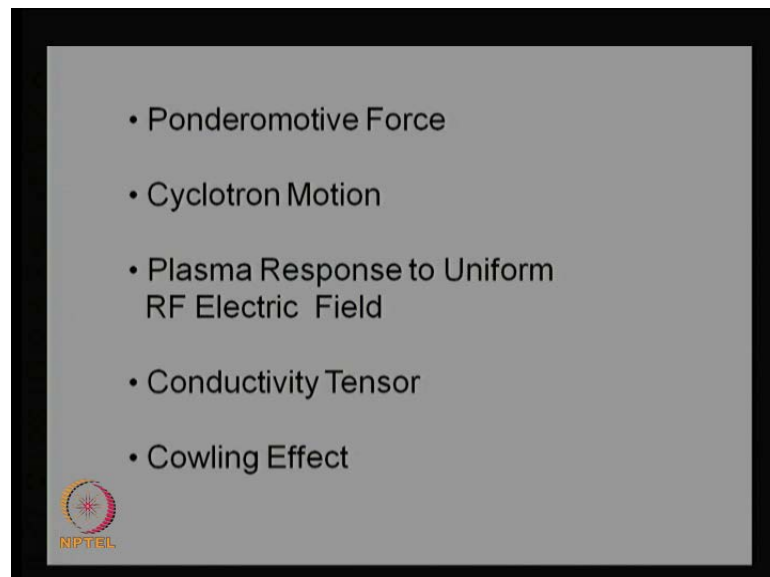


Plasma Physics
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Lecture No. # 05
RF Conductivity of Plasma (contd.)

Well friends. Last time we discussed the response of plasma to radio frequency electric field, in the absence of any other field like a magnetic field or anything else. Today, I like to extend that study to electric field that has space dependence, and also study the response of plasma to an electric field in the presence of a DC magnetic field. Because most plasmas, if you have to confine them from expansion stops them from expansion, you employ a static magnetic field. And hence, a realistic plasma response would have to be investigated in the presence of a DC magnetic field, but before I go over to discuss the response of plasma in the presence of a DC magnetic field.

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Let me consider the problem of plasma response to a RF field, which depends on a space. And then I will go over to discuss the electron motion in a magnetic field alone, and then plasma response to uniform radio frequency electric field. I will reduce an

expression for conductivity tensor. And Then I will discuss an important effect called Cowling effect.

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$$\vec{E} = \vec{A}(\vec{r}) e^{-i\omega t}$$

$$m \left(\frac{\partial \vec{v}}{\partial t} + \vec{v} \cdot \nabla \vec{v} \right) = -e \vec{E} - e \vec{v} \times \vec{B}$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} = i\omega \vec{B}$$

$$\vec{B} = \frac{1}{i\omega} \nabla \times \vec{E}$$

$$m \frac{\partial \vec{v}}{\partial t} = -e \vec{E} - \underbrace{m \vec{v} \cdot \nabla \vec{v} + e \vec{v} \times \vec{B}}_{\vec{F}_P \text{ Ponderomotive force}}$$

So, in order to understand the response of plasma to an electric field, which varies, whose amplitude varies in position, say r in general and time dependence of this form. This is very important issue, understand, that the electric field gives rise to velocity to particles, velocity means their position changes. So, when position changes r will change it time. So, you cannot take this to be constant.

So, there is an inherent or hidden dependence of an on time. As explicitly, it does not depend on time, But because the position of a particle that, say this electric field varies from time to time due to the velocity it acquires, due to the RF field, this A has hidden time dependence.

So, this is a serious problem and I will have to employ the technique of iteration to solve the equation of motion. Let me write down the equation of motion, for the sake of simplicity, I will ignore collisions; I will ignore pressure term anything.

Just consider, simple response of plasma to an electric field of this form, the equation of motion is $m \Delta v$ by Δt plus $v \cdot \nabla v$ is equal to minus $e E$. However, there is a problem here, if there is time varying electric field with amplitude a function of position, if you take curl of e , it may be 0, and it may not be 0. If it is not 0, then third Maxwell

equation will tell that, there is a magnetic field also should be there. And hence, I must add a term here, minus $e \mathbf{v} \times \mathbf{B}$. So, this \mathbf{B} is the magnetic field associated with a time dependent electric field, whose amplitude depends on r .

Now, this is a little complicated equation. This is the electric force; this is the magnetic force which involves the response \mathbf{v} also. My goal is to find \mathbf{v} due to e , what is the particle velocity due to the electric field e ? Well, first let me calculate, what is \mathbf{B} ? From third Maxwell equation, curl of \mathbf{E} is equal to minus $\Delta \mathbf{B}$ by Δt in mts units. And because electric field is changing in time like this. In the steady state, \mathbf{B} will also change with time like this.

So, I can replace this by $i\omega$, $\Delta \Delta t$ by $-i\omega$, already minus is there. So, simply this becomes like this. So, you see here, that \mathbf{B} is equal to 1 upon $i\omega$ curl of \mathbf{E} . If my electric field is small, \mathbf{B} is also small and the response \mathbf{v} will also be small.

So, how do we solve this equation? The terms that depend on the product of a small quantity are called the small terms. Like, \mathbf{v} if I multiply with another \mathbf{v} here, this will be product of very small quantities; it will be much smaller than this term. Similarly, \mathbf{v} is a small response, small velocity, if my e is not large. And \mathbf{B} also will be small, because e is not large. So, this is a product of two small quantities and you can ignore them.

