

Basic Quantum Mechanics
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Module No. # 08
Angular Momentum - II
Lecture No. # 02
Angular Momentum Problem (Contd.)

In our last lecture, we were discussing the operator algebra involving the angular momentum operators. We will continue our discussions on that. We first said that the operators J^2 and J_z commute with each other and since they are observables, we can have a set of simultaneous Eigen vectors of the operator J^2 and J_z .

So, we wrote down $J^2 \ket{\lambda m} = \lambda(\lambda+1) \ket{\lambda m}$ and $J_z \ket{\lambda m} = m \ket{\lambda m}$. Then, we had defined two operators J_+ and J_- . J_+ was $J_x + i J_y$ and J_- was $J_x - i J_y$ and now, let me consider the commutation of J_+ and J_z . So, this is $J_+ J_z - J_z J_+$. So, this is equal to $J_x J_z + i J_y J_z - J_z J_x - i J_z J_y$. So, this will be $J_x J_z - J_z J_x + i(J_y J_z - J_z J_y)$. Now, you know that $J_y J_z - J_z J_y$ commutator is just $i J_x$. So, this will be i times i is minus J_x and this is in the wrong order. So, $J_z J_x$ as you would recall from the previous lecture is $i J_y$. So, $J_x J_z$ is the minus of that. So, this is minus $i J_y$.

So, if I take the minus sign, so you will get $J_x + i J_y$ and so, this is equal to minus J_+ . I hope I have done it correctly. So, the left hand side is $J_+ J_z - J_z J_+$ is equal to so much. What we do is we operate this on $\ket{\lambda m}$, that is we operate these both sides on $\ket{\lambda m}$. So, you will get, you had here. So, let me put a plus sign here, plus sign here and minus sign here and take this to the other side. So, we will have, please see this $J_z J_+$. Let me do this once more.

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$$\begin{aligned}
 [J_+, J_z] &= -J_+ \\
 J_+ J_z |\lambda m\rangle - J_z J_+ |\lambda m\rangle &= -J_+ |\lambda m\rangle \\
 -m J_+ |\lambda m\rangle + J_z J_+ |\lambda m\rangle &= +J_+ |\lambda m\rangle \\
 J_z \{ J_+ |\lambda m\rangle \} &= (m+1) \{ J_+ |\lambda m\rangle \} \\
 J_z |p\rangle &= (m+1) |p\rangle \\
 J_z \{ J_+ |p\rangle \} &= (m+2) \{ J_+ |p\rangle \}
 \end{aligned}$$

So, we have J_+ , J_z the commutator this was equal to $-J_+$ we derive that so we had $J_+ J_z$. Let me leave some space minus $J_z J_+$ let me leave some space. This is equal to J_+ .

What we do is as we have mentioned operate each term on ket λm operate this on ket λm operate this on ket λm . Then, you can see that J_z operating. Did I do it correctly? $J_+ J_z$ is equal to this was minus J_+ minus. So, this is minus, I am sorry, there is a minus sign here. So, J_z operating on ket λm is m times J_+ ket λm minus $J_z J_+$ ket λm minus J_+ ket λm .

So, I multiply all sides by a minus sign. So, I get a minus sign here, plus sign here and a plus sign here. Now, I take this term on the right hand side. So, I get J_z operating on J_+ ket λm is equal to $m+1$, $m+1$ operating on J_+ ket λm . Thus, if ket λm is an Eigen ket of the operator J_z belonging to the Eigen value m , then J_+ ket λm is also an Eigen ket of the same operator J_z , but now belonging to the Eigen value $m+1$.

So, let me denote this file, say ket p . So, J_z ket p is equal to $m+1$ ket p . Now, I do the same thing. Instead of operating this by ket λm , I operate this on ket p . So, I will obtain $J_z J_+$ ket p will now be equal to $m+2$ J_+ ket p of course, provided ket p is not a null ket. Similarly, if ket p is not a null ket, then J_+ ket p is also an Eigen ket of the operator J_z belonging now to the Eigen value $m+2$, provided this is not a null ket.

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Handwritten equations on a blue background:

$$J^2 [J_{\pm} |\lambda m\rangle] = \lambda [J_{\pm} |\lambda m\rangle] \quad |\lambda m\rangle$$

$$J_z [J_{\pm} |\lambda m\rangle] = (m \pm 1) [J_{\pm} |\lambda m\rangle]$$

$$J_{\pm} |\lambda m\rangle = C |\lambda, m \pm 1\rangle$$

$$J^2 [J_{\pm} J_{\pm} |\lambda m\rangle] = \lambda [J_{\pm} J_{\pm} |\lambda m\rangle]$$

$$J_z [J_{\pm} J_{\pm} |\lambda m\rangle] = (m \pm 2) [J_{\pm} J_{\pm} |\lambda m\rangle]$$

$$\lambda \geq m^2$$

$m_{\max} = j$
 $|\lambda j\rangle$ is not a null-ket
 $J_{+} |\lambda j\rangle$ is a null ket

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We had seen that J , therefore we had seen that J plus ket lambda m . In the last lecture, we had shown that is also an Eigen ket of the operator J square belonging to the same Eigen value J plus ket lambda and $J_z J$ plus ket lambda m is equal to m plus 1 J plus ket lambda m . Therefore, if ket lambda m of course, a non-null ket is a simultaneous Eigen vector of the operator J square and J_z . Then, J plus ket lambda m is also a simultaneous Eigen ket of J square and J_z belonging to the same Eigen value of J square, but 1 raised Eigen value of J_z .

