

Introduction to LASER
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Lecture - 10
Laser Rate Equations: 4-Level System

Welcome to this MOOC on LASERs. In the last lecture we discussed about Rate Equations in a 3-level system and today, we will take it up further to the 4-Level System.

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Recap : 3-Level System

Saturated Gain Coefficient

- $\gamma(\nu) = \frac{\gamma_0(\nu)}{1 + I_\nu/I_s}$
- $I_s = \frac{h\nu}{\sigma(\nu)\tau_s}$
- $\tau_s = \frac{(3W_p + 2T_{32})}{T_{32}(W_p + T_{21})}$
- $W_p \sim T_{21}, T_{32} \gg T_{21}$
- ∴ $\tau_s \approx \frac{2}{(W_p + T_{21})}$
- ∴ As $W_p \uparrow, \tau_s \downarrow \Rightarrow I_s \uparrow$
- Also, $\gamma_0(\nu) = \left(\frac{W_p - T_{21}}{(W_p + T_{21})}\right) N$
- ∴ As $W_p \uparrow \Rightarrow \gamma_0(\nu) \uparrow$

Scheme of 3-level Amplification

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First, a very quick recap of the 3-level system. So, the scheme of amplification is shown here, that is the 3-level amplification scheme where you consider 3 atomic energy levels of the system. So, the ground state and excited state, upper excited state.

A pump of energy $h\nu_p$ which raises atoms from the ground state to the third excited, second excited state and from there they make rapid transitions to the level e_2 with population N_2 here and then this is the level which is the meta stable state or which is the level that has a longer lifetime, because the level has a long lifetime.

Atoms tend to accumulate here and as qualitatively shown at steady state the number of atoms N_2 can be greater than number of atoms N_1 ; if we pump it sufficiently hard, because the transition T_{32} is a fast usually non radiative transition. The number of atoms at level 3 or N_3 is much smaller usually very small, because as soon as the atoms reach there they go to level 2 by a fast non radiative transition. So, the scheme illustrates and we had also obtained an expression for the gain coefficient and the expression for the saturated gain coefficient is here, γ_{ν} is equal to $\gamma_0 \frac{I_{\nu}}{I_s}$ where γ_0 is the small signal gain coefficient.

So, this is the small signal gain coefficient divided by $1 + \frac{I_{\nu}}{I_s}$, where I_{ν} is the intensity of the radiation which is being amplified and I_s is the saturation intensity. The expressions for I_s are given here $h\nu$ divided by σ_{ν} , σ_{ν} is the cross section and τ_s is the saturation time defined in this way. So, we have made this derivation and we have seen this.

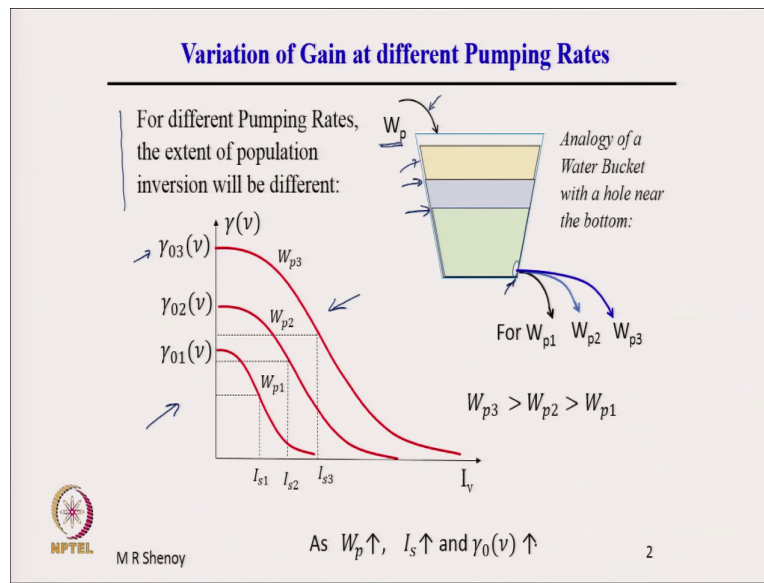
Now, since T_{32} is much larger than T_{12} so, here T_{21} , that is here T_{32} is much larger than T_{21} , this is fast, this is slow, T_{21} is slow then we can simplify this. Further, W_p is of the order of T_{21} we have seen that at threshold W_p is equal to T_{21} and above threshold W_p may be slightly greater than T_{21} , but it is of the same order as T_{21} .

