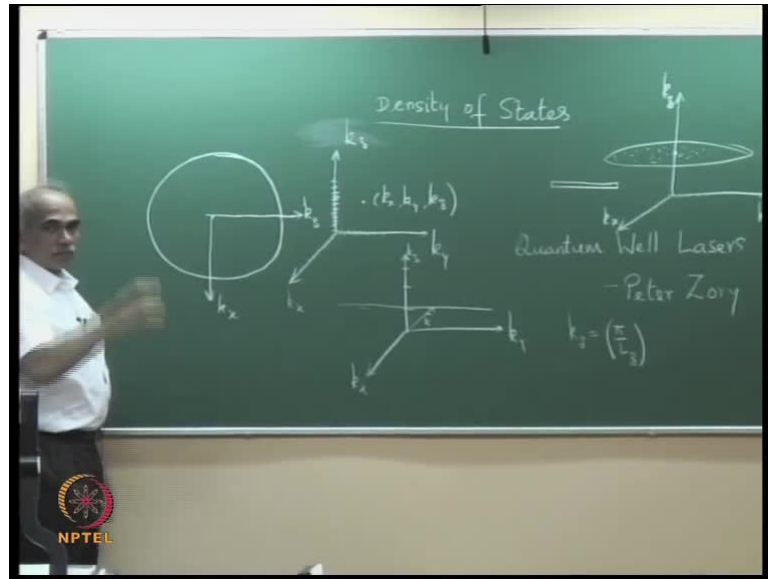


Semiconductor Optoelectronics
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Lecture - 6
Density of States in a Quantum Well Structure

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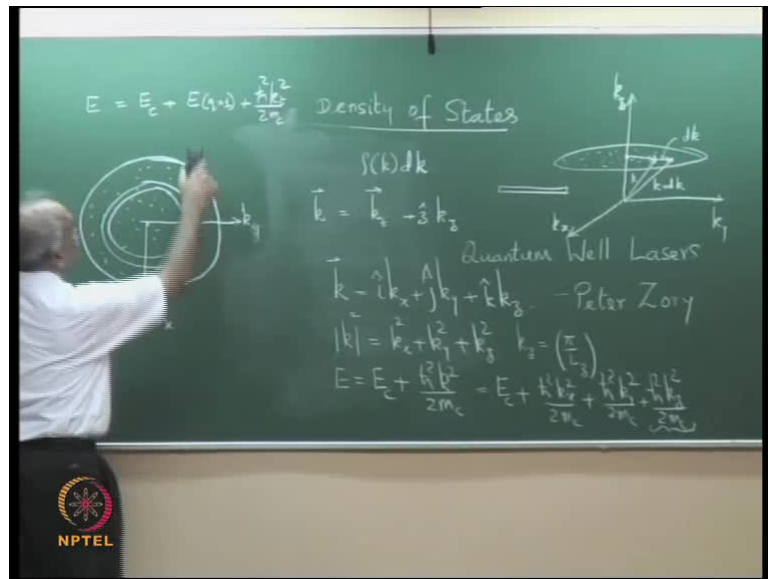
So, we continue with the density of states discussion and density of states in a quantum well structure. In the last class, we have discussed density of states for bulk semiconductors; and in the k space, in the k space I had shown that if this is k_x , k_y and k_z ; then in the k space, the k components can take very large number of values in general for a bulk semiconductor and therefore, the density of states available density of states is very large. And if you every point in the k space corresponding to a k_x , k_y , k_z represents this state, then we have very large number of such states. But if I reduce the dimension of the semiconductor in one direction, let us say I reduce the dimension in the z direction to make it into a thin sheet, then correspondingly the density of states change. And what we have is k_x and k_y still are in very large number of values, however the allowed values for k_z become highly discretised.

And since every k must have these three components; in this sheet an electron wave propagating through the sheet has to have three components of this; the first component of k non zero component of k_z is here. So this is k_z is here, which means every k value that you can think of must have one value k_z here, which means the tip of the k vector,

the tip of the k vector in the k space must lie on a plane k_z equal to constant. So for the first value k_z equal to constant is simply k_z equal to π/L_z into m equal to 1, q equal to 1. So k_z the first value is π/L_z therefore, all points the vector k can take large number of values still, but the tip of the vector must lie in this plane.

Therefore, in calculating the density of states, so we had this picture here, where this is k_z , k_x and k_y we considered discs with k_z equal to constant, the first one. And in this because k_x and k_y take very large number of values or permitted to take large number of values corresponding to each k_x and k_y , we have a point a permitted point in this disc, which means if I slip the disc and show the disc like this k_z is here central axis k_z ; this is k_x and this is k_y .