Therefore, J plus ket lambda m must be something like some constant times m plus 1 m increases by 1, but lambda value remains the same. I can go on doing that. Then, we will have J square operating on J plus J plus ket lambda m . That will belong to the same Eigen value of J square J plus J plus ket lambda m , but for this, we say will now be m plus 2. That is why, they are known as ladder operators provided this is not a null ket. So, the value lambda remains the same on m increases, but I have the inequality that lambda must be greater than equal to m square. So, at some m value, we must have a maximum value of m because this process cannot be on indefinitely. Otherwise, it will violate the inequality lambda is greater than or equal to m square and therefore, at some maximum value of m .

So, let there must exist a maximum value of m and let the maximum value of m be equal to J . Then, the max, this Eigen ket is a non null ket, is not a null ket, but J plus ket

λJ is a null ket because J m value with the maximum value of m is J . So, you cannot have m plus 1 and therefore, J plus ket λJ must be a null ket.

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$$\begin{aligned}
 J_+ J_- |\lambda j\rangle &= 0 & j = m_{\max} \\
 J_+ J_- &= (J_x - iJ_y)(J_x + iJ_y) & J_+ J_- |\lambda j\rangle &= 0 \\
 &= J_x^2 + J_y^2 + i(J_z) & J_z |\lambda j\rangle &= j |\lambda j\rangle \\
 &= J^2 - J_z^2 - J_z & J_z^2 |\lambda j\rangle &= j^2 |\lambda j\rangle \\
 (J^2 - J_z^2 - J_z) |\lambda j\rangle &= 0 & \lambda - j^2 - j &= 0 \\
 (\lambda - j^2 - j) |\lambda j\rangle &= 0 & \Rightarrow \boxed{\lambda = j(j+1)} &
 \end{aligned}$$

So, therefore, we must have that J plus ket λJ . That what is J . This is a null ket that J is the maximum value of m . I operate this by J minus and I had shown that you remember that J minus J plus. We had evaluated this. This is J_x minus $i J_y$ multiplied by J_x plus $i J_y$. So, this is J_x square plus J_y square and this is plus $i J_x J_y$ minus $J_y J_x$. So, that is $i J_z$. So, this is equal to J square minus J_z square minus J_z . So, I operate J minus J plus this thing.

So, you must understand the equation that J plus ket λJ is a null ket. This is not a null ket. Of course, this is not a null operator, but operating on a ket gives you a null ket. So, I operate this by J minus. So, I get J square minus J_z square minus J_z operating on λJ must be a null ket. This is a non null operator, this is a non null ket, but this operating on this is a null ket.

Now, this is an Eigen ket of J square. So, the J square operating on that is λ . The Eigen value of J_z is J of J_z square. If J_z ket λJ is equal to J λJ , if this is the Eigen value of J_z and therefore, if I operate this again by J_z , so J_z square λJ will be equal to J operating J_z λJ . So, that will be J square λJ . So, this will be λ minus J square minus J operating on λJ must be a null ket.

Now, here this quantity is a number, this is an operator. It is a non null operator, but this is a non null ket by definition. So, this number must be 0. So, lambda minus J square minus J. This is really beautiful is 0 which gives me the beautiful result that lambda must be equal to J square plus J or J into J plus 1.

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$$\begin{aligned}
 J_- &= J_x - iJ_y \\
 [J^2, J_-] &= [J^2, J_x] - i[J^2, J_y] = 0 \quad J^2 |\lambda m\rangle = \lambda |\lambda m\rangle \\
 J^2 J_- |\lambda m\rangle &= J_- J^2 |\lambda m\rangle \\
 J^2 \{ J_- |\lambda m\rangle \} &= \lambda \{ J_- |\lambda m\rangle \} \\
 \underbrace{J^2 \{ J_- |\lambda m\rangle \}}_{|A\rangle} &= \lambda \{ J_- |\lambda m\rangle \} \\
 J^2 \{ J_- |A\rangle \} &= \lambda \{ J_- |A\rangle \} \\
 |B\rangle &= J_- |A\rangle = J_- J_- |\lambda m\rangle
 \end{aligned}$$

Now, if you have not been able to follow the algebra, let me do it all over again for the J minus operator. So, you have the J minus operator. Let us forget about everything, but keep this result somewhere. So, the J minus operator is equal to J x minus i J y. So, J square commutes with J x, J square commutes with J y. So, J square commutes with J minus. So, J square, J minus is equal to J square, J x minus i times J square, J y. Both are 0, both vanish. So, this is zero.

