

**Advanced Condensed Matter Physics**  
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**Lecture – 11**  
**S – Matrix and free electron Green's Function**

So, in the context of Green's function at 0 temperatures we have introduced the 3 representations that we would be working in, namely the Schrodinger representation Heisenberg representation. And interaction representation and we say that we would be working in the interaction representation mode and in that context we have computed the time evaluation of the wave function psi. And have introduced another operator you which was defined and it was found that the time evaluation of this operator U and the wave functions psi are similar and then finally the solution was written in an integral form.

Let us carry this discussion forward and introduce another important quantity called as the S-Matrix.

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S- Matrix

$$\psi(t) = \underbrace{e^{iH_0 t/\hbar} e^{-iH t/\hbar}}_{U(t)} \psi(0) = U(t) \psi(0) \quad H = H_0 + H'$$

Define a S-matrix

$$\psi(t) = S(t, t') \psi(t') = S(t, t') U(t') \psi(0)$$

$$= U(t) \psi(0)$$

$$S(t, t') U(t') = U(t) \Rightarrow S(t, t') = U(t) U^\dagger(t')$$

Some properties of the S-matrix (1)  $S(t, t) = \mathbb{1} = U(t) U^\dagger(t)$

Proof:  $\left( e^{iH_0 t/\hbar} e^{-iH t/\hbar} \right) \left( e^{iH t/\hbar} e^{-iH_0 t/\hbar} \right) = \mathbb{1}$

So, let us see what S-Matrix is and how it simplifies our discussion? So, to remind ourselves the wave function in the interaction representation was written as psi of t which is equal to exponential i H naught t by H cross psi of sorry this is exponential i H t

minus  $iHt$  by  $H$  cross and  $\psi$  of  $0$  can refer to the discussion earlier and this was actually written as  $U$  of  $t$ . Hence this was written as  $U$  of  $t$  and  $\psi$  of  $0$ . Now so, this is how the wave function would evolve from the wave function at time  $t$  equal to  $0$ . So,  $U$  is to operate by this  $U$  of  $t$  operator, which is written as exponential  $iH$  naught  $t$  by  $H$  cross and exponential minus  $iHt$  by  $H$  cross, where  $H$  is the total Hamiltonian equal to  $H$  naught plus  $H$  prime and  $H$  naught is a non interacting part of the Hamiltonian and  $H$  prime is the interaction term.

So, now one can also write it as. So, we can define in  $S$  matrix like this and by this which connects the wave function at  $t$  and the wave function at  $t$  prime. So,  $\psi$  at  $t$  is connected to  $\psi$  at  $t$  prime by a matrix, which is given by  $S$   $t$   $t$  prime. So,  $S$   $t$   $t$  prime connects the wave function at  $t$  and  $t$  prime. And this can still be written as  $S$   $t$   $t$  prime  $U$   $t$  prime  $\psi$  of  $0$ .

Now this  $\psi$  of  $t$  as we know that can be written as  $U$  of  $t$   $\psi$  of  $0$  hence this  $S$   $t$   $t$  prime  $U$   $t$  prime equal to  $U$  of  $t$ . So,  $S$   $t$   $t$  prime  $U$   $t$  prime equal to  $U$   $t$  we can write the same thing as  $S$   $t$   $t$  prime, equal to  $U$   $t$   $U$  dagger  $t$  prime. So, this is the definition of  $S$  matrix which is constructed from the  $U$  operators, at time  $t$  and  $t$  prime it is a  $U$   $t$   $U$  dagger  $t$ . And now since we know the solution of  $U$  of  $t$  we can find out  $S$   $t$   $t$  prime see the solution of  $U$  of  $t$  is in terms of the interaction term  $H$  prime which we have shown earlier.

So, now let us some properties of the  $S$  matrix and what are the properties one of them is  $S$   $t$   $t$  which is equal to  $1$ , because this is equal to  $U$   $t$   $U$  dagger  $t$  which is equal to  $1$ . So, this  $1$  means it is a unit matrix. So, that is that can be easily shown let us show this, because  $U$  of  $t$  is like exponential  $iH$  naught  $t$  by  $H$  cross exponential minus  $iHt$  by  $H$  cross that is my  $U$  of  $t$  and  $U$  dagger  $t$  would be exponential  $iHt$  by  $H$  cross exponential minus  $iH$  naught  $t$  by  $H$  cross.

