 and total momentum is a conserved quantity okay. So if Lagrangian is independent of  $q_j$  corresponding generalized momentum  $\frac{\partial L}{\partial \dot{q}_j}$  over  $q_j$  dot will be conserved quantity and in this case it just turns out that quantity is the component of the total momentum along  $\hat{n}$  cap vector and that is the quantity that could be conserved in this case.

So this is the general conservation law for linear momentum of the entire system. We can do exactly the same thing for the angular momentum of the entire system think in terms of you know there is a particular coordinate which causes the entire system to turn okay then we will be able to derive a similar law for the angular momentum. Now I will not do that here but I will ask you people to try that at home and see for yourself that you can actually derive same similar condition for the angular momentum case okay. Now I will do one more conservation theorem which we get from the Newton's equations and that is energy conservation.

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*Energy Conservation.*

$$L(q, \dot{q}, t) \text{ ind of } t = \frac{\partial L}{\partial t} = 0$$

$$\frac{dL}{dt} = \sum_i \left( \frac{\partial L}{\partial q_i} \dot{q}_i + \frac{\partial L}{\partial \dot{q}_i} \frac{d}{dt} \dot{q}_i \right) + \frac{\partial L}{\partial t}$$

$$= \sum_i \frac{d}{dt} \left( \dot{q}_i \frac{\partial L}{\partial \dot{q}_i} \right) + \frac{\partial L}{\partial t}$$

$$\Rightarrow \frac{d}{dt} \left( \underbrace{\dot{q} \frac{\partial L}{\partial \dot{q}} - L}_{h} \right) + \frac{\partial L}{\partial t} = 0$$

*h : Energy Function.*

$$\Rightarrow h = \text{Constant}$$

Okay remember you have a system whose Lagrangian is given by  $q$  and Lagrangian of course you know whenever you have a path  $q$  of  $t$  Lagrangian is of course is a function of time  $t$  and it is a implicit function of time  $t$  through variables  $q$  and  $\dot{q}$  and it also may be explicit function of

time  $t$  okay. Let me assume that given Lagrangian is independent of  $t$  you see we will have identity which is similar to Beltrami identity.

So let me write that here if I calculate  $dL/dt$  this is nothing but because there is an implicit dependence  $\frac{\partial L}{\partial q} \dot{q}$  sorry  $\frac{\partial L}{\partial \dot{q}} \ddot{q} + \frac{\partial L}{\partial q} \dot{q} + \frac{dL}{dt}$ . So of  $\dot{q}_i$  okay and  $+$  they will be a partial derivative with respect to at time  $t$ . Now this we know how to this can be rewritten as so  $d/dt$  remember this is similar to what we did in case of Beltrami identity. So the first factor here  $\frac{\partial L}{\partial \dot{q}}$  can be replaced by  $d/dt$  of  $\frac{\partial L}{\partial q}$  from the Lagrange's equation.

And this of course can be written as  $\dot{q} \frac{\partial L}{\partial \dot{q}}$  okay and then summed over  $i$   $+ \frac{dL}{dt}$  and put it together this will become  $\frac{d}{dt} (\dot{q} \frac{\partial L}{\partial \dot{q}} - L) = 0$ . And the quantity which goes here in the bracket this is usually denoted by  $h$  and is called energy function of the system. And remember if  $L$  is independent of time  $t$  that means there is no explicit dependence on time  $t$  that means partial derivative is 0 this immediately gives us that  $h$  is a constant or is a conserved quantity okay.

Now how does this turn out to be an energy conservation that is not immediately clear here but if I look at one example it probably will immediately become clear.

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$$\begin{aligned} \text{Ex } L &= \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) - V(x, y) \\ h &= \dot{x} \frac{\partial L}{\partial \dot{x}} + \dot{y} \frac{\partial L}{\partial \dot{y}} - L \\ &= \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) + V(x, y) = \text{total Energy} \end{aligned}$$

$$\text{If } L = L_0 + L_1 + L_2$$

$$L_2 = \sum_{i,j} c_{ij} \dot{q}_i \dot{q}_j \quad c_{ij}(q_1, \dots, q_n, t)$$

$$L_1 = \sum_i b_i \dot{q}_i$$



Again take an example suppose the Lagrangian of the system is of a single particle which is travelling in two dimension and the influence of some potential field okay and the potential field here depends only on the sorry only on the coordinates and not on the velocities okay. In this case I can now calculate the h function and h function here would become so this would be  $x \dot{y} - y \dot{x}$  I will first write this explicitly  $+ y \dot{x} - x \dot{y}$ .