So, you have J square J minus is equal to J minus J square. I operate both sides by lambda m. So, we have this is equal to lambda. So, this will be lambda J minus ket lambda m. So, you will have J square J minus. Thus, if ket lambda m is an Eigen ket of the operator J square belonging to the Eigen value lambda as expressed by this equation, then J minus ket lambda m is also an Eigen ket J square belonging to the same Eigen value lambda. Similarly, let me denote this by say, ket a. So, then, we can write down J square exactly in a similar way J square J minus ket a will be equal to lambda J minus ket a. So, if I denote this by ket B which is J minus ket A which is equal to J minus J minus ket lambda m, then we say.

Listen, please listen carefully. If ket lambda m is an Eigen ket of J square belonging to the Eigen value lambda, then J minus J minus is also an Eigen ket of this same operator J square belonging to the same Eigen value lambda provided of course, a is not a null ket and J minus a and b is also not a null ket. Therefore, once again we must expect the genre is and this is an Eigen ket, this is an Eigen ket and this is also an Eigen ket. Then, we consider the commutation relation of J z.

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$$\begin{aligned}
 [J_z, J_-] &= [J_z, J_x - iJ_y] = [J_z, J_x] - i[J_z, J_y] \\
 &= 0 - i(-iJ_x - J_y) = -J_x - J_y = -J_- \quad J_z |\lambda m\rangle = m |\lambda m\rangle \\
 J_z J_- |\lambda m\rangle &= J_- J_z |\lambda m\rangle - J_- |\lambda m\rangle \\
 J_z \{ J_- |\lambda m\rangle \} &= (m-1) \{ J_- |\lambda m\rangle \} \\
 J_z \{ J_- J_- |\lambda m\rangle \} &= (m-2) \{ J_- J_- |\lambda m\rangle \}
 \end{aligned}$$

So, let me consider the commutation relation of J z of J minus. This is J z, J x J J minus is J x minus i J y. So, this is J z, J y. So, J z, J x is equal to i J y. I J y and J z J y will be equal to minus. If you look up the first, we wrote down that J z, J y was minus i J x. So, minus minus plus and i square is minus. So, this becomes minus J x. So, if I take minus sign outside. So, I get J x minus i J y, sorry i J y. So, this will be minus J minus. So, I had J z J minus. Please leave some space minus J z J minus J z, leave some space is equal to minus J minus.

Now, what I will do is operate this on lambda m, operate this on lambda m and operate this on lambda m and you have here J z ket lambda m is m lambda m. So, I get this first term becomes J z J minus ket lambda m and if you take it to the other side, then it will become m minus 1 J minus ket lambda m. Therefore, if ket lambda m is an Eigen ket of the operator J z, if ket lambda m is an operator is an Eigen ket of the operator J z

belonging to the Eigen value m , then $J_- | \lambda m \rangle$ is also an Eigen ket of the same operator J_z .

Now, belonging to 1 less Eigen value $m - 1$ provided of course, this is not a null ket and I can continue this again, so J_- operating on $J_- | \lambda m \rangle$. I hope you can now immediately say that this must be $m - 2$ J_- $J_- | \lambda m \rangle$.

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$| \lambda m \rangle$ is a ket of $J^2(\lambda)$ & $J_z(m)$
 $J_- | \lambda m \rangle$ is also a " ket of $J^2(\lambda)$ & $J_z(m-1)$
 $J_- J_- | \lambda m \rangle$ ket of $J^2(\lambda)$ & $J_z(m-2)$
 \vdots
 \vdots
 $\lambda \geq m^2$
 $m_{\min} = J'$
 $| \lambda J' \rangle$ is not a null ket
 $J_+ J_- | \lambda J' \rangle = 0$
 $J_+ J_- = (J_x + i J_y)(J_x - i J_y)$

So, therefore, if ket, let me summarize If ket λm is a simultaneous Eigen ket of J^2 belonging to the Eigen value λ and J_z belonging to the Eigen value m , then $J_+ | \lambda m \rangle$ is also a simultaneous Eigen ket of J^2 belonging to the same Eigen value λ and of J_z belonging to the Eigen value $m + 1$. No, I will put minus sign here, but now this will become $m - 1$ provided this is not a null ket and then, $J_- | \lambda m \rangle$ will also be a simultaneous Eigen ket of J^2 belonging to the same Eigen value λ , but of J_z belonging to the Eigen value $m - 1$ and so on provided we do not hit a null ket, but I have. So, this is an arithmetical preparation $m, m - 1, m - 2, m - 3$. If it goes on indefinitely, then this quantity will become much greater than λ and we have the inequality that λ must be greater than m^2 .

