

Introduction to LASER
Prof. M. R. Shenoy
Department of Physics
Indian Institute of Technology, Delhi

Lecture - 08
Part-II: Scheme of Light Amplification
Laser Rate Equations: 2-Level System

Welcome to this MOOC on lasers. So, today we will take up part 2 in which, we will discuss the Scheme of Light Amplification. In particular today's lecture we will discuss Laser Rate Equations. So, let me first recap what we have done in part 1.

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Recap: PART- I

- Interaction of Radiation with Matter
- Einstein Coefficients
- Atomic Lineshape Function
- - Rates of Emission and Absorption
- - Amplification by Stimulated Emission:

“*Population inversion* is the necessary condition for Amplification by Stimulated Emission”

→ Now, Q: How to achieve ‘*Population inversion*’ ?

Ans: By ‘*Pumping*’ → i.e. by exciting the atoms.

M R Shenoy2

In part 1, we discussed about the Interaction of Radiation with Matter. We saw the Einstein Coefficients, then the we introduced the Atomic Lineshape Function and discussed the Rates of Emission and Absorption. So, rates of emission and absorption and then the condition for

Amplification by Stimulated Emission. And we have seen that Population inversion is the necessary condition for amplification by stimulated emission.

Now, in this part therefore, the next question is how to achieve population inversion? We have seen that the population inversion is the necessary condition and therefore, how to achieve population inversion in atomic system. We have already seen, that at thermal equilibrium, there is no population inversion the number of atoms in the excited state is much much lesser than the number of atoms in the ground state and therefore, how to achieve population inversion.


The answer is by Pumping, that is by exciting the atoms by means of an external pump and that is called Pumping. So, we will discuss this in more detail in this part.

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In PART- II

- **SCHEME OF AMPLIFICATION**
- • Schemes of Pumping: 3-level and 4-level Systems
- • Rate Equations → Requirements to achieve *Population Inversion*
- • Rare-earth Doped Laser Amplifiers
- • Examples: Nd:YAG Amplifier and the EDFA → Amplifier Configurations

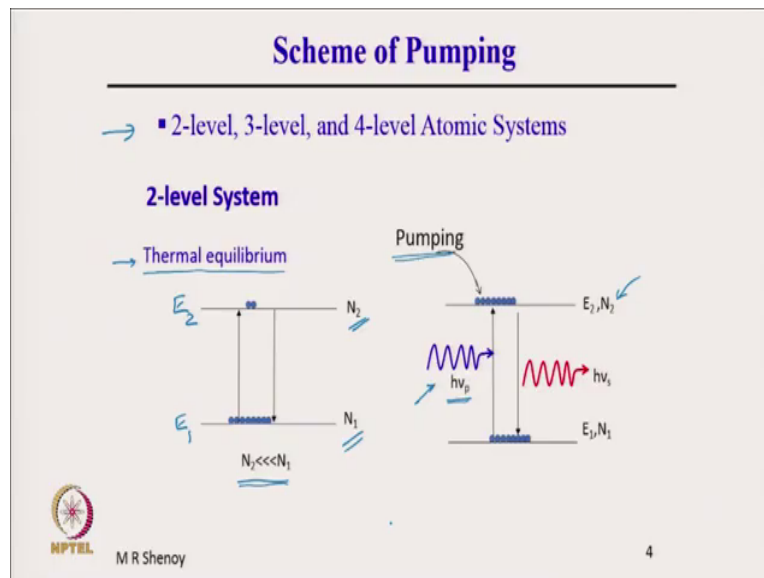
Er-doped Fiber Amplifier

 NPTEL M R Shenoy 3

So, in this part that is Scheme of Amplification, we will first see the Scheme of pumping, possible schemes of pumping in 3 level and 4 level atomic systems, we will see what are these 3 level and 4 level systems. And then we will write down the Rate Equations, which describe the rate of change of population of the levels. And they will tell us, they will give us the requirements to achieve Population Inversion, the conditions or the requirements of the atomic system, which can lead to Population Inversion.

And then we will discuss Rare earth Doped Laser Amplifiers with specific examples of Nd YAG Amplifier, Neodymium doped YAG Itanium Aluminum garnet amplifier is a widely used amplifier and laser Nd YAG laser, and the EDFA, Erbium doped fiber amplifier. So, Erbium is a rare earth atom. So, Erbium doped fiber amplifier most widely used amplifier in optical fiber communication systems amplifier. So, EDFA we will see the amplifier configurations and typical characteristics of these amplifiers.

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So, first the Scheme of Pumping and let us see first what are 2 level, 3 level and 4 level atomic systems. First 2 level system, what is a 2 level system? A 2 level system is an atomic system; we know that atoms are characterized by large number of discrete energy levels. So, there are not just 2 levels there are many many levels, but a 2 level system when we say we are discussing about a 2 level system, then we are referring to an atom, where two atomic energy levels one is the ground state and an excited state, which is interacting with the radiation.