So, we solve this equation iteratively, by first ignoring this term and this term. Rather than ignoring them, let me rewrite this equation like this; $m \Delta \mathbf{v}$ by Δt is equal to minus $e \mathbf{E}$, So, take this term on the right hand side, it becomes; minus $m \mathbf{v} \cdot \nabla \mathbf{v}$ minus $e \mathbf{v} \times \mathbf{B}$. These two terms are called non-linear terms, because they are the product of small quantities and they have the dimensions of force. So, I can combine them together from here to here and give a name to them, in literature, the name given to them is F_p ; non-linear force also called ponderomotive force.

So, I would like to calculate the value of this iteratively. First of all, I will ignore the ponderomotive force, obtain the response of electrons to the electric field alone and then, use the value that I will calculate here, in these terms and evaluate those terms. Let us do that.

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To the zeroth order

$$m \frac{d\vec{v}}{dt} = -e\vec{E}, \quad \begin{aligned} \vec{E} &\sim e^{-i\omega t} \\ \vec{v} &\sim e^{-i\omega t} \end{aligned}$$

$$\frac{d}{dt} = -i\omega$$

$$\vec{v} = \frac{e\vec{E}}{m i \omega}$$

$$\vec{F}_p = -m \vec{v} \cdot \nabla \vec{v} - e \vec{v} \times \vec{B}$$

$$= -\frac{m}{2} \text{Re} [\vec{v} \cdot \nabla \vec{v}^*] - \frac{e}{2} \text{Re} [\vec{v} \times \vec{B}^*]$$

So, through the zeroth order, what do I get, $m \Delta v$ by Δt is equal to minus $e E$ and because this has simple time dependence of the form, exponential minus $i \omega t$.

So, velocity also, I would expect to have time dependence of the form exponential minus $i \omega t$. Hence, Δt , I will replace by minus $i \omega$ as I discussed earlier.

Put this in here, you will get the velocity of the electron is equal to $e E$ upon $m i \omega$. So, now my goal is to calculate the ponder motive force, F_p which is equal to minus $m \nabla \cdot v$ minus $e v \times B$. I will have to use the complex number identity to evaluate these products. Because here really, real part of v is to be placed here, a real part of v has to be used here. Similarly, real part of v has to be used there. So, it is a product of real part of v into real part of this quantity. Similarly, real part of v is to be multiplying with the real part of capital B there.

What will I do? I can write this as, minus m real part of, so, real part is implied; I will not write explicitly the symbol Re . what is implied? Or let me write this, half there and $v \nabla \cdot v$; I am not interested in time dependent part of ponder motive force, for the moment, I want time independent part of ponder motive force. So, I will take only complex conjugative.

I just explained, in my previous lecture, that when you have to multiply terms, then, there are always two terms; one without complex conjugate, the other one is complex

conjugate. The term with complex conjugate, the exponentials cancel out. So, this becomes a time independent term. So, I will write only that.

And then, here it becomes e by 2 real part of \mathbf{v} cross \mathbf{B} star. So, this is time average ponder motive force otherwise I should write two more terms, I will not write them.

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$$\begin{aligned} \frac{d}{dt} &= -i\omega \\ \vec{v} &\sim e^{-i\omega t} \\ \vec{v} &= \frac{e\vec{E}}{m i \omega} \\ \vec{F}_p &= -m \vec{v} \cdot \nabla \vec{v} - e \vec{v} \times \vec{B} \\ &= -\frac{m}{2} \text{Re} [\vec{v} \cdot \nabla \vec{v}^*] - \frac{e}{2} \text{Re} [\vec{v} \times \vec{B}^*] \\ &= \frac{e^2}{2m\omega^2} \text{Re} [\vec{E} \cdot \nabla \vec{E}^*] - \frac{e^2}{2m\omega^2} \text{Re} [\frac{1}{i} \vec{E} \times \vec{B}^*] \end{aligned}$$

Let me evaluate this number, this expression, this force. What you get here, v if I substitute from here, and takes complex conjugate there, this becomes simply equal to e square, one e is there. So, e square m square comes in the denominator, one m will cancel out with this, so e square upon $2 m \omega$ square will be common, and will be real part of $\mathbf{E} \cdot \nabla \mathbf{E}^*$ minus if I put the value of v again here, it becomes e square upon $2 m$.

Now, $i \omega$ is there, but \mathbf{B} also has... So, let me just remind you well, I think let me put this, to in two steps; $i \omega$ is because of this. I think I should cancel this here, because real part I have to take. So, this becomes e square upon $2 m \omega$ outside, and this becomes 1 upon i into $\mathbf{E} \times \mathbf{B}^*$. So, I just substituted the value of v from this expression in here, and taken the constant terms outside, imaginary term I should keep inside, this is what you get.