Remember there is summation over  $q_i$  so here it would be  $x \dot{y} - y \dot{x} + y \dot{x} - x \dot{y}$  into  $\frac{d}{dt} (x \dot{y} - y \dot{x})$  and  $-$  the Lagrangian and can you immediately see that this simplifies to  $\frac{1}{2} m (\dot{x}^2 + \dot{y}^2) - V(x, y)$  there you go this is nothing but the total energy. This is the total energy of this so at least in Cartesian coordinates if  $V$  is the only function of coordinates and not of velocities and time then you have a conservative system the total energy is conserved.

So this is the conservation theorem for total energy but remember normally  $h$  will not be same as the total energy of the system that can be different. Now if your Lagrangian is a form like this I will write this as  $L_0 + L_1 + L_2$  okay now the piece  $L_2$  is say only quadratic in velocities. So I will write this as  $\sum_{i,j} C_{ij} \dot{q}_i \dot{q}_j$  okay summation is over  $i$  and  $j$  okay and there will be double counting which will take care of when we write this  $C_{ij}$ . Remember  $C_{ij}$  themselves will be function of  $q_1$  to  $q_n$  and probably also time  $t$  and similarly  $L_1$  is a linear term so I will write this as  $\sum_i b_i \dot{q}_i$  okay it is only a linear term in  $\dot{q}_i$ .

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$$\begin{aligned} \text{If } L &= L_0 + L_1 + L_2 \\ L_2 &= \sum_{i,j} C_{ij} \dot{q}_i \dot{q}_j \quad C_{ij}(q_1, \dots, q_n, t) \\ L_1 &= \sum_i b_i \dot{q}_i \\ L_0 &= L_0(q, t) \\ h(q, \dot{q}, t) &= L_2 - L_0 \\ \vec{r}_i &\rightarrow q \quad \vec{r}_i = \vec{r}_i(q_1, \dots, q_n) \\ \Rightarrow &\text{transformations are ind of } t \\ L &= L_2 + L_0 \\ h &= L_2 - L_0 = T + V \end{aligned}$$

And  $L_0$  this is a term which does not depend on the generalized velocities okay it is just the function of  $q$  and  $t$ . In this case if I calculate  $h$  this is the exercise for you people from the assignment so that this will be equal to  $L_2 - L$  okay. So  $h$  function just turns to be this much in this special case what happens is this if I have suppose you have the position vector  $r_i$  and you make a transition to generalized coordinates  $q$ .

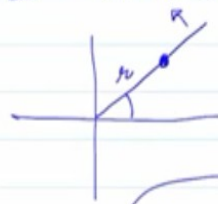
So when I write this transformation equation  $r_i$  as  $r_i$  function of  $q_1$  to  $q_n$  and in this transformation are independent of time  $t$  then what you get is a homogeneous so when you take  $r_i$  dot you get a homogeneous function of the generalized velocities and it is linear in generalized velocities. So when you take squares you only get  $L_2$  term so this remember if the transformations are independent of  $t$  then you get  $L_2$  term in Lagrangian and you will get  $L_0$  term okay.

The first linear term in generalized velocities will be absent now what would be  $L_0$  term?  $L_0$  term is your potential term in the potential term if we assume that the potential depends only on coordinates and not on generalized velocities or on time then your  $h$  function in fact will simply become  $L_2 - L_0$  and this would be same as kinetic energy + potential energy. So in this special case the  $h$  function looks like the total energy or = total energy and in this case if your Lagrangian is not explicitly dependent on time the total energy is concerned.

So that is how we get the energy conservation rule remember  $h$  energy function even though it is called as energy function it is not equal to total energy and it may happen that  $h$  is a conserved quantity but our total energy may not be conserved quantity. I will give you one example of this or even it may happen that  $h$  is actually = total energy but the total energy is not conserved. Remember this  $h$  function so Lagrangian even though there is a unique Lagrangian the  $h$  function actually depends on the coordinates so which coordinates you choose will keep on changing your  $h$  function. Now also the total energy function so here is an example which illustrates this.