Remember that when you take a dagger you have to write the operators so, this is like taking a transpose and that has to be taken in the reverse order. So, we have written this as a so, it is a transpose and then it is taken as a conjugate and then it is written backward. So, this is off course equal to  $1$ , because these  $2$  will become equal to one then farther these  $2$  will become equal to  $1$ . So, that is  $1$  property.

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$$\begin{aligned}
 (2) \quad S^\dagger(t, t') &= U^\dagger(t') U(t) = S(t', t) \\
 (3) \quad S(t, t') S(t', t'') &= S(t, t'') \\
 \text{Proof: } \psi(t) &= S(t, t') \psi(t') = \underbrace{S(t, t') S(t', t'')}_{= S(t, t'')} \psi(t'') \\
 &= S(t, t'') \psi(t'') \\
 \text{EOM (Equation of Motion)} \\
 \frac{\partial}{\partial t} S(t, t') &= \frac{\partial}{\partial t} U(t) U^\dagger(t') = -i H'(t) S(t, t') \\
 S(t, t') &= T \exp \left[ -i \int_{t'}^t dt_1 H'(t_1) \right]
 \end{aligned}$$

Let us see the second property of this matrix. So, property number 2, which is equal to  $S^\dagger(t, t')$  equal to  $U^\dagger(t') U(t)$ , which is equal to  $S(t', t)$ . So, which means that  $S^\dagger(t, t')$  is same as  $S(t', t)$  and the third one is  $S(t, t') S(t', t'')$  is equal to  $S(t, t'')$ . So, this says that when we multiply these  $S$  matrices the resultant matrix will be the first one and the last the first time index and the last time index. And this can be you know extended to more than 2 products. And the reason that this is true is that let us give a short proof of this one as well.

So, your  $\psi(t)$  is equal to  $S(t, t') \psi(t')$  and this is equal to  $S(t, t') S(t', t'') \psi(t'')$  and  $\psi(t'')$  can be related to the  $\psi(t')$  by simply using  $S(t', t'') \psi(t'')$ . And if you compare the this one and this one it is easy to see that what is written in 3 follows from that.

So, these are the properties and now the equation of motion every time will use this word and write it as EOM. So, the equation of motion of  $\psi$  is  $\frac{\partial}{\partial t} S(t, t') \psi(t')$  equal to  $\frac{\partial}{\partial t} U(t) U^\dagger(t')$ , which is equal to  $-i H'(t) S(t, t')$ . Now you see that the equation of motion exactly takes the same form as  $U(t)$  or  $\psi(t)$ . So, the  $S$  matrix also has a solution which looks like it is a  $t$  exponential minus  $i$   $t'$  to  $t$   $H'(t)$ . The only difference is that the  $S$  matrix is function of 2 time variables  $t$  and  $t'$ , that is why this integral does not go from 0 to  $t$  rather it goes from  $t'$  to

$t$  and it is the same interaction term and it is exponential related and with the time ordering operator.

This is exactly the same as we are written for the  $U$  of  $t$  or the  $\psi$  of  $t$ . So, far we have introduced the 3 representations and have computed closed form for the wave functions  $\psi$  the operator  $U$  and the  $S$  matrix and now the whole problem.

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$\psi(0)$  is unknown  $\rightarrow$  genuine problem.  
 Eigensolutions of  $H_0$  are known.  $H_0$ : noninteracting  $H$ .  
 $\epsilon_0, \phi_0$  are known.  
Gellmann - Low Theorem  
 $|\psi(0)\rangle = S(0, -\infty)|\phi_0\rangle$   
 Also known,  
 $\psi(t) = S(t, 0)\psi(0)$ .  
 Operate by  $S(0, t) \Rightarrow \psi(0) = S(0, t)\psi(t)$

Makes sense because the interacting many body ground state must have evolved from the non-interacting ground state in the distant past.

So, far is that that our  $\psi$  of 0 is unknown and it is actually a genuine problem, because we need to know the  $\psi$  of 0 because everywhere it is we are getting that the  $\psi$  of  $t$  is generated from  $\psi$  of 0 if  $\psi$  of 0 is not known, then we cannot get  $\psi$  of  $t$  even if we get the  $U$  of  $t$  or  $S$   $t$   $t$  prime etcetera. However, there is one thing that is known is the Eigen solutions of  $H_0$  are known  $H_0$  to remind you is the non-interacting Hamiltonian. So, let us write that because it will help you. So, it is a non-interacting Hamiltonian and it is assumed that we can solve it exactly and we know the wave functions and the ground state energy.