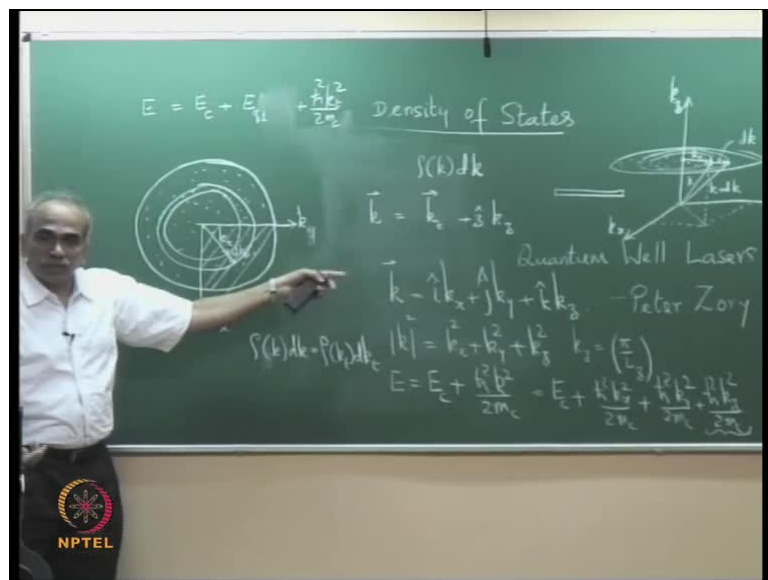
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So, k_x , k_y and k_z , k_x , k_y , k_z that is this disc in the k space I am showing like this, so k_z is coming up k_x , k_y , k_z in this you see large number of allowed values. Our objective is to find out $\rho(k)$ which is defined through $\rho(k) dk$ is the number of states between k and $k + dk$, so if this is some value k and the next value $k + dk$ if it stands here, so this is k , this is $k + dk$ which means this is the vector dk here, this is dk then our objective is to find the number of allowed states between k and $k + dk$ if you see in this the k will correspond to a value which is here please see this is the disk the k can take any value like this between k and $k + dk$.

So, this will correspond to a circle here a circle corresponding to k and the next circle outside corresponding to $k + \Delta k$, this transverse component here you see k is equal to $k_x \hat{i} + k_y \hat{j} + k_z \hat{k}$ and $k^2 = k_x^2 + k_y^2 + k_z^2$, we have energy E is equal to $E_c + \frac{\hbar^2 k^2}{2m}$, for this disc q is equal to 1 the first value of k_z corresponds to $q = 1$ therefore, I write this as $E_c + \frac{\hbar^2 k_x^2}{2m} + \frac{\hbar^2 k_z^2}{2m}$ I am just substituting this and $\hbar^2 k_x^2 + \hbar^2 k_y^2$ by $2m(E - E_c)$ where $E - E_c$ is this component $\hbar^2 k_z^2$ by $2m$ if I put $q = 1$ k_z takes the value $k_z = \frac{\pi}{L}$.

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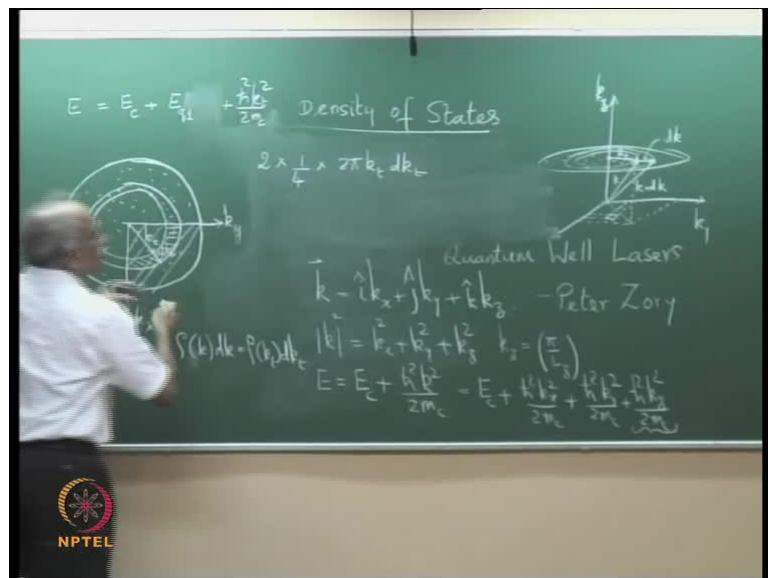


So, I can write this as $E_c + E_{q=1}$ this E stands for that component only here or sometimes it is denoted as $E_{q=1}$ either you can write this as $E_c + E_{q=1} + \frac{\hbar^2 k_x^2 + \hbar^2 k_y^2}{2m}$ our objective is to find the number of points in this plane, so this is k_t and this additional vector here is Δk_t please see here this is k , this is $k + \Delta k$ corresponding to k there is a k_t which has only 2 components

x component and y component you can just drop down here and you find out what is the x component what is the y component and this is the z component, so what is k_t ? k_t is just this that vector which has x component and y component do you follow, so it is these 2 components make this transverse vector, so this you see if I drop from here that is k_t and dk_t is the incremental volume, incremental vector, so that is dk_t . So we have to find out the number of points therefore, the number of points between k and k plus dk is the same as number of points between k_t and k_t plus dk_t do you agree because k can sit anywhere around this circle and k plus dk can sit anywhere here the number of points is the same between k and k plus dk in fact it is ρ_k , ρ of k , dk is equal to ρ of k_t , dk_t the number of points as far as the number of points is concerned it is the same.

So, to find out the number points, point means every point represents an allowed state number of points between k and dk plus and remember we have to take only the positive k_x , positive k_y and positive k_z values because the negative values are already taken into account in our boundary condition. And therefore, I need to consider only this quadrant, positive quadrant like in the case of the sphere I have to consider only the octant, positive octant.