And in a 2 level system therefore, we are considering the ground state which is here energy E_1 and an excited state usually the first excited state E_2 , in the presence of radiation. So, at thermal equilibrium, first at thermal equilibrium here there is no external radiation at thermal equilibrium atoms can get excited to a higher state.

And they would also come down and excite to the ground state. So, there is a dynamic equilibrium atoms keep going up and coming down, and at any instant or at any time there is a dynamic equilibrium with certain number of atoms N_2 in the upper state and N_1 in the ground state, and we have already seen that N_2 is much much less than N_1 .

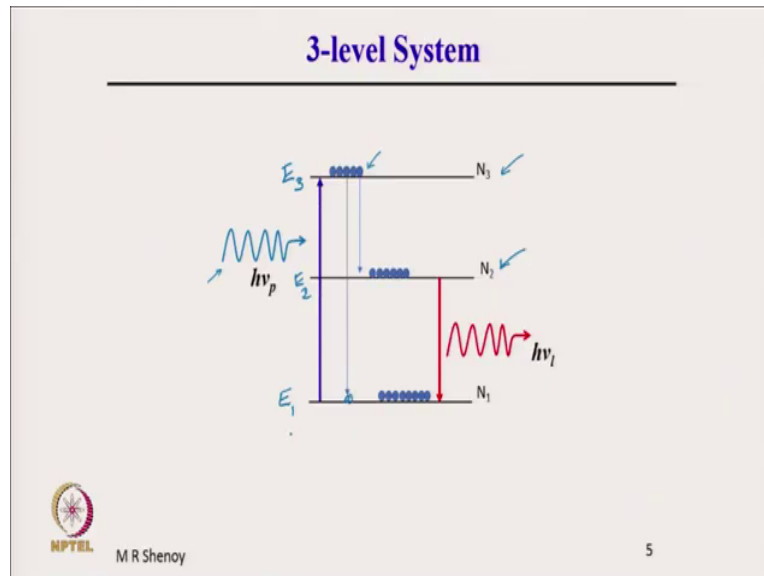
Therefore, there is no question of population inversion, but if you pump that is by some external means for example, I have shown here $h\nu_p$, means photons of energy $h\nu_p$ corresponding to the energy difference $E_2 - E_1$, then atoms will go to excited state. It need not be excitation by light, it could be some other means that is why I have shown here pumping by some other means like electric discharge, or any one of the means to excite atoms to the excited upper state.

Then there is radiation emission and radiation which is exciting, there is a dynamic equilibrium now, but in the presence of a pump. And naturally because of the additional excitation mechanism

Thermal equilibrium is the normal excitation mechanism because of the finite temperature of the system, but in the presence of pump it is like an additional pump temperature is one pump, but an external excitation is an additional pump this would always lead to a higher value of N_2 , but N_2 may still be not greater than N_1 , but N_2 will be more in the presence of an additional pumping mechanism.

And let us we will see the rate equations, which will tell whether it is possible to get N_2 greater than N_1 or not. So, that is a 2 level system.

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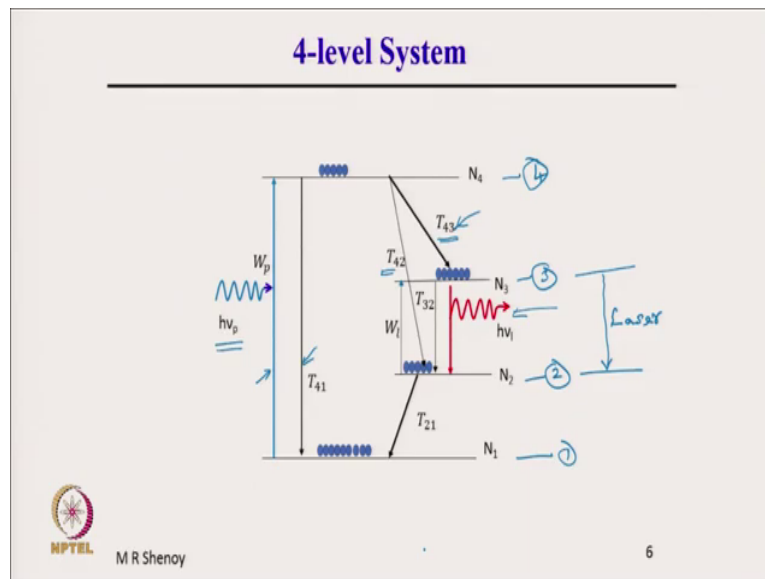


And then let us see, what is a 3 level system? So, here is a 3 level system. Now we are considering three energy levels of the atomic system, this one being the ground state E_1 , the upper state could be E_2 or some other higher state also, but the lowest level is the ground state and this could be E_3 . Let us say 3 the ground state and 2 excited states.

In this case if we have an external pumping mechanism, which is shown here by the blue sinusoid, $h\nu_p$ is the energy of the pumping photons, there is the pump photons. They will raise excite atoms from the ground state to an upper state and then from there atoms would make a downward transition to the second state in between and from there they could make a downward transition to the ground state. The atoms in the upper state here could also make a direct transition to the ground state here. So, they could also make a direct transition here to the ground state.