So, W_p is of the same order as T_{21} , but T_{32} is much greater than T_{21} and therefore, this W_p here, $3W_p$ here can be neglected in comparison to T_{32} and then T_{32} and this T_{32} cancels and therefore, we have τ_s is equal to is approximately equal to 2 by W_p plus T_{21} .

This clearly shows that as W_p increases that is the pumping rate increases τ_s drops down and I_s goes up, because in the expression for I_s saturation intensity τ_s is in the

denominator. So, similarly for γ_0 of ν is given by this expression and as W_p increases, γ_0 would also go up and therefore, the variation of gain for different pumping rates would be as shown in this diagram here.

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So, this clearly shows that as the pumping rate increases the small signal gain coefficient γ_{01} , γ_{02} , γ_{03} will be higher and the corresponding saturation intensities are also higher. So, they are qualitatively schematically shown in this diagram how the gain curve changes with the pumping rate. So, for three different pumping rates we have shown that the extent of population inversion will be different at different pumping rates and therefore, you have different curves.

Now, this extent of population inversion will be different can be easily recognized or can be easily identified if we take a simple example of with analogy with a water bucket with a hole

near its bottom. So, there is a water bucket here. So, this is the water bucket in which there is a small hole near the bottom, if you pour water into the bucket water would go out from the hole and when the rate at which water is poured into the bucket becomes equal to the rate at which water goes out of the bucket we will have a steady state.

For example if we pump at a rate of W_p that is a certain rate at which water is pumped into this bucket then the level will build up to a certain level, the level will go up to a certain height so that the pressure with which water goes out of this hole will be such that the quantity of water going out rate at which water goes out will be equal to the rate at which water is poured into the bucket.


If we start pouring at a higher rate then naturally the level has to go to a higher level now, higher height so that water goes out with higher pressure, so that the rate at which water goes out increases. So, the point is depending on the pumping rate, the water level will be different in the bucket to reach steady state that is the rate at which water goes out will be equal to the rate at which water comes in then we have a steady level in the bucket.

Same thing is true here. In the atomic system if you pump harder then the extent of population inversion will be different. The extent of population inversion can be compared with the level of water here and therefore, the small signal gain coefficient will be different and then the gain varies as per the expression for the saturated gain coefficient and as we have discussed as W_p increases I_s also increases and γ_0 of ν increases alright.

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Example: Ruby Laser Amplifier

$\sim 400 \text{ nm} \rightarrow$ • Ruby: $\text{Cr:Al}_2\text{O}_3 \sim 10^{19}/\text{cc} \rightarrow N$ (total no. of atoms)
 $\sim 550 \text{ nm} \rightarrow$ • $\lambda_p \sim 500 \text{ nm}; W_{pt} = T_{21}$
 • $T_{21} = \frac{1}{\tau_l}, \tau_l \approx 3 \text{ ms} \rightarrow T_{21} = \frac{1}{3 \times 10^{-3}} \text{ sec}^{-1}$

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Let us take example of the ruby laser amplifier. We had mentioned that ruby laser is a 3-level example of a 3-level system and let me take the example and calculate what kind of pumping power is required to have lasing action in a ruby laser.

So, what is shown schematically is the ruby laser. There is a ruby rod ruby is a chromium doped alumina. The color, the pinkish red color of ruby comes from the chromium doping and typical doping concentration is 10 to the power of 19 per cc and that is our N that is the total number of atoms.

So, these are the atoms which are participating in the interaction. It is the three level which we are discussing are the levels of the chromium ion. So, in the presence of a pump, pump is generally a flash lamp. So, this could be a xenon flash lamp or krypton flash lamp and ruby has pumping bands of blue and green around 400 nanometer or so 400 nanometer or 550

nanometer so approximately, 500 nanometer so that is why we have taken lambda P of the order of 500 nanometer.