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The image shows a chalkboard with the following handwritten equations:

$$\vec{B} = \frac{1}{i\omega} \nabla \times \vec{E}, \quad \vec{B}^* = -\frac{1}{i\omega} \nabla \times \vec{E}^*$$

$$\vec{F}_p = -\frac{e^2}{2m\omega^2} \text{Re}[\vec{E} \cdot \nabla \vec{E}^*] - \frac{e^2}{2m\omega^2} \text{Re}[\vec{E} \times \nabla \times \vec{E}^*]$$

$$\text{Re} \vec{C} \equiv \frac{1}{2} (\vec{C} + \vec{C}^*)$$

$$\vec{F}_p = -\frac{e^2}{4m\omega^2} [\vec{E} \cdot \nabla \vec{E}^* + \vec{E}^* \cdot \nabla \vec{E} + \vec{E} \times (\nabla \times \vec{E}^*) + \vec{E}^* \times (\nabla \times \vec{E})]$$

$$= -\frac{e^2}{4m\omega^2} [\vec{A} \cdot \nabla \vec{A}^* + \vec{A}^* \cdot \nabla \vec{A} + \vec{A} \times \nabla \times \vec{A}^* + \vec{A}^* \times \nabla \times \vec{A}]$$

An NPTEL logo is visible in the bottom left corner of the chalkboard image.

Let me, remind you that B, I had written as, 1 upon i omega curl of E. So, use there; B star will be minus 1 upon i omega into curl of E star, complex conjugate of e.

So, if I use this in the expression for F p, becomes, equal to, I forgot the first term was a negative sign also, in the previous case it was actually had a negative term, minus e square upon twice m omega square and real part of E dot del E star minus e square upon twice m omega square and you will get real part of E cross curl of E star. This factor is common in both the terms. And you always know that, real part of any complex number C is the same thing as, half C plus C star. If there is any complex number, if you add its complex conjugate with it, the imaginary parts will cancel out and this is always true.

So, now this is a complex number in general, what I can do? I can write this expression as, this expression plus its complex conjugate divided by 2. So, F p becomes minus e square upon 4 m omega square and this becomes E dot del E star plus complex conjugate of this quantity, which is E star dot grad E, just reverses the stars.

Similarly here, this is plus E cross curl of E star plus E star cross curl of E. This is a very general expression for ponderomotive force due to a RF field, which does not depend, whose amplitude does not depend on time, but depends on space and this could be arbitrary displacement dependence. So, this is the general expression.

Now, put this, e in terms of A , then, this can be written simply as, minus e square upon $4 m \omega^2$, all the exponential factors will cancel out, because they are complex conjugate terms and they will be simply, $A \cdot \nabla A^*$ plus $A^* \cdot \nabla A$ plus $A \times \nabla A^*$ here plus $A^* \times \nabla A$. This is a vector quantity, it has x component, y component and z component. I evaluate only one component, $F_p x$ and then, I will generalize to. Similarly, I will write down y, z component. So, let me write down the x component, for a general case. $F_p x$ will be how much? Let us see.

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$$\begin{aligned}
 F_{px} &= -\frac{e^2}{4m\omega^2} \left[\vec{A} \cdot \nabla A_x^* + \vec{A}^* \cdot \nabla A_x + (\vec{A} \times \nabla A^*)_x + (\vec{A}^* \times \nabla A)_x \right] \\
 &= -\frac{e^2}{4m\omega^2} \left[A_x \frac{\partial A_x^*}{\partial x} + A_y \frac{\partial A_x^*}{\partial y} + A_z \frac{\partial A_x^*}{\partial z} \right. \\
 &\quad + A_x^* \frac{\partial A_x}{\partial x} + A_y^* \frac{\partial A_x}{\partial y} + A_z^* \frac{\partial A_x}{\partial z} \\
 &\quad + A_y \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) - A_z \left(\frac{\partial A_z}{\partial x} - \frac{\partial A_x}{\partial z} \right) \\
 &\quad \left. + A_y^* \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) - A_z^* \left(\frac{\partial A_z}{\partial x} - \frac{\partial A_x}{\partial z} \right) \right] \\
 &= -\frac{e^2}{4m\omega^2} \frac{\partial}{\partial x} [A_x A_x^* + A_y A_y^* + A_z A_z^*]
 \end{aligned}$$

$F_p x$ would be minus e square upon $4 m \omega^2$ out. Now, the first term was $A \cdot \nabla A^*$. So, I am writing x component, this will be $A \cdot \nabla A^*$, x component, I should write here, forgot this sign here. Because $A \cdot \nabla$ is a scalar quantity. So, when I had written a star, I should write like this, plus second term was $A^* \cdot \nabla A$, which is a scalar quantity, so total quantity has to be written and A was there, so, put $A \times$ there.

Similarly, $A \times \nabla A^*$, I want x component. Let me write this simply, $A \times \nabla A^*$, I want it is x component. Then, I want x components quantity, $A^* \times \nabla A$, x component of this quantity I want.

Let me simplify this, this is equal to minus e square upon $4 m \omega^2$, $A \cdot \nabla$ has three terms; this is $A_x \frac{\partial}{\partial x} A_x^* + A_y \frac{\partial}{\partial y} A_x^* + A_z \frac{\partial}{\partial z} A_x^*$

$\mathbf{A} \cdot \nabla (\mathbf{A} \times \mathbf{A}^*)$. $\mathbf{A} \cdot \nabla$ is a scalar product. So, it has $\mathbf{A} \cdot \nabla \mathbf{x}$ plus $\mathbf{A} \cdot \nabla \mathbf{y}$ plus $\mathbf{A} \cdot \nabla \mathbf{z}$ and all the three components have to be added, but this they all operate over $\mathbf{A} \times \mathbf{A}^*$.