So, by Eigen solutions what we mean is that  $\epsilon_0$  and  $\phi_0$  are known where  $\epsilon_0$  is the energy of the non-interacting system and  $\phi_0$  is the wave function for the non-interacting system. These are known's these are the known quantities. And now we have to generate  $\psi$  of 0 from these  $\epsilon_0$  and  $\phi_0$  in particular  $\phi_0$  because  $\psi$  of 0 is the  $t$  equal to 0 wave function; for the interacting system which is not known; however, the non-interacting ground state or the energy the wave function is known.

So, we need to get that now this is facilitated by a theorem called as the Gellmann low theorem. So, the Gellmann low theorem says that the interacting wave function at  $t = 0$ . Now when I say wave function at the back of my mind I have ground state wave function, because that is the lowest energy state and in condensed matter physics as we have said earlier also that, we are only interested in the low energy states and which is particularly the ground state wave function and the ground state energies. So, while  $\psi(0)$  can actually mean any state at  $t = 0$ , but in we actually mean that it is a ground state.

So, this is written in including the S matrix as  $S(0, -\infty)$  and  $\psi(0)$ . So, it is clear that  $\psi(0)$ , which is the wave function of the ground state wave function for the interacting system can be generated from the wave function of the non-interacting system, which is known by this S matrix, which are where the 2 time variables are 0 and minus infinity that is the theorem statement of the theorem. Why does the theorem makes the sense? It make sense because the interacting many body ground state must has evolved from the non-interacting ground state in the distant past.

So, also it is known that let me write down because this is important makes sense because the interacting many body ground state must have evolved from the non-interacting ground state in the distant past, which means at  $t = -\infty$  that is how this S matrices is written as  $S(0, -\infty)$ .

Also it is true that  $\psi(t) = S(t, 0)\psi(0)$  if we operate by  $S(0, t)$  and in that case we have a  $\psi(0)$ , which is equal to  $S(0, t)\psi(t)$ . So, my interacting many body ground state at  $t = 0$  is obtained from the from the many body ground state at a finite time  $t$  using this S Matrix.

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$$\begin{aligned}
 &\text{For } t = -\infty \\
 &|\psi(0)\rangle = S(0, -\infty) |\psi(-\infty)\rangle = S(0, -\infty) |\phi_0\rangle \\
 &|\psi(-\infty)\rangle \equiv |\phi_0\rangle \\
 &\text{At } t = +\infty \\
 &|\psi(\infty)\rangle = S(\infty, 0) \psi(0). \\
 &|\psi(\infty)\rangle \text{ and } |\phi_0\rangle \text{ may be same or differ by} \\
 &\text{a trivial phase factor } \chi. \\
 &\phi_0 e^{i\chi} = \psi(\infty) = S(\infty, 0) \psi(0) = S(\infty, -\infty) \phi_0. \\
 &\boxed{e^{i\chi} = \langle \phi_0 | S(\infty, -\infty) | \phi_0 \rangle}
 \end{aligned}$$

And why is this done is because for t equal to minus infinity my psi of 0 it is equal to S 0 minus infinity psi of minus infinity i just put t equal to minus infinity there, which is nothing, but it is equal to S of 0 minus infinity and phi of 0 you want to write it with the brackets. So, psi of minus infinity is identical with phi of 0 which means that at t equal to minus infinity the interaction is not even switched on. So, the interaction is switched on only after t equal to minus infinity.

And so, this is the non-interacting ground state now there is an additional property that we would be needing at this in order to define the Green's function. So, at t equal to plus infinity we have psi of infinity it is equal to S of infinity 0 and psi of 0. Now it can also be assumed that the interaction is switched off at t equal to plus infinity. So, the interaction is switched on after t equal to minus infinity and is switched off after t equal to plus infinity. And in which case the psi infinity. So, psi infinity and phi 0 may be same and even if they are not same they should only differ by a trivial phase factor or differ by a trivial phase factor say a chi.

So, we can write which is considering the case in which it differs is phi 0 exponential i chi it is equal to psi of infinity, which is equal to S of infinity 0 psi of 0 which is equal to S of infinity minus infinity and a phi 0. I am not writing the (Refer Time: 19:57) always you can if you wish you can write it.