So, there are actually two absorption bands as we have seen earlier. So, the pump wavelength could be of this order and near threshold W_{pt} is equal to T_{21} and T_{21} we know, because it is T_{21} is equal to $1/\tau_{21}$. Lifetime is a measurable quantity. So, for ruby it is typically 3 millisecond. So, we know the value of T_{21} which means we know the pumping rate, the threshold pumping rate W_{pt} .

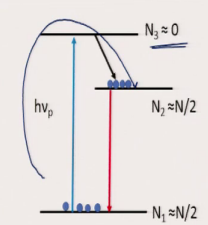
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Example: Ruby Laser Amplifier

$N_2 - N_1 = 0$ is the threshold for Amplification

→ $W_{pt} = T_{21}$, is the "threshold pumping rate"

→ Assuming $N_3 \approx 0$, $N_2 \approx N_1 \approx \frac{N}{2}$




• Pumping rate at threshold = $W_{pt} N_2 \approx W_p \frac{N}{2}$ ← Threshold Pumping

• Pumping Power at threshold $\approx W_{pt} \frac{N}{2} \times h\nu_p$ ← Power

$$= \frac{1}{3 \times 10^{-3}} \times \frac{10^{19}}{2} \times 6.6 \times 10^{-34} \times \frac{3 \times 10^8}{500 \times 10^{-9}}$$

$$= 660 \text{ Watts} \sim 1 \text{ kW}$$

→ Matches with practical laser system

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So, let us see what kind of power is required to reach the threshold. So, $N_2 - N_1 = 0$ is the threshold for amplification that is just before N_2 becomes greater than N_1 and corresponding to that the pumping rate is W_{pt} is the threshold pumping rate. Assuming N_3 is equal to 0, N_3 is the upper level.

Please see the mechanism that the atoms are excited to this level, but from here they rapidly go down here so that at any instant the number of atoms in this excited level, third excited level or the pump level N_3 is approximately 0 or much smaller compared to these and if N_3 is 0 then N_2 nearly equal to N_1 equal to N by 2.

N is the total number of atoms at threshold N_2 minus N_1 is 0 therefore, N_1 must be nearly equal to N_2 must be equal to N by 2 that is 50 percent of the atoms will be here and 50 percent of the atoms will be here, neglecting the presence of any atom at N_3 . Therefore, the pumping rate at threshold is given by $W_{p \rightarrow 2} N_2$. $W_{p \rightarrow 2} N_2$ gives us the number of atoms to be raised per unit time to achieve this kind of, achieve the kind of a population inversion or just threshold when N_2 is equal to N_1 .

Therefore, the corresponding pumping power is given by multiplied by, so this is the number. So, this is the number of atoms to be raised multiplied by the energy of photon raised per second therefore, energy per second will be the power. So, this is the power. So, this is the threshold pumping power. Threshold pumping rate multiplied by the photon energy of 1 photon gives you threshold pumping power maybe, I have written somewhere pumping power.

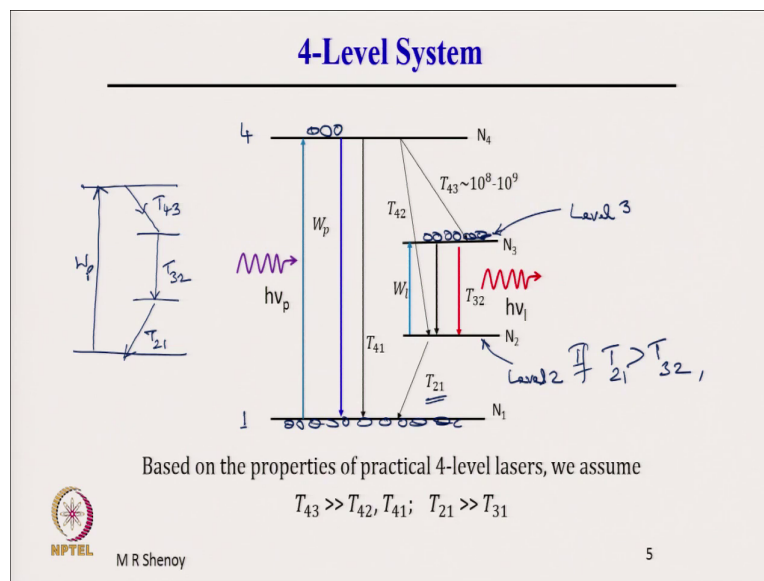
So, we are now substituting $W_{p \rightarrow 2}$ we had shown that is equal to T_{21} which is equal to 1 by τ_{21} this is N_1 by 2 is N_2 multiplied by h Planck's constant into ν_P the pump wavelength which is c by λ ; λ we have assumed that 500 nanometer and therefore, this comes out to be 660 and 60 Watts or of the order of 1 kilo Watt. So, this matches with that of practical systems.

