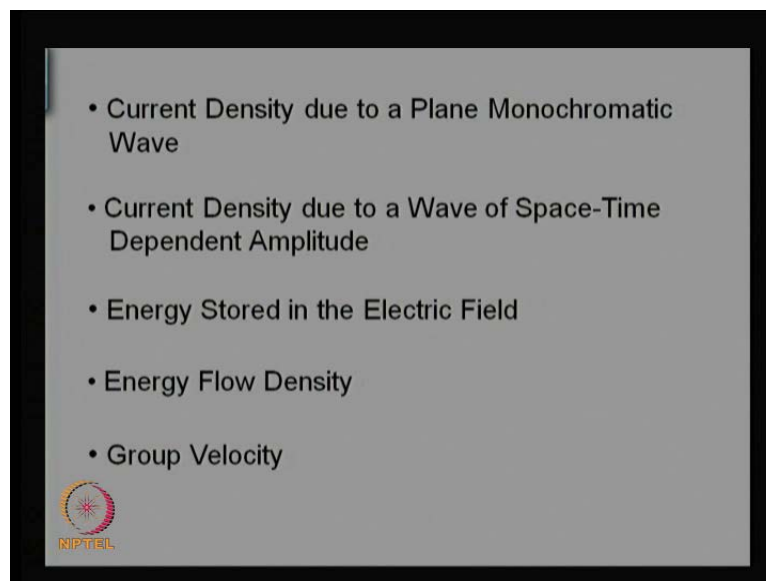


Plasma Physics
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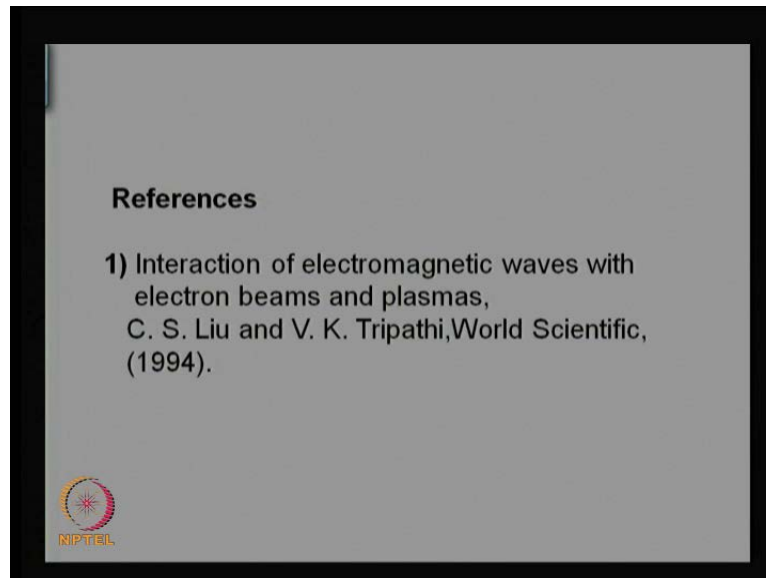
Lecture No. # 11
Energy Flow With an Electrostatic Wave

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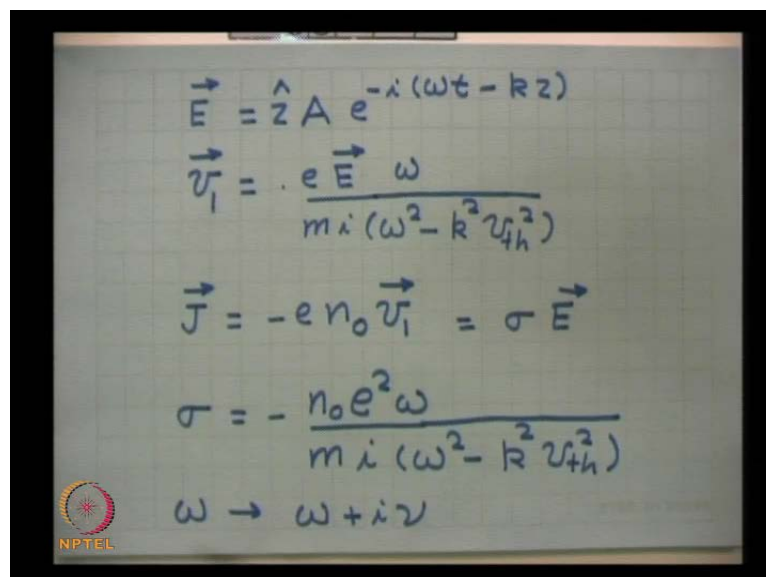
Well friends. Today I would like to discuss the energy flow with an electrostatic wave. We shall discuss current density due to a plane monochromatic wave, then current density due to a wave whose amplitude is a function of space and time, and then deduce an expression for energy stored in the electric field of a wave. And deduce the expression for energy flow per unit area per unit time with the electro static wave, and we shall also discuss the group velocity.