So, this phase factor is written as just the expectation value of the infinity minus infinity operator between the non-interacting ground state. So, this is an interesting and important property or rather an equation that you would need while defining the Green's function. So, let us finally get to the Green's function at  $t$  equal to 0.

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Green's function at  $T=0$

At  $T=0$ , in Heisenberg representation,

$$G(\lambda, t-t') = -i \langle \Psi_0 | T [c_\lambda(t) c_\lambda^\dagger(t')] | \Psi_0 \rangle \quad (1)$$

$\lambda$ : a relevant quantum mechanical parameter.

For an electron gas,  $\lambda = (\vec{k}, \sigma)$

$|\Psi_0\rangle$  is the interacting many body ground state.

$$c_\lambda(t) = e^{iHt} c_\lambda(0) e^{-iHt} \quad (\hbar=1)$$

Heisenberg representation.

For  $t' > t$ , the Green's function is defined as,

$$G(\lambda, t-t') = i \langle \Psi_0 | c_\lambda^\dagger(t') c_\lambda(t) | \Psi_0 \rangle$$

And will write it in the Heisenberg representation. So, at  $t$  equal to 0 in Heisenberg representation there is a reason to write it initially in Heisenberg representation and then will switch over to the interaction representation.

So, the Green's function is defined as  $G$  and some parameter  $\lambda$  which is a good property of the problem or the system under consideration, this is  $t$  minus  $t'$  equal to minus  $i$  and between the interacting ground state and it is a time order  $c_\lambda(t) c_\lambda^\dagger(t')$ . So, that is the definition which you should keep in mind  $\lambda$  is any parameter it could be energy, it could be momentum, it could be some other property, which is like spin or something which is equal to minus  $i$  times the ground state expectation value of  $c_\lambda(t) c_\lambda^\dagger(t')$  with the time ordering involved there.

That is the definition and so,  $\lambda$  will quantum mechanical parameter. Now it could be for an electron gas for an electron gas  $\lambda$  could be set of 2 quantum numbers namely  $k$  and spins  $\sigma$  or for other problems it could be something else. And once again to remind ourselves that  $|\Psi_0\rangle$  is the interacting many body ground state and a

$c_\lambda$  and  $C_\lambda(t)$  and  $C_\lambda(t')$  are written in the Heisenberg representation as  $C_\lambda(t) = \exp(iHt) C_\lambda(0) \exp(-iHt)$  and this is in the Heisenberg representation, this we have discussed at length.

So, this is my operator. So, for  $t > t'$  the Green's function is defined as  $G_\lambda(t, t') = i \psi_0^\dagger C_\lambda(t') C_\lambda(t) \psi_0$ .

Since we have resorted to a particular case where  $t'$  is greater than  $t$  the time ordering says that we have to write down  $t'$  on the left, in order to do that I gained an additional negative sign, because of the anti-commutation relation of the Fermi operators, which absorbs this negative sign that written in equation one and also I now do not need the time ordering operator because I have sorted the particular case this just an example that we are wanted to site.

So, let us just look at things such as.

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$$\begin{aligned}
 |\psi_0\rangle &= S(0, -\infty) |\phi_0\rangle \\
 c_\lambda(t) &= e^{iH_0 t} c_\lambda(0) e^{-iH_0 t} \\
 c_\lambda(0) &= e^{-iH_0 t} c_\lambda(t) e^{iH_0 t} \\
 \text{These operators have to be converted to interaction representation,} \\
 c_\lambda(t) &= e^{iH_0 t} \underbrace{e^{-iH_0 t} c_\lambda(t) e^{iH_0 t}}_{c_\lambda(t)} e^{-iH t} \quad \left. \begin{array}{l} e^{iH_0 t} \\ c_\lambda(0) \\ e^{-iH t} \end{array} \right\} \\
 &= U^\dagger(t) c_\lambda(t) U(t) \quad U(t) = e^{iH_0 t} e^{-iH t} \\
 &= S(0, t) c_\lambda(t) S(t, 0)
 \end{aligned}$$

So, what is my  $\psi_0$ , once again it is equal to  $S(0, -\infty) \phi_0$  it is already written earlier. And now we have to go to the interaction representation, because this is what we had planned now in order to go to the interaction representation we write down the



operators in the interaction representation as this. From now on we would take  $H$  cross equal to 1 so, not put that in exponent the denominator of the exponent.