Typically, ruby lasers require a threshold pumping power of the order of kilo Watts. So, a very simple calculation tells us that the numbers are alright it is of the order of 1 kilo Watt. We will see that if we go to a 4-level system the pumping power required is much smaller that is, because we are creating population inversion between two excited levels. It is normally we know that the number of atoms in the ground state is much-much higher than those in the excited states.

So, if you want to create population inversion between a ground state and an excited state more than 50 percent of the atoms from the ground state have to be raised to the excited state and therefore, the pumping power required is very high in a 3-level laser.

So, in a 4-level laser the population inversion is created between two excited states therefore, even if you have a number like 100 in one level and 90 in the lower level. You still have population inversion, that is the advantage of a 4-level system. So, will the practical lasers such as helium neon lasers, Nd YAG lasers are 4-level lasers and we will take up the 4-level system.

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So, the 4-level system; the scheme is this that this is the ground state from where atoms are raised to an excited state level 4. So, level 1. So, this is level 1, 2, 3, and level 4. From there the atoms rapidly come down to level 3. So, this is level 3 and if the lifetime of level 3 is

large or if level 3 happens to be and a meta stable state then atoms can get accumulated in this level here.

So, we have large number of atoms here; the ground state and then atoms are raised to the upper state here, from there they rapidly come down to this state and if the level has a long lifetime then atoms can accumulate here. When they come down to level 2, if T_{21} is a fast transition that is whatever comes here, immediately goes down to the ground state that is if T_{21} is greater than or much greater than T_{32} .

Then we will have net accumulation here and we will have population inversion between level 3 and level 2. So, this is level 2. This is what I mentioned that you can create population inversion between two excited state by raising much smaller number of atoms.

So, this is the scheme of amplification in a 4-level system. So, although there are transitions from 41, T_{41} , T_{42} , and T_{43} there are 3 transitions which 3 types of transitions spontaneous transitions that can take place from level 4, but in practice T_{43} in practical lasers, most of the practical lasers T_{43} is much larger, the rate is much larger compared to T_{42} and T_{41} .

T_{42} T_{41} are of the order of 10^4 , 10^5 , 10^6 , but 10^8 , T_{43} is much larger of the order of 10^8 or 10^9 and therefore, in practice we can neglect T_{41} and T_{42} and therefore, if we neglect those then our transition, our levels, our transitions will remain only this.

So, we are raising atoms here, then they come down by this way only. The other two transitions T_{42} and T_{41} we can neglect and then this comes down here and finally, goes back. So, this is T_{21} , this is T_{32} , we are showing spontaneous transitions and this is T_{43} and this is the pump W_p . So, this is the 4-level scheme, it does not mean there are no other transitions, but other transitions can be neglected in view of the practical numbers alright.

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Rate Equations: 4-Level System

At steady state:

→ $\frac{dN_4}{dt} = W_p(N_1 - N_4) - T_{43}N_4 \dots(1)$

→ $\frac{dN_3}{dt} = W_l(N_2 - N_3) + T_{43}N_4 - T_{32}N_3 \dots(2)$

→ $\frac{dN_2}{dt} = W_l(N_2 - N_3) + T_{32}N_3 - T_{21}N_2 \dots(3)$

→ $\frac{dN_1}{dt} = W_p(N_4 - N_1) + T_{21}N_2 \dots(4)$

→ And $N = N_1 + N_2 + N_3 + N_4 = N_1 \left(1 + \frac{N_2}{N_1} + \frac{N_3}{N_1} + \frac{N_4}{N_1} \right)$

• At steady state

$$\frac{dN_3}{dt} = 0 = \frac{dN_1}{dt} = \frac{dN_2}{dt} = \frac{dN_4}{dt}$$

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So, now we straight away come to writing rate equations for such a 4-level. You can see here I have shown only T 43 and T 21 the other two namely T 41 and T 42 I have not shown at all and therefore, we can write the rate equations at steady state $\frac{dN_4}{dt}$ that is the population of level 4 $\frac{dN_4}{dt}$ rate of change of the population is $W_p(N_1 - N_4)$. We can now directly write, because we know that W_p is the pumping rate here.

