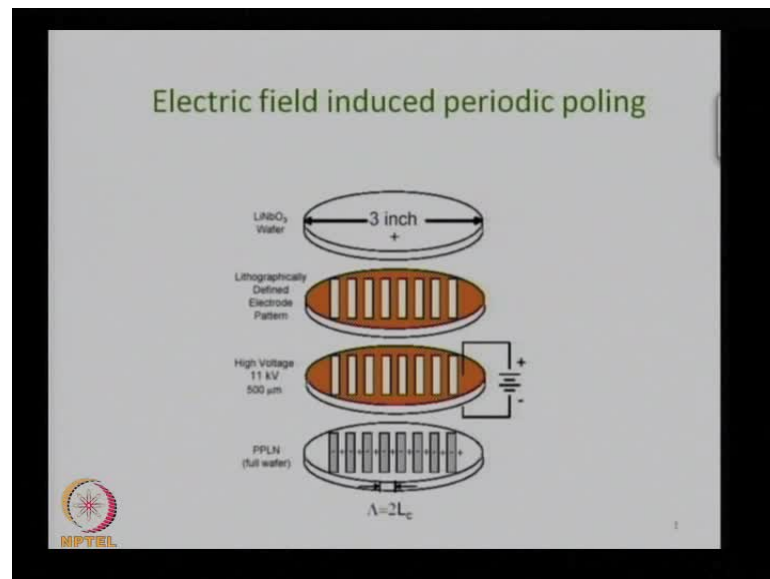


Quantum Electronics
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Module No. # 03
Second Order Effects
Lecture No. # 11
Non - Linear Optic

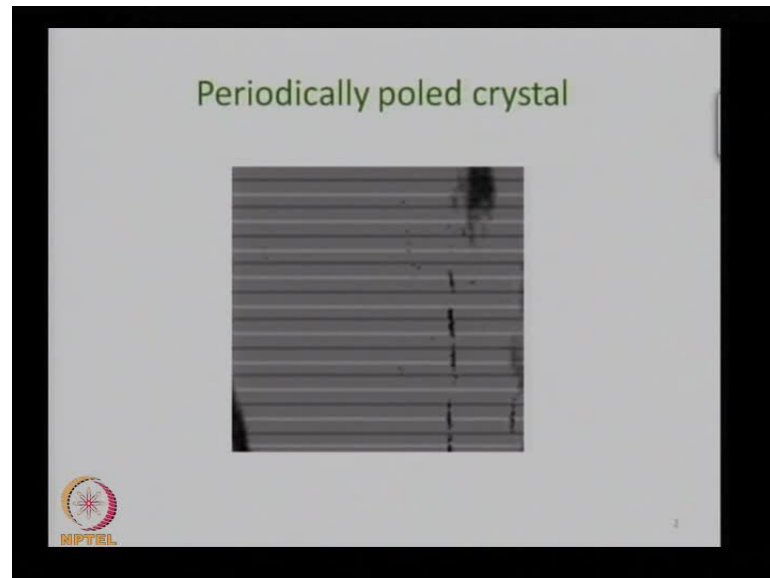
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Before we continue with discussion of the mathematical formulation, I will show you some slides on this poling. How to generate periodically domain reverse structures in lithium niobate? So, as I was mentioning to you, lithium niobate is a ferroelectric material and the direction of polarization. There is a permanent polarization in the material. So, one takes a lithium niobate wafer, which is the plus is written there, that means the z axis is pointing up, its optic axis is pointing up, and then as I was telling you one lithographically defines an electrode pattern; and then applies a high voltage typically about as is written here, 11 kilovolts and 500 micrometers. So, that is 22 kilovolts per millimeter the breakdown strength is much higher.

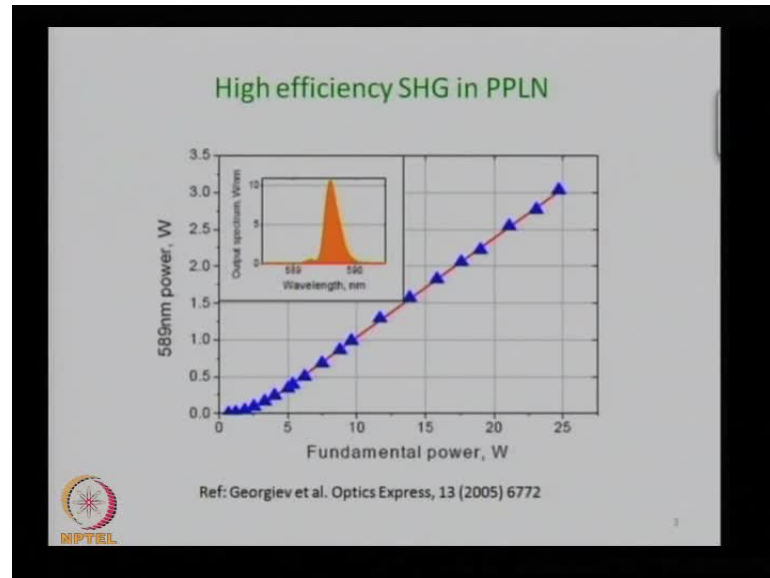
So, what happens is where the electric fields are applied? Domain switch from vertically up to vertically down. So, you get a periodic poling of the crystal. And you can actually, have a more complex electrode structures to achieve this.

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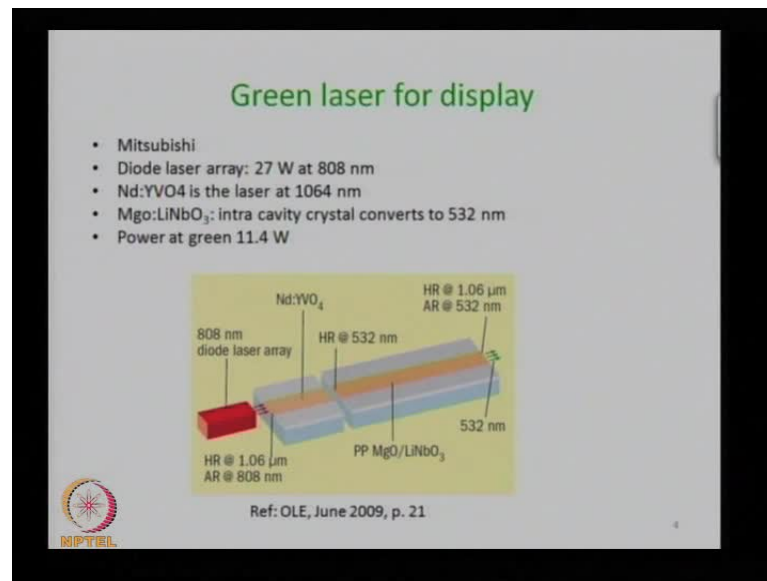
So, for example - let me show you, the top surface of the substrate after poling. And what happens, is you can actually show the poling, by etching the z plus and z minus surfaces get etched at different rates. So, this structure that you see is actually, the periodically poled substrate; with this in one of them and the other one is this, z minus z plus z minus z plus. Its periodic you can, you can, see it is very clearly periodic with a certain period. And the propagation direction is like this vertical perpendicular. So, these domains are like this and you are propagating perpendicular to the domains. So, you get a periodic poling which is possible in lithium niobate by periodically applying electric fields and so, this is now a standard technique, which is used to pole crystals of lithium niobate.

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Let me show you some examples, of high efficiency generation in using periodically poled substrates; here is a report from a journal optics express, which is a journal published by the optical society of America and you can see here for about 25 watts of fundamental power, you can generate, something like 3 watts of second harmonic power. So, that is a good efficiency and as you can see that shows the change of the fundamental power versus the second harmonic power versus fundamental power. And this is just to show you that you can achieve very high efficiencies by using periodically poled crystals.