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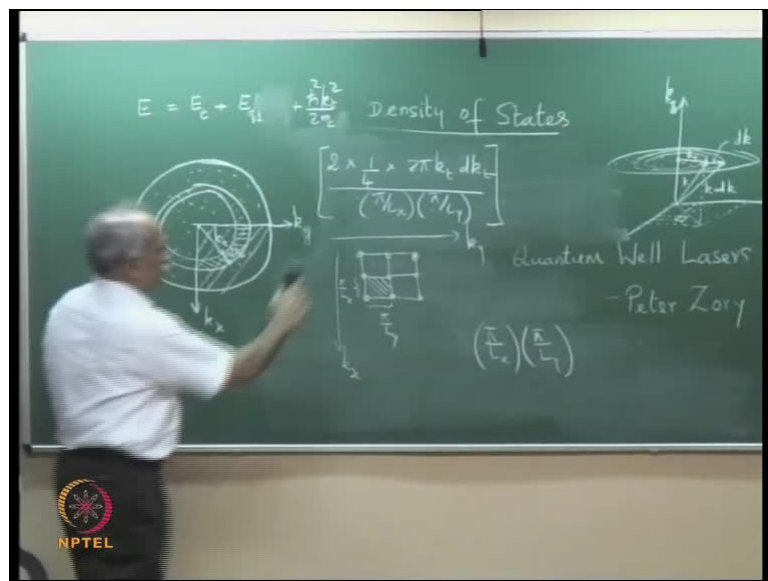


So, here of this circle I need to consider only the quadrant and therefore, what is the density of h the number of points always remember it has to come from definitions the number of points between k and k plus dk is in the positive octant or quadrant. So one

forth into what is the number of points to find out the number of points I have to find out the area here now it is disc not the sphere in sphere what did we do we found out the volume between that shell and then divided by the volume occupied by one state.

Now it is the area it is a disc therefore, area between k and $k + dk$, what is the area? One fourth into $2\pi k dk$, $2\pi r dr$ the area of the, so the area of this is $2\pi r dr$, so $2\pi k dk$, one fourth of it and of course, we need a factor to take care of the electron spin this is in terms of area. Now we need to know the number of points which means what is the area occupied by 1 point in this case what do you think is the area occupied by 1 point please see corresponding to every value here and every value. Here there is a point, the next value here there is a point, next value here there is a point, so what do you have you have a square actually rectangle if L_x is not equal to L_y it is a rectangle, so you have a rectangle with the 4 points at the corner, but the every point is shared by 4 other rectangles so if I enlarge the view.

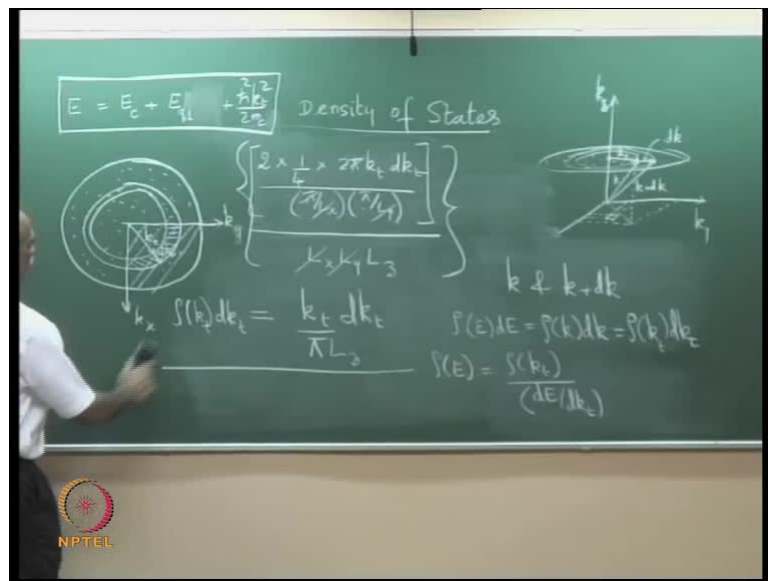
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We have a situation, so we have points like, so this is x direction this is, so this is k_x direction and this is k in the k space these points are here they are very closely packed points. I have just add 2 points, but actually if you zoom in then you will see that the points are sitting like this and what is disc separation π by L_x , what is this separation π by L_y . So the area of a rectangle here is π by L_x into π by L_y in the k space area of this rectangle where points represent k states allowed states, but you see that every point

is surrounded by 4 such rectangles which means number of point per rectangle is one fourth one fourth because each contributes one fourth to this one fourth to this and one fourth to this like this. Therefore, the number of points per rectangle or area per point is π by L_x divided by π by L_y , area corresponding to one allowed state is this much in the k space the area in the first quadrant here is this divided by area corresponding to one state will give me number of states is this clear the area in the k space divided by area occupied by one state will give me number states.