So, there must be a minimum value of m and let that minimum value be equal to J' . That means that ket $\lambda J'$ is not a null ket. However, $J_- | \lambda J' \rangle$ is a null ket. Even then J' is not a minimum value. So, once again this is not

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So, this is equal to $J_x^2 + J_y^2 + J_z^2 - J^2$ and then, you will have $J_x J_y - J_y J_x$. So, this is $i \hbar J_z$. Actually, $i \hbar$ cross J_z . So, this will be $i^2 \hbar^2 J_z^2$, so plus. So, this will become $J_x^2 + J_y^2 + J_z^2 - J^2 + i \hbar^2 J_z^2$. So, this is my J^2 minus, but we are seeing that J^2 plus $i \hbar^2 J_z^2$ operating on J^2 prime. So, this I operate on ket. This whole thing I operate this on J^2 prime. So, this will be a null ket. So, J^2 operating on this, this is an Eigen ket of J^2 . So, this will become J^2 minus J^2 prime plus J^2 prime operating on J^2 prime must be 0.

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Handwritten notes on a blue grid background showing the derivation of the minimum value of m . The notes include the following equations and text:

- $m_{\max} = j$
- $m_{\min} = j'$
- $m_{\min} = -j$
- $\lambda = j(j+1)$
- $\lambda = j'(j'-1)$
- $j'(j'-1) = j(j+1)$
- $j' = -j$ (boxed)
- $j = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots$
- $j = \frac{1}{2}, 1, \frac{3}{2}, 2, \dots$
- A vertical list of values: $+j, \dots, -j, 0, -1$

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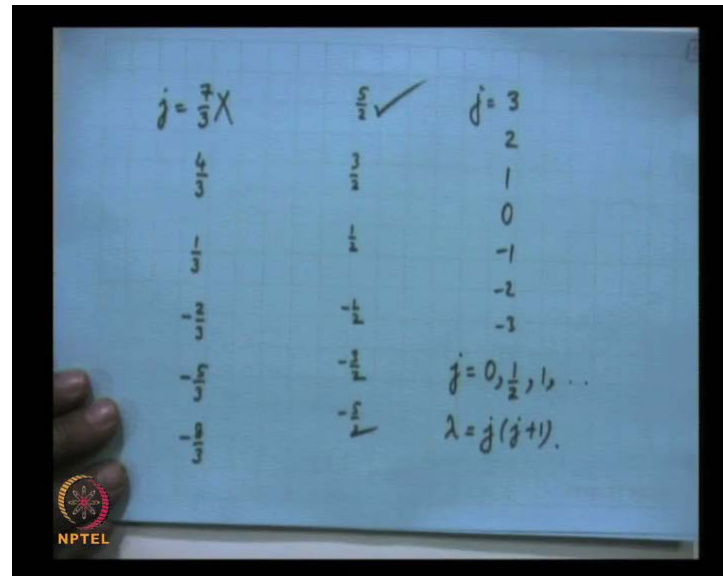
So, please remember two things, ket λm is a simultaneous Eigen ket of J^2 and J_z . m is the Eigen value of J_z and λ is the Eigen value of J^2 . First, we derive that λ must be greater than or equal to m^2 and we said that what the maximum value is. So, we said that since this, because of this inequality the maximum value of this is J and we found that λ must be equal to $J(J+1)$.

Then, we use ladder operators J_{\pm} . We found that m must have a minimum value and let that be J' and we just now found that λ must be equal to J' into $J' - 1$. So, J' into $J' - 1$ is equal to J into $J + 1$ and J' . This is a quadratic equation in J' , but since $J' < J$, so the root is J' is equal to $-J$ because if I put $-J$ here, so you will get $-J$ into $-J - 1$. So, that is J into $J + 1$.

So, the solution of this equation is J' must be equal to $-J$. So, the maximum value is J . The minimum value is $-J$ and now, see that I start with a value J and I go down in steps of 1 $J - 1$, $J - 2$, $J - 3$, $J - 4$, $J - 5$, $J - 6$ and then, I must hit $-J$. That is only possible if J itself is equal to $0, \frac{1}{2}, 1, \frac{3}{2}, 2$ etcetera not for any other number because we have said that if m is an Eigen value, then $m + 1$ is also an Eigen value. If m is an Eigen value, $m - 1$ is an Eigen value, it can hop in units of 1 and it has to land up on $-J$. Otherwise, you will never get a null ket and therefore, as you start with J , you hop in units of 1 and you go to $-J$. That will only be possible if J takes these values. So, if you have J is equal to 1, then you start with

1 decrease by 0. They are the m values. You start with J is equal to 3 by 2, so you have 3 2 half minus half and minus 3 by 2.

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Let me say that I take a, let us suppose I have J is equal to 7 by 3 to 7 by 3 minus 1 is 4 by 3, 4 by 3 minus 1 is 1 by 3, 1 by 3 minus 1 is minus 2 by 3 minus 2 by 3 minus 1 is minus 5 by 3 and then, minus 8 by 3 and so on. So you have not hit minus 7 by 3 to 7 by 3 is not possible, but if you use 5 by 2, so 5 by 2 minus 1 is 3 by 2, 3 by 2 is half minus, half minus 3 by 2 and minus 5 by 2. So, that is possible. So, J is equal to 3, then 2, 1, 0, minus 1, minus 2, minus 3.