So, this is how the time evolution of the operators are written with the  $H$  naught and the minus  $H$  naught  $t$ . And so, if i take if i want to write the  $c$  lambda of 0 so, that then i can write it as  $i H 0 t$ ,  $C$  lambda of  $t$  exponential  $i H 0 t$ . So, what I did is that i left multiply a exponential  $i H 0 t$  and right multiply a exponential  $i H$  naught  $t$ .

And that gives  $C$  lambda at  $t$  equal to 0 in terms of  $C$  lambda at  $t$ . Now these operators' operators have to be converted to interaction representation. And in order to do that we can write down a  $C$  lambda  $t$ , which is exponential  $i H t$  exponential minus  $i H 0 t$   $c$  lambda  $t$ , exponential  $i H 0 t$ , exponential minus  $i H t$ . Remember the time evolution of the operators in the Heisenberg representation is given by exponential  $i H t$  and  $C$  lambda 0 exponential minus  $i H t$ . And now my  $C$  lambda 0 is same as this, which is from the step earlier. So, I will have to write a in a interaction representation my  $C$  lambda as this and this can be written as  $U$  dragger  $t$   $C$  lambda  $t$   $U$  of  $U$  of  $t$  and remember that write it neatly.

So, these  $U$  of  $t$   $S$  are the operators that we have written earlier where  $U$  of  $t$  is written as exponential  $i H 0 t$  exponential minus  $i H t$ . So, this is how the interaction that operators in interaction representations are written and this can be written as  $S$  of 0  $t$   $C$  lambda  $t$   $S$  of  $t$  is 0.

So, I introduce the  $S$  matrices which we have learned and written down the operators and this was precisely the reason that we have introduced the  $S$  matrices and  $U$  in matrices, and now we know the complete solution of the  $S$  matrices and the  $U$  matrices which would be used in order to define the Green's function.

So, let us write down. So, this will be slightly bigger expressions, but nothing to worry these are well established and they can how they can be used to deal with physical problems we shall see.

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$$\begin{aligned}
 G(\lambda, t-t') &= -i \Theta(t-t') \langle \phi_0 | S(-\infty, 0) S(0, t) c_\lambda(t) S(t, 0) S(0, t') c_\lambda^\dagger(t') \\
 &\quad S(t', 0) S(0, -\infty) | \phi_0 \rangle \\
 &\quad + i \Theta(t'-t) \langle \phi_0 | S(-\infty, 0) S(0, t') c_\lambda^\dagger(t') S(t', 0) S(0, t) c_\lambda(t) \\
 &\quad S(t, 0) S(0, -\infty) | \phi_0 \rangle \\
 \text{Extreme left end of the terms, can be replaced by,} \\
 \langle \phi_0 | S(-\infty, 0) &= e^{-i\lambda} \langle \phi_0 | S(\infty, -\infty) S(-\infty, 0) \\
 &= \frac{\langle \phi_0 | S(\infty, 0)}{\langle \phi_0 | S(\infty, -\infty) | \phi_0 \rangle}
 \end{aligned}$$

So,  $G(\lambda, t-t')$  is what we have learned that if  $t$  is greater than  $t'$  the theta function. So, this is the theta function which will take a value equal to 1 and if not the it will take a value equal to 0. Now I am introducing the non-interacting ground states and writing it as minus infinity  $S(0, t) C(\lambda) S(t, 0) S(0, t')$   $C(\lambda) S(t', 0)$ .

Now the entire thing has to be written for the right part, which is  $S(t', 0) S(0, -\infty)$  and a  $\psi(0)$  of this only term there is another term, which will have to write for  $t'$  greater than  $t$ , which will come by sign opposite to the first one and hence will have to write it as plus  $i \Theta(t' - t)$ . The now it will be  $\psi(0)$ , now everything that we have written on the other side will come here it will be  $S(0, -\infty) S(0, t')$  and  $C(\lambda) S(t', 0)$  and  $S(0, t)$  and  $C(\lambda) S(t, 0)$ . So, everything we have written on the right now will come in the left, but in reverse order and  $S(0, -\infty) S(0, t)$  and  $S(0, -\infty)$  and a  $\psi(0)$ .

So, the total Green's function it is big no doubt, but you do not have to worry because ultimately we are going to bring it a very simple form simple looking form and the utility to physical systems will be shown. So, these are the 2 terms which are taken for  $t > t'$   $t < t'$ . Now what we can do is that the extreme left and back head for the first term left end of the first term or rather it is both the terms of the terms can be replaced by  $\psi(0) S(-\infty, 0) S(0, -\infty) = e^{-i\lambda} \psi(0)$

infinity minus infinity by the property of the S matrix that we have written is this. And now exponential i chi as we have shown that this is equal to the this expectation value of S infinity minus infinity.