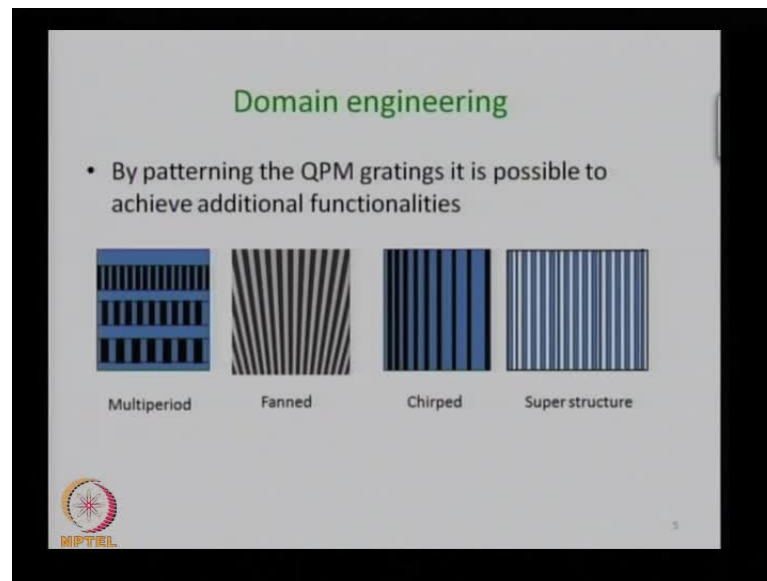
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In fact, in a recent publication in 2009 June, they have actually shown, you can generate green light using this second harmonic generation. In fact, there are some laser pointers-green laser pointers they have actually a red laser, a second harmonic generator within the green laser pointer.

So, you can actually convert 808 nanometers to 404 nanometers actually it is 808 nanometers is used to pump a lutetium. Why we opt for crystal, that generate 1064 nanometer. The 1064 nanometer then is converted to 532 nanometers using a magnesium oxide doped lithium niobate crystal and you can see, the powers coming out at the green 11 watts. The argon laser, which is a huge laser in the laboratory that generates a few watts of power, so that is a huge power and these are now being seen for application for - I mean - laser display, you have you need to have the three primary colors and this is one of the primary colors green laser.

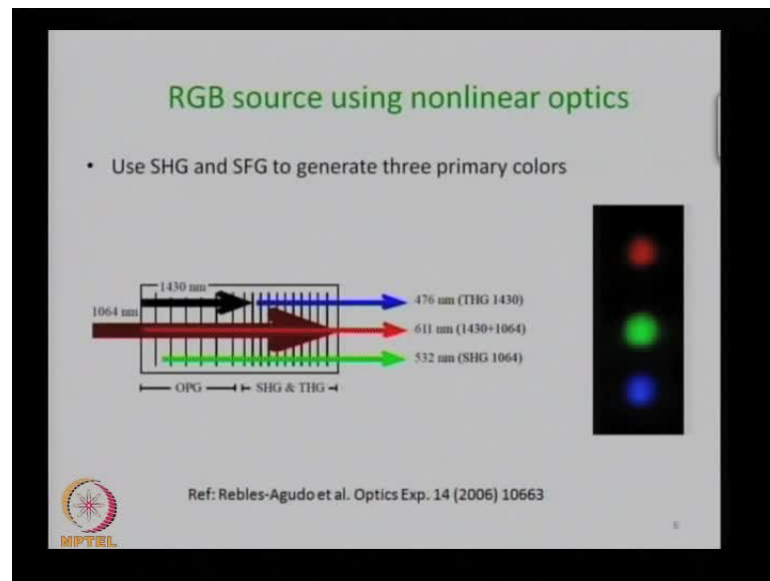
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Let me, show you some interesting configurations, you can do by making crystals with multi periods as you have one period, then second period, third period. So, depending on which period you use, you can achieve phase matching for different frequencies, it is remembered that the phase matching period, depends on the frequency of the fundamental waves. So, you can actually make a crystal with multiple periods and use this multiple periods to achieve phase matching at different fundamental frequencies or you can actually, fan the grating like this. So, depending on the vertical position where you are coming. The period that will be encountered by the light will be different.

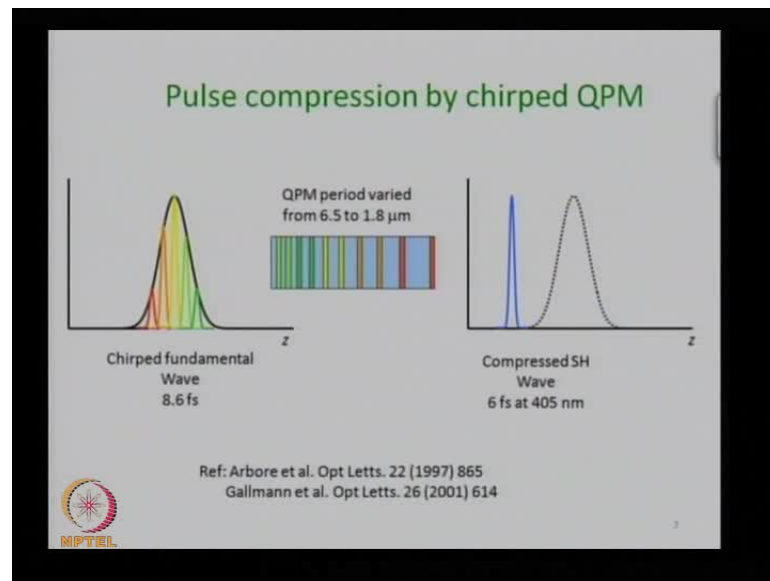
So, you can achieve tunable phase matching essentially. So, you can actually move the crystal and depending on where you are focusing the laser beam, you can have different phase matching condition satisfied; you can have a chirped, where the spatial period is changing with position. You can have super structure. So, there is a one period, which is the internal period and then there is another period, which is this is being repeated again and again the super period gratings. So, you can actually do a lot of interesting engineering to achieve different functions multifunctional, as you will see later on for difference frequency generation the periods are different. So, you can do second harmonic generation and then difference frequency etcetera, in a single crystal.

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So here is an example of, **how** in a single crystal you launch 1064 nanometer light, use optical parametric generation to generate a new frequency. Then use second harmonic and third harmonic generation to generate multiple wavelengths simultaneously. This is a single crystal in which you generate multiple domains of different periods, each period is meant to achieve different nonlinear effects essentially. So, finally you can generate all the three primary colors. So, there is a lot of activity in using this nonlinear interaction to achieve new wavelengths, which are otherwise not possible, because conventional lasers have a fixed wavelengths which are emitted by the laser.