So, depending in such an atomic system that which is called a 3 level atomic system, we will see that it is possible depending on the lifetime of these levels. It is possible to have Population Inversion between the second level here E_2 and E_1 , it is possible under certain conditions that N_2 can be greater than N_1 . And this we will the conditions we will achieve by writing the rate equations for this 3 level system. So, we will discuss this in detail.

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Now let us see what is 4 level system.

So, the 4 level system here refers to, 4 participating levels. So, 1, 2, 3 and 4. So, the 4 participating levels, these are not necessarily the 1st excited, 2nd excited, 3rd excited, we will see subsequently that the levels could be different, but there are 4 predominantly participating levels in this interaction.

So, atoms from the ground state are excited to the 4th state here, as I mentioned this 4th could be actually a 8th state, 8th level, 8th excited level, but we call it a 4th state, this will become more clear when we discuss the scheme in more detail.

So, the pump atoms, pump photons, pump photons excite atoms from level 1 to level 4, from there of course, they can make a downward transition as shown here. This transition rate we call as T_{41} , that is level 4 to level 1, this is excitation upward excitation and T_{41} is d excitation from the 4th level to the ground state.

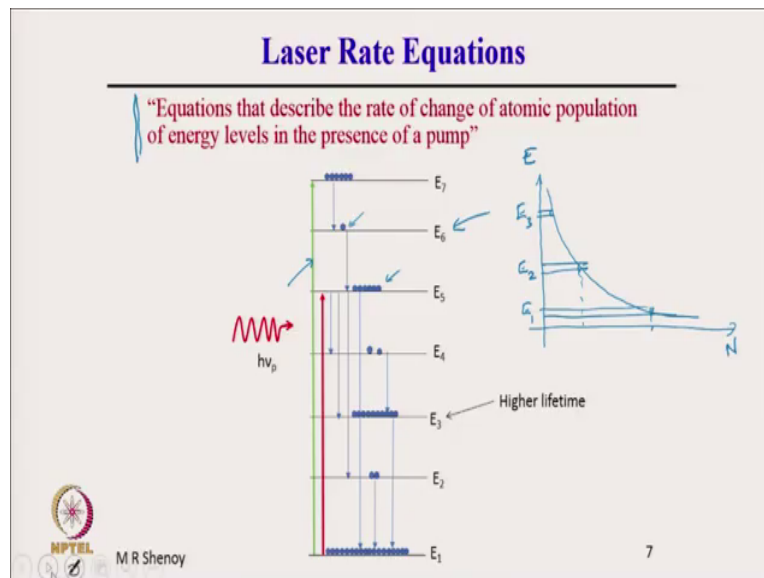
Atoms could also come down to the 3rd level which is represented by T_{43} , T_{43} is the rate at which it would come down, and T_{42} is the rate at which atoms would come down from level 4 to level 2. We will see some of the some of the transitions I have shown with bold arrows in a practical 4 level laser system.

We will see that the T_{43} , which is shown here by the bold arrow is predominant, that is it is a fast transition and therefore, it is possible under certain circumstances if T_{43} and T_{21} are much faster compared to T_{32} , then it is possible to have population inversion between this level 3 here and the level 2..

So, between these two levels it is possible to have population inversion. And then we can have laser action as shown here, that the laser action could take place laser action I refer to the generation of coherent laser radiation because of stimulated emission.

So, laser action could take place between level 3 and level 2 and this is called a 4 level system, we will write the rate equations for this 4 level system as well and see and also compare how the 3 level and 4 level lasers work and which one is preferable and why, ok.

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So, I have been mentioning about Laser Rate Equations. So, what are Laser Rate Equations? Laser Rate Equations are equations that describe the rate of change of atomic population of the energy levels in the presence of a pump.

The equations which describe the rate of change of atomic population of the energy levels in the presence of a pump, here pump normally we refer to an external excitation mechanism. So, here a general atomic system with several atomic energy levels of course, they are not equally separated it is just qualitatively shown as E_1, E_2, E_3 I shown here 7 levels.

For example, there is a pump, which is raising a pump may have more than one frequencies. For example, we will see that later on when we see the flash lamps, krypton flash lamp or

xenon flash lamp we will see that there are several excitation lines, which can excite atoms from the ground state to an upper state excited state as shown here.

For example, this photon has sufficient energy to raise atom from E_1 to E_7 , from there it makes downward transition and depending on the lifetime of the levels, the steady state number in different levels may be different. As you can see I shown here large number of atoms here very small number of atoms just and again large number of atoms small number.

The number of atoms which are accumulating in any energy level depends on its lifetime. We will discuss this again in detail that the lifetime of the levels will determine the population in the presence of an external pump. And we will exploit this to achieve Population Inversion between two excited levels, not necessarily ground state and an excited level it could be between any two excited levels to get laser action alright.

So, where the atoms are accumulating; that means the transition from for example, the transition from E_3 to ground here must be much slower that is why atoms are getting accumulated. And if you remember in Thermal Equilibrium we have the atomic population dropping exponentially like this.