So, it raises atoms from N_1 and therefore, $W_p(N_1 - N_4)$ is the number of atoms which are reaching here $W_p(N_1 - N_4)$ is the number of atoms which go out and therefore, the net will be, $W_p(N_1 - N_4)$ net addition will be $W_p(N_1 - N_4)$ plus atoms are also going down through spontaneous transitions here which is $T_{43}N_4$. The negative sign says that N_4 is decreasing, because of this transition and positive sign wherever we have says the number is increasing, because of that transition.

If we write $\frac{dN_3}{dt}$ then we have W_p of course so we have W_1 here. So, N_3 will increase, because of W_1 N_3 will increase, because of this N_3 will decrease, because of W_1 into N_3 here and N_3 will decrease, because of the spontaneous transitions. So, we have two positive terms and two negative terms. So, one positive term, one positive and two negative terms. So, similarly we can write the rate equations for all the levels.

Further, we also have N is equal to N_1 plus N_2 that is the total number N is equal to N_1 plus N_2 plus N_3 plus N_4 . In practice of course, most of the atoms are in the ground state therefore, at any time N_1 is much-much greater than N_2 , N_3 , N_4 . In fact, even if you add all of these this will also be much smaller compared to N_1 and this can be written as N_1 into the same way that we had done for the 3-level system. At steady state as before $\frac{dN_3}{dt}$ each one of them is equal to 0.

So, recall the bucket where at steady state although water is getting poured into the bucket, water is going out of the hole, there is a steady state and therefore, there is a certain level which is maintained in the bucket. Exactly like that there is a certain number of atoms maintained at any given instant at any time whereas, it is a dynamic equilibrium photons are exciting atoms continuously to the upper state from where they get de excited, but there is a steady state which is maintained alright. From equation 1 that is the 1st equation here.

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Rate Equations: 4-Level System (contd.)

- From Eqn. (1)

$$T_{43}N_4 + W_p N_4 = W_p N_1 \Rightarrow \frac{N_4}{N_1} = \frac{W_p}{W_p + T_{43}}$$

$$\text{Since } T_{43} \gg W_p, \quad \frac{N_4}{N_1} \approx \frac{W_p}{T_{43}} \quad \dots(a)$$

- From Eqn. (4)

$$W_p \left(\frac{N_1 W_p}{T_{43}} - N_1 \right) = -T_{21} N_2$$

$$N_1 \left(W_p - \frac{W_p^2}{T_{43}} \right) = T_{21} N_2$$

$$\Rightarrow \frac{N_2}{N_1} = \frac{W_p (T_{43} - W_p)}{T_{21} T_{43}}$$

$$\text{Since } T_{43} \gg W_p, \quad \frac{N_2}{N_1} \approx \frac{W_p}{T_{21}} \quad \dots(b)$$



So, if we put dN_4/dt equal to 0 then this equal to this one and this contains only N_4 and N_1 therefore, we can get the expression for N_4 by N_1 . So, that is what is written here N_4 by N_1 is equal to W_p divided by $W_p + T_{43}$ and as indicated earlier, T_{43} is much faster, because it is a high speed non radiative transition and therefore, we can write N_4 by N_1 is equal to W_p by T_{43} these are approximations which are practically correct approximation.

So, instead of retaining large number of terms, we use these approximations right at the beginning so that we get a small compact expression which is not very different from the rigorously correct expression. Similarly, from equation 4 now, we look at equation 4 that is the fourth equation here this also has only two terms therefore, if you put dN_1/dt equal to 0 we can get the ratio as shown here. W_p into this equal to this one which means it is here written we have substituted for N_4 by N_1 from here W_p by T_{43} and we have an expression for N_2 by N_1 is equal to W_p into T_{43} minus W_p into this.