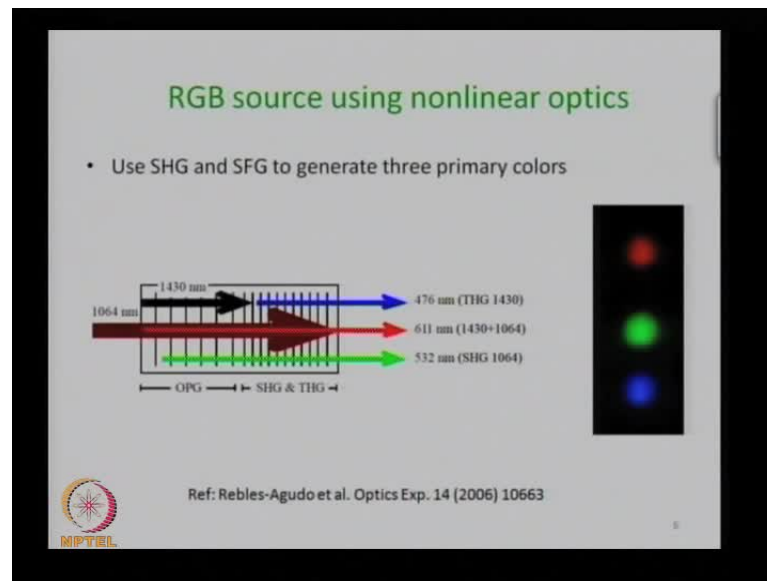
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You can actually do more interesting things of compressing the pulse, where it generates a second harmonic and so here is a very interesting period paper in which you have a period varying from 6.5 to 1.8 microns. And that you can show is actually a fundamental waves, comes in with a duration of 8.6 femtoseconds; and it gets compressed into second harmonic to about 6 femtoseconds; see you can do all of change that is by essentially engineering the domains of the crystal, which is not possible if you have birefringence phase matching. So, quasi phase matching gives you a lot of flexibility in achieving multi-functionalities.

Sir, in the figure, that you showed that different colors were generated, three colors are generated. So, there what happens is that first the incident wave comes and its second harmonic successor gets generated? Then that second harmonic of that second harmonic is generated.

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No, you are putting in 1064 nanometers; you are also putting at 1430 nanometers. So, it generates a difference frequency that is the optical parametric generation. So, we will discuss this difference frequency generation little later. Then there is the second harmonic of 1064 is 532 nanometers, which comes out. Then the 1430 and 1064 mix to generate a sum frequency generation at 611 nanometers and then, you have a 476 nanometers coming from the 1430 nanometer.

So, you actually multiple processes nonlinear process all depending on χ^2 . Please remember, χ^2 process can be second harmonic sum frequency, difference frequency. So, the generated frequency is so powerful. Now, it is not very low power. Now, the efficiency is quite high that you can use the generated light to further undergo; further, nonlinear processes for which you require different periods and so the periods can be different at different position of crystal to compensate, for the phased mismatch for that interaction process.

Is it possible that the efficiency for second harmonic generation is that high, I mean if you use chirped gratings, these chirped grates, then is it possible that the efficiency for the first, second harmonic generation in so high; that further we can use a period that second harmonic of that second harmonic can be generated, I mean the efficiency **can be...**

no that is not chirped, I will have to generate two different gratings; one grating for fundamental second harmonic, another grating from second harmonic to the fourth harmonic of the fundamental .

2 omega to 4 omega

So we cannot do it in one?

No, because the period required are very different. One period is depends on $n \omega$ and $n 2 \omega$, the another depends on $n 2 \omega$ $n 4 \omega$ and because of strong dispersion, as you go to the higher frequencies, the periods required may be very different, you cannot chirp it so much, there is no point in chirping, because chirping I need for a to extend in the bandwidth of the interaction process.

But otherwise, I know the period to convert from λ_1 to λ_2 ω_1 to 2ω into 2ω . Once I convert ω into 2ω then I put another period to 2ω to 4ω .

So, you will send in ω and out will come 4ω , you would not, you would not, care what is happening inside? So, it is possible in general to have multiple processes taking place simultaneously.

Point was that we can incident a light, which is at a given frequency – low.

At a lower frequency, which is on the lower end of the visible line? So, we can generate two ω from that. And then we get a net four ω , from that so we can reach from violet to red. From the in that process so using

No, this is the other way round red to violet.

Red to violet,

In only one input,

Sure, this is what I said, it is possible to go from ω to 4ω , by going from ω into 2ω and then 2ω to 4ω .

Essentially it comes from the outputs of the 2 omega also.

Surely, it is coming out here 532 is coming out here and 532, which is generated can also be mixed with other wavelengths at input to generate new wavelengths, for that I need to calculate, what is the QPM required? period required, and I need to make that period in microscope.

So, what we will now do is a continue our discussion, on the second harmonic generation that we were doing the other day, we were actually discussing the situation where the efficiency could be very high that is when I do not neglect the pump depletion.

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$$\frac{dE_1}{dz} = i \frac{\omega d}{c n_1} E_2 E_1^*$$
$$\frac{dE_2}{dz} = i \frac{\omega d}{c n_2} E_1^2$$
$$E_1(z) = u_1 e^{i\phi_1}$$
$$E_2 = u_2 e^{i\phi_2}$$
$$\Delta k = 0$$
$$n_1 = n_2$$
$$\kappa = \frac{\omega d}{c n_1}$$

Remember, we started with these equations dE_1/dz is equal to $i \omega d / c n_1 E_2 E_1^*$ and dE_2/dz is equal to $i \omega d / c n_2 E_1^2$. So, I am assuming Δk is equal to 0, I am assuming Δk is equal to 0.

Anyway let me lighten it, I forgot to switch.

Is that the maximum brightness can you see?

So, we had assumed Δk is equal to 0 phase dispatched is equal to 0 and let me call this coefficient as kappa. So, because Δk is equal to 0 n_1 is equal to n_2 and let me call this as kappa is equal to $\omega d / c n_1$.

So, what we did was we substitute at E 1 of z is equal to u 1 exponential i phi 1 E 2 is equal to u 2 e to the power i phi 2 into these two equations, then where u 1 and e 2 u 2 and phi 1 and phi 2 are all real quantities. We can equate, the real imaginary parts on both sides and got the following four equations.

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$$\frac{du_1}{dz} = -\kappa u_1 u_2 \sin \theta$$

$$\frac{du_2}{dz} = \kappa u_1^2 \sin \theta$$

$$u_1 \frac{d\phi_1}{dz} = \kappa u_1 u_2 \cos \theta$$

$$u_2 \frac{d\phi_2}{dz} = \kappa u_1^2 \cos \theta$$

$$\theta = \phi_2 - 2\phi_1$$

$$\frac{d\theta}{dz} = \frac{d\phi_2}{dz} - 2 \frac{d\phi_1}{dz}$$

$$= \frac{\kappa u_1^2 \cos \theta}{u_2} - 2 \kappa u_1 \cos \theta$$

$$= \kappa \cos \theta \left(\frac{u_1^2}{u_2} - 2u_1 \right)$$

$$\frac{d}{dz} (\ln u_1^2 u_2) = \frac{1}{u_1^2 u_2} \left(u_1^2 \frac{du_2}{dz} + 2u_1 u_2 \frac{du_1}{dz} \right)$$

$$= -\frac{d}{dz} (\ln \cos \theta)$$

Let me rub these equations. So, what we got was du 1 by dz is equal to minus kappa u 1 u 2 sine theta du 2 by dz is equal to kappa u 1 square sine theta u 1 d phi 1 by dz is equal to kappa u 1 u 2 cos theta and u 2 d phi 2 by dz is equal to kappa u 1 square cos theta, where theta was defined as phi 2 minus 2 phi 1.

So, what I do is I substitute, the complex fields E 1 and E 2, which I write as u 1 times exponential i phi 1 and u 2 times exponential i phi 2 substitute into those equations and equate the real and imaginary parts on both sides; and I get instead of two complex equations I get four real equations now. What I will do is I can actually solve these equations, by the following procedure.

So, let me calculate, d theta by d E z this should be equal to d phi 2 by dz minus 2 d phi 1 by dz theta is phi 2 minus 2 phi 1. So, I substitute for d phi 2 by dz and d phi 1 by dz from here so, this is equal to kappa u 1 square by u 2 cos theta minus 2 kappa u 2 cos theta.