Similarly, second term will give me, plus $\mathbf{A} \times \mathbf{A}^* \cdot \nabla (\mathbf{A} \times \mathbf{A}^*)$ plus $\mathbf{A} \times \mathbf{A}^* \cdot \nabla \mathbf{y}$ plus $\mathbf{A} \times \mathbf{A}^* \cdot \nabla \mathbf{z}$ of $\mathbf{A} \times \mathbf{A}^*$. Then, the x component of this product which is little complicated, but it is not that bad.

So, I treat this as one vector, this is another vector; $\mathbf{A} \times \mathbf{B}$ kind of thing. So, I will get x component of this quantity means, y here z there minus z here y there. So, this will be, x component will be, A_y multiplied by z component of this quantity, z component means, $x y$ minus $y x$. So, $\nabla \cdot (\mathbf{A} \times \mathbf{A}^*) - \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$, this is A_y into z component of this quantity, minus z component of this and y component of this. So, y component of this quantity would be $\nabla \cdot (\mathbf{A} \times \mathbf{A}^*) - \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$. I am very explicit on this, because this is something, which is very useful, very important in non-linear plasma theory.

So, the derivation of this, I am really doing step by step, so, that you have some feeling for this, plus I have to write down the x component of this quantity, which is very similar to this one, only the stars are interchanged.

So, this becomes $A_y \nabla \cdot (\mathbf{A} \times \mathbf{A}^*) - \nabla \cdot (\mathbf{A} \times \mathbf{A}^*) - A_z \nabla \cdot (\mathbf{A} \times \mathbf{A}^*) + \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$ like this. I have been very explicit, but this is useful.

Now, you may note here, that only two terms survive, others cancel out. Not to me, let this, most of the terms cancel out. For instance, this term $A_y \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$, you will find $A_y \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$ here as well, $A_y \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$; this term will cancel with this term. This $A_z \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$ term, $A_z \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$ term, this will cancel out with this. Then this term, this will survive, $A_y \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$, $A_y \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$, this is negative, this is positive, they will cancel.

Then this, $A_z \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$, $A_z \nabla \cdot (\mathbf{A} \times \mathbf{A}^*)$, this will cancel out with this one. So, lot of terms have canceled out, what you are left with, this is $\mathbf{A} \cdot \nabla (\mathbf{A} \times \mathbf{A}^*)$

A x star and this is, these two terms can be combined, become delta delta x of x star A x A x star.

Similarly, this term delta delta x is common, so it becomes, this term can be combined with A y term, this one. So, these two can be combined, so, this can be combined with this term, this can be combined with this term and this term can be combined to the last one. They have same sign. So, this can be combined with this term and when you combine them, you get a simple expression minus e square by 4 m omega square delta delta x of A x A x star plus A y A y star plus A z A z star.

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$$\begin{aligned} \vec{F}_p &= - \frac{e^2}{4m\omega^2} \frac{\partial}{\partial x} |A|^2 \\ \vec{F}_p &= - \frac{e^2}{4m\omega^2} \nabla |A|^2 \\ &= - \frac{e^2}{4m\omega^2} \nabla |E|^2 \end{aligned}$$

$\begin{array}{c} \uparrow \text{low field} \\ \text{---} \rightarrow \text{high field} \\ \downarrow \end{array}$

Let me rewrite this expression, this becomes, so, the x component of ponder motive force is simple, minus e square upon 4 m omega square delta delta x of modulus of A square.

So, this is, this is not vector, this is a scalar quantity; x component. Similarly, if you calculate the y component, z component, then, you will write F p vector is equal to minus e square upon 4 m omega square del operating over A square modulus of A square. But modulus of A is the same as modulus of the electric field, so, minus e square upon 4 m omega square gradient of e square.

This quantity is very important in plasma theory, plasma physics. What does it tell? This says that, if any region of space electric field has non-uniform amplitude, suppose the electric field has a larger value here and a smaller value there, then, the gradient of

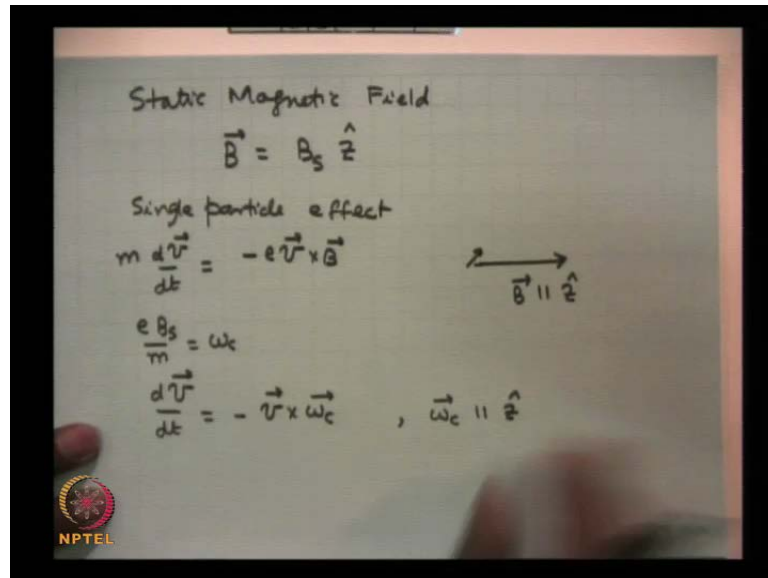
electric field is in this direction, the force will be opposite. Means, this force compels the electron to go from high field region, this is high field region for instance, to low field region.