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Ex Bead on a rotating wire



$$\theta = \omega t$$

$$L = \frac{1}{2} m (\dot{r}^2 + r^2 \omega^2) \leftarrow \text{ind of } t$$

$$E = \frac{1}{2} m (\dot{r}^2 + r^2 \omega^2)$$

$$h = \dot{r} \frac{\partial L}{\partial \dot{r}} - L = \frac{1}{2} m (\dot{r}^2 - r^2 \omega^2) \leftarrow \text{conserved quantity}$$

N is doing work  
at  $t=0$   $r=r_0$  and  $\dot{r}=0$

$$h = -\frac{1}{2} m r_0^2 \omega^2$$

$$\Rightarrow (m r^2 \omega^2 - m r_0^2 \omega^2)$$

This is a bead on rotating wire okay so you have a wire and there is a bead on the wire and this wire is in fact rotating about origin with the constant angular speed omega okay. So here because there are two coordinates to begin with but there is one constraint and because of this constraints all I really need is one generalize coordinate which is r. The other coordinate theta is in fact just equal to omega times t okay.

Now let us write down the Lagrangian for this will be half m r dot square + r square theta dot square but that we already know is just omega square okay. This is the Lagrangian for the system and what is the total energy of the particle? So total energy of the particle here because there are other than constraint forces they no other forces acting on the particles. So total energy is simply half m r dot square + r square omega square okay and what is the h function?

Let us calculate h function so h function will be energy function will be equal to r dot del L over del r dot - L and this is half m r dot square - r square omega square. See here the h function is not same as energy function now let me ask you a question is this a conservative system? That means is the energy consort quantity here? Surely not see if I had a free rotating rod like a merry go round which is not driven.

And if you sit on it and if you move outwards what would normally happen is the angular speed decreases but in this case even if the particle moves away from the origin the external system constraint system is actually driving the rod or the wire with the same constant angular speed

which means the external forces or the normal reaction here is doing work. And that work is actually reflected in the energy  $e$  so what remains conserved here?

Remember Lagrangian again independent of  $t$  that means this is conserved quantity what does this mean? How do I interpret this? In fact we can immediately see that if I hold the particle at  $t = 0$  at some initial instant and let go that means of the initial instant  $\dot{r}$  is 0 and your  $h$  function is in fact is nothing but so at  $t = 0$   $r = r_{\text{naught}}$  and the speed the radial speed of the particle is 0. See we are letting go at one instant so this is in fact  $-\frac{1}{2} m r_{\text{naught}}^2 \omega^2$  and what it says is this.

If the particle moves out to some other location then its new energy will be same as this quantity + the amount of work done by the constraint forces okay that is why the difference between  $h$  and  $e$  is exactly equal to  $m r_{\text{naught}}^2 \omega^2$  okay. So the difference the gain in the energy could  $m r_{\text{naught}}^2 \omega^2$  this would be total amount of work done by the constraint forces.

So in this example what we see is at the total energy is not conserved but as long as the Lagrangian is independent of time  $t$  there is some conserved quantity which is called as energy function and that energy function is given by this in this case. Okay now these ideas in fact are also related to each one of this conservation principle is related to the symmetries of the Lagrangian and this going to be focus of my next session.

So the conservation laws are intimately connected to symmetry of the Lagrangian and this is going to be focus of this small section today. If you see that if you make a coordinate transformation this is a great facility given by Lagrangian equation for very simple reason that the general form of the Lagrangian equations it remains same what that means is this when I make a transformation the Lagrange's equation as we write them.

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$$q = (q_1, \dots, q_n) \rightarrow s = (s_1, \dots, s_n)$$

$$q_i = q_i(s_1, \dots, s_n) \quad \forall i$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0 \Rightarrow \leftarrow$$

$$\frac{d}{dt} \left( \frac{\partial L_s}{\partial \dot{s}_i} \right) - \frac{\partial L_s}{\partial s_i} = 0 \leftarrow$$

$$L_s(s, \dot{s}, t) = L(q(s, t), \dot{q}(s, \dot{s}, t), t)$$
  

Ex  $(x, y) \rightarrow (r, \theta) \quad L = \frac{1}{2} m(\dot{x}^2 + \dot{y}^2)$   
 $L_s = \frac{1}{2} m(\dot{r}^2 + r^2 \dot{\theta}^2)$