So, this is equal to $\kappa \cos \theta$ into u_1^2 by u_2^2 minus $2 u_1 u_2$. So, it is a mathematical trick that I am going to do, to just solve these equations. So, I get $d\theta$ by dz is so much. Now, let me calculate, so you can also calculate d by dz of $\log u_1^2 u_2^2$,

This can be calculated, this will be in terms of so... this is actually 1 by $u_1^2 u_2^2$ into $u_1^2 du_2$ by dz plus $2 u_1 u_2 du_1$ by dz and I have expressions for du_1 by dz and du_2 by dz here which I can substitute here and this is simple exercise, you can show that this is actually happens to be equal to minus d by dz of $\cos \theta$.

So, if you substitute for du_2 by dz and du_1 by dz here, you will see that the right hand side is essentially related to $d\theta$ $d \cos \theta$ by dz . Due to this...

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The whiteboard contains the following handwritten equations:

$$\frac{du_1}{dz} = -\kappa u_1 u_2 \sin \theta$$

$$\frac{du_2}{dz} = \kappa u_1^2 \sin \theta$$

$$\frac{d}{dz} \left(\ln(u_1^2 u_2^2 \cos \theta) \right) = 0$$

$$\Rightarrow u_1^2 u_2^2 \cos \theta = \text{Constant} = C$$

On the right side of the whiteboard, there are additional notes:

at $z=0$
 $E_2(z)=0$
 $\Rightarrow u_2=0$
 $\theta = +\frac{\pi}{2} = \frac{\pi}{2} - 2\phi$

So, this is actually minus 1 by $\cos \theta$ into minus sine θ $d\theta$ by dz . So, if you expand this you will get in terms of $d\theta$ by dz , you can use this equation and show that these two are equal, which implies essentially that d by dz of... So, if I bring it to the left hand side I will get d by dz of $\log u_1^2 u_2^2$ plus $\log \cos \theta$ is equal to 0 .

So, what does it imply \log of $u_1^2 u_2^2 \cos \theta$ is equal to 0 . So, this implies $u_1^2 u_2^2 \cos \theta$ is a constant of motion independent of z . As the waves propagate u_1 and u_2 and θ everything changes, but $u_1^2 u_2^2 \cos \theta$ remains constant.

It is just simple mathematical equation that I am using to get this condition. Now, you note that for the case of second harmonic generation, what is the initial condition, I have only the fundamental incident on the crystal. So, at z is equal to 0 E_2 of z is equal to 0 there is no second harmonic incident.

So, I have a crystal in which I launch wave at ω frequency. So, this implies u_2 is equal to 0 because E_2 is u_2 times exponential $i\phi_2$ and u_2 is equal to 0 implies C must be 0 because at this is a constant of motion; so, at z is equal of 0 also it is valid. So, at z is equal to 0, I have u_1^2 of 0 u_2 of 0 $\cos\theta$ of 0 and because u_2 of 0 is 0 the constant is 0 and as it changes u_1 will change u_2 will change. So, how will I maintain this equation?

θ has to be $\cos\theta$ must be 0, which means θ must be plus minus $\pi/2$. So, θ will become and remember θ is ϕ_2 minus $2\phi_1$.

Sir, we have taken that n_1 is equals to n_2 .

Phase matched

We are solving this equation, only under phase matching, it is possible to solve in the original paper by Angstrom. I will give you the reference in 1962 physical review a. This equation is solved in more general terms, which in terms of elliptical integrals and so on. So, I am not going to do in details, but I just want to show you that it is possible to solve this equation, when you have phase matched situation, it is much easier. So, Δk is equal to 0; I get those two equations and then I am trying to solve this equation under this condition Δk is equal to 0.

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$$\frac{du_1}{dz} = -\kappa u_1 u_2 \sin \theta$$

$$\frac{du_2}{dz} = \kappa u_1^2 \sin \theta$$

$$\frac{du_1}{dz} = -\kappa u_1 u_2$$

$$\frac{du_2}{dz} = \kappa u_1^2$$

$$\frac{du_2}{dz} = \kappa (u_{10}^2 - u_2^2)$$

$$\frac{d}{dz} (u_1^2 + u_2^2) = 0$$

$$u_1^2(z) + u_2^2(z) = u_1^2(0)$$

$$= u_{10}^2$$

$$u_1^2 = u_{10}^2 - u_2^2$$

So, once theta becomes defined like this, these two equations will give me. Now, these two equations gets simplified to the following two equations. So, du_1 by dz . Let me, assume theta is equal to plus pi by 2, I leave the problem to you, to do for theta is equal to minus pi by 2 for theta is equal to plus pi by 2 minus kappa $u_1 u_2$ and du_2 by dz is equal to kappa u_1 square sine theta becomes 1. Now, again I want to leave for you to show these equations, satisfy the following condition d by dz u_1 square plus. What is this equation? Energy conservation, please remember this is under n 1 is equal to n 2 otherwise, you will have $n_1 \text{ mod } E_1 \text{ square by } 2 c \mu_0$ plus $n_2 \text{ by } 2 c \mu_0 \text{ mod } E_2 \text{ square}$ must be constant.

Because n 1 is equal n 2 is simply becomes $E_1 \text{ square } E_2 \text{ square}$ or $u_1 \text{ square is } u_2 \text{ square}$. So, this implies $u_1 \text{ square of } z \text{ plus } u_2 \text{ square of } z$ must be equal to $u_1 \text{ square of } 0$ because at z is equal to 0 there is no second harmonic incident, I am assuming that the crystal there is only fundamental wave incident in the crystal. So, let me call this u_{10} .

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$$\int_0^{u_2} \frac{du_2}{u_{10}^2 - u_2^2} = \int_0^z \kappa dz$$

$$u_2(z) = u_{10} \tanh(u_{10} \kappa z)$$

$$P_2(z) = \frac{n_2}{2c\mu_0} u_2^2 S$$

$$P_2(z) = P_1(0) \tanh^2 \left(\sqrt{\frac{2c\mu_0 P_1(0)}{n_1 S}} \kappa z \right)$$

① For low η
 η same as before

② $P_1(z) = ?$

So, I need to solve this equation, so, u_1 square becomes u_{10} square minus u_2 square and so, this equation simplifies to du_2 by dz κ times u_{10} square minus. So, I can actually integrate this equation, I will have integral du_2 by u_{10} square minus u_2 square is equal to integral κdz from 0 to z from u_2 is 0 to from u_2 at z is equal to 0 u_2 is 0 at z the value of u_2 is so much, that is an integral function.

So, I will give you the solution, the solution comes out to be u_2 of z is equal to \tan hyperbolic function, you can actually integrate this equation by transforming variables and show the solution of that equation that integral is in terms of \tan hyperbolic inverse and so, you get this expression for u_2 of z from here, you can calculate P_2 of z is equal to n_2 by $2c\mu_0 u_2^2$ into area. n_2 by $2c\mu_0$ mod E_2 square into area mod E_2 square is u_2 square into area. So, $u_2 z$ square is given here and n_1 is equal to n_2 so this you can show again that this is $P_1(0) \tan$ hyperbolic square of you can this is actually square root of $2c\mu_0 P_1(0)$ by n_1 into area into I have replaced u_{10} , by the power in the fundamental and κ is a known quantity ωd by $c n_1$.

So, the power in the second harmonic close with z like this, and if you all know **time** \tan hyperbolic square function is a monotonically increasing function. So, from z is equal to **to...** if you plot as a function of z it will go like this, the maximum power you can get is $P_1(0)$ this is P_2 of z versus z .

Now, I leave the following problem that if the efficiency is small then you can convert the tan hyperbolic function into.... So, if the efficiency is small that means $P_2 z$ must be very very small that means, this argument must be small then you can except the tan hyperbolic function in terms of the argument. And I would like you to show that the efficiency comes out be exactly, what you have calculated before under the low pump depletion approximation.