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Well, this is based on this book interaction of electromagnetic waves with electron beams and plasmas by Professor C.S. Liu and myself. However, you can find discussion of today's theme in any book in plasma physics.

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Well, let me begin with the response of a plasma to a monochromatic plane electrostatic wave whose electric field is E is equal to A to be simple or a specific lets the wave is longitudinally polarized in the z direction, and amplitude is constant in the time variation is $i \omega t$, and space variation is like $k z$. This is the electric field of an electrostatic wave, we have already

solved the equation of motion for electrons including the pressure term, and we deduced an expression for electron velocity from there we wrote that the perturbed velocity of electrons due to this waves is equal to $\frac{e E}{m i \omega^2 - k^2 v_{\text{thermal}}^2}$ into ω .

If thermal velocity is not there this is the thermal velocity of electrons then, this is the same expression for drift velocity as due to an electromagnetic wave, but the electrostatic wave is difference because of the thermal effects the denominator is modified and from there if you write down the perturbed current density due to the wave which is a product of electron charge equilibrium electron density, and perturbed velocity or drift velocity of electrons.

If you substitute this there then, this is expressible as σ into E , and σ turns out to be $\frac{-n_0 e^2}{m i \omega^2 - k^2 v_{\text{thermal}}^2}$ and ω in the denominator upon $m i \omega^2 - k^2 v_{\text{thermal}}^2$ in this derivation I have ignored the effect of collisions. If collisions are finite then wherever ω occurs you have to replace ω by $\omega + i \nu$ then this equation is valid for collisional plasma also. Usually when we are talking of Linear waves ω is around ω_p , and collision frequency is orders of magnitude is smaller than ω_p , hence this effect is small though it may be significant, but it is small. This is the current density due to an electrostatic wave whose amplitude is constant in time as well as in space. In order to deduce an expression for energy flow with the wave, we like to examine the current density of plasma response to a wave whose amplitude has a slow dependence, but significant dependence on time and position. So, let me go over to that stage.

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Response when wave amp.
is a function of time
and space

$$\vec{E} = \hat{z} A(t, z) e^{-i(\omega t - kz)}$$
$$\frac{\partial \vec{v}_1}{\partial t} = -\frac{e\vec{E}}{m} - \frac{v_{th}^2}{n_0} \nabla n_1$$
$$\frac{\partial n_1}{\partial t} + \nabla \cdot (n_0 \vec{v}_1) = 0$$

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Now, I would like to discuss the electron response when wave amplitude is a function of time and space. I think we shall learn something very significant from this deduction. I am considering that my system has an electric field which is directly along z because it is a longitudinal electrostatic wave propagating on z axis, but amplitude I will permit to have a dependence of time and z Exponential minus i omega t minus kz . So, besides having fast phase dependence like this the way has a slow dependence on time and amplitude in terms of it is amplitude. Let us see what the consequence of this. If you write down the equation of motion is if I divide the equation of motion by mass it reads as Δv by Δt is equal to minus eE upon m minus pressure term which reads as v_{th}^2 which is te upon m upon n_0 and gradient of n .

Here I have used the linearized equation of motion. Where $v \cdot \text{Del } v$ term has been dropped means the products of perturbed quantities have been dropped. This is what we call as a linearized equation of motion that we had written last time. And actually this is for the perturbed velocity and perturbed density. So, I have written like this. And the equation of continuity for the perturbed density is Δn_1 by Δt plus divergence of $n_0 v_1$ is equal to 0. I wish to solve these equations when n_1 is not a constant, but a slowly varying function of t and z . I will slowly following a procedure of iteration solve these equations and derive some a result and then I will point out a general procedure which is very simple, but in order to appreciate that simple procedure. I need to do some deduction from these equations. I will be little slow and elaborate in deriving or rather in obtaining the solution to these equations.

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$$\vec{v}_1 = \hat{z} a(t, z) e^{-i(\omega t - kz)}$$

$$n_1 = b(t, z) e^{-i(\omega t - kz)}$$

$$\frac{\partial a}{\partial t} < \omega a, \quad \frac{\partial b}{\partial t} < \omega b$$

$$\frac{\partial a}{\partial z} < k a, \quad \frac{\partial b}{\partial z} < k b$$

$$-i\omega a + \frac{\partial \vec{a}}{\partial t} = -\frac{eA}{m} - \frac{v_{th}^2}{n_0} (ikb + \frac{\partial b}{\partial t})$$

$$S_1 = -\frac{\partial a}{\partial t} - \frac{v_{th}^2}{n_0} \frac{\partial b}{\partial t}$$

What I do I express my perturbed electron velocity n in terms of some quantity a and I am taking all quantities vector quantities along z because it is a longitudinal wave. Suppose this is slowly varying amplitude of velocity which is upon t and z exponential minus i omega t minus k z . I have written the past time and z dependence the same way as the electric field because there is a source there is a response, but I am permitting my amplitude to have some dependence on t and z . Then I can write down my perturbed density also as say b which is a function of t and z exponential minus i omega t minus k z . Now what I do, I substitute these two expressions in the equation of motion continuity and treat Δa by Δt to be small as compared to minus i omega and Δb by Δt also to be small.