So, we get the result that the value of J can only take 0 to half, 1 and lambda is equal to J into J plus 1. These are the Eigen values of J square operator. Now, you see a remarkable result that we have obtained. Earlier, we had said we had used differential operator representation.

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$$L^2 Y_{lm}(\theta, \phi) = l(l+1)\hbar^2 Y_{lm}(\theta, \phi)$$

$$L_z Y_{lm}(\theta, \phi) = m\hbar Y_{lm}(\theta, \phi)$$

$$l = 0, 1, 2, \dots$$

$$J^2 |\lambda, m\rangle = j(j+1)\hbar^2 |\lambda, m\rangle$$

$$J_z |\lambda, m\rangle = m\hbar |\lambda, m\rangle$$

$$J^2 |j, m\rangle = j(j+1)\hbar^2 |j, m\rangle \quad j = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots$$

$$J_z |j, m\rangle = m\hbar |j, m\rangle \quad m = -j \text{ to } +j$$

We had found that the Eigen functions are the spherical harmonics and the Eigen values were $l^2 Y_{lm}(\theta, \phi)$ were equal to $l(l+1)\hbar^2 Y_{lm}(\theta, \phi)$. In fact, these were simultaneous Eigen functions of L^2 and L_z , but here l is equal to 0, 1, 2, 3, but here I will get that $J^2 |\lambda, m\rangle = j(j+1)\hbar^2 |\lambda, m\rangle$. This is equal to $J^2 |\lambda, m\rangle = j(j+1)\hbar^2 |\lambda, m\rangle$ and $J_z |\lambda, m\rangle = m\hbar |\lambda, m\rangle$ is equal to $m\hbar$ cross, but the difference between the two is that here J can take half integral values.

As we know that the electron and the proton and the neutron are endowed with a intrinsic angular momentum which is half \hbar cross and that is predicted by quantum mechanics that are now derived from quantum mechanics and just by using commutation relations and nothing else. So, we finally obtain actually instead of λ , the convention is to label the Eigen kets as $|j, m\rangle$. This is not the Eigen value. The Eigen value is really λ , but the convention is $J^2 |j, m\rangle = j(j+1)\hbar^2 |j, m\rangle$ and $J_z |j, m\rangle = m\hbar |j, m\rangle$ I have put. So, these and these values of J that we can take are 0, half, 1, 3 by 2 etcetera and m can take from minus J to plus J .

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$$j_{\min} = -j \quad \lambda = j(j+1) \quad j - j$$

$$J_+ |j, m\rangle = C_+ |j, m+1\rangle$$

So, I had shown that J_{\min} which is the minus value of which is the minimum value of J was equal to minus J . Therefore, these are the m values for a given value of J . The m values go from J to minus J . I had now proved that just by using commutation relation. So, all the previous results are proved that because the differential operators also satisfy the same commutation relations.

So, therefore, I have proved that for a given value of J , the m values go from J to minus J . Now, let me do now one more algebra and that is we have shown that $J_+ |j, m\rangle$ is an Eigen ket of J^2 and J_z , then $J_+ |j, m\rangle$ is also an Eigen ket, but now belonging to the Eigen, same Eigen value λ which is $J(J+1)$ λ is equal to $J(J+1)$, but m value is raised. So, let me write this down and this must be a multiple of $|j, m+1\rangle$. May be you put a , here. What we must always remember on the right hand side, I may write it down that $J^2 |j, m\rangle$.

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$$J^2 |j m\rangle = j(j+1) |j m\rangle$$

$$J_z |j m\rangle = m |j m\rangle$$

$$|p\rangle = J_+ |j m\rangle = c_+ |j, m+1\rangle$$

$$\langle p| = \langle j m| J_- = c_+^* \langle j, m+1|$$

$$\langle p|p\rangle = \langle j m| J_- J_+ |j m\rangle = |c_+|^2$$

$$J_- J_+ = (J_x - i J_y)(J_x + i J_y)$$

$$= J_x^2 + J_y^2 + i(J_z) = J^2 - J_z^2 - J_z$$

$\langle j' m'| j m\rangle = \delta_{jj'} \delta_{mm'}$
 $\langle j m| j m\rangle = 1$
 $J_+ = J_x + i J_y$
 $J_- = J_x - i J_y$

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Let me write it on a fresh piece of paper. J^2 ket $J m$ is equal to J into J plus 1 \hbar cross square is there, but we assume that in a system of units where \hbar cross is 1 and J_z $J m$ is equal to m \hbar cross is 1. We are working in a system of units of this.