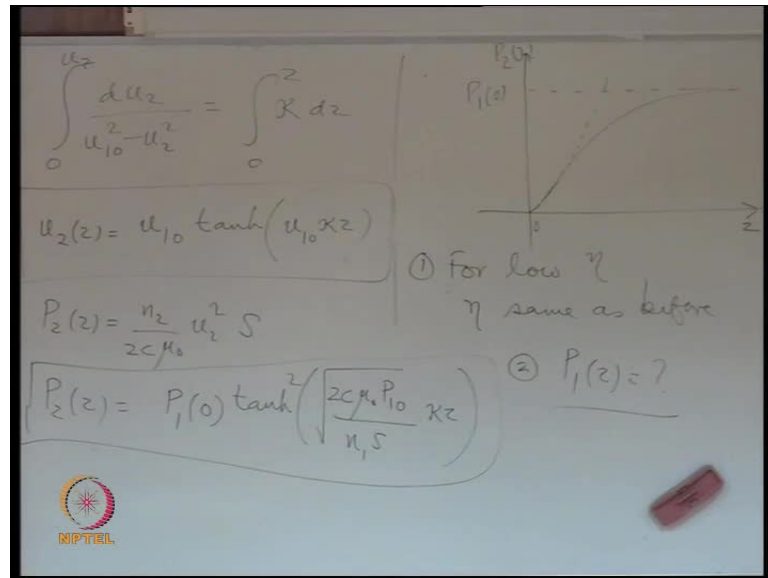
That means, here the expression, which we calculated will go like this, it will be overlapping with this variation, until about a efficiency of a few may be 5 6 percent or 10 percent and then it will diverge up, because an increases quadratically with z the increase is got quadratic it is in the form tan hyperbolic function.

So, two things I want you to show, first thing is for low efficiencies P_2 of z for low efficiencies η same as before and the second thing, I want you to calculate is P_1 of z . This is the power of the second harmonic, what is the corresponding power in the fundamental as z increases? Note that in this calculation, you can never convert all the power from the fundamental second harmonic, because that is a tan hyperbolic function it a synthetically approaches $P_1 0$, it is not a periodic exchange between ω and 2ω .

So, ω you launch ω , you generate 2ω continue to generate 2ω and the 2ω , power keeps on increasing. And if you have perfectly phase matched, it actually approaches the full efficiency of 100 percent at a essentially in finite lines.

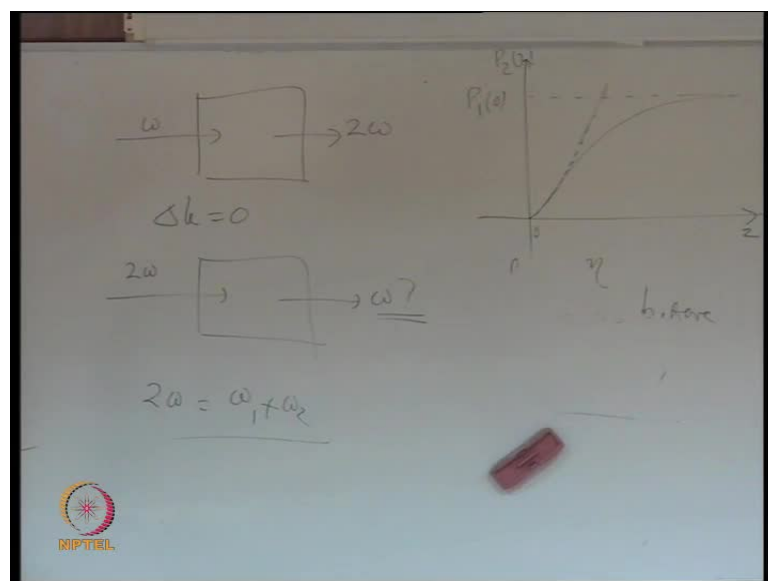
Sir what is that dotted curve?

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This dotted curve is for low efficiencies, what we had calculated sine square delta k z by 2 by delta k by 2 whole square in terms of that function. So, we had actually calculated earlier P 2 of z. Assuming, E 1 was a constant. So, on that approximation efficiency will be low and this is the low efficiency expression and this match for low efficiencies for higher efficiencies that equation cannot be used. And actual equation is this and the efficiency goes to hundred percent, as you increase the length of interaction.

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So, now let us look at another problem, and that is as I was mentioning to you before, generation of ω from 2ω . So, what we have discussed is I have a crystal, I launch ω and I generate 2ω , second harmonic generation. So, for this Δk is equal to 0. Now I take the same crystal and I launch 2ω generate ω ; the sub harmonic generation.

Remember, we have, **we have**, two equations d by dz and dE^2 by dz . we solve the dE^2 by dz equation under low, under low pump depletion approximation and got a variation of we got certain efficiency; and we got an efficiency for complete conversion, for where you do not neglect the depletion. Can I reverse the problem and I say I launch on 2ω and I want to generate ω .

So, here what is happening is 2 photons are frequency ω , merged to form 1 photon at frequency 2ω ; here 1 photon at 2ω is suppose to split into 2 photons that frequency ω . This is also called parametric down-conversion, you are decreasing the frequency, you are generating a lower frequency from a higher frequency, it is called parametric down-conversion.

And in fact, this 2ω can be written as a sum of a infinite sets of two frequencies, 2ω is ω plus ω 2ω is ω_1 plus ω_2 some ω_1 plus ω_2 is also 2ω .

So, in fact, there are infinite combinations of this equation, of course, you need to satisfy phase matching condition to generate the new frequencies. So, we will go to this little later, but first thing, I want to address the problem is if I do not if I launch only 2ω can I generate ω .

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Handwritten mathematical derivations on a whiteboard:

$$\frac{dE_1}{dz} = i \frac{\omega d}{c n_1} E_2 E_1^*$$

$$E_2(z) = \text{Constant}$$

$$E_1(0) = 0$$

$$\frac{d^2 E_1}{dz^2} = i \frac{\omega d}{c n_1} E_2 \frac{dE_1}{dz}$$

$$= i \frac{\omega d}{c n_1} E_2 \left(-i \frac{\omega d}{c n_1} E_2^* E_1 \right)$$

$$= \frac{\omega^2 d^2}{c^2 n_1^2} |E_2|^2 E_1$$

Diagram of a crystal with length $2a$ and incident waves at ω and 2ω .

$$\frac{dE_1}{dz} = i E_1$$

$$E_1(z) = A \sinh \beta z + B \cosh \beta z$$

$$B = 0$$

$$A = 0$$

$$E_1(z) = 0$$

So, let me go back to this equation and let me assume no pump depletion approximation, which means I assume E_2 is a constant. So, the equation, which I need to solve now is just first equation $i \omega d / c n_1$; with a condition that E_2 of z constant and E_1 of 0 that this I have at the input of a crystal, a second harmonic with a certain electric field amplitude E_2 and no incidence of any fundamental wave. There is no ω incident; there is only 2ω incident.

In the first case, there was only ω incident. There was no 2ω incident, I generated ω 2ω from ω . Now, I am trying to generate, ω from 2ω by launching only 2ω .

Sir E_2 represents 2ω and represents ω .

I continue to represent, the second harmonic by the electric field E_2 the fundamental by this electric field E_1 . Now first thing you already notice here that, if E_1 of 0 is 0 $d E_1 / dz$ at z is equal to is also 0; this equation because this is proportional to E_1^* , so, what you expect?