$(x, y) \quad m\ddot{x} = 0, m\ddot{y} = 0$   
 $(r, \theta) \quad m(\ddot{r} - r\dot{\theta}^2) = 0 \quad m(r\ddot{\theta} + 2\dot{r}\dot{\theta}) = 0$

So if we start from coordinates  $q$  which are given by  $q_1$  to  $q_n$  and make a transformation to  $s$  which is  $s_1$  to  $s_n$ . So the transformation here would be  $q_i$  is written in terms of  $s_1$  to  $s_n$  for all for each  $q_i$  okay under this transformation the form of the Lagrangian equation remains same that means  $d/dt$  of  $\partial L / \partial \dot{q}_i - \partial L / \partial q_i$  this is equal to 0 this is if it is Lagrange's equations for  $q$  coordinates then for  $s$  coordinates it is a same equation except that sorry  $\dot{s}_i - \partial L_s / \partial \dot{s}_i = 0$ .

The only thing is the Lagrangian is of course changes the Lagrangian will be given by the new Lagrangian which is function of  $s$ ,  $\dot{s}$  and  $t$  this must be obtained from the previous Lagrangian which was written in terms of  $q$  but remember  $q$  is a function of  $s$  and possibly also  $t$  then  $\dot{q}$  will be function of  $s$ ,  $\dot{s}$  and  $t$  and the explicit dependence on  $t$ . So this is how the new Lagrangian is calculated but the general form remains same.

This is not true for the coordinates though even though the general form is same the individual equation when you substitute coordinates will be different I mean you cannot simply you know take this equation for  $q_j$  and substitute  $s_j$  in that equation to get a new equation here what I mean by this is that if you start with Cartesian coordinate let me take simple example where you go from  $xy$  coordinates to  $r\theta$  coordinates.

And the Lagrangian is that of a free particle so you have half  $m\dot{x}^2$  and  $+ y\dot{y}^2$  so when you make a transformation to the new Hamiltonian sorry new Lagrangian  $L_s$  which is

function of  $r$   $\theta$   $\dot{r}$   $\dot{\theta}$  and so on this would come to half  $I \dot{\theta}^2 + r^2 \dot{\theta}^2$ . So the Lagrangian looks different what about equation of motion or they also looked very different.

Because in first case the equation simply turns out to be  $m\ddot{x} = 0$  and  $m\ddot{y} = 0$  but so these are the equations in  $x$  and  $y$  but in  $r$  and  $\theta$  we already know that  $m\ddot{r} - r\dot{\theta}^2$  must be equal to 0 and the second equation for the  $\theta$  becomes  $m r \ddot{\theta} + 2\dot{r}\dot{\theta}$  must be equal to 0. You see I cannot simply take the  $x$  equation  $m\ddot{x}$  substitute  $r$  in it and write down the equation for  $r$  that is not true.

So the even though the Lagrange's equation in the form are same the actual equation in terms of coordinates they will be of course different. Now when will they be same they will be same if the Lagrange's Lagrangian has the same form.

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<u>Ex</u>	$(x, y) \rightarrow (r, \theta)$	$L = \frac{1}{2} m(\dot{x}^2 + \dot{y}^2)$
		$L_S = \frac{1}{2} m(\dot{r}^2 + r^2 \dot{\theta}^2)$
$(x, y)$	$m\ddot{x} = 0, m\ddot{y} = 0$	
$(r, \theta)$	$m(\ddot{r} - r\dot{\theta}^2) = 0$	$m(r\ddot{\theta} + 2\dot{r}\dot{\theta}) = 0$
<u>Ex</u>	$(x, y) \rightarrow (x', y')$	
	$x' = \cos\alpha x - \sin\alpha y$	
	$y' = \sin\alpha x + \cos\alpha y$	
	$L(x, y, \dot{x}, \dot{y}, t) = \frac{1}{2} m(\dot{x}^2 + \dot{y}^2)$	
	$L_S(x', y', \dot{x}', \dot{y}', t) = \frac{1}{2} m(\dot{x}'^2 + \dot{y}'^2)$	$\curvearrowright$
	$\Rightarrow m\ddot{x}' = 0, m\ddot{y}' = 0$	