As before since T_{43} is much greater than $W_p N_2$ by N_1 is nearly equal to W_p by that is this is much greater. So, we have neglected then T_{43} T_{43} cancels and we have W_p by T_{21} .

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Rate Equations: 4-Level System (contd.)

→ Eq.(2): $\frac{dN_3}{dt} = W_l(N_2 - N_3) + T_{43}N_4 - T_{32}N_3 = 0$

gives $W_l N_2 + T_{43} N_4 = (W_l + T_{32}) N_3$

$N_1 \left(W_l \frac{W_p}{T_{21}} + T_{43} \frac{W_p}{T_{43}} \right) = (W_l + T_{32}) N_3$

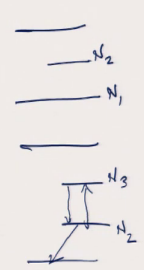
$W_p N_1 \left(\frac{W_l + T_{21}}{T_{21}} \right) = (W_l + T_{32}) N_3$


Or $\frac{N_3}{N_1} = \frac{W_p (T_{21} + W_l)}{T_{21} (T_{32} + W_l)} \dots (c)$

The required 'population inversion is $N_3 - N_2 > 0$

From (c) and (b), $\frac{N_3 - N_2}{N_1} = \frac{W_p}{T_{21}} \left[\frac{(T_{21} + W_l)}{(T_{32} + W_l)} - 1 \right] = \frac{W_p (T_{21} - T_{32})}{T_{21} (T_{32} + W_l)}$

or $\frac{\Delta N}{N_1} = \frac{W_p (T_{21} - T_{32})}{T_{21} (T_{32} + W_l)}$ $\Delta N > 0, \text{ if } T_{21} > T_{32}$





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Now, let us take equation 2 that is the 2nd equation there we had written 4 equations and the 2nd equation dN_3 by dt is equal to this. So, let us look at this. From equation 2, which is written here when you put that equal to 0 you can write this way and simplify it further and we get an expression for N_3 by N_1 . Now, we have an expression for N_4 by N_1 , N_3 by N_1 , and N_2 by N_1 . Why did we want this? Because we wanted to use the equation that N capital N that is the total number is equal to N_1 into 1 plus N_2 by N_1 plus N_3 by N_1 plus N_4 by N_1 .

So, we have got all these terms now so that we can relate the population inversion to the total number of atoms N alright. So, from c and b; so c equation here and b in the previous here, so

equation b, we can write. If you subtract one is N_3 by N_1 other is N_2 by N_1 . So, N_3 minus N_2 by N_1 comes out to be this and which is equal to this. There are no approximations, only approximation we have made is T_{43} is much greater than W_p or ΔN .

So, N_3 minus N_2 now we call as ΔN . In the 4-level, in the 3-level system we so in the 3-level system we called N_2 minus N_1 as ΔN , in the 4-level system we call this as ΔN . So, this is N_3 minus N_2 as ΔN , because ΔN here, represents the extent of inversion, the magnitude of inversion, N_3 minus N_2 is ΔN . So, in N_3 minus N_2 by N_1 is equal to so, this expression.

So, we have got the expression here. Note that ΔN for ΔN to be greater than 0 we need T_{21} to be greater than T_{32} . All other numbers are positive and therefore, ΔN that is greater than 0 or population inversion is possible if T_{21} is greater than T_{32} that is the rate at which atoms go from here to here, if this is faster than the rate at which atoms come from here to here there will be population inversion. Mathematically now, we see that yes we have got the same expression. So, this is what we discussed qualitatively that the requirement is T_{21} must be greater than T_{32} .

So, we have to look for a laser medium where the lifetimes or the transition rates are such that T_{21} that is from the lower excited states the transition rate is much faster compared to the upper excited state or if the upper excited state is a meta stable state ok. So, that I have rewritten here.

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Rate Equations: 4-Level System (contd.)