So, I am presuming that Δa by Δt is less than omega into a and Δb by Δt less than omega v . This is called iteration approximation. What I am going to do, first of all I will ignore the dependence of a on time and b on time and similarly I am assuming there z derivative also to be small as compared to k , means Δa by Δz is less than k a and Δb by Δz less than k b . This is the approximation I am going to make. What I do first of all when I substitute this I ignore the derivatives of a and b in the equation of motion continuity then those equations read like this. Let me actually substitute this and see how those equations look when you substitute this expansion for v_1 in the equation of motion you get minus i omega a plus Δa by Δt . This is Δa by Δt term exponential factor is

common and will cancel out in both sides is equal to minus e into A minus the pressure term which is v thermal square upon $n_0 \Delta z$ of n_1 .

So, when you m also is here. When you evaluate that term there are two terms $i k$ into b plus Δb upon Δt . This is the term and this is term. These two terms together are called perturbed terms and I can call them if I take this term on the right hand side and add these two terms together I will call them as S_1 . Then this equation becomes, let me define S_1 is equal to minus Δa by Δz minus b thermally square upon $n_0 \Delta v$ by Δt . These two terms must I ignore and obtain the value of b in terms of A ; obviously, there is b here.

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$$-i\omega b + \frac{\partial b}{\partial t} + n_0 (i k a + \frac{\partial a}{\partial z}) = 0$$

$$S_2 = -\frac{\partial b}{\partial t} - n_0 \frac{\partial a}{\partial z}$$

$$-i\omega a + i k \frac{v_{th}^2}{n_0} b = -\frac{eA}{m} + S_1$$

$$i k n_0 a - i\omega b = S_2$$

$$b = \frac{k}{\omega} n_0 a, \quad a = + \frac{eA\omega}{m i (\omega^2 - k^2 v_{th}^2)}$$

So, I should write down the equation of continuity as well and the equation of continuity is, if you substitute the expansion for density it gives you minus $i \omega b$ plus Δb by Δt plus n_0 divergence of Δz of v_1 which gives me $i k a$ plus Δa by Δz is equal to 0. It looks little complicated, but simple substitution of v_1 and n_1 in the equation of motion and continuity. Here also there are two small terms Δv by Δt and Δa by Δz . I can take these two terms on the right hand side and define the sum of these 2 as S_2 which is equal to minus Δb by Δt , because I take on the right hand side and then this becomes and there is another term minus $n_0 \Delta a$ by Δz .

This is the second term S_2 that I get what you should have is in terms of S_1 and S_2 the equation of motion then is minus $i \omega a$ plus $i k v$ thermal square upon n_0 into b is equal to minus $e A$ upon m plus S_1 . The equation of motion takes this form and the equation of

continuity this one, when I transfer these partial derivative terms from the right hand side takes the form $i k$. This term I will write first $n_0 a$ then minus $i \omega b$ is equal to S_2 . What I am suggesting is that first ignore S_1 and S_2 , when you ignore S_1 and S_2 , you can solve these equations to obtain b is equal to from this equation k upon ω into $n_0 a$ and when you use this value of b in this equation and ignore this one you will get a is equal to plus $e A$ ω upon $m i \omega^2 - k^2 v_{th}^2$.

This is the same velocity. If you multiply the exponential term as we had obtained in the local approximation when the amplitude was a constant and this is the density perturbation when you multiply the phase term. So, as before, we obtain the same thing when you ignore this S_1 and S_2 , but one use this value of b and a in S_2 and S_1 . S_2 is put this value here and this value of a here and solve these two equations you can get in the new case when S_1 and S_2 are retained. This is when S_1 and S_2 are ignored. If S_1 is equal to 0 and S_2 is equal to 0 these are the values of b and a . Then I am suggesting that now evaluate S_1 and S_2 using these values of b and a and then solve these two equations returning these two terms and you will get the following. I just give the result, this algebra is simple and I am pretty sure you can easily simplify this. The net result is this.