Now that also I will explain using an example so you start with  $xy$  coordinates and go to new coordinates which are  $x'$  and  $y'$  and these  $x'$  and  $y'$  are basically obtained by rotating  $x$   $y$  prime by an angle  $\alpha$ . So the connection would be  $x'$  will be equal to  $\cos\alpha x - \sin\alpha y$  and  $y'$  will be equal to  $\sin\alpha x + \cos\alpha y$  and the original Lagrangian here which is function of  $xy$ ,  $\dot{x}$   $\dot{y}$  and  $t$  this is half  $m \dot{x}^2 + \dot{y}^2$  and what is the new Lagrangian?

Great now the new Lagrangian also has exactly the same form which is equal to half  $m \dot{x}^2 + y \dot{y}^2$  and you see this is what is called as the invariance of the Lagrangian under coordinate transformation. Even if you go from here to here the Lagrangian form of the Lagrangian is same. And in that case of course because the form of general form of Lagrange's equation is same here I can find out the equation for  $x$  prime simply by taking equation for  $x$  and substitute by  $x$  prime and that immediately gives you this equation.

And similarly  $m \ddot{y}$  will be equal to 0 so these are the new equations obtain new Lagrangian so I will call this as  $L_s$  so this was the new Lagrangian. So this is the idea that we want to follow under which transformation the Lagrangian remains invariant.

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Ex  $(x, y) \rightarrow (x', y')$

$$\begin{aligned} x' &= \cos \alpha x - \sin \alpha y \\ y' &= \sin \alpha x + \cos \alpha y \end{aligned} \quad \left. \vphantom{\begin{aligned} x' &= \cos \alpha x - \sin \alpha y \\ y' &= \sin \alpha x + \cos \alpha y \end{aligned}} \right\} \begin{array}{l} \checkmark \text{ invariance transformation} \\ \text{for system.} \end{array}$$

$$L(x, y, \dot{x}, \dot{y}, t) = \frac{1}{2} m(\dot{x}^2 + \dot{y}^2)$$

$$L_s(x', y', \dot{x}', \dot{y}', t) = \frac{1}{2} m(\dot{x}'^2 + \dot{y}'^2)$$

$$\Rightarrow m \ddot{x}' = 0, \quad m \ddot{y}' = 0$$

$$q = (q_1, \dots, q_n) \rightarrow s = (s_1, \dots, s_n)$$

def  $s_i = s_i(t_1, \dots, t_n, t)$  - point transformations.

$$L_s(s, \dot{s}, t) = L(q(s, t), \dot{q}(s, \dot{s}, t), t)$$

A Lagrangian is invariant under a point transformation if  $L_s(s, \dot{s}, t) = L(q, \dot{q}, t)$

Invariance transformation.

So make a general case we have a system with  $n$  degrees of freedom with one set of variables which are  $q_1$  to  $q_n$  and we making transformation to another system another set of coordinate which is  $s_1$  to  $s_n$  and the transformations here is  $s_i$  are given in terms of  $q_1$  to  $q_n$  and possibly  $t$ . Now of course when you make this coordinate transformation all the other conditions of you know differentiability and the invertibility must hold here and will assume that.

These transformation are called as point transformation okay and then we get new Lagrangian  $L_s$  which is function of  $s$ ,  $\dot{s}$  and  $t$  so remember this is in some sense definition which must be equal to  $L$  at  $Q$  but now written as function of  $s$  and  $\dot{q}$  which is written as function of  $s$ ,  $\dot{s}$

and  $t$  and  $t$  so you have to be basically take this transformation find the inverse transformations and in the Lagrangian wherever there is a  $q$  you substitute the form in terms of  $s$ .

So this is what you get a new Lagrangian and then we will say the Lagrangian is invariant under a point transformation if the new Lagrangian has exactly same form as old Lagrangian remember this means you take the form of old Lagrangian wherever there is  $q$  is substitute by  $s$  where is this  $q$  dot is substitute by  $s$  dot okay. This is called as and then the point transformation is called as invariance transformation.