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The Green's function becomes,

$$G(\lambda, t-t') = \frac{-i}{\langle \phi_0 | S(\alpha, -\infty) | \phi_0 \rangle} \left[ \begin{aligned} & \Theta(t-t') \langle \phi_0 | S(\alpha, t) c_\lambda(t) S(t, t') c_\lambda^\dagger(t') S(t', -\infty) | \phi_0 \rangle \\ & - \Theta(t'-t) \langle \phi_0 | S(\alpha, t') c_\lambda^\dagger(t') S(t', t) c_\lambda(t) S(t, -\infty) | \phi_0 \rangle \end{aligned} \right]$$

First term can be simplified as,

$$\Theta(t-t') \langle \phi_0 | S(\alpha, t) c_\lambda(t) S(t, t') c_\lambda^\dagger(t') S(t', -\infty) | \phi_0 \rangle$$

$$= \Theta(t-t') \langle \phi_0 | T c_\lambda(t) c_\lambda^\dagger(t') S(\alpha, -\infty) | \phi_0 \rangle$$

$S(\alpha, -\infty)$  involves operators in 3 time intervals,  $(\alpha, t)$ ,  $(t, t')$ ,  $(t', -\infty)$ .

So, this can be written as a phi a S infinity 0 these 2 can be combined as S infinity 0 and then we can write it with the denominator replacing the exponential i chi, which is equal to sorry there is a exponential minus i chi with this, minus i divided by this exponential i chi term S infinity minus infinity phi 0 and there is this theta function with the t minus t prime and a phi 0 S infinity t C lambda t S t t prime C lambda dragger t prime and S t prime and also there is another term in the bracket with minus, theta function t prime minus t phi 0 S infinity t prime C lambda dragger t prime S t prime C lambda t t minus infinity phi 0 and the bracket.

So, these are the 2 terms and we have introduced the S matrices properly and so, the let us see the first term. So, this term is theta t minus t prime phi 0 S infinity t C lambda t S t t prime, C lambda t prime dragger S t prime minus infinity and a phi of 0 this can be written as theta t minus t prime phi 0 t C lambda t C lambda dragger t prime S infinity minus infinity phi 0, this is an important step have a look at this and this says that C S infinity and minus infinity involves operators in 3 time intervals. So, intervals and these are infinity to t t to t prime in this case it is t is greater than t prime and it is t prime to

minus infinity. So, the time ordering operator automatically will sort these time intervals, because if we put the time that is largest will to the left will put it to the left.

So, proper sequencing of the time will be done by the time ordering operator and in which case we can write it and the second term also can be written it in the similar manner. So, this term also can be written it for the similar manner and again it will be 3 intervals, but; however, we have t prime is greater than t. So, there will be t t prime instead of that will have t prime t and there will be a t minus infinity.

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$$G(\lambda, t-t') = \frac{-i \langle \phi_0 | T c_\lambda(t) c_\lambda^\dagger(t') S(\infty, -\infty) | \phi_0 \rangle}{\langle \phi_0 | T S(\infty, -\infty) | \phi_0 \rangle}$$

$$G(\lambda, t-t') = -i \langle \phi_0 | T c_\lambda(t) c_\lambda^\dagger(t') | \phi_0 \rangle$$

Final form of the  $T=0$  Green's function.

So, if we combine both my Green's function will be written as lambda t minus t prime minus i phi 0 T C lambda t C lambda dragger t prime and S infinity minus infinity and a phi 0, and divided by this phi 0 S infinity minus infinity phi 0. In fact it is not required to put the S infinity minus infinity, because the time ordering will take care of the time that are appearing in the time variables that are appearing in the problem.

And in which case we can drop the S infinity and minus infinity and we can also sort of write the wave function without the write the Green's function without the denominator as G lambda t minus t prime, which is equal to minus i phi 0 T C lambda t C lambda, it is just pure applications there is no comma there. So, it is C lambda t prime dragger a phi 0 there. So, this is the final form of the T equal to 0 Green's function.