Electrons will go like this, they will flow out. It is a very important consequence; it has a very important consequence in plasmas, because if you are sending for instance, a high power laser in plasma, the intensity of the laser is large on the axis and small outside. So, when the laser goes into plasma, the electric field is large on the laser axis and less outside. So, laser repels the electrons from the central region or actual region to outside, creating a plasma channel. This non-linear force has been recognized, has been measured and has been observed in numerous experiments and has been found to play very important role in laser guiding and many other non-linear phenomena.

Here, I have calculated only the time average part of this ponder motive force. If you had calculated the time dependent part of ponder motive force, this is responsible for harmonic generation, which is another important non-linear effect.

So, in our response of plasma to radio frequency field, we have learnt that the radio frequency response at low fields can be described in terms of a quantity called plasma conductivity. But when the fields are strong and when they become time space dependent, in that case, an additional response term of force \times on the particles, which is called ponder motive force. What is the response of plasma to this force? We shall discuss later on, when we discuss the phenomenon of shell focusing, may be a few lectures later. So, we conclude our discussion on radio conductivity or radio frequency response of plasma in the absence of a magnetic field.

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Now, I would like to take you to the plasma response in the presence of a static magnetic field. A static magnetic field in the easiest form can be written as, B static which is independent of position and time. And suppose the direction, I choose to be in the z direction. First of all, I would like to see, what an individual electron will do in the presence of a magnetic field. And then, later on, I will discuss, if you have millions and trillions of electrons. So, electron is a family, electron is a community, how do they respond to this magnetic field or in the presence of magnetic field, how do they respond to the electric field.

So, first of all, let me do the single particle effect. Well suppose, this is my z axis, where I have applied my dc magnetic field. If I launch an electron in the system, which is moving with some velocity here for instance, you know that the velocity of the particle can be resolved into two components; one along the magnetic field and one perpendicular. The electron will experience a force. The force will be charge of the electron into velocity of the electron cross B , this is the force on the electron. And the equation of motion would be $m \frac{d\vec{v}}{dt}$ for a single electron, this is the equation of motion.

So, what will happen? The electron, physically let us understand what happens, the electron resolve in two components; one parallel to magnetic field, one perpendicular. The component of velocity which is parallel to magnetic field, the force vanishes. Only

because of the perpendicular velocity, this will experience a force. And this will try to; the force will be perpendicular to both the velocity and magnetic field. So, perpendicular to the board and electron gyrates about the line of force.

And that motion can be easily understood from this equation. If I divide this equation by m and I define a quantity $e B$ upon m is equal to ω_c for instance. So, this equation becomes $d \mathbf{v}$ by $d t$ is equal to minus \mathbf{v} cross ω_c . Well, ω_c is a vector, parallel to z axis. Because static magnetic field in the z direction.

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z comp
 $\frac{dV_z}{dt} = 0$, $V_z = \text{const}$

x comp
 $\frac{dV_x}{dt} = -V_y \omega_c$

y comp
 $\frac{dV_y}{dt} = V_x \omega_c$

$\frac{d^2 V_x}{dt^2} = -\omega_c \frac{dV_y}{dt} = -\omega_c^2 V_x$

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So, I want to solve this equation, this simple equation. But this has very rich physics. First of all, I will write down the z component of this equation, it becomes, $d v_z$ by $d t$ is equal to 0, which means that the particle velocity in the z direction will be constant.

Secondly, I write down the x component, that would give me, $d v_x$ by $d t$ is equal to v_y cross ω_c ; if you write down the x component will be, minus $v_y \omega_c$. And then, write down the y component the equation of motion, $d v_y$ by $d t$ is equal to $v_x \omega_c$. The x and y components of velocity equations are coupled. So, what you should do, take the time derivative of the first equation, you will get $d^2 v_x$ by $d t^2$ is equal to minus, ω_c is already there, $d v_y$ by $d t$. Now, use of value of $d v_y$ by $d t$ from here, it becomes, minus $\omega_c^2 v_x$. And a general solution of this equation can be written. Let me, write down that solution.

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$$\begin{aligned}v_x &= a \cos(\omega_c t + \delta) \\v_y &= a \sin(\omega_c t + \delta) \\v_x^2 + v_y^2 &= a^2 = v_{\perp}^2 \\ \frac{dx}{dt} &= v_x = v_{\perp} \cos(\omega_c t + \delta) \\x &= x_g + \frac{v_{\perp}}{\omega_c} \sin(\omega_c t + \delta) \\y &= y_g - \frac{v_{\perp}}{\omega_c} \cos(\omega_c t + \delta)\end{aligned} \quad \left| \begin{aligned}(x-x_g)^2 + (y-y_g)^2 \\ &= \frac{v_{\perp}^2}{\omega_c^2} = \rho^2\end{aligned}\right.$$

That solution is, v_x is equal to some constant, let me call this constant as, $a \cos \omega_c t$ plus another constant, δ . A second order differential equation has two constants. So, this is one constant, this is another constant. And if, v_x is like this and if you use the x component equation of motion, then, v_y can be obtained to be equal to $a \sin \omega_c t$ plus δ .