So when you have invariance transformation the Lagrangian remains form invariant where is the example so there it is. In the previous two examples this one here this transformation is not a invariance transformation because if you take a new Lagrangian function of  $r$  and  $\theta$  wherever there is a  $r$  substitute by  $x$  whereas there is a  $\theta$  substituted by  $y$  will not work okay. So this one is not an invariance transformation but this one is here this one is a invariance transformation for this particular system okay.

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Cyclic Co-ordinates:  $q_j$  Cyclic if  $L$  is ind of  $q_j$

$$s_i = q_i \quad i \neq j$$

$$s_j = q_j + \alpha$$

$$L_0(s_1, \dots, s_n, \dot{s}_1, \dots, \dot{s}_n, t) = L(s_1, \dots, s_n, \dot{s}_1, \dots, \dot{s}_n, t) \quad \forall$$

if  $q_j$  is cyclic then  $L$  is invariant under translation along  $q_j$ .  $\Rightarrow P_j = \text{Conserved}$ .

Now what is the connection with the conservation laws so there it so you get conservation laws I will first consider the simplest possible case of cyclic coordinates remember coordinate  $q_j$  is cyclic if  $L$  is independent of  $q_j$  and we already know that if this happens then quite obviously the corresponding generalized momentum or conjugate momentum is conserved. Now we will see



what is connection of this statement with the symmetry of the Lagrangian think of this the transformation where  $s_i = q_j$  if  $i$  is not equal to  $j$  and  $s_j$  is  $q_j + \sum \alpha$  okay.

And you can immediately see that the new Lagrangian here which is  $s_1$  to  $s_n$   $s_1$  dot to  $s_n$  dot  $t$  this is in fact exactly = Lagrangian the whole Lagrangian at  $s_1$  to  $s_n$ ,  $s_1$  dot to  $s_n$  dot and  $t$ . That is because each  $s_i$  is in fact same a  $q_i$  so that is what go into the Lagrangian and  $s_j$  is different from  $q_j$  but wait your Lagrangian does not even have the term  $q_j$  so it would not have  $s_j$ . So this in that case holds that means if  $q_j$  is cyclic then  $L$  is invariant under translation along  $q_j$ .

See what we have done is in fact keeping all other coordinate same we have just translated the system along the coordinate  $q_j$  okay. So under the translation transformation or the system here the Lagrangian here is invariant under the translation transformation for or along the direction of  $q_j$  or along the coordinate  $q_j$ . And this results in the conservation law that  $p_j$  is conserved will see exactly which quantity is conserved in this case okay.

The connection right now we have this segment  $q_j$  cyclic  $p_j$  is conserved  $q_j$  is cyclic Lagrangian is invariant now we want to put these two together and say if the Lagrangian is invariant under a particular transformation there is conservation law okay alright so let me do that here.

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$$\begin{aligned}
 &L(q, \dot{q}, t) \\
 &q_i \rightarrow q_i + \delta q_i \\
 &\dot{q}_i \rightarrow \dot{q}_i + \delta \dot{q}_i \\
 &\text{Lagrangian is invariant} \rightarrow \\
 &L(q, \dot{q}, t) = L(q_i + \delta q_i, \dot{q}_i + \delta \dot{q}_i, t) \quad \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) \\
 &= L(q_i, \dot{q}_i, t) + \sum_i \left( \frac{\partial L}{\partial q_i} \delta q_i + \frac{\partial L}{\partial \dot{q}_i} \delta \dot{q}_i \right) \\
 &\Rightarrow \frac{d}{dt} \left( \underbrace{\sum_i \frac{\partial L}{\partial \dot{q}_i} \delta q_i}_{\text{conserved quantity}} \right) = 0
 \end{aligned}$$

So let me start by taking a Lagrangian of  $q$   $\dot{q}$  and  $t$  okay and we will say that we have this infinitesimal transformation which takes each  $q_i$  to  $q_i + \text{some } \delta q_i$  okay. With this

transformation what would happen to the Lagrangian if  $q_i$  is change like this then  $\dot{q}_i$  are also  $\dot{q}_i + \delta \dot{q}_i$  okay and let us put this together. So if under this transformation Lagrangian is invariant implies that the Lagrangian as a function of coordinate  $q$ ,  $\dot{q}$  and  $t$  also must be equal to Lagrangian substituting the new coordinates now  $q_i + \delta q_i$ ,  $\dot{q}_i + \delta \dot{q}_i$  and  $t$ .