So, it says that the operators  $C_{\lambda t}$  and  $C_{\lambda t'}$  are taken the expectation is taken the time ordered expectation is taken with respect to the non-interacting ground state, which are known. And one has to multiply it with a minus i factor and the physical consequence or the physical situation that would it would correspond to is that if we create a particle given by  $C_{\lambda t'}$  and later take out the particle from the assembly of this many particle system at  $T$  equal to  $t$ .

The particle are rather the system will have the information that what kind of interactions are gone into the inside the system and the Green's function will provided information about that. It is more important, now to relate physical quantities to this Green's function which is what we will do. We will now do one example of this Green's function and the most simple case that is the non-degenerate electron gas at which means it is a non-interacting system.

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Example  
Degenerate, non-interacting electron gas  
 Simple metal at  $T=0$ . All electronic states for  $\epsilon_k (= \frac{\hbar^2 k^2}{2m}) > \mu$  are empty and  $\epsilon_k < \mu$  are filled.  
 Define  $\xi_k = \epsilon_k - \mu$   
 $\langle \phi_0 | C_k^\dagger C_k | \phi_0 \rangle = \langle n_k \rangle = \theta(k_F - k) \stackrel{\beta \rightarrow \infty}{=} \frac{1}{e^{\beta \xi_k} + 1}$   
 $\langle \phi_0 | C_k C_k^\dagger | \phi_0 \rangle = 1 - \langle n_k \rangle = \theta(k - k_F)$

So, example so, it is a degenerate non interacting electron gas and let us take a simple metal at  $T$  equal to 0. So, all electronic states for  $\epsilon_k$ , which is equal to  $\frac{\hbar^2 k^2}{2m}$  and being the electronic mass for  $\epsilon_k$  greater than  $\mu$  are empty and for  $\epsilon_k$  less than  $\mu$  are filled.

This you must have studied in the first course of solid state physics, which is the model for a metal and we do not need to have  $t$  equal to 0 we will relax this condition. That

since we are talking about the chemical potential we are restricting ourselves to  $t$  equal to 0, when we rather we can talk about the Fermi energy when we talk about  $t$  equal to 0, but in principle this the illustration will carry over to  $t$  naught equal to 0.

Now since all the energies are measured from the Fermi level let us define a variable  $\epsilon_k$  which is equal to  $\epsilon_k - \mu$  and we have for a uniform spherical Fermi surface, we have  $\phi_0 = \sum_k c_k^\dagger c_k$ . Now your  $\lambda$  is equal to simply equal to  $k$  we do not need to put the spin in the less there is term let us the spin. So, this is equal to nothing, but equal to  $n_k$  and this is equal to  $k_f - k$ , which is limit  $\beta$  tending to infinity, which means  $t$  tending to 0 of this exponential  $\beta \epsilon_k + 1$  that is the Fermi distribution.

So, this is the definition that we have the ground state expectation non interacting ground state expectation value of the  $c_k^\dagger c_k$  is nothing, but expectation value of  $n_k$  which is nothing, but the theta functions are all  $k$  is greater than  $k_f$  that is  $k < k_f$  which is nothing, but the Fermi function.

And for the other term because we have 2 terms in the Green's function, because of the time ordering and this is equal to  $c_k^\dagger c_k$   $\phi_0$ , which is nothing, but equal to  $1 - n_k$  and that is equal to theta of  $k_f - k$ , if the states which are all empty at  $t$  equal to 0 or  $t$  tends to 0.

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So the Green's function (unperturbed) is

$$G_{\lambda}^0(t-t') = -i \theta(t-t') \langle \phi_0 | c_{\lambda}(t) c_{\lambda}^{\dagger}(t') | \phi_0 \rangle + i \theta(t'-t) \langle \phi_0 | c_{\lambda}^{\dagger}(t') c_{\lambda}(t) | \phi_0 \rangle$$

$$= -i \left[ \theta(t-t') \theta(\epsilon_k) - \theta(t'-t) \theta(-\epsilon_k) \right] e^{i\epsilon_k(t-t')}$$

Fourier transform of this,

$$G^0(k, \omega) = -i \left[ \theta(\epsilon_k) \int_0^{\infty} dt e^{it(\omega - \epsilon_k + i\eta)} - \theta(-\epsilon_k) \int_0^{\infty} dt e^{it(\omega - \epsilon_k - i\eta)} \right]$$

$$= \frac{\theta(\epsilon_k)}{\omega - \epsilon_k + i\eta} + \frac{\theta(-\epsilon_k)}{\omega - \epsilon_k - i\eta}$$

So, the so the Green's function, now this Green's function, because we are talking about a non-interacting problem will call it an unperturbed Green's function this word will be clarified more as we go along is  $G(\lambda, t - t')$ , it is equal to  $\theta(t - t')$  and a  $\phi_0 C(\lambda, t)$  this I am writing it in the form that we have introduced and  $\theta(t - t')$  there is a general form and the plus  $\theta(t - t')$  and  $\phi_0$  the sign change occurs, because the operators swap the positions. And this is equal to  $\theta(t - t')$  and a  $\theta(x_k - x_k')$  and exponential  $i x_k t - x_k' t'$ .