So, the probability distribution given by Boltzmann distribution, if you remember the population diagram the numbers have to drop like this this is the exponential drop. So, what I have plotted here is energy E versus the number. So, this is the number, number of atoms in level of energy E_1 , E_2 , E_3 and so on.

But what I have shown in this diagram you see that E_3 has much larger number of atoms than E_2 and similarly E_5 has larger number of atoms than E_4 . So, this is possible depending on the lifetime of the levels we will discuss this mathematically writing the rate equation.

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Photon Flux $\phi(\nu)$

Recall: $\frac{A}{B} = \frac{8\pi h \nu^3}{(c/n)^3}$, $A = \frac{1}{t_{sp}}$, $u_\nu = \frac{I_\nu}{(c/n)}$


$\rightarrow \Gamma_{21} = W_{21} N_2 = B u_\nu g(\nu) N_2$

$= \frac{(c/n)^2}{8\pi t_{sp}} \frac{g(\nu) I_\nu}{\nu^2 h \nu} N_2 = \sigma(\nu) \left(\frac{I_\nu}{h\nu}\right) N_2$

$\rightarrow \frac{I_\nu}{h\nu} = \phi(\nu) \rightarrow \text{photon flux}$

$\therefore \Gamma_{21} = \sigma(\nu) \phi(\nu) N_2$ and $W_{21} = \sigma(\nu) \phi(\nu)$

Note: Rate Equations may be written in terms of W_{ij} ,
or in terms of the photon flux $\phi(\nu)$.



M R Shenoy

8

Before we write down the rate equations I would like to bring this concept of photon flux, because several authors you will see in several books that the rate equation is also written in terms of the Photon Flux, ϕ of ν that is the number of photons per unit time per unit area incident or crossing a plane is called the photon flux.

So, recall that we had the relation between Einstein coefficients A by B is equal to $8\pi h \nu$ cube divided by c by n cube. And A is equal to 1 divided by t_{sp} the spontaneous emission lifetime and u_ν the energy density at the frequency ν is equal to the intensity I_ν divided by c by n .

And we have written this rate of stimulated emission from level 2 to 1 was written as W_{21} into N_2 where W_{21} is the rate per atom multiplied by N_2 gives you the total rate of

stimulated emission, which was equal to $B_{21} u_{\nu}$ into the atomic line shape function into N_2 , which is then written in this form. So, we have substituted for B in terms of A .

And therefore, we get this expression here and u_{ν} if you substitute I_{ν} by c by n we get I_{ν} by $h \nu$ into N_2 , which was written as the cross section σ_{ν} of ν . This we have discussed in detail, but I am just recalling into I_{ν} by $h \nu$ into N_2 , this quantity I_{ν} by $h \nu$ is the photon flux.

I_{ν} is the intensity which is power per unit area power is energy per unit time per unit area divided by energy of one photon gives us the number of photons per unit area per unit time, which is called the photon flux, ϕ_{ν} you. Therefore, γ_{21} is equal to σ_{ν} of ν here. So, this is the photon flux ϕ_{ν} . So, ϕ_{ν} and therefore, γ_{21} is equal to this and because our W_{21} is equal to σ_{ν} of ν into ϕ_{ν} .

Therefore the rate equations maybe return in terms of these W_{ij} that is W_{21} , W_{12} and 3 , 4 and so on or in terms of the photon flux ϕ both are related. So, the rate equations may be written in terms of W_{ij} or the photon flux. We will write the rate equations in terms of the transition rates W_{ij} ok.

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RATE EQUATIONS – 2-level System

→ Consider a 2-level atomic system in the presence of a pump:

As $W_{12} = W_{21} = W_p$,

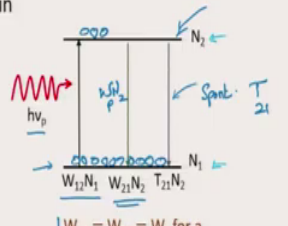
→ $\frac{dN_2}{dt} = W_p N_1 - W_p N_2 - T_{21} N_2$

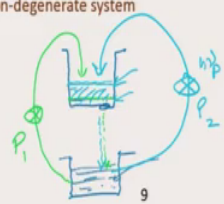
→ $\frac{dN_1}{dt} = -W_p N_1 + W_p N_2 + T_{21} N_2$

$= -\frac{dN_2}{dt}$

→ $N = N_1 + N_2 \rightarrow$ total no. of atoms, so,
 $\frac{dN}{dt} = 0 \Rightarrow \frac{dN_1}{dt} = -\frac{dN_2}{dt}$

→ At steady state: $\frac{dN_1}{dt} = 0 = \frac{dN_2}{dt}$





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Let us take up the rate equation.

First rate equations for a 2 level system. So, consider a 2 level atomic system in the presence of a pump that is illustrated here. There is an external pump of photon energy $h \nu_p$ which leads to absorption of atoms from the ground state to the upper state here. And that rate is W_{12} absorption here stimulated absorption W_{12} rate multiplied by N_1 , the number of atoms in that level and atoms once they reach here.