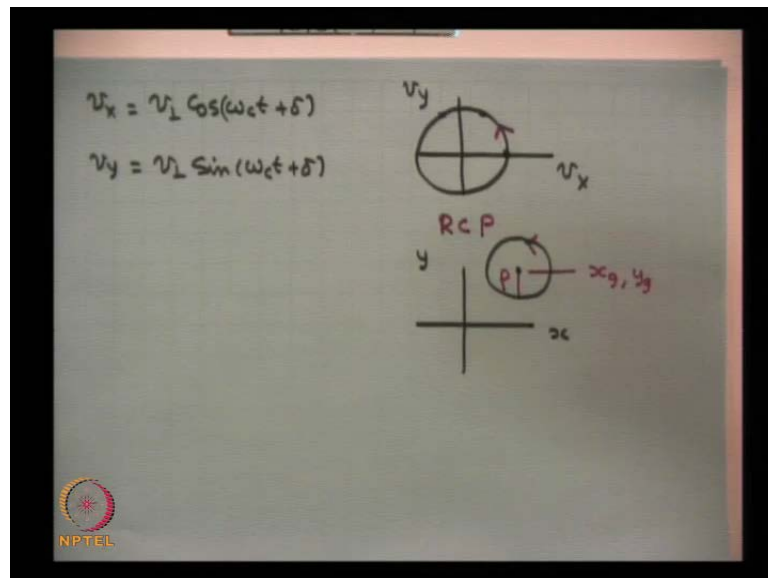
So, I have substitute the value of v_x in the x component of equation of motion, I get the value of v_y . So, these two I obtained. And you may note one thing in here, that $v_x^2 + v_y^2$, if I add, they give me a square. So, a is a constant, means $v_x^2 + v_y^2$ is a constant; this is usually given a symbol v_{\perp}^2 .

So, rather than writing a , I will write down a symbol v_{\perp} . Then, this equation becomes, v_x is equal to $v_{\perp} \cos(\omega_c t + \delta)$, but v_x is the same thing as $\frac{dx}{dt}$. So, I can integrate this equation to obtain x as a function of time, because v_{\perp} is a constant and this time dependence is simple.

So, x turns out to be some constant, let me call this constant as, $x_g + \frac{v_{\perp}}{\omega_c} \sin(\omega_c t + \delta)$. And similarly, if you solve this equation for v_y , you will get, y is equal to $y_g - \frac{v_{\perp}}{\omega_c} \cos(\omega_c t + \delta)$. This is a very general expression. But the beauty here is, that if I take, $(x-x_g)^2 + (y-y_g)^2$ and add y minus y_g square to it, you will find any constant. So, these two

equations give me, $x - x_g$ whole square plus $y - y_g$ whole square is equal v_{\perp}^2 upon ωc^2 , this quantity is given a symbol called rho or normal radius. This is the equation of a circle whose center is at x_g, y_g and whose radius is rho, which is v_{\perp} upon ωc .

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And so, the electron really rotates on a circle, if you draw a graph of, suppose I plot here, v_x and v_y here and v_x I have taken is equal to $v_{\perp} \cos \omega c t$ plus delta and v_y is equal $v_{\perp} \sin$ of $\omega c t$ plus delta.

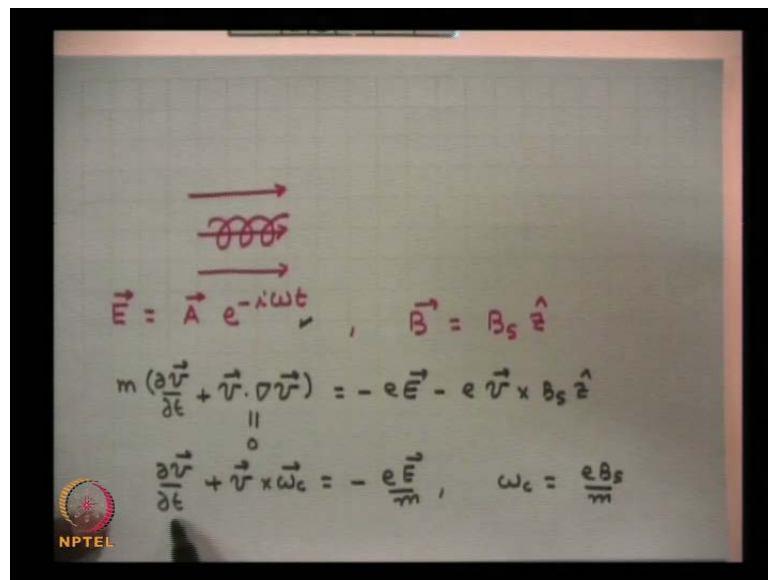
Now, if you plot for the sake of specificness, suppose I choose delta equal to 0, really does not matter, what delta you choose? Delta equal to 0, time t equal to 0, v_x will be v_{\perp} and v_y will be 0, so, you are somewhere here.