It will continue to say 0, let me calculate. So, let me differentiate this equation, $\frac{d^2 E_1}{dz^2}$ is equal to $i \omega d / c n_1 E_2 d E_1^* / dz$. I substitute from here, so $i \omega d / c n_1 E_2$ into minus $i \omega d / c n_1 E_2^* E_1$, which is equal to $\omega^2 d^2 / c^2 n_1^2 \text{mod } E_2^2$ into E_1 .

Now, let me call this, γ^2 its $\omega^2 d^2 / c^2 n_1^2$ into $\text{mod } E_2^2$. So, this equation is simply $\frac{d^2 E_1}{dz^2}$ is equal to $\gamma^2 E_1$. What is the solution? E_1 of z is equal to or in terms of hyperbolic functions. So, what will it be $A \sinh \gamma z$ plus $B \cosh \gamma z$.

Because, E_1 of 0 is 0. What should be 0?

B B

Because E_1 of 0 is 0 constant d must be 0, because E_1 of 0 is 0 plus B . so, B is 0 $d E_1$ by dz is also 0 at z is equal to 0 that means A is also 0.

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The whiteboard contains the following handwritten notes and equations:

- Top left: $\frac{dE_1}{dz} = i \frac{\omega d}{c n_1} E_2 E_1^*$
- Top right: $\gamma^2 = \frac{\omega^2 d^2}{c^2 n_1^2} |E_2|^2$
- Middle left: $E_2(z) = \text{Constant}$
- Middle left (circled): $E_1(0) = 0$
- Middle left (with diagram): $\frac{d^2 E_1}{dz^2} = i \frac{\omega d}{c n_1} E_2 \frac{dE_1^*}{dz}$. A diagram shows a rectangular block with length $2a$ and width a , with arrows indicating wave propagation.
- Middle left (simplified): $= i \frac{\omega d}{c n_1} E_2 \left(-i \frac{\omega d}{c n_1} E_2^* E_1 \right)$
- Middle left (final simplified): $= \frac{\omega^2 d^2}{c^2 n_1^2} |E_2|^2 E_1$
- Middle right: $\frac{d^2 E_1}{dz^2} = \gamma^2 E_1$
- Middle right: $E_1(z) = A \sinh \gamma z + B \cosh \gamma z$
- Middle right: $B = 0$
- Middle right: $A = 0$
- Middle right (underlined): $E_1(z) = 0$

So, E_1 of z , so this gives me B is equal to 0, because E_1 of 0 is 0 and A is equal to 0, because $d E_1$ by dz is 0 at z is equal to 0. So, E_1 of z is equal to 0, no generation of sub harmonic.

Omega can get converted in 2 omega. Those equations predict, but those same equations tell me that I cannot convert 2 omega at omega, there is nonlinearity, there is nonlinear effects. The powers are large there are strong electric fields, but no generation of omega.

When I do an experiment, I find omega and the explanation is quantum mechanical. Note that according to quantum mechanics, I cannot have complete vacuum in terms of no electric fields and no magnetic fields. I cannot have a situation, where the electromagnetic energy is 0; there is always 0 point energy, 0 point fluctuations and it is so 0 point fluctuations, which will induced the second harmonic to down converter.

For example, if I do not put this condition and I as, if I assume a finite electric field to be incident at frequency omega, I can solve this equations, Because, then those both of them do not become 0 and I will get a solution to the problem.

Let me try to analyze, for example ,so let me assume, that instead of this, so the first thing is that without this noise, without this quantum noise, there is no sub harmonic generation, I cannot generate omega from 2 omega. I cannot generate. This is called, what is called as spontaneous parametric fluorescence or Spontaneous Parametric Down-Conversion SPDC?

It is a spontaneous process in which the 2 omega photons spontaneously split into 2 omega photons, it is parametric down-conversion. It is a down-conversion process, because it goes from a higher frequency to lower frequency and spontaneously it has to happen.

So, SPDC this Spontaneous Parametric Down-cNversion has a completely quantum mechanical origin and this situation is called degenerate parametric fluorescence, because the two frequencies which are coming out are the same.

At a split a 2 omega photon into two different frequency photons, whose sum of frequency is equal to 2 omega. So, that will be non-degenerate parametric fluorescence. This is called degenerate parametric fluorescence, degenerates means the output 2 photons have the same frequency.

The equations, we have derived is only for degenerate, because we have only consider two frequencies omega and 2 omega in the process, after we finish second harmonic

generation and this discussion, we will going to more general case, where I can have three frequencies interacting with each other and that will be used for understanding more things about parametric processes.

Now, let me try to show you mathematically some very interesting features of this equations.

Sir, I seen that because use this same equations to that it is possible to generate omega 2

No, **we have to go** then we go to quantum mechanical description, then only I can show you but these equations can be used, when E_{10} is not equal to 0, that means in addition to 2ω , if I put a small omega here.

Sir, but this is a similar to the condition, when omega **1 got the when the electric field at omega got depleted and there was no (())**

Exactly, so, it is in between condition, I launch omega in the case after some distance, I have both omega and 2ω .

So, but according to it which the energy only gets pumped into E_2 and not and from E_1 once **this no**.

You are right, but I will show you that depending on the phase difference between these two waves energy, can move from omega into 2ω or from 2ω to omega.

This is in the first case that we did where **(())**

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Because that says different number appear here, I will show you depending on the phase difference between these two waves, I can either go from omega to 2 omega to 2 omega 2 omega.

You said that quantum mechanically electromagnetic energy cannot be 0, you said that.

You cannot have a situation where electric and magnetic fields are 0.

If they are finite, so, how can we know that there must be some component at frequency omega presents that some harmonics are generated.

Because they are presented in all frequencies. In free space, all frequencies are allowed and so there is fluctuations corresponding to all frequencies, if you take a cavity a cavity only allows certain frequencies of oscillation and then noise is present at those frequencies; the other frequencies cannot exist in where inside the cavity.

I mean in theory assuming that these frequencies are very well defined and individual frequencies. So, this situation is where I have a strong 2 omega beam and a weak omega beam got again but this is very weak but not 0.

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$$E_1(0) = u_{10} e^{i\phi_{10}}$$

$$B = u_{10} e^{i\phi_{10}}$$

$$\frac{dE_1}{dz} = A\gamma \cosh \gamma z + \gamma u_{10} e^{i\phi_{10}} \sinh \gamma z$$

$$\left. \frac{dE_1}{dz} \right|_{z=0} = A\gamma = i \frac{\omega d}{c n_1} E_2(0) E_1^*(0)$$

$$E_2(0) = u_2 e^{i\phi_2}$$

$$E_1(0) = u_{10} e^{i\phi_{10}}$$

So, let me solve this equation. So, what is an expression, I will get now, so let me call this, this is not 0. So, I call this let me write here again, $E_1(0)$ is E_{10} . So, I get B is equal to E_{10} from here.

So, $\frac{dE_1}{dz}$ is equal to $A\gamma \cosh \gamma z + \gamma E_{10} \sinh \gamma z$ so $\frac{dE_1}{dz}$ at z is equal to 0 is equal to $A\gamma$ is equal to $i \frac{\omega d}{c n_1} E_2(0) E_1^*(0)$.

And substituting the expression for $\frac{dE_1}{dz}$ that is $i \frac{\omega d}{c n_1} E_2(0) E_1^*(0)$. So, let me assume that $E_2(0)$ is $u_2 e^{i\phi_2}$ and $E_1(0)$ is $u_{10} e^{i\phi_{10}}$ and exponential $i\phi_{10}$.

There is a sine or sine hyperbolic

Sine hyperbolic, I am sorry, sine hyperbolic. So, I get an expression for A . A is equal to $i \frac{\omega d}{c n_1}$. Remember, I have called this κ so $\kappa = \gamma u_1 u_2$ this is a let me symbol.