So, here are the ground state atoms large number of atoms in the ground state. Once the atoms is there they can come down to the lower level by stimulated emission because the same radiation would also stimulate atoms to come down or by spontaneous emission.

So, this is spontaneous emission which is given by the rate T_{21} . So, spontaneous transition rate is T_{21} and W_{12} because it is a 2 level system it is a non-degenerate system, for a non-degenerate system here W_{12} is equal to W_{21} is equal to W_p , the subscript p referring to the pump for a non-degenerate system and therefore, we can write the rate of change of population of level N₂, that is this here, what is the rate of change of population N₂ will increase because of absorption here atoms making upward transition, N₂ will decrease, because of spontaneous emission and stimulated emission, spontaneous emission here and stimulated emission.

So, dN_2/dt is equal to $W_p N_1$, please see $W_p N_1$ is the rate of stimulated absorption that is the number of atoms absorbed per unit time per unit volume that is dN_2/dt , which leads to a change in the atomic number atomic population of level 2. Hence this is called a rate equation, dN_2/dt is equal to $W_p N_1$ minus $W_p N_2$ because here $W_p N_2$. So, this is $W_p N_2$ it is written here. So, W_p times N₂ and the 3rd term it is also decreasing the number N₂ is decreasing because of spontaneous transitions T_{21} into N₂.

If you write the rate of rate equation for N₁, then dN_1/dt equal to N_1 reduces because of this upward transition. So, minus W_p times N₁, N₂ was increasing, but N₁ is decreasing because of this transition and N₁ is increasing because of these two because atoms are being added to the lower level and therefore, dN_1/dt is equal to it is quite obvious.

But I am explaining it because we will write more complicated rate equations for 3 level and 4 level systems and we can see that this is equal to minus dN_2/dt , because that is true because we have a 2 level system which means the total number of atoms N is shared between the 2 levels. So, we have $N_1 + N_2$ total number dN by $d t$ equal to 0, because the total number of atoms is constant therefore, it is 0 which implies dN_1/dt is equal to minus dN_2/dt .

Now, at steady state in the presence of pump at steady state dN_1/dt equal to 0 and it is equal to d. So, this can be imagined very easily like this that you have a bucket, let us say you have a bucket with a hole here and there is another bucket this bucket example that I give is.

So, this is the ground state where there is let us say this is a lower bucket which has plenty of water. So, full of water here and there is very little water here because of the thermal pump that is thermal energy atoms get excited, let me change the color because of the thermal pumps atoms get excited and water is being lifted it is like a pump.

So, we have a pump here. So, pump 1, P_1 atoms get lifted here into the upper level or upper bucket, but the bucket has a small hole because atoms also come down spontaneously. So, atoms will come down from here.

But at any state depending on the pumping rate you will have a certain steady state level water is being pumped water also comes down here, there will be a level is reached at which time the amount of water being pumped is equal to the amount of water which goes down. For example, if this hole is bigger then water will come down much faster the steady state level will be different.

Now, if I use another pump. So, let me change the color again. So, if I use another pump to pump water from here to here, another pump a second pump and this pump is like the external mechanism $h\nu_p$ we are using a pump. Now, because of a second pump now the n_2 level will go up like this to a n_2 height because water is being pumped because of temperature.

Water is being pumped because of the external radiation the second pump. So, this I want to call a second pump, water goes down and now the level has to go higher because a steady state will come when the amount of water going here will be equal to the amount of water this plus this. And therefore, the level will go up, but there will be a steady state.

So, long as the pumping rates are constant we will get a new steady state, this was original steady state now you have a new steady state level in this water and pump scheme. The same thing happens here that because of thermal equilibrium we had a certain distribution of atoms N_2 and N_1 , but because of the presence of an additional radiation of energy $h\nu_p$ the numbers will change now, but there will be a steady state. And at steady state dN_1 that is the

rate of change of the numbers in both the levels will be equal to 0, because steady state it is if we put these dN_2 equal to 0 for example, here.

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2-level System (contd.)

$\rightarrow N_1 W_p = (W_p + T_{21}) N_2$ or $\frac{N_2}{N_1} = \frac{W_p}{W_p + T_{21}}$


Since $W_p, T_{21} > 0 \Rightarrow N_2 < N_1$

\Rightarrow *population inversion is not possible in a 2-level system for any pumping rate. (at steady state)*

Now,

$$\frac{N_2 - N_1}{N_2 + N_1} = \frac{\frac{N_2}{N_1} - 1}{\frac{N_2}{N_1} + 1} = \frac{-T_{21}}{2W_p + T_{21}} = \frac{\Delta N}{N}$$

where $\Delta N = N_2 - N_1$, and $N = N_2 + N_1$



$\Rightarrow \Delta N = \frac{-T_{21}}{2W_p + T_{21}} N$

10

M R Shenoy

Then we have N_1 into W_p . So, this is dN_2 is by dt equal to 0 gives you $N_1 W_p$ is equal to this or N_1 by N_2 is equal to W_p divided by W_p plus T_{21} .