Now, we say that J_+ $J m$ and these are orthonormal ket. These are simultaneous Eigen ket, but they are orthonormal, that is J prime, m prime, $J m$ here. All delta, $J J$ prime delta, m m prime, they are all orthonormal. Both of them and this is normalized, that is $J m$ $J m$ is 1. We assume that because they belong to (\odot) operators, so J_+ ket $J m$ is therefore a multiple of this ket $J m$ plus 1.

So, let this be c_+ , let this be ket p . Then, $\langle p|$ is equal to $\langle j m| J_+$ bar. This will be equal to this is just a number, so $c_+^* \langle j m|$ plus 1. So, I now write $\langle p|$ ket p , this is equal to $\langle j m| J_+$ bar is J_+ plus is equal to J_x plus $i J_y$. So, J_+ bar is J_- because this will be J_x minus $i J_y$. So, this is $J_- J_+$ $J m$ is equal to $c_+^* \langle j m|$ mod square and this is 1 this times this is 1. So, once again $J_- J_+$ is equal to J_x minus. This you must remember by now J_x plus $i J_y$. So, this is J_x^2 plus J_y^2 plus i times $J_x J_y$ minus $J_y J_x$. So, that is i times J_z . So, this is J_x^2 plus J_y^2 is J^2 minus J_z^2 minus J_z . You have derived this earlier.

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$$\langle P | P \rangle = \langle j m | J^2 - J_z^2 - J_x^2 | j m \rangle = |C_+|^2$$

$$(\hbar^2(j(j-1) - m^2 - m)) \langle j m | j m \rangle = |C_+|^2$$

$$(\hbar^2(j(j+1) - m^2 - m)) \langle j m | j m \rangle = |C_+|^2$$

$$C_+ = \sqrt{(j-m)(j+m+1)}$$

So, we have therefore, bra p ket p is $J_m J^2 - J_z^2 - J_x^2$ is $J^2 - J_z^2 - J_x^2$ operating on J_m is equal to $\hbar^2 C_+^2$. So, this is operating on this, this is an Eigen ket of J^2 , J_z^2 and J_x . So, J^2 that will be $J(J+1)\hbar^2$ minus $m^2\hbar^2$ minus $m\hbar^2$. This is C_+^2 . I am sorry this will be $J(J+1)$ minus m^2 minus m . This is not the Eigen value. The Eigen value is $J(J+1)$ as we had done here $J(J+1)$.

So, this will be $J(J+1) - m^2 - m$ bra J_m ket J_m which is 1 because they are orthonormal ket. This is equal to C_+^2 . So, this is equal to 1 and therefore, we obtain C_+ is equal to the under root of that. So, this will be $\sqrt{J(J+1) - m^2 - m}$. Therefore, J_+ ket λ_m , I am sorry let me do it again.

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$$J_+ |j, m\rangle = \sqrt{(j-m)(j+m+1)} |j, m+1\rangle$$

$$J_- |j, m\rangle = \sqrt{(j+m)(j-m+1)} |j, m-1\rangle$$

$$m=j \quad |j, j\rangle \quad \left\{ \begin{array}{l} J_+ |j, j\rangle = 0 \\ J_- |j, -j\rangle = 0 \end{array} \right.$$

$$j = \frac{1}{2} \quad ; \quad m = \frac{1}{2}, -\frac{1}{2}$$

$$\begin{cases} |1\rangle = |\frac{1}{2}, \frac{1}{2}\rangle \\ |2\rangle = |\frac{1}{2}, -\frac{1}{2}\rangle \end{cases}$$

$$J^2 |1\rangle = \frac{3}{4} \hbar^2 |1\rangle$$

$$P_{11} = \langle 1 | P | 1 \rangle$$

$$P_{12} = \langle 1 | P | 2 \rangle$$

$$P_{22}$$

So, therefore, we will have $J_+ |j, m\rangle$ will be equal to $c_+ |j, m+1\rangle$ and that is square root of this. You have to remember $J_- |j, m\rangle = \sqrt{(j+m)(j-m+1)} |j, m-1\rangle$. These are ladder operators. It raises the value of m to $m+1$ and similarly, I leave it as an exercise exactly in a similar way you show that the operator J_- is proportional to $J_- |j, m\rangle = \sqrt{(j+m)(j-m+1)} |j, m-1\rangle$. These two relations we had to remember and we see that when m is equal to j , this is 0. So, $J_+ |j, j\rangle$ is the maximum value of m .

So, $|j, j\rangle$ is a non null ket. It is a well defined ket but, $J_+ |j, j\rangle$ is a null ket. Similarly, $|j, -j\rangle$ is a non null ket, but $J_- |j, -j\rangle$ will be a null ket. That comes out automatically from here. Now, let me do an example. Let me consider the case where this spin half j is equal to half. This is an extremely important problem. So, j is equal to half. So, m can be only half and half minus half, sorry m can be half and minus half. So, we can only have two states. State 1 which I denote as half, half. This is the value of j . So, the Eigen value of J^2 is $j(j+1) \hbar^2 = \frac{1}{2}(\frac{1}{2}+1) \hbar^2 = \frac{3}{4} \hbar^2$, that is $3/4 \hbar^2$.