At a later time, this quantity will decrease, this will increase and this electron will rotate like this. So, this is the way the electron will move, the electron was initially here, it went from here to here, then here and everywhere. So, this is rotating, in a sense that if I take a right handed screw and rotate along the path of this like this, if I rotate it like this, then, this will advance in the direction of z axis, this electron motion is known as Right Circularly Polarized motion. Right circularly polarized or it is rotating in the right handed sense.

Similarly, if I plot, x here particle coordinate x there and particle coordinate y there, then, around a point x g, y g, the electron will rotate over a circle and the coordinates are like this. This coordinate are x g, y g and this radius is r, rho rather and electron is rotating like this.

So, what is happening that, electrons rotate about a line of force, they can move along the line of force with velocity v z, but they cannot live the line of force. So, whenever there is a magnetic field in the plasma, the electrons in the absence of any other force, I am considering no collisions, no temperature variation, nothing except a dc magnetic field, which is uniform.

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So, if you have lines of force like this, then, the electrons move along the line of force like this, but they do not move perpendicular to the line of force.

So, anyway the electrons are confined to the lines of force and magnetic field provides a confinement to the plasma. Obviously, there is problem that what would you do, either you find some sort of reflection of electrons along the lines of force by some mechanism or do something or close the line of force, only then, you can confine the plasma, that is a separate issue. So, but we have really found something very significant, that a single electron in the absence of any other force will not move perpendicular to line of force.

Now, I would like to examine the effect of a dc magnetic field, RF electric field on the electrons in the presence of a magnetic field.

So, my issue is, I have a plasma in which I am applying a electric field of this form whose amplitude is constant in time, the field is independent of position coordinate also. And I would like to examine the response of the plasma, when there is a static magnetic field also, which is equal to B_s in the z direction.

So, the electric field of the RF can have a component along z axis and it can also have components perpendicular to B_s . Let us see, what is the response? Because we are ignoring any space dependence. So, the equation of motion, then, in this case would be, $m \frac{dv}{dt} + v \cdot \nabla v$ is equal to $-eE - e v \times B_s$, which is in the z direction. This is my equation of motion, I take this term to be 0, because v does not depend on x y z.

So then, if I divide this equation by m, my equation becomes, $\frac{dv}{dt}$, I will take this from the left hand side, so it becomes, $v \times \omega_c$ is equal to $-\frac{eE}{m}$. Remember, ω_c I am defining as, $\frac{eB_s}{m}$. The same ω_c was appearing in the equation of motion in a pure electric magnetic field and this is known as the electron cyclone frequency, because the electron rotates with this frequency around the line of force.

So, I have to solve this equation. Because my time dependence of the electric field is like this, in the steady state, I expect the v also to be having same time dependence. Hence, I will replace $\frac{d}{dt}$ operator by $-i\omega$ and let us see, what do I get.

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Handwritten equations on a whiteboard:

$$\frac{\partial}{\partial t} = -i\omega$$

$$-i\omega \vec{v} + \vec{v} \times \vec{\omega}_c = -\frac{e\vec{E}}{m} - \nu \vec{v}$$

↑ collision term

$$(\nu - i\omega) \vec{v} + \vec{v} \times \vec{\omega}_c = -\frac{e\vec{E}}{m}$$

Z component

$$v_z = -\frac{e E_z}{m(\nu - i\omega)}$$

x-comp.

$$-i(\omega + i\nu)v_x + v_y \omega_c = -\frac{e E_x}{m}$$

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So, by replacing $\frac{\partial}{\partial t}$ by $-i\omega$, my equation of motion becomes, $-i\omega v + v \times \omega_c$ is equal to $-\frac{eE}{m}$. If there were collision, then I should use a collision term also which will be $-\nu v$ is the collision term, I can add. And this can be combined easily with this v term here. So, it becomes this is the collision term, in many cases you can ignore this, but for the sake of completeness one can include this here. So, you get, $(\nu - i\omega)v + v \times \omega_c$ is equal to $-\frac{eE}{m}$.

I will write down different components of this equation. Let me begin with the simplest one, which is z component, because this term vanishes. So, v_z turns out to be, $\frac{eE_z}{m(\nu - i\omega)}$ with a negative sign. Then, write down this x component of equation, I can write down this x component is equal to $-\frac{eE_x}{m}$, actually take $-i$ out, $(\omega + i\nu)v_x + v_y \omega_c$ is equal to $-\frac{eE_x}{m}$, this is my x component equation.

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$$-i(\omega + i\nu)v_y - v_x\omega_c = -\frac{eE_y}{m}$$

$$v_x = \frac{e[-i(\omega + i\nu)E_x - \omega_c E_y]}{m[(\omega + i\nu)^2 - \omega_c^2]}$$

$$v_y = \frac{e[-i(\omega + i\nu)E_y + \omega_c E_x]}{m[(\omega + i\nu)^2 - \omega_c^2]}$$

$$\vec{J}_e = -ne\vec{v} = \underline{\underline{\sigma}}_e \cdot \vec{E}$$

$$J_{ex} \equiv \sigma_{exx}E_x + \sigma_{exy}E_y + \sigma_{exz}E_z$$

Similarly, I will write down the y component, which turns out to be, equal to minus i omega plus i nu into v y minus v x into omega c is equal to minus e E y upon m. These equations are coupled algebraic equations for v x and v y, which can be solved to obtain v x and v y. And I will simply write the value of v x and v y separately. This turns out to be, is equal to e upon m omega plus i nu whole square minus omega c square multiplied by minus i omega plus i nu into E x minus omega c E y, this sort of expression you get.