Note that, W_p is a pumping rate and T_{21} is a transition rate, both are positive since W_p and T_{21} are greater than 0, N_2 must be always less than N_1 , whatever be the pumping rate unless W_p becomes infinity. Only when W_p becomes infinity then N_2 by N_1 will become 1. N_2 could become equal to 1 at best, but otherwise N_2 is always less than N_1 .

In other words it means population inversion is not possible in a 2 level system for any pumping rate at steady state. I want to add this at steady state because in transient mode it is

possible to momentarily achieve population inversion, but you cannot get a steady state population inversion in a 2 level system.

So; obviously, we will have to go to 3 level and 4 level system, but before going let us discuss a little bit more. So, here therefore, N_2 by N_1 . So, if you do N_2 by N_1 minus 1 divided by N_2 by N_1 plus 1, then that will be N_2 minus N_1 divided by N_2 plus N_1 which will be equal to minus T_{21} .

So, here you simply do minus one divided by plus one we get minus T_{21} divided by $2W_p$ plus T_{21} . Now N_2 minus N_1 we call as ΔN . So, it is shown here ΔN is equal to N_2 minus N_1 and N_2 plus N_1 is the total number of atoms that is N . So, we have therefore, we have ΔN is equal to T_{21} minus T_{21} divided by $2W_p$ plus T_{21} into N . Now what are we reaching it let us see next.

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2-level System (contd.)

Thus, $\gamma(\nu) = \sigma(\nu)\Delta N = -\sigma(\nu) \frac{T_{21}}{2W_p + T_{21}} N = -\alpha(\nu)$ (say)

where $\alpha(\nu) = \sigma(\nu) \frac{N}{1 + \frac{2W_p}{T_{21}}} = \sigma(\nu) \frac{N}{1 + 2W_p\tau_1}$ $\left| \frac{1}{\tau_{21}} = \gamma_{21} \right.$


→ $\alpha(\nu)$ → Absorption Coefficient

Now, $2W_p\tau_1 = 2\sigma(\nu) \frac{h\nu}{h\nu} \tau_1 = \frac{I\nu}{I_s}$ (say),

with $2\sigma(\nu) \frac{\tau_1}{h\nu} = \frac{1}{I_s}$; $I_s = \frac{h\nu}{2\sigma(\nu)\tau_1}$ → Saturation Intensity

→ $\alpha(\nu) = \frac{\sigma(\nu)N}{1 + \frac{I\nu}{I_s}} = \frac{\alpha_0(\nu)}{1 + \frac{I\nu}{I_s}}$ → Saturated loss coefficient

$\sigma(\nu)N = \alpha_0(\nu)$ → Small-signal absorption coefficient



MPTEL M R Shenoy 11

The gain coefficient we have derived this earlier that, the gain coefficient γ of ν is given by σ of ν into ΔN , where σ of ν was the cross section interaction cross section therefore, for ΔN if we substitute from the previous expression here. So, if we substitute for ΔN from this expression in here, then we get γ of ν is equal to minus σ of ν into T_{21} divided by twice W_p plus T_{21} into N , which we call the keeping this minus sign here minus α of ν . So, say where α of ν is all the rest. So, everything which contains to this 1 is α of ν .

So, α of ν is equal to σ of ν into N . So, we have divided by T_{21} to the denominator. So, we get 1 plus twice W_p by T_{21} . And 1 by T_{21} we have already seen this is nothing, but equal to lifetime of that level τ_1 .

And therefore, we write α of ν is equal to σ of ν cross section into N total number of atoms divided by 1 plus W_p pumping rate per atom into τ_1 , where α of ν is called the absorption coefficient. So, γ of ν is minus α of ν γ of ν is the gain coefficient. Therefore, α of ν if it is negative then we call it α of ν and then that represents the absorption coefficient.

So, absorption coefficient of a 2 level atomic system. Now twice W_p into τ_1 is equal to we have substituted for W_p σ of ν into I_ν by $h\nu$ into τ_1 . So, this is W_p and this if we call this as I_ν by I_s where I_s is all of this. So, all of these if we call as I_s that is 1 by I_s , then we have 1 by I_s .

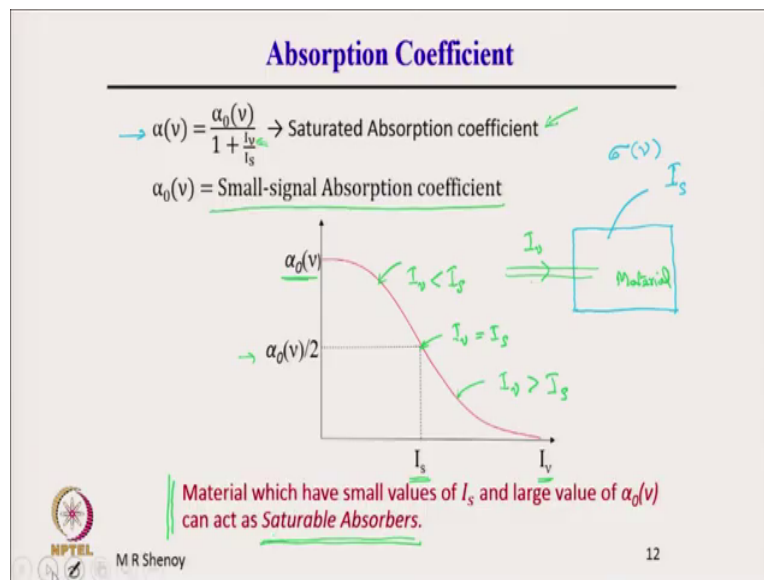
So, here with 1 by I_s equal to this then we have twice W_p into τ_1 is equal to I_ν by I_s twice W_p into τ_1 here this term is equal to I_ν by I_s . And therefore, we can write α of ν is equal to σ of ν into N divided by 1 plus I_ν by I_s , which again we write as α_0 of ν that is this term σ of ν into N in the numerator we call as α_0 of ν .