And you can of course expand this so this would become Lagrangian at  $q_i$ ,  $\dot{q}_i$  and  $t + \sum$  over  $i$  and then you have  $\frac{\delta L}{\delta q_i} \delta q_i + \frac{\partial L}{\partial \dot{q}_i} \delta \dot{q}_i$  okay. And then of course we can immediately say that you remember the trick that we did in case of Beltrami identity this is nothing but replace  $\frac{\delta L}{\delta q_i}$  here by  $\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i}$  and then this just looks like a product rule for differentiation and I will put this.

So Lagrangian of course will cancel on both the sides and what you get is  $\frac{d}{dt} \sum \frac{\partial L}{\partial \dot{q}_i} \delta q_i$  and this is a conserved quantity. So there you go for every infinitesimal transformation which more accurately are called as continuous transformation for infinitesimal transformation there is a conserved quantity will immediately look at the examples of this.

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$$\begin{aligned} \text{Ex. } L &= \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) + \frac{1}{2} m \omega^2 (x^2 + y^2) \quad \text{2D Harmonic Oscillator} \\ x' &= x - \alpha y \Rightarrow \delta x = -\alpha y \quad \alpha, \text{ infinitesimal} \\ y' &= \alpha x + y \Rightarrow \delta y = \alpha x \\ L &= \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) + \frac{1}{2} m \omega^2 (x^2 + y^2) + O(\alpha^2) \\ \frac{\partial L}{\partial x} \delta x + \frac{\partial L}{\partial y} \delta y &= \text{const.} \\ -\alpha m \dot{x} \dot{y} + \alpha m \dot{y} x &= \text{const} \\ \Rightarrow \alpha m (\underbrace{x \dot{y} - y \dot{x}}_{L_2}) &= \text{const} \end{aligned}$$

So let me take a example of 2D Harmonic oscillator here the Lagrangian is  $\dot{y}^2 + \frac{1}{2} m \omega^2 (x^2 + y^2)$  okay. And now I am going to do a transformation to new coordinates. So if I take  $x'$  which is remember this is same as the rotation of the system by an angle  $\alpha$  but now I am going to assume that the angle  $\alpha$  is infinitesimal. So if angle

alpha is infinitesimal then I can write  $x'$  as  $x - \alpha y$  and  $y'$  will be equal to  $\alpha x + y$  okay.

This immediately gives us  $\Delta x$  to be equal to  $-\alpha y$  and this gives us  $\Delta y = \alpha x$  and if is this particular Lagrangian invariant under such transformation yes if you substitute all this there then the new Lagrangian will be equal to  $\frac{1}{2} m \dot{x}'^2 + \frac{1}{2} m \dot{y}'^2 + \frac{1}{2} m \omega^2 (x'^2 + y'^2)$  and + the term which are quadratic in  $\alpha$  square. So all the linear terms in  $\alpha$  actually have vanished leaving the Lagrangian invariant at least infinitesimal invariant under this transformation.

And what is this conserved quantity then so conserved quantity will be  $\frac{\partial L}{\partial \dot{x}}$  into  $\Delta \dot{x} + \frac{\partial L}{\partial \dot{y}}$  into  $\Delta \dot{y}$  and this is the conserved quantity and we can immediately identify this. What is  $\frac{\partial L}{\partial \dot{x}}$  that is nothing but  $m \dot{x}$  and what is  $\Delta \dot{x}$ ?  $\Delta \dot{x}$  is nothing but  $-\alpha \dot{y}$  and for the second term this is  $m \dot{y}$  into  $\alpha x$  and this must be constant.

And this immediately gives you  $\alpha m (\dot{x} y - y \dot{x})$  is constant and what does this immediately tell you that this quantity also must be constant of motion and this is in fact nothing but z component of the angular momentum. So in case very obvious that there is a symmetry circular symmetry or symmetrical symmetry in xy plane in which case the z component of the angular momentum is conserved. So this is how the conservation principles emerge from the invariance of Lagrangian under point transformation.