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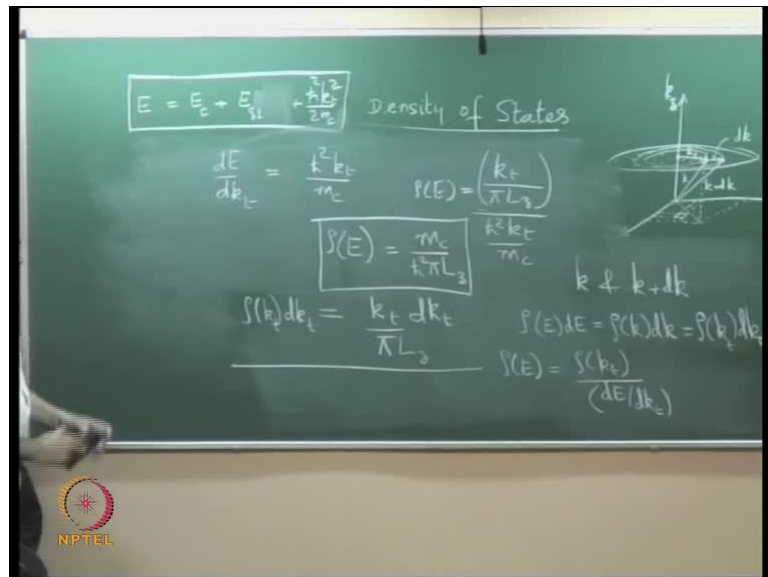
So, this divided by π by L_x into π by L_y this is the number of points between k and k plus $d k$, disc per unit volume of the material will give me density of states, so this what is the volume of the material L_x into L_y into L_z , so this gives me the density of state. So please see it is just in one line or one expression the entire derivation is there 2 standing for the spin one fourth of the quadrant of area $2 \pi k_t d k_t$ divided by area occupied by one state will give me the numerator gives me number of states between k_t and $d k_t$ or $k_t k_t$ plus ρk_t plus k_t plus $d k_t$ between k and k plus $d k$ this is the number of points per unit volume that is the definition, so number of states per unit volume, so simplify this what will you get what do you get, so $L_x L_y L_z$ cancels.

And L_z remains, so ρ of k_t , $d k_t$ equal to ρ of k , $d k$ is equal to k_t , $d k_t$, πL_z we see this you can cancel these though one π goes here, so one π remains in the

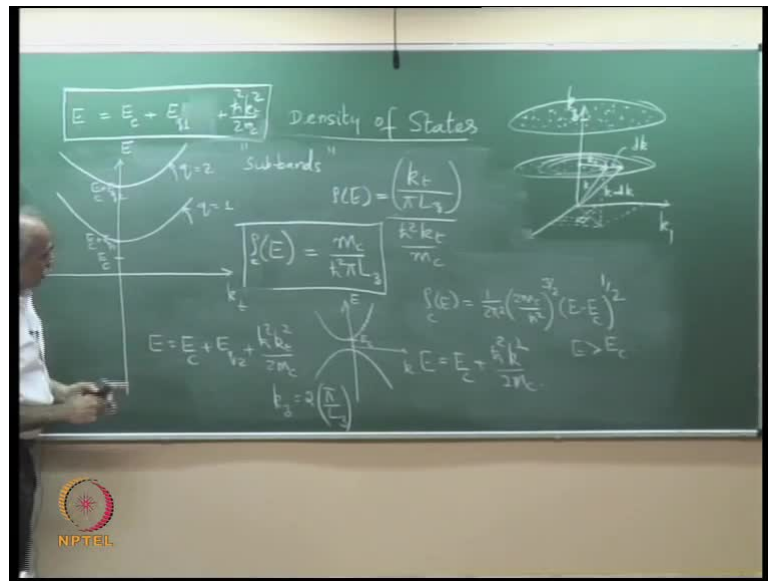
denominator, L_z remains in the denominator here, L_z per unit volume, π by L_x divided by that is right, so in the k space we have got the density of states.

Now, as I indicated in the last class we would like to know the density of state in the E space then how do we go for E space we will use the expression relation between k and E and use the fact that $\rho(E) dE$ is equal to $\rho(k) dk$ is equal to $\rho(k) dk$ and therefore, $\rho(E)$ finally, we need the density of states in the E space, so $\rho(E)$ is equal to $\rho(k) dk$ divided by dE by dk , I have used this equal to this and we have the expression for $\rho(k)$ already.

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So if I may erase this, so if you have a clearer picture in one expression the entire derivation of the density of states is present. So dE by dk_{\parallel} let me do this simple algebra complete it, so dE by dk_{\parallel} . So this is a constant this is a constant for that plane this is a constant therefore, we simply will have \hbar^2 cross square into kt by m_c^2 , 2 cancels \hbar^2 cross square kt by m_c this is dE by dk_{\parallel} , so $\rho(E)$ therefore, is equal to $\rho(E)$ is equal to $\rho(k_{\parallel})$ which is kt divided by π into L_z that divided by dE by dk_{\parallel} , so divided by \hbar^2 cross square kt , $\rho(k_{\parallel})$ divided by dE by dk_{\parallel} , \hbar^2 cross square kt divided by m_c , so m_c goes to the top, kt , kt cancels, so we have $\rho(E)$, $\rho(E)$ is equal to in a quantum well, $\rho(E)$ is equal to m_c divided by \hbar^2 cross square π into L_z where L_z is the thickness of the quantum.

Something interesting that you see here that the density of states is independent of the energy right hand side there is no dependence on energy. Recall the density of state in a bulk $\rho_c(E)$ this is for the conduction band actually