→ $N_3 - N_2 = \Delta N = N_1 \left[\frac{W_p (T_{21} - T_{32})}{T_{21} (T_{32} + W_l)} \right]$
 For Population Inversion, we must have $T_{21} > T_{32}$


→ As before, by using $N = N_1 \left(1 + \frac{N_2}{N_1} + \frac{N_3}{N_1} + \frac{N_4}{N_1} \right)$

And with $W_l = \frac{I_\nu \sigma(\nu)}{h\nu}$, $I_s = \frac{h\nu}{\sigma(\nu)\tau_s}$

Where, $\tau_s = \frac{(2W_p + T_{21})}{T_{21}(W_p + T_{32})} \approx \frac{1}{(W_p + T_{32})}$

SHOW $\Delta N = \frac{\Delta N_0}{1 + I_\nu/I_s}$ $\gamma(\nu) = \sigma(\nu) \Delta N$

$\Rightarrow \gamma(\nu) = \frac{\gamma_0(\nu)}{1 + I_\nu/I_s} \rightarrow$ **saturated gain coefficient**



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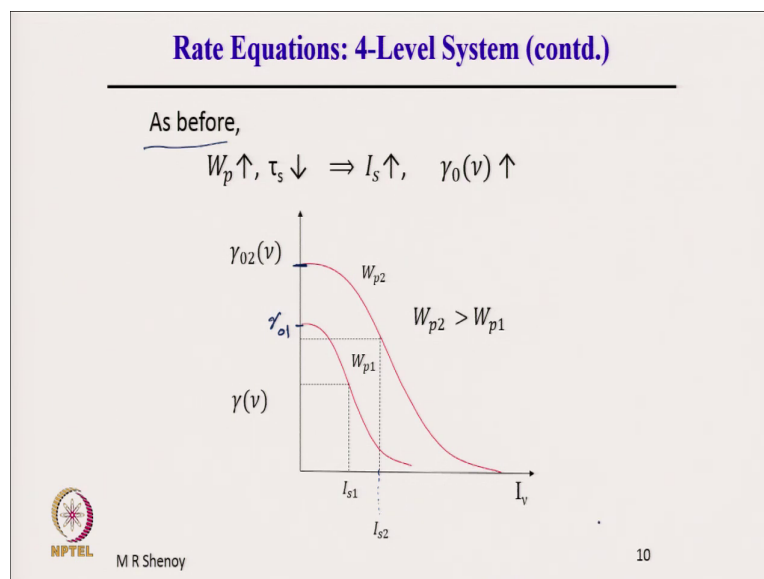
$N_3 - N_2$ is equal to ΔN into this therefore, for population inversion we must have T_{21} greater than T_{32} . Note that there is no condition on W_p , the pumping rate. In a 3-level system we had a condition on W_p that is the pumping rate. There was a pumping rate has to satisfy this condition, but now there is no, whatever be the pumping rate so long as you have T_{21} greater than T_{32} you will be able to achieve population inversion.

So, as before we can use this expression we have got a, b, c, N_2 by N_1 , N_3 by this expressions we have in terms of the rates and W_l is equal to I_ν into σ the cross section by $h\nu$, this is the definition of W_l similarly, for W_p and I_s is $h\nu$ divided by cross section into τ_s where τ_s is this. So, with these expressions we can show as before. In the 3-level case I had shown that ΔN we can be written as ΔN_0 which depends only on the pumping rate W_p divided by $1 + I_\nu$ by I_s .

We can write exactly similar expression here only the expression becomes slightly lengthier, because there are many more transition rates in the 4-level system, but you can get the same expression as ΔN_0 divided by $1 + I_\nu / I_s$ and therefore, γ_0 of ν is equal to we know that γ_0 of ν is equal to σ_0 of ν into ΔN . The gain coefficient is given by this and therefore, once ΔN is in this form then you will have γ_0 of ν is equal to γ_0 of ν divided by $1 + I_\nu / I_s$ which is called the saturated gain coefficient.

So, we have got almost similar expressions, same procedure, but note that in a 3-level system there was a condition on the minimum pumping rate which is required to achieve population inversion. So, we calculated that in the case of ruby laser that threshold pumping rate is of the order of kilo Watt. Here, we do not have such condition, what the requirement is T_{32} should be slower or T_{21} should be faster than T_{32} for population inversion alright.