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$$\underline{\text{Ex}} \quad \mathcal{L} = \frac{1}{2} m (\dot{v}^2 + r^2 \dot{\theta}^2) - V(r)$$

$$\theta' = \theta + \alpha \quad \text{small } \alpha.$$

$$\delta\theta = \alpha$$

$$r' = r \Rightarrow \delta r = 0$$

$$\Rightarrow \delta r \cdot \frac{\partial \mathcal{L}}{\partial \dot{r}} + \delta\theta \cdot \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = \text{constant}$$

$$\Rightarrow 0 + \alpha \underbrace{(mr^2 \dot{\theta})}_{L_\theta} = \text{const}$$

And another example in this example I will take a simpler case or here the Lagrangian is half  $m \dot{r}^2 + r^2 \dot{\theta}^2$  and let me assume that the potential only depends on  $r$ . This situation is similar to the previous one okay situation is exactly same because in the previous case the potential I had explicitly given it as the potential of the Harmonic oscillator was proportional to  $x^2 + y^2$  and that is that means it is proportional to  $r^2$ .

Here what we are saying is that the potential here does not depend on the angle  $\theta$  which also makes  $\theta$  a cyclic coordinate here remember and the transformation that we are going to is  $\theta' = \theta + \alpha$  okay. Small  $\alpha$  this is like a translation along the  $\theta$  direction which eventually means a rotation which is exactly same as the previous example but here the transformation looks like when you look at it from the coordinate  $\theta$  point of view it is like a translation by small amount  $\alpha$ .

And in this case again the conserved quantity will be so here  $\delta\theta = \alpha$  what is  $\delta r$  we are not change in  $r$  okay so the new coordinate  $r'$  is exactly equal to  $r$  which gives us  $\delta r = 0$ . So this is the translation where all the coordinates are fixed except one which is being translated the conserved quantity then would be equal to  $\delta r \frac{\partial \mathcal{L}}{\partial \dot{r}} + \delta\theta \frac{\partial \mathcal{L}}{\partial \dot{\theta}}$  and in this case this will be equal to constant but this of course is  $0 + \delta\theta \frac{\partial \mathcal{L}}{\partial \dot{\theta}}$  and  $\delta\theta$  is nothing but  $\alpha$  and  $\frac{\partial \mathcal{L}}{\partial \dot{\theta}}$  is nothing but  $mr^2 \dot{\theta} = \text{constant}$ .

This is again the z component of angular momentum so there is an intimate connection between the invariance of the Lagrangian or symmetry of Lagrangian to the conservation principles. So every continuous invariance transformation there is a conservation law and the conserved quantity is also given by.

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For every infinitesimal transformation  $\delta q$ , there is a conservation law:

$$\left( \sum_i \delta q_i \cdot \frac{\partial L}{\partial q_i} \right) = \text{const.}$$

So for every infinitesimal transformation given by  $\delta q$  okay there is so conservation law which states that  $\sum_i \delta q_i \frac{\partial L}{\partial q_i}$  is a conserved quantity. This theorem is called as Noether's theorem this is by a mathematician Emmy Noether's.

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Noether's Theorem:

$$L' = L(q, \dot{q}, t) + \frac{d\Lambda}{dt}(q, t)$$

So it is in honor of her work this is called as Neother's theorem right. Now only one last point before i stop here remember the invariance of Lagrangian the way we posed it here. In fact it is little bit restrictive if you notice and this is also problem in assignment is that if you take a Lagrangian you get equations of motion. Now if I have set of equation motion do I have a unique Lagrangian for it may be possible that different Lagrangian may give you the same set of equation.

And that actually happens we can immediately or easily prove that if you have a Lagrangian  $L$  you construct new Lagrangian  $L'$  prime by taking  $L$  which is function of  $q$ ,  $\dot{q}$  and  $t$  and act to it  $t$  of  $q$  and  $t$  both  $L$  and  $L'$  prime they in fact will give you exactly same set of equations in coordinates to. What that means? Is when I was definition my invariance earlier when I make a point transformation I get a new Lagrangian and as long as new Lagrangian is equal to old Lagrangian + a total derivative of sum function of  $q$  and  $t$  I still get a invariance okay.

So the new definition of invariance should also include this. So  $L'$  prime at  $q$ ,  $\dot{q}$  and  $t$  if it is equal to old  $L$  old Lagrangian at  $q$ ,  $\dot{q}$  and  $t$  + another total different derivative of sum function then also your Lagrangian is considered as invariant.