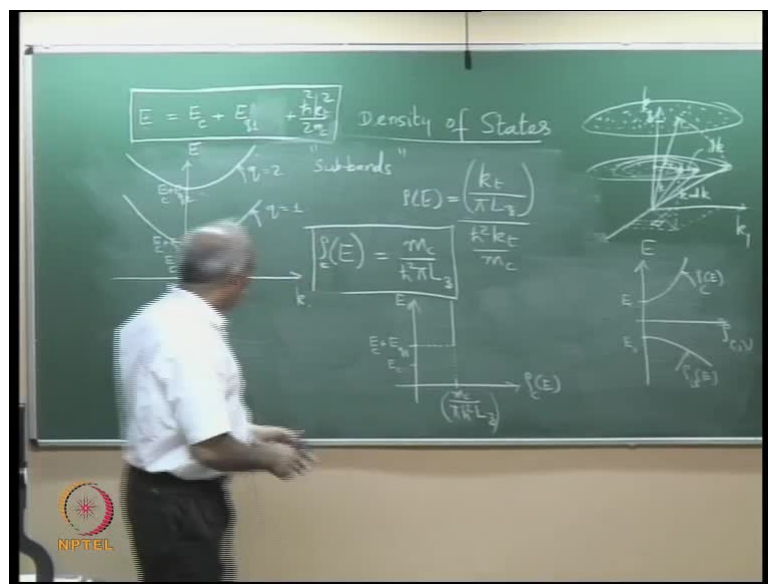
Therefore this is also ρ_c , $\rho_c(E)$ is equal to 1 over $2\pi^2$ square, $2m_c$ by \hbar^2 cross square to the power E by 2 into $E - E_c$ to the power half conduction band which means E for E greater than E_c , there is a energy dependence, but there is no energy on the right hand side which means the density of state is constant, first point.

Now, let us come to some additional discussion about this energy, let us see what is this? This tells us that recall the $E-k$ diagram how did we have the $E-k$ diagram we had E , so E

varying parabolically, so E here, so E was equal to E_c plus $\frac{h^2 k^2}{2m_c}$ divided by twice m_c right from at k is equal to 0, so this is k here, E is equal to E_c , so this is E_c , what are these? These are simply nothing, but points of allowed states allowed states here, so you have large number of allowed states for k greater than 0, it is varying parabolically.

In this case you plot the density of states the $E-k$ diagram, k versus E , up to E_c you have no states, up to E_c if you put k equal to zero I have E_c plus E_q , E is equal to E_c plus E_q , what is E_q ? E_q is the allowed energy value of electrons. So this is E_c the next level here is E_c plus E_q beyond this as k almost takes continuous values it varies parabolically in other words the band varies like this parabolically for q is equal to 1. We have now plotted with respect to k not the total k with respect to k , but what you need to see is earlier we had allowed states right from E_c , now there are no states above E_c up to E_c plus E_q and your band starts from here, so this is a parabolic variation of the band what about when q is equal to 2 I have energy E is equal to E_c plus E_q plus $\frac{h^2 k^2}{2m_c}$, if I put q equal to 2 which means my k_z is now 2 times π by L_z originally for q equal to 1, k_z was π by L_z , now q equal to 2, 2 times π by L_z that value is suddenly a jump here, so this is E_c plus E_q at E_c equal to E_q I again have a parabolic variation, so we have the next sub band these are called energy sub bands, sub bands.

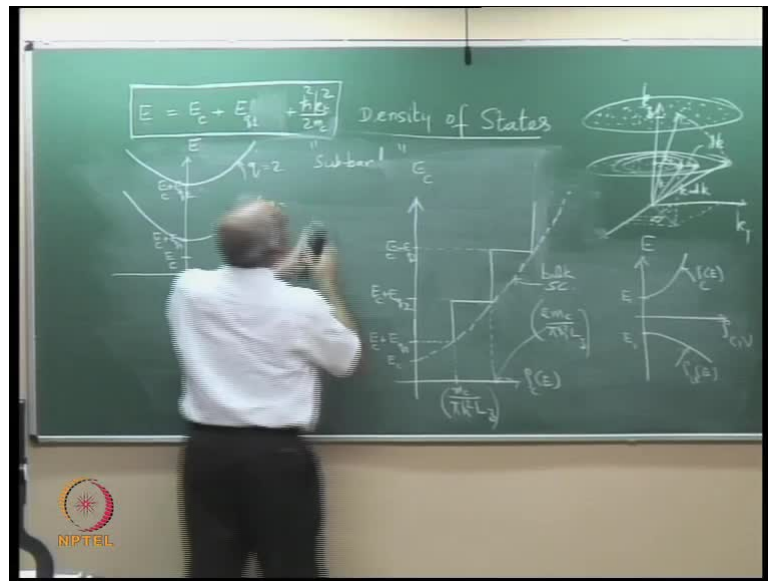
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The E_k diagram in a bulk has a parabolic variation here, in the case of a quantum well you have energy sub bands, these are bands because almost continuum k_x and k_y take almost continuous values because k_x and k_y take very large number of values, but k_z has quantum jumps because π/L_z is a very large number because L_z is a very small number. So it just jumps here for one value of k_z for this value of k_z you have large number of k_x, k_y values permitted next it goes to the next level and again large number of k_x, k_y values that is why in the case of quantum wells the energy band E_k diagram is characterized by sub bands and corresponding to each sub band you have a density of states. Now why did I bring this here because if you see this, this was for q equal to 1 at q equal to 2, I have a second disc, now a second disc is available where again large number of points are present, if I change k_z energy changes, E changes if I have to plot this how would it look like let me first plot this $\rho(E)$ you remember for the bulk I keep it here this is E and this is ρ , ρ of c comma v , we had $E_c, E - E_c$ square root dependence here.

So this was ρ_c of E we had shown this in the last class, so E_c and similarly, for E_v we had density of state varying like this where this variation is ρ_v of E we want to plot the corresponding variation here the density of states, let me show first only the conduction band at E is equal to E_c we have no state because the first value of k_z starts E is equal to $E_c + E_q$, So E_q that is $E_c + E_q$, so this is $E_c + E_q$, E_q is only that component $\hbar^2 k_z^2 / 2m_c$, so $E_c + E_q$ is the total energy at this value what is on this axis ρ_c of E . So up to this there were no density of states when we came here there is density of states and can you tell me what is the value here m_c by m_c divided by $\pi \hbar^2 L_z$ the density of states is constant this is the value what does that mean that means beyond E is equal to $E_c + E_q$, we have density of state, but density of state is constant, density of state continuously increases with energy here, density of state is constant repeatedly I keep telling that why am I ...