Because this one the expression which we have here for α of ν is called the saturated loss coefficient and the quantity I_s here is known as saturation intensity. We can see the dimensions for example, $h\nu$ is here energy and σ of ν is area and τ_1 is time. So,

energy per unit time per unit area is intensity and the dimensionally that is why we have put it as I_s , we called it as I_s and that term represents the saturation intensity, we will see the importance what is the saturation intensity.

So, the absorption coefficient can be written as $\sigma \nu N$. So, this is the cross section total number of atoms divided by $1 + I \nu / I_s$, I_s is a characteristic parameter of the material, which is the saturation intensity for a given material at a given wavelength we have a saturation intensity given by $h \nu / 2 \sigma \nu \tau$.

(Refer Slide Time: 32:59)



So, let us see this a little bit more. So, α of ν this absorption coefficient is a saturated absorption coefficient with the numerator called small signal absorption coefficient. What is the meaning of this? And what is its importance? if you consider a material, then this

material is characterized by a σ of ν at any at a given frequency it has a cross section and a saturation intensity I_s .

If we have radiation of some intensity I_ν . So, I_ν is passing through this material in the absence of any pumping the material will absorb this radiation, but the absorption by this material will depend on the intensity I_ν , that is what is being discussed here.

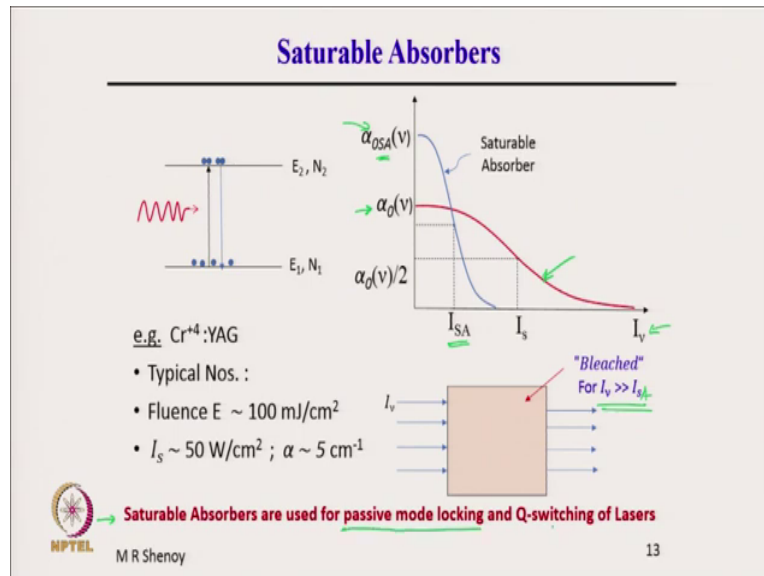
So, what we have plotted is I_ν intensity of the radiation as versus the absorption coefficient, when I_ν is very small compared to I_s , when I_ν in the denominator is very small, then we can neglect it compared to I_s . So, this is 0. So, we simply have the last coefficient equal to α_0 of ν and that is why. So, what is α_0 of ν ? α_0 of ν is the absorption coefficient of this given material this is a material for an atomic system. This material when the intensity of the radiation is very small hence the name small signal absorption coefficient.

If the intensity of the radiation I_ν increases, then in the denominator I_ν becomes more and therefore, as I_ν becomes comparable for example, when I_ν equal to I_s then we have α of ν is equal to α_0 by 2 that is half. So, this also is the definition of I_s the saturation intensity is the intensity of the radiation at which the absorption coefficient drops down to half of its value. That is the saturation intensity if the intensity increases further that is when. So, this portion here is for I_ν greater than I_s and this portion is I_ν less than I_s and this is when I_ν equal to I_s .

So, this is the absorption coefficient when I_ν equal to I_s intensity equal to the saturation intensity. And when the intensity of the radiation is very small then we have the small signal absorption coefficient this is important. Similarly, we will get small signal gain coefficient, which leads to saturation. So, that is why this expression is called saturated expression for absorption coefficient, saturated absorption coefficient and α_0 of ν is the small signal absorption coefficient.

What is the important? Materials here, materials which have small values of I_s and large values of α_0 of ν can act as Saturable Absorbers, Saturable Absorbers are widely used in laser physics. So, let us discuss about Saturable Absorbers.