So, the Fourier transform of this of this is  $G$  now let us write it because we are talking about the unperturbed principle let us write it with 0. So,  $G_0(k, \omega)$  it is equal to  $\theta(x_k - x_k')$  and using the definition of the theta function for the time  $x$  equal to 0 to infinity and the  $\int dt \exp(i t, \omega - \epsilon_k + i \eta)$  and for the other one will write it with  $\theta(x_k - x_k')$  to infinity  $\int dt \exp(i t, \omega - \epsilon_k - i \eta)$  and this tells me that this becomes equal to this  $x_k \omega - x_k + i \eta$  and the plus theta minus  $x_k$  divided by this minus sign.

Let me write it a little more clearly and it is a  $\omega - \epsilon_k - i \eta$ .

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In a compact form,

$$G_0(k, \omega) = \frac{1}{\omega - \epsilon_k + i\eta_k}$$

where  $\eta_k = \eta \operatorname{sgn}(\epsilon_k - \mu) = \eta \operatorname{sgn}(\epsilon_k - \mu)$   
 $= +\eta$  when  $\epsilon_k > \mu$  (1st term)  
 $= -\eta$  when  $\epsilon_k < \mu$  (2nd term).

Introducing additional properties, such as, spin etc.

$$G_{\alpha\beta}^0(k, \omega) = \sum_{\alpha\beta} \left[ \frac{\theta(k - k_f)}{\omega - \epsilon_k + i\eta} + \frac{\theta(k_f - k)}{\omega - \epsilon_k - i\eta} \right]$$

So, this can be written in a compact form as a  $G_0(k, \omega)$  equal to 1 divided by  $\omega - \epsilon_k + i \eta_k$ , where  $\eta_k$  equal to some  $\eta \operatorname{sgn}(\epsilon_k - \mu)$  where this is equal to  $\eta \operatorname{sgn}(\epsilon_k - \mu)$

Sign of  $\epsilon_k - \mu$ . Which says that whether  $\epsilon_k$  is greater than  $\mu$  or  $\epsilon_k$  is less than  $\mu$   $\eta$  will take form, which is plus or minus where  $\eta$  is positive number.

So, will have a  $i\eta$  or  $-i\eta$  depending on whether we are talking about the particle or the hole will clarify this more as we go along, but this term the first term that is the term on the left is called as the electron propagator and whereas the term on the right would be called as the hole propagator, because this corresponds to filled states and this corresponds to empty states.

So, this is equal to plus  $\eta$  when  $\epsilon_k$  is greater than  $\mu$ . So, that is the first term and  $x$  equal to minus  $\eta$ , when  $\epsilon_k$  is less than  $\mu$  that is the second term. So, finally, the Green's function if we introduced spin so, introducing additional properties such as spin etcetera. One can write the  $G_{\alpha\beta}(k, \omega)$  will be equal to  $\delta_{\alpha\beta} \theta(k - k_f)$  divided by  $\omega - \epsilon_k + i\eta$ , and plus  $\theta(k_f - k)$  divided by  $\omega - \epsilon_k - i\eta$ . So, in fact, this correspond to the electron particle term the second term, sorry for messing that up earlier and the first term corresponds to the hole term which are all empty where  $k$  is greater than  $k_f$ .

So, at  $t = 0$  none of these states are occupied and all the states are occupied, which are called as the particle states for  $k > k_f$  and hole states for  $k < k_f$ .

So, this is the Green's function that we compute for a degenerate free electron gas and so, it gives the entire energy spectrum of the poles of the Green's function gives the energy spectrum for both the particles and the holes and they just differ by of course, the numerator and in the poles there will fine, whether the complex term actually pushes the pole up of the real line or it pushes down of the real line we will clarify all these things in our subsequent discussion.