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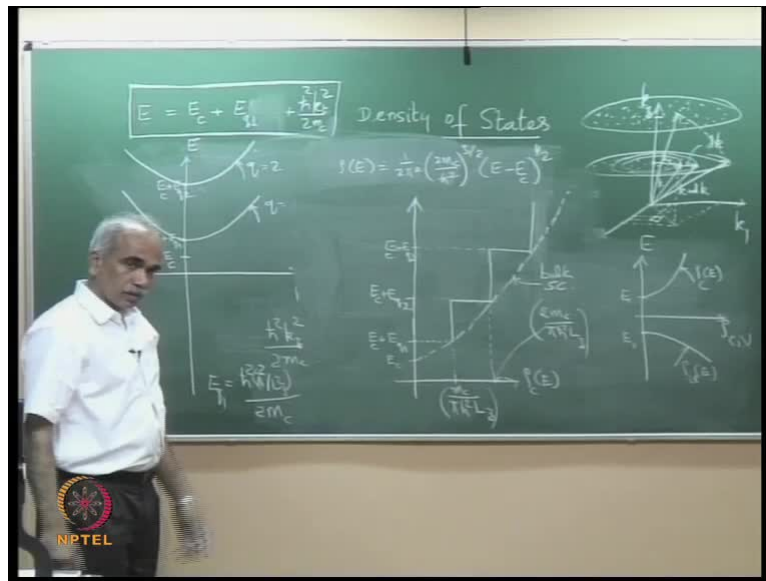


Density of state, so much many will think why after density of states for the last three class. Let me assure you that it is the density of states which will determine the device characteristics and the device performance. It is very important to understand the density of states, so its constant here what happens when energy as energy exceeds $E_c + E_{q2}$ we have a second sub band equivalently we have a second disc, so we have another amount of area that is if my k value becomes large. So let me show this k like this, this length is large if k is large, which means energy is large the tip of the vector can either lie in this disc or it can also lie E_c the length is large it can also lie here, so when the energy is large you have the second disc available which means the density of state doubles the density of state for disc is what we have calculated here m_c divided by $\pi \hbar^2$ cross square into L_z when you have a second disc available the density of states doubles and therefore, the density of states here for E , as E increases. That we continue on the same diagram E increases to $E_c + E_{q2}$ we have a second disk available and density of state simply jumps to this value.

What is this value here? 2 times m_c by $\pi \hbar^2$ cross square L_z , so this value here is $2 m_c$ by $\pi \hbar^2$ cross square L_z when you reach you can extend this $E_c + E_{q3}$, so $E_c + E_{q3}$ there is a third sub band available here which means there is a third disc available and you have the density of states making a third jump. What we have seen is this step like behavior in the density of states it so happen that if you draw the density of states my diagram is not very good. So what I have drawn here is the density of state

corresponding to the bulk, bulk semi conductor this one if you draw the density of states corresponding to the bulk semi conductor you will get this step kind of variation. I leave this as an exercise to you to show that indeed in the bulk if you plot the density of state to quantum wells and bulk corresponding to E is equal to E_c plus E_q 1 the density of states that you get for the bulk is the same as this value. And if you put E_c plus E_q 2 in the expression for this you will get 2 times m_c by π h cross square L_z .

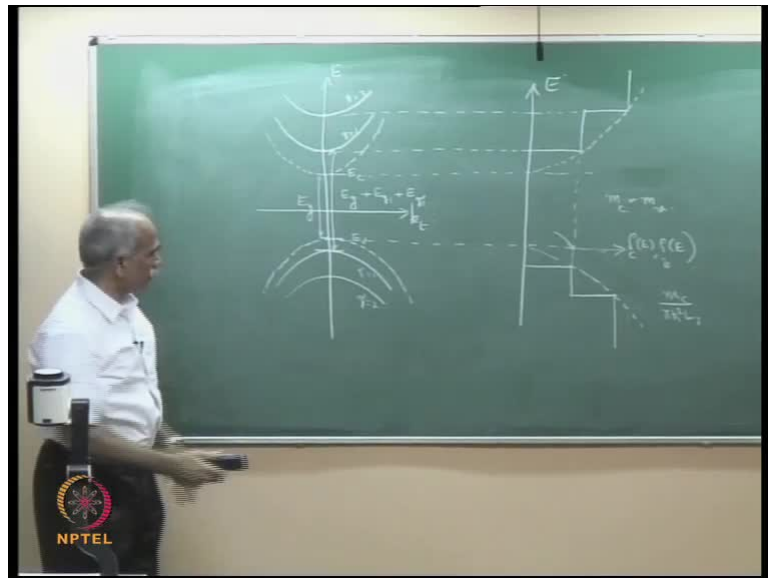
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You just verify this which expression the expression for bulk where we have density of states $\rho(E)$ is equal to $\frac{1}{2\pi^2} \left(\frac{2m_c}{\hbar^2} \right)^{3/2} (E - E_c)^{1/2}$. If you substitute for E this value here you will get this result as this, if you substitute for E this value E_c plus E_q 2 means E_c plus $4\pi^2 \hbar^2 m_c^{-1} L_z^{-2}$, so what is E_q 2, E_q 2 is $\hbar^2 \pi^2 E_z^2$ divided by twice m_c , so E_q 1 is equal to $\hbar^2 \pi^2$ by L_z^2 divided by $2m_c$, π^2 by L_z^2 divided by $2m_c$, if you put E_q 2 it will be 4 times because q equal to 2, so q square will become 4. So 4 times this if you substitute that value here you will get this number that is why the density of states is shown like that. Otherwise how do you know that this is touching this you substitute and see with indeed touches that value, so what we have seen is the density of states has a step like variation we will see its implications later when we go to the devices. My discussion has been focused only on the conduction band exactly similar discussion is applicable for the valance band and you will see the same step like behavior in the valance band.