So, $J^2 |1\rangle = \frac{3}{4} \hbar^2 |1\rangle$ and $J^2 |2\rangle = \frac{3}{4} \hbar^2 |2\rangle$. Now, so this I take as the base vectors. I make a representation. I will try to make a representation of the operators J_x, J_y, J_z . Now, when we make a representation, you have to first tell me what your base vectors are like in an I want to describe a vector. If I want to describe a vector and in terms of its component, I will first ask you tell me your x and y axis. So, this is the

base vectors. So, in terms of the base vectors 1 and 2, I will write the operators as a matrix. So, for example, since this is a two-dimensional space, there are only two base vectors. We may have 2 by 2 matrices.

If I had J is equal to 1 where m can take 1, 0 and minus 1, so there will be three independent vectors and it will span a three-dimensional space. We will have 3 by 3 matrices. So, of any operator, the matrix representation will be let us suppose of operator p , there will be four numbers p_{11} which is $\langle 1|p|1\rangle$, that is $\langle 1|p|2\rangle$ and $\langle 2|p|1\rangle$ and $\langle 2|p|2\rangle$. These are the four numbers which will give you a representation of the operator. So, let me take these as my base vectors.

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Handwritten mathematical derivations for angular momentum operators J_x , J_y , and J_z acting on basis states $|j, m\rangle$.

Given $j = \frac{1}{2}$, $m = +\frac{1}{2}$ and $-\frac{1}{2}$.

Basis states:

$$|1\rangle = |\frac{1}{2}, \frac{1}{2}\rangle$$

$$|2\rangle = |\frac{1}{2}, -\frac{1}{2}\rangle$$

Operators and their actions:

$$J^2 |1\rangle = \frac{3}{4} \hbar^2 |1\rangle$$

$$J^2 |2\rangle = \frac{3}{4} \hbar^2 |2\rangle$$

$$J_z |1\rangle = \frac{1}{2} \hbar |1\rangle$$

$$J_z |2\rangle = -\frac{1}{2} \hbar |2\rangle$$

Commutator relations and specific actions:

$$J_+ |1\rangle = J_+ |\frac{1}{2}, \frac{1}{2}\rangle = 0$$

$$J_- |1\rangle = J_- |\frac{1}{2}, \frac{1}{2}\rangle = |2\rangle$$

$$J_+ |2\rangle = |1\rangle$$

$$J_- |2\rangle = 0$$

General formulas for raising and lowering operators:

$$J_- |j, m\rangle = \sqrt{(j+m)(j-m+1)} \hbar |j, m-1\rangle$$

$$J_+ |j, m\rangle = \sqrt{(j-m)(j+m+1)} \hbar |j, m+1\rangle$$

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So, I will take ket 1. So, I consider J is equal to half. Then, m is equal to plus half and minus half. So, the base kets are half half and the second is half minus half. These are simultaneous Eigen kets of J^2 and J_z , that is J^2 ket 1 is equal to half into half plus 1. That is $\frac{3}{4} \hbar^2$ cross square ket 1. J^2 ket 2 is also $\frac{3}{4} \hbar^2$ cross square ket 2 and J_z ket 1 is equal to this is the Eigen value of J_z half \hbar cross ket 1 and J_z ket 2 is equal to minus half \hbar cross ket 2.

Now, let me write down what is J_+ ket 1. J_+ ket 1 will be equal to J_+ half half. So, this will be square root of $J - m$. So, this is 0. So, remember that J_+ ket $J - m$ was let me write it below. J_+ ket $J - m$ was under root of $J - m$ into $J + m + 1$ of $m + 1$. So this is 0 and J_- ket 1 will be J_- half half. So, this will be

please see this J minus half h will be under root of J plus m into J minus m plus half h J, m minus 1. So, J is half m is half. So, half plus half is 1 square root of 1 half minus half is 0, 1. So, this will be half minus half. So, that is ket 2.

Now, let me write down what is J plus ket 2. So, here J is half m is minus half. So, here J minus m half minus minus half half plus half 1 half minus half is 0. So, this is equal to 1 and J minus ket 2 J minus ket 2 means m will go down, but that cannot be minus 3 by 2. So, this will be a null ket. Ok, so we all have the relations now.

(Refer Slide Time: 53:02)

Handwritten notes on a whiteboard showing the derivation of J^2 and J_z matrices for a spin-1 system.