So, let me call this, This is E_{10} actually because $E_1(0)$ is E_{10} . This is complex please just correct this. This is $E_{10} u_{10} e^{i\phi_{10}}$ amplitude and phase so this is B is equal to $u_{10} e^{i\phi_{10}}$.

And this is actually gamma, anyway the second harmonic is constant, so there is no z dependence on the E 2. So, this is A is equal to i kappa u 2 u 10 exponential i phi 20 minus phi 10 minus phi 10, because you have an E 10 star E 10 star is u 10 exponential minus i phi 10 and so that is the expression for A.

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The whiteboard shows the following derivations:

$$E_1(z) = u_{10} e^{i\phi_{10}} \left[\cosh \gamma z + e^{i(\phi_{20} - 2\phi_{10} + \frac{\pi}{2})} \sinh \gamma z \right]$$

$$E_1(z) = A \sinh \gamma z + B \cosh \gamma z$$

$$= \left(i \frac{\kappa u_2}{\gamma} u_{10} e^{i(\phi_{20} - \phi_{10})} \right) \sinh \gamma z + \left[u_{10} e^{i\phi_{10}} \right] \cosh \gamma z$$

So, let me substitute into the solutions and I get a complete solution. As, so let me write this first, one first so what is what was E 1 of z A sine hyperbolic gamma z plus B cos hyperbolic gamma z. so, A is I kappa u 2 by gamma u 10 exponential phi phi 20 minus phi 10 sine hyperbolic gamma z plus u 10 e to the power i phi 10 cos hyperbolic and substituting the values of A and B.

So, I can write this as E 1 of z, Let me take this common now, remember, what was gamma square gamma square was omega square d square by c square n 1 square mod E 2 square. This is equal to kappa square into u 2 square and you find gamma square as omega square d square by c square n 1 square into mod E 2 square omega B by c n 1 is kappa and mod E 2 square is u 2 square so kappa times u 2 at is equal to gamma. So, these cancelled out and I write this as u 10 exponential i phi 10 into cos hyperbolic gamma z plus exponential I have i phi 20 minus 2 phi 10 plus pi by 2 the i factor multiplying here I have written as exponential i pi by 2 and then I have sine hyperbolic gamma z.

What I have done is I have used the initial conditions at z is equal to 0 to calculate the two constants A and B substituted into the solution and explained the solution like this now you see $E_1(z)$ is not 0.

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The whiteboard shows the following derivations:

$$E_1(z) = u_{10} e^{i\phi_{10}} \left[\cosh \gamma z + e^{i(\phi_{20} - 2\phi_{10} + \frac{\pi}{2})} \sinh \gamma z \right]$$

$$\phi_{20} - 2\phi_{10} + \frac{\pi}{2} = 0 \quad \left| \quad \phi_{20} - 2\phi_{10} + \frac{\pi}{2} = \pi$$

$$E_1(z) = u_{10} e^{i\phi_{10}} e^{\gamma z} \quad \left| \quad E_1(z) = u_{10} e^{i\phi_{10}} e^{-\gamma z}$$

At the top right, there is a small equation: $\lambda = \frac{\gamma \kappa}{\gamma} u_2 u_{10} e^{i(\phi_{20} - \phi_{10})}$

Now, I can have two situations suppose this coefficient – suppose $\phi_{20} - 2\phi_{10} + \frac{\pi}{2}$ was equal to 0. The initial phase of the fundamental initial phase of the second harmonic. Remember ϕ_{20} , what is ϕ_{20} ? ϕ_{20} was u_{20} exponential, ϕ_{10} , this is the phase of the second harmonic 2ω wave at the input this is the wave of the fundamental wave at the input the phase of fundamental wave.

Suppose I could satisfy this condition, what will happen E_1 of z ? What happens to the sum? This becomes exponential, becomes 1 sum of cos hyperbolic and sine hyperbolic.

Becomes exponential γz , becomes exponential γz cos hyperbolic x plus sine hyperbolic x is exponential x cos hyperbolic x minus sine hyperbolic x is exponential minus x.

If I had this is one case, so, if I had $\phi_{20} - 2\phi_{10} + \frac{\pi}{2}$ is equal to π or π then $E_1(z)$ will be $u_{10} e^{i\phi_{10}}$.

If I choose this exponential factor to this phase to be minus $\pi/2$. So, this is 0. So, this happens with 0 that means $\phi_{20} - 2\phi_{10}$ is minus $\pi/2$, then the fundamental

wave grows exponentially. This assuming, no pump depletion that means I am solving the equation assuming E^2 is a constant.

So, there is an exponential growth, but you cannot extend this to large values of that so, the ω frequency grows exponentially, if you choose this phase of $\phi_0 - 2\phi_0$ is plus $\pi/2$, if this difference is plus $\pi/2$ then the fundamental decrease and this decay, which means that the energy is transferring from ω to 2ω here. And here it is transferring from 2ω to ω , what is this amplification, u_0 was the input field at the ω frequency as z increases the ω frequency keeps on increasing in amplitude. This is an optical amplifier. This is called an optical parametric amplifier it uses the nonlinear property to amplify the signal and interesting thing is this amplifier depends on the phase of the input signal.

If your input signal phase ϕ_0 satisfy this condition, the signal will get amplified if the input signal phase, satisfy this condition it will get deamplified. This is not like laser amplifiers, if you take a laser amplifier irrespective of what you do to the input phase it is always amplified it uses population inversion. This uses a different principle completely it uses a nonlinear effect and this is called a parametric amplifier and this is a phase sensitive parametric amplifier.

Later on, I will show you the parametric amplifier can operate in phase sensitive fashion of phase insensitive fashion and you can use this phase sensitive nature of this amplification process to very interesting features in generating light with noise properties which are better than, what you can achieve with lasers.

You can use this process to what is called as squeezed vacuum to generates squeezed vacuum has certain fluctuations, you can reduce the fluctuations below, what standard vacuum has by using this kind of a process and that is called squeezed vacuum. And squeezing is a very important technique that is used today to reduce the noise of signals and this is used in many applications including in gravitational wave observatory. For example, there is a light orb interferometer which is which is being built to detect gravitational waves.

So, you need waves, which whose noise levels are very low and the... So, this squeezed light which comes out from this kind of a process can be used to for applications where you want very low noise signals, we cannot have the quiz now.

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Handwritten equations on a whiteboard:

$$\lambda = e^{i\frac{\pi}{2}} u_2 u_{10} e^{i(\phi_{20} - \phi_{10})}$$

$$E_1(z) = u_{10} e^{i\phi_{10}} \left[\cosh \gamma z + e^{i(\phi_{20} - 2\phi_{10} + \frac{\pi}{2})} \sinh \gamma z \right]$$

$$\phi_{20} - 2\phi_{10} + \frac{\pi}{2} = 0$$

$$E_1(z) = u_{10} e^{i\phi_{10}} e^{\gamma z}$$

$$\phi_{20} - 2\phi_{10} + \frac{\pi}{2} = \pi$$

$$E_1(z) = u_{10} e^{i\phi_{10}} e^{-\gamma z}$$

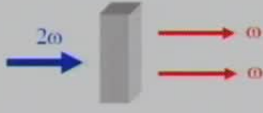
NPTEL logo is visible in the bottom left corner of the whiteboard image.

So, there is there is a very interesting feature here, so let me show you some slides before I stop. So, will have the quiz may be on thursday itself.


So, this is a very interesting feature of an amplifier, which does not use population inversion but which uses nonlinear effects and this the amplifier is phase sensitive. So, we will come back with this phase sensitive amplifiers a little later but let me finish the lecture by showing you some slides.