So, let me give the final results for the density of states in a quantum well because the energy expression that I have written is for energy in the conduction band exactly similar discussion in the quantum in the valance band will give you similar results that I want to write the results.

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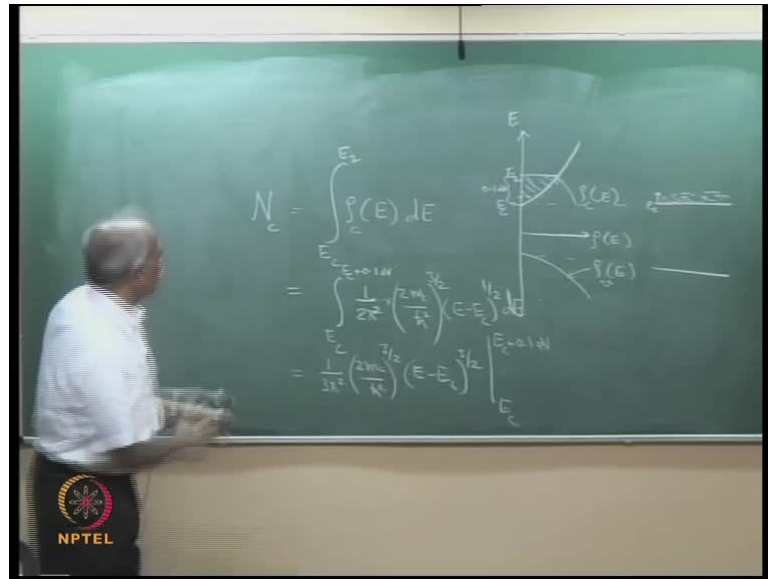


So, the $E-k$ diagram for a bulk $t-d-L$ for a quantum well, so this is q dash equal to 1, q dash equal to 2, q dash is the corresponding number for the valance band, q equal to 1, q equal to 2 this is $E-k$ diagram. So E versus k and here is E_c and I have drawn this by dotted line because if you make a quantum well structure the first allowed states starts from here and immediate implication is the band gap of a semiconductor is E_c minus E_v this is E_g if you make a quantum well structure the first allowed state is here and the first allowed whole state lowest energy whole is here the lowest the effective band gap is this that is original E_g plus E_{q1} plus E_{qdash1} , this is E_{qdash1} .

So what you see is by changing the structure into a quantum well structure you have effectively changed the band gap. We will discuss more about this later when we discuss about band gap engineering and we will discuss more about this, but this is the corresponding $E-k$ diagram and you have that corresponding density of states diagram. Everywhere vertical axis is E to the density of states for bulk and corresponding to, so we have this similar density of states, so this is ρ_c of E and ρ_v of E , so the dotted line corresponds to bulk and solid line corresponds to the quantum well structure. And I

have assumed in showing that this point is the same as this point, I have assumed that m_c equal to m_v because recall that this value here was m_c divided by $\pi \hbar^2$ cross square into L_z , if I am showing this value the same it means I have assumed m_c equal to m_v if m_c is not equal to m_v both the curvature here and the first the first value corresponding to the density of states would be different in the valance band which is in general to...