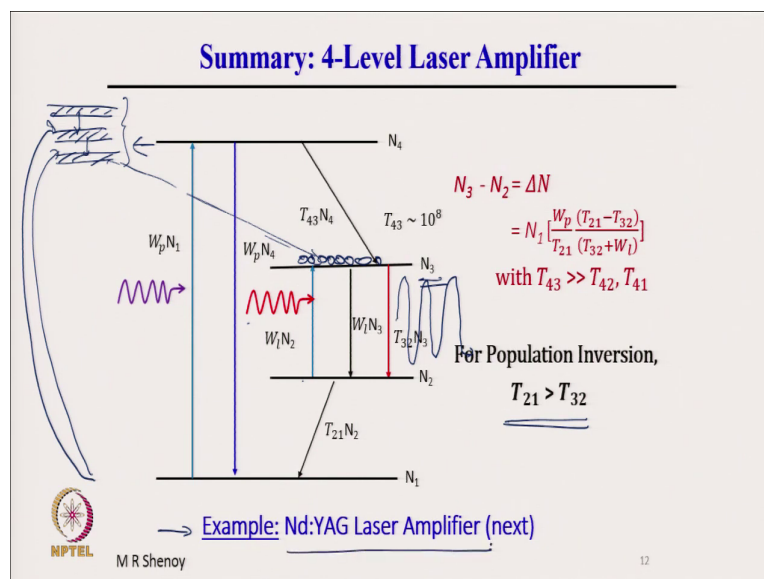
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So, as before written, as before the discussion is the same. W_p , if W_p increases τ_s decreases, we can see the expression τ_s is here. So, if W_p increases τ_s decreases. This W_p is negligible compared to the T_{21} here, T and therefore, as W_p so, its written here as if W_p increases τ_s decreases and that implies I_s saturation intensity increases and also γ_0 of ν increases.

So, just as in the case I discussed it earlier, in this lecture that for the 3-level system, for different pumping rates we had different gain curves that is the small signal gain coefficient increases. So, this is for W_p 1. So, γ_{01} here, this value here and this is γ_{02} small signal gain coefficient, which goes on dropping down as the intensity increases, the saturation intensity is shown here. This is the saturation intensity for the second case. So, the saturation intensity also increases.

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So, now we come to the summary of this 4-level laser amplifier. So, we have a 4-level system which comprises of 3 excited states and the ground states. This is the schematic that we show, but as I explained in the case of a 3-level system, the excited state this itself may comprise of several states or maybe bands, we will see this, this may be several bands. We will see this in next, we will see the example of a Nd YAG laser where I will show that if the 4th-level which I am indicating here as level 4 may comprise of 2 or 3 bands and atoms are pumped to these bands.

So, atoms may be raised to these bands, because if you take usually if we will see in the next lecture that Nd YAG laser is pumped by krypton flash lamp. These rays of course, diode pump solid state lasers will also be discussed later, but if you use krypton flash lamp it has several lines which will raise atoms to these bands and these bands within themselves have very rapid transitions from here to here and then from this level it comes down to the lower level which is here. In other words, the level 4 itself may comprise of several levels, but we do not call this as a 5 level or a 6-level laser.

This is still a 4-level laser and the pump band all of them are represented by level 4 from where atoms make rapid transition to a meta stable state. This level 3 is a meta stable state where the atoms have to accumulate. I am repeating this part in this summary that atoms are excited to an upper state or states from where they rapidly decay to a lower state, a lower excited state and if that state happens to be a meta stable state with the long lifetime then atoms tend to accumulate at that level.

The next excited state; if it has a fast transition to the ground state then we can always achieve population inversion for any pumping rate between these two levels. This is the mechanism of achieving population inversion in a 4-level system and a signal which corresponds to the energy difference between level 3 and level 2 will get amplified. Once population inversion is created then this signal will get amplified by stimulated emission.

So, the necessary condition for population inversion is T_{21} , this transition must be faster compared to T_{32} . We will take up in the next lecture the example of a Nd YAG laser

amplifier which is widely used and which is a very nice amplifier to discuss. The picture will become more clear.

Thank you.