Left side (Matrix elements of J^2):

$$J^2|1\rangle = \frac{3}{4}\hbar^2|1\rangle$$

$$J^2|2\rangle = \frac{3}{4}\hbar^2|2\rangle$$

$$(J^2)_{11} = \langle 1|J^2|1\rangle = \frac{3}{4}\hbar^2$$

$$(J^2)_{12} = \langle 1|J^2|2\rangle = 0$$

$$(J^2)_{21} = \langle 2|J^2|1\rangle = 0$$

$$(J^2)_{22} = \frac{3}{4}\hbar^2$$

Right side (Matrix representation of J^2 and J_z):

$$(J^2) = \frac{3}{4}\hbar^2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$J_z|1\rangle = \frac{1}{2}\hbar|1\rangle$$

$$J_z|2\rangle = -\frac{1}{2}\hbar|2\rangle$$

$$(J_z)_{11} = \frac{1}{2}\hbar$$

$$(J_z)_{21} = \langle 2|1\rangle \cdot \frac{1}{2}\hbar = 0$$

$$(J_z) = \frac{1}{2}\hbar \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

An arrow labeled σ_z points to the matrix (J_z) .

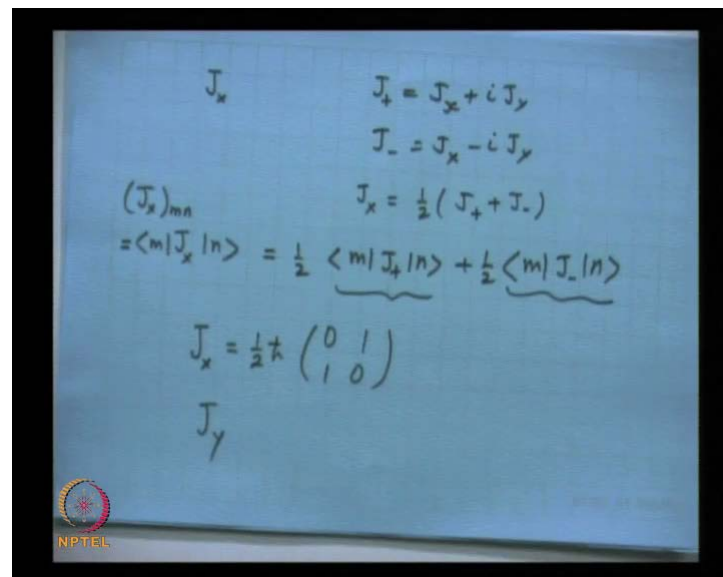
We can write down the matrices. First, let me write down for J square. So, we had first we wrote down J square ket 1 is equal to 3 by 4 h cross square ket 1. J square ket 2 is equal to 3 by 4 h cross square ket 2, so J square. There will be four elements 1 1, that is bra 1 J square ket 1, so bra 1. So, this will be equal to 3 by 4 h cross square. Then, J square 1 2 will be bra 1 J square ket 2.

So, this will be bra 1 ket 2. So, that will be 0. So, J square 2 1 will be 2 J square ket 1. That will be also 0 and J square 2 2 will be equal to, I leave it as an exercise. So, 3 by 4 h cross. So, the matrix representation of the operator J square is 3 by 4 h cross square 1 0 1. Similarly, I have J z ket 1 is equal to half h cross ket 1 because m value is half and J z ket 2 is equal to minus half h cross ket 2. So, again since these are Eigen kets, so if you write J z 1 1, so this will be multiplied by bra 1. So, this is half h cross, but J z 1 2 or 2 1

J_z will be $\hbar m$. So, this will be multiplied by \hbar . So, this will be $\hbar m$ and J_z will be minus half \hbar .

So, the operator representation, the matrix representation of J_z will be \hbar cross $\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$. This is the sigma z, the Pauli matrix. We have derived that what are the Eigen values of J^2 at this matrix 1 . Therefore, what are the Eigen values of J^2 ? That is just 3 by $4 \hbar^2$. What are the Eigen values of J_z ? The Eigen value of these matrix are 1 and minus 1 . So, the Eigen values of J_z are half \hbar and minus half \hbar .

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Handwritten derivation on a blue background:

$$J_x = \frac{1}{2} (J_+ + J_-)$$

$$(J_x)_{mn} = \langle m | J_x | n \rangle = \frac{1}{2} \langle m | J_+ | n \rangle + \frac{1}{2} \langle m | J_- | n \rangle$$

$$J_x = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Below the matrix, J_y is written.

So, I have just given you a hint as to I mean, rest we will do in our next lecture. I want now to make a representation of J_x . So, we have $J_+ = J_x + i J_y$ and $J_- = J_x - i J_y$. So, I add them up and get $J_x = \frac{1}{2} (J_+ + J_-)$. So, J_x bra m ket n , this is the m, n th matrix element is given by this. So, this will be half \hbar plus J_+ ket n plus half \hbar minus J_- ket n and I leave this as an exercise for you. We will try to do this next time, but if you can work this out, this will be very good and we will, you should find that the matrix representation of J_x will come out to be half \hbar cross $\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ and similarly, for J_y .

So, with this we conclude today's lecture. What we have achieved is that we have been able to obtain the Eigen value spectrum, simultaneous Eigen kets of the operator, J^2

square and J_z . In our next lecture, we will complete the analysis for obtaining the matrix representation of J^2 , J_x , J_y and J_z . Thank you.

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