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So, if you have a 2 level system then we have the expression for the absorption coefficient and we know the small signal absorption coefficient α_0 of ν and the red curve is the curve, which I have plotted in the previous slide. So, we can see the previous slide here. So, this is the red curve the absorption coefficient and discussion of variation of absorption coefficient with the intensity I_ν materials in which have small saturation intensity.

So, I_{SA} is the saturation intensity for a saturable absorber. So, saturable absorbers are absorbers which have a small value of saturation intensity and a large value of absorption

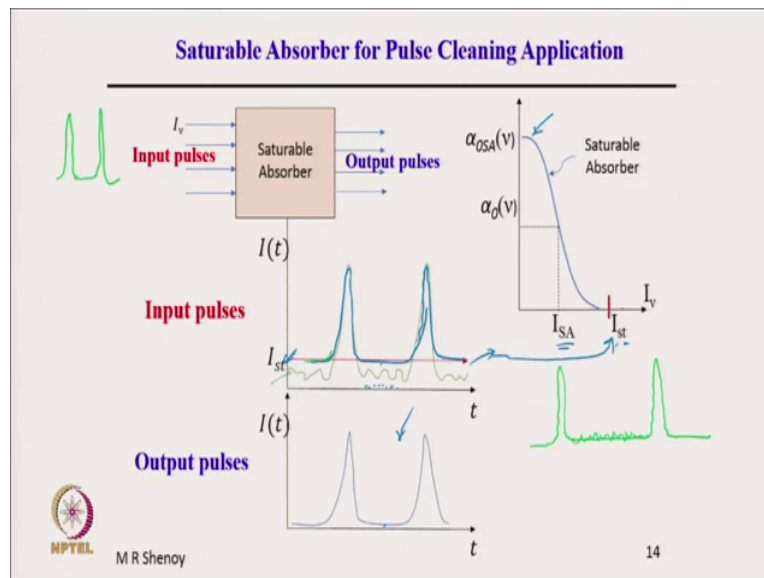
coefficient. So, this is the small signal absorption coefficient of the saturable absorber, α_0 of the saturable absorber and saturation intensity of the saturable absorber.

So, graphically we can show qualitatively, this is a qualitative explanation that saturable absorbers are materials, which have relatively large absorption coefficient at small intensities. And it completely goes down to 0 or we call it it gets bleached for $I \gg I_s$ much greater than I_s for intensities which are much greater than I_s , here I_s it gets completely bleached, bleached means what there is no absorption the absorption completely drops and the material behaves as if it is a transparent material.

At small intensities the absorption coefficient is very high, but at higher intensities it is almost transparent. We call the material gets bleached as the intensity becomes higher and higher such materials are called saturable absorbers.

Saturable absorbers are very important in the study of lasers, because saturable absorbers are used for passive mode locking and Q switching of lasers. We will study both of these in detail at a later stage, but at this point I just wanted to bring in this idea of a saturable absorber from the discussion of a 2 level system.

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One the, one application of the saturable absorber is for pulse cleaning, which can be easily understood that is why we have introduced this here. So, what you have is input pulses. So, there are pulse. So, this the input pulses we have pulsed input and output pulses here what is shown is the input pulses here, that is the intensity varying as a function of time. So, pulses have high intensity over a small duration of time and then there is no intensity ideally it should be 0.

But usually when pulses are generated there are small noisy fluctuations in the intermediate time period. Ideally we should have had a pulse, which is like this and then completely 0 till the next pulse comes, but usually in the generator pulse generator there are some noisy mechanisms and you have some small intensity fluctuations. So, there are applications where you need clean pulses with 0 fluctuation in between or very little fluctuation in between.

So, one of the ways that can be that can be used to clean the pulses is by passing through the input pulses through the saturable absorber. So, it is illustrated here that beyond a certain intensity it is completely transparent. So, you have all the intensity coming. So, the effectively what has happened is this pulse has become like this. Ok let me show in the blue color.

So, the pulse after if I want to show in the same graph, if I want to show the pulses then it has become like this after sorry, after passing through the saturable absorber that is what is shown in the graph below.

Because the pedestal which contained these low frequency, low intensity variations has been eliminated because at low intensities the absorption coefficient of the saturable absorber is very high, but at higher intensity. So, our laser as the intensity which is more than this. So, the laser intensity if it is here. So, this level that we have set here is the threshold level some threshold level here.

I s t saturable the intensity threshold value, this is the saturation intensity of the saturable absorber and this is the threshold value, that we have set. And you can see beyond this threshold the absorption coefficient is literally 0, which means all the intensity of the pulse, which is above this threshold value here completely passes through because absorption coefficient is 0.

Below that the absorption coefficient is high and therefore, it gets attenuated and therefore, the output would look like this basically we have shifted the baseline here. So, this level has now come to here. So, that is an application of a easily understandable simple application of a saturable absorber.

But more important applications of use of saturable absorbers in mod locking and Q switching will be discussed at a later stage. So, we will stop here with the 2 level system, in the next lecture we will take up 3 level system.

Thank you.