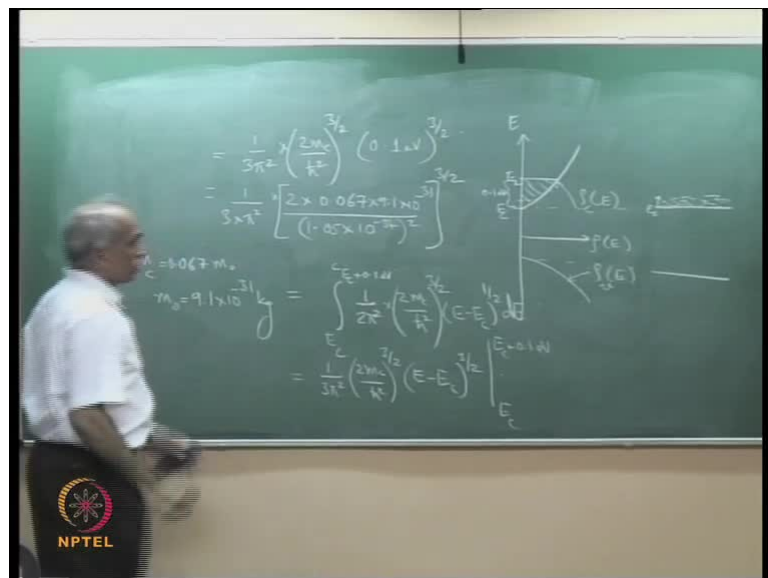
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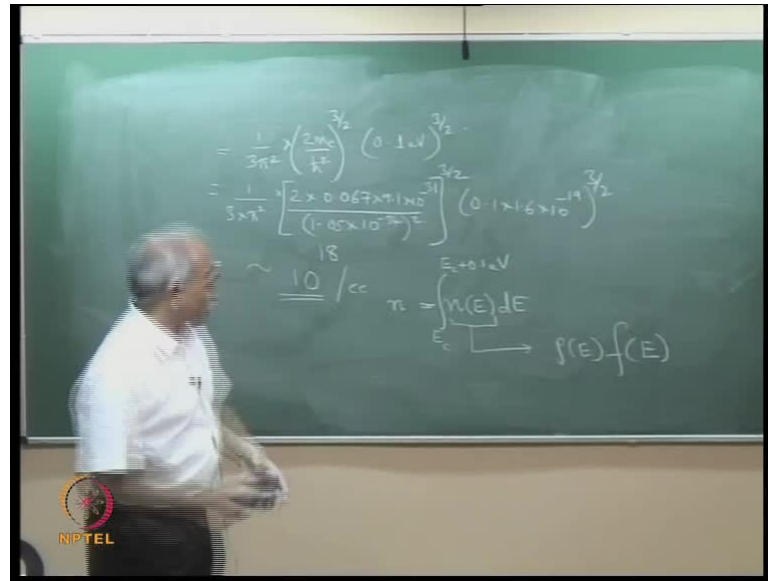
In general the density of states and the $E-k$ diagram variation is different in the valance, but if the effective masses are the same. Then they come at the same dilemma is this all right, we will discuss more about the quantum well and how to make etcetera at a later stage. So this is the $E-k$ diagram and this is E versus the density of stage why do we need density of states. Let us answer this question why do we need density of states because as I have already said that we want to know the carrier concentration, the carrier concentration is determined by the available states multiplied by the probability of occupation. How do we calculate the number of available states? The number of available states N is equal to ρ_c of E number of available states in the conduction band is equal to ρ_c of E dE integrate over E_c to some value let us say up to some value E_2 . You need the number of states up to E_2 what is this E_2 , so E_2 I am plotting again this is E , this is ρ_c of E ρ_c of E and ρ_v of E , this tells me the density of the states therefore, I want to calculate the actual number of states.

So, the number of states this is E_c and I want to know up to some energy E_2 . How many states are present here the dense the number of states is ρ_c of E , dE from E_c to E_2 , if you want to calculate the density of states in the entire band, so you write from E_c up to the top of the band, so if you want to see up to some value E_2 substitute here and let us see what do we get, so this is equal E_c for example, E_2 is equal to let us say this height is $0.1 E_v$ let us put some numbers here. Now $0.1 E_v$ I want to calculate from the you are familiar old band diagram which does not tell anything except energy gap, so if you want to calculate the density of state the number of states available here from E_c this bottom up to some height how many states are available here. So I can calculate only through the density of states, so now we are coming to numbers. Let us see from E_c up to this which means I calculate it from E_c plus $0.1 E_v$ substitute the values very quickly 1 over 2π square into twice m_c by h cross square to the power 3 by 2 into E minus E_c to the power half dE , integrate this what do you get this is constant.

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This is constant $E - E_c$ to the power half integrate you get $3/2$ and a factor 2 by 3 2 2 cancels, so you are left with 1 over 3 π square into 2 m c by h cross square to the power of 3 by 2 into $E - E_c$ to the power 3 by 2 from E_c to $E_c + 0.1$ E v if you put E equal to E_c it is 0 , so it is simply $E_c + 0.1$ E v which means substitute the values 1 divided by 3 π square into twice m c by h cross square to the power 3 by 2 into $E - E_c$ is simply 0.1 E v the 0.1 E v to the power 3 by 2 . So this is equal to 1 over 3 into π square substitute value for π into 2 into m c let us say gallium arsenate m c is equal to 0.67 times m_0 where m_0 equal to 9.1 into 10 to the power of minus 31 k g , so substitute m c equal to 0.067 , 0.067 into 9.1 into 10 to the power of now you are like engineers calculate, this divided by h cross square h cross square is h , how much is h cross? 1.01 01 05 .

Student: (())

05 all right.

So, 1.05 into 10 to the power of minus 34 h cross this is h cross, so h cross square and the whole to the power 3 by 2 into 0.1 E v it is in electron Volts whereas, all the others are s i units, so you have to convert this into Joules, so it is 0.1 into 1.6 into 10 to the power of minus 19 to the power 3 by 2 , simplify this and find out what is the answer I think this will be of the order of 10 to the power of 18 the number is of the order of 10 to the power of 18 per cc , so we have find out how many states are available from the

bottom of the conduction band up to some height $0.1 E_v$ this is the available states, but if you want to know the carrier concentration then we have to have $N(E)$, dE I want to get the carrier concentration N then it is $N(E) dE$ you integrate again from E_c to E_c plus $0.1 E_v$, but what is $N(E)$ density of carriers, carrier density? Carrier density is given by $\rho(E)$ density of states multiplied by the probability of occupation.

The probability of occupation is given by the fermi function density of states multiplied by the probability. So now to know carrier concentration we need probability of occupation occupation probability. So this will be our next topic is occupation probability if I substitute the occupation probability multiplied by the density of states and you integrate you will get carrier concentration, now on we will get the numbers. So once the basic physics and the picture is clear then calculating numbers is not a problem at all, so our next topic will be probability of occupation, so I will stop here.