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Parametric fluorescence



- One photon at 2ω splits spontaneously into two photons at ω
- Explanation is quantum mechanical
 - Vacuum fluctuations lead to this process
- The generated photon pairs can have esoteric properties of **ENTANGLEMENT**
- Application in quantum information science



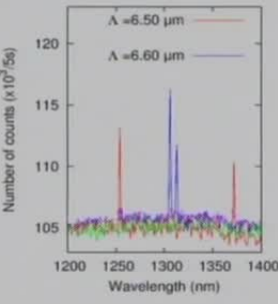
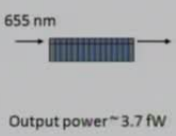
1

So, this is what is called parametric fluorescence. This is a degenerate parametric fluorescence where photons at 2ω splits into 2 photons at ω and these 2 photons, which are coming out, actually are generated from the same photon 2ω and they can possess certain peculiar features of what are called as entanglement properties of entanglement? So, there are correlations between these 2 photons, which are beyond what classical physics predicts and this process is very much used in in quantum information science.

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Parametric fluorescence

Input laser at $\lambda \sim 655$ nm



655 nm

Output power ~ 3.7 fW


Number of counts ($\times 10^7/5s$)

Wavelength (nm)

$\Lambda = 6.50 \mu m$

$\Lambda = 6.60 \mu m$

Ref: Martin et al., Opt Exp 17 (2009) 1033

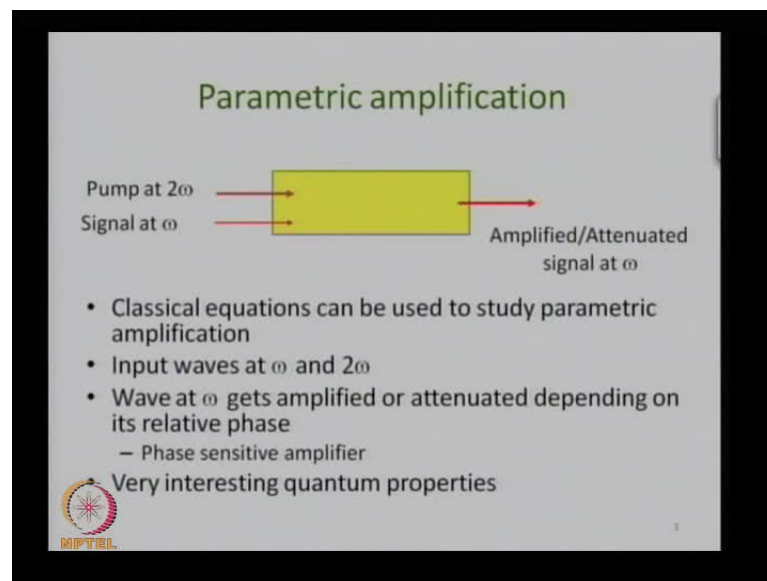


2

So, here is a typical example of light being pumped at 655 nanometers and out comes light at longer wavelengths, which is down-conversion and depending on the quasi phase matching period. The pairs of light wavelengths coming out are such that the ω_1 plus ω_2 is equal to the ω of the input light. 655 nanometers photons are splitting into 2 photons of lower frequencies or higher wavelengths and the sum of these two frequencies is equal to 655. The sum of these two is also 655 but in this case, these are coming out because you satisfy the phase matching condition for that pair of frequencies.

So, if you change the period of this quasi phase matching, you can generate different pairs of photons, which are coming out please store the power levels control watts very very low generation powers. **These are this is called** this is called fluorescence parametric fluorescence, you shine light at this wavelength and out comes light at new wavelengths, which are longer wavelengths than the input light.

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What we discussed today is a parametric amplifier, you can amplify input signal and this is phase sensitive. So, it has very interesting quantum properties, which we will come back when we discuss quantum mechanics at that time I will bring back this problem and we look at some quantum properties of this.

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Parametric amplifiers: History

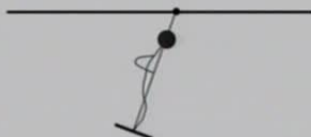
- First observations made by Michael Faraday in 1831:
 - oscillations of one frequency excited by forces of double the frequency observed on the surface in a wine glass excited by moving a moistened finger around its circumference
 - Arrived at the conclusion experimentally that there were two complete vibrations of the support for each complete vibration of the liquid.
- Melde's experiment in 1859 of exciting a string attached to a tuning fork when the fork is vibrating at twice the frequency and parallel to the string
- Lord Rayleigh analyzed these in 1883-1887
- Classic example: child on a swing



Incidentally, parametric amplifiers were investigated about Michael Faraday in 1831: so you take a wine, if you rub your finger on a wine glass, the wine started to vibrate and he found out that there were two complete vibrations of the support for one vibration of the liquid 2ω .


Many of you may not have done this experiment - Melde's experiment where you can excite a string into oscillation by forcing it at the twice the frequency and not really as analyzed this, and let me show you classical example of a child on a swing.

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The diagram shows a child on a swing. A horizontal line represents the pivot point. A vertical line connects the pivot to a black dot representing the seat. A stick figure is attached to the seat, leaning forward. The swing is shown in a slightly displaced position from the vertical.

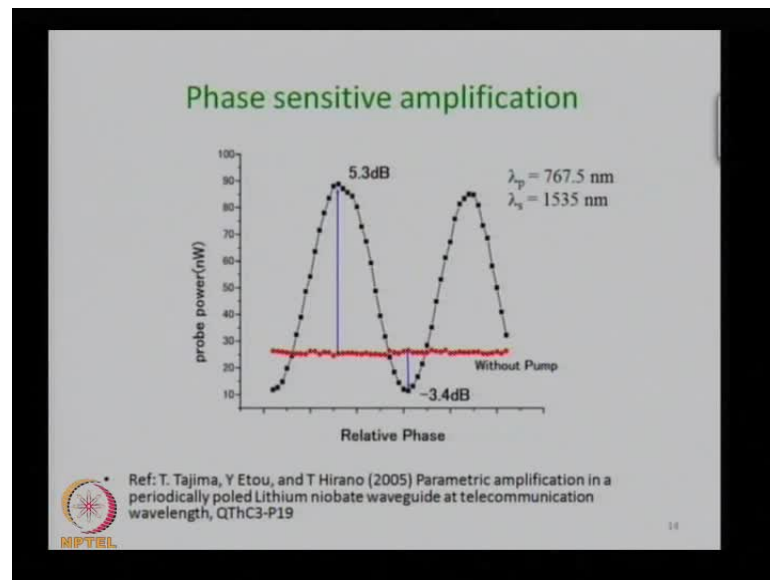
- A child can amplify the amplitude of the swing by periodic up and down motion at twice the natural frequency of the swing if pumped at the right phase



So, child sits then stands up, again sits back and stands up, two frequencies of the child sitting and standing for one period of the swing. The child is pumping the swing at 2ω and the swing is oscillating at ω frequency.

And so, also note that the child has to pump the swing at the right phase, if it gets lot of phase, the swing will come down to rest and also nobody needs to push. The child can generate some noise in the swing and then start to amplify that then, the moment it starts to swing a little bit you sit and stand at the right time, the right phase and you start to amplify the signal. So, you can actually generate from noise amplify that noise and generate the oscillations.

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And here is some experimental results showing, how the amplification depends on the relative phase, which is what this equation predict and you can these are some experimental results showing, phase sensitive amplifiers and so, these are very interesting. So, there is a lot of work going on in phase sensitive amplifiers and as I said one of the most interesting application of this is in quantum noise reduction or generating non-classical states of light and so on.

Thank you.