And v y, can simply be written, v x is equal to this expression, I have correctly written and v y is the same similar expression, e upon m omega plus i nu whole square minus omega c square into minus i omega plus i nu into E y plus omega c E x. Once, we know the velocities of the particles, drift velocities of the electrons, you can write down the current density. Current density due to electrons, I will write as minus n e v. You may note one thing, v x involves not only E x, it involves E y also. V y involves not only E y, it involves E x also. Means, one component of J will involve not only the electric field that component, but other components of electric field also. So, J x will involve J y besides E x, J y will involve E x besides E y and so on.

So, in general, if I substitute this v in this one, I can write down, sigma e electron conductivity dot E. It is a symbolic representation; sigma is a tensor called conductivity tensor. And this implies that, if I have to write down, J e electron current density x component, then, I must write down the x component here. Now, this tensor has two

indices, sigma e so, I should write down the first index, is the same as of x of J and other is running. So, like this is, J x x E x plus sigma e x y E y plus sigma e x z E z. This is the meaning of the y definition. So, in tensor representation, J x is a scalar quantity, so I should not put a vector sign there. Like this similarly, you can write down J e y. So, if you compare the coefficients of E x in x component of current density, you will write down the value of sigma e x x, sigma e x y, sigma e x z and if you evaluate these quantities, you get this expression.

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$$\sigma_e = \begin{pmatrix} \sigma_{exx} & \sigma_{exy} & 0 \\ -\sigma_{exy} & \sigma_{exx} & 0 \\ 0 & 0 & \sigma_{ezz} \end{pmatrix}$$

$$\sigma_{exx} = \frac{ne^2 i (\omega + i\nu)}{m [(\omega + i\nu)^2 - \omega_c^2]} \approx \frac{ne^2 i \omega}{m (\omega^2 - \omega_c^2)}$$

$$\sigma_{exy} = \frac{ne^2 \omega_c}{m [(\omega + i\nu)^2 - \omega_c^2]}, \quad \sigma_{ezz} = \frac{ne^2}{m (\nu - i\omega)}$$

$$= \frac{ne^2 \omega_c}{m (\omega^2 - \omega_c^2)}$$

The expression is, sigma is a tensor electron conductivity, which is equal to, the following components are finite, sigma e x x, this is sigma e x y, sigma e x z is 0. The sigma e y x component is equal to sigma e x y with a negative sign. The sigma e y y is equal to sigma e x x. This is 0, this is 0, this is 0 and this is sigma e z z. Where these values are, sigma e x x turns out to be equal to n e square upon m and here is, omega plus i nu whole square minus omega c square multiplied by i omega plus i nu. This is correct expression.

And sigma e x y is equal to n e square upon m omega plus i nu whole square minus omega c square into omega c. And sigma e z z is equal to n e square upon m nu minus i omega. So, this is a tensor and the tensor has non-diagonal terms, finite and they are of opposite sign. The diagonal terms are, two of them are equal, but the third one is not

equal. So, this is a characteristic, this anisotropy is a characteristic of a magnetic field or a magnetized plasma.

You may note one thing more here, if you forgot collisions, because collisions are usually a very small quantity. If you forgot collisions, then, this can be written as, $n_e^2 \frac{\omega^2}{m \omega^2 - \omega_c^2}$. There is a cyclotron resonance effect here, when the frequency of the electric field is close to the cyclotron frequency of electrons, in that case, this quantity could be very huge. And similarly, this quantity, $n_e^2 \frac{\omega_c^2}{m \omega^2 - \omega_c^2}$ could also be huge. So, this is one important consequence of applying a magnetic field, so magnetic field not only confines the plasma, it can enhance the conductivity immensely, in the vicinity of cyclotron frequency.

Secondly, if ω is less than ω_c , much less than ω_c , then, σ_{xx} goes as m into ω_c^2 . In the case of ions, if you write ion conductivity, these expressions are valid. Just replace mass of the electrons, by mass of the ion, and ω_c by ion cyclotron frequency, which is very small quantity. And it is possible that ion conductivity could be comparable or even bigger than the electron conductivity that is a very important consequence of magnetic field. And this, presence of the cyclotron resonance effect at electron cyclotron frequency, as well as, at ion cyclotron frequency gives rise to very important wave phenomena, that are useful in exploring space plasmas, as well as, laboratory plasmas. It has tremendous consequence on wave propagation, there are lots of waves in a magnetized plasma, that are possible because of these cyclotron effects.

So, plasma response is very strongly influenced by the presence of magnetic field and ion dynamics cannot be ignored in general. Obviously, at high frequencies, when ω is comparable to ω_c , much bigger than the ion cyclotron frequency, you can ignore the ion motion, but at lower frequencies you cannot. So, we have to study a little more carefully the consequence of magnetic field in plasmas.

This is one thing, the derivation of conductivity tensor. It will be very useful, if you can go step by step and appreciate this. Because when we study the structures, which have non-uniform magnetic field, then, the particle dynamics becomes a little more

complicated, but the understanding of this tensor will help us in understanding more complex geometries.