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$$a = \frac{e \omega A}{m i (\omega^2 - k^2 v_{th}^2)} - \frac{e}{m (\omega^2 - k^2 v_{th}^2)}$$

$$\left[(\omega^2 + k^2 v_{th}^2) \frac{\partial A}{\partial t} + 2 k \omega v_{th}^2 \frac{\partial A}{\partial z} \right]$$

$$J_2 = - e n_0 a e^{-i(\omega t - k z)}$$

$$= - \frac{n_0 e^2 \omega E_2}{m i (\omega^2 - k^2 v_{th}^2)} + \frac{n_0 e^2}{m (\omega^2 - k^2 v_{th}^2)} \left[(\omega^2 + k^2 v_{th}^2) \frac{\partial A}{\partial t} + 2 k \omega v_{th}^2 \frac{\partial A}{\partial z} \right] e^{-i(\omega t - k z)}$$

So, to the next approximation when you evaluate a and b . They turn out to be a is equal to $e \omega A$ upon $m i \omega^2 - k^2 v_{th}^2$ where v_{th} is the electron thermal velocity this term as before. Now the additional terms are these minus e

upon $m\omega^2 - k^2 v^2$ thermally square multiplied by I will write down this in different ink because these are the 2 important contributions. Because of S_1 and S_2 which is $\omega^2 + k^2 v^2$ thermally square into Δa by Δt plus twice $k\omega v$ thermally square Δa by Δz . This is an important addition to the solution, this is if you ignore these two terms this last two terms then a is the same as due to a uniform amplitude wave. And these are the contributions that take into account that time variation of amplitude and space variation of amplitude. These two terms together. They are important terms and we shall learn in little while the role played them in energy flow.

Well, if you know the velocity the current density can be written and current density J_z in this case will be $-\epsilon_0 n_0 v$. Velocity is $a \exp(-i\omega t - kz)$. This is the value of J_z and when you substitute the value of a here you get an expression which is equal to $-\epsilon_0 n_0 e^2$ just multiply $n_0 e$ here. Because $-\epsilon_0 n_0 e^2$ square upon $m(i\omega^2 - k^2 v^2)$ thermally square into ω this ω into a into this quantity is simply $e E_z$. This is first term and second term is plus because plus $\epsilon_0 n_0 e^2$ square upon $m(\omega^2 - k^2 v^2)$ thermally square multiplied by these two terms in the bracket. Let me write them with red ink $\omega^2 + k^2 v^2$ thermally square Δa by Δt plus twice $k\omega v$ thermally square Δa by Δz into exponential term this exponential. So, I do not have room here to write exponential. So, I will write e and then put dot there this exponential factor phase factor is to be there.

Now the beauty here is that if you call the coefficient of e_z as σ then, you can easily see that this quantity this term which is multiplying Δa by Δt is simply $\Delta \sigma$ by $\Delta \omega$ with an i and similarly this quantity is exactly same as $\Delta \sigma$ by Δk you can just verify. Let me write this, I will define this coefficient from here through here this quantity is conductivity. If I called σ then these other coefficients can be easily shown and you must show yourself If must check yourself.

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The whiteboard shows the following equations:

$$J_z = \sigma E_z + i \frac{\partial \sigma}{\partial \omega} \frac{\partial A}{\partial t} e^{-i(\omega t - kz)}$$

$$+ i \frac{\partial \sigma}{\partial k} \frac{\partial A}{\partial z} e^{-i(\omega t - kz}$$

$$\vec{E} = \hat{z} A(t, z) e^{-i(\omega t - kz)}$$

$$\frac{\partial \vec{E}}{\partial t} = (-i\omega A + \frac{\partial A}{\partial t}) \hat{z} e^{-i(\omega t - kz)}$$

Below these equations, there are substitutions for the frequency and wave number:

$$-i\omega \rightarrow -i\omega + \frac{\partial}{\partial t} \quad \text{or} \quad \omega \rightarrow \omega + i\frac{\partial}{\partial t}$$

$$k \rightarrow k - i\frac{\partial}{\partial z}$$

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That J can be written as σE_z plus $i \frac{\partial \sigma}{\partial \omega} \frac{\partial A}{\partial t}$ exponential minus $i\omega t$ minus kz , and another term plus $i \frac{\partial \sigma}{\partial k} \frac{\partial A}{\partial z}$ exponential minus $i\omega t$ minus kz .

This is a very simple comprehensive way of writing an expression for J, because only in terms of a single quantity called conductivity and its derivatives. I had been able to write J due to an electrostatic wave whose amplitude is a function of time and z. Of course, I have assumed that σ depends slowly on time and slowly on z. This is the approximation I have made, but in that limit this is good. Well this was a very special case of a wave when there was no static magnetic field when I did not consider collisions and so on. But in general I like to point out a general procedure. Whenever you are encountered with waves of time varying amplitude and position what to do if you know the response to a monochromatic wave whose amplitude is constant you can easily deduce the response due to time varying amplitude and time varying and the space varying amplitude the procedure is as follows.

You know if you have electric field is equal to say some function of t and z exponential minus $i\omega t$ minus kz , suppose this is my electric field. If I differentiate this partially with time I will get $\frac{\partial E}{\partial t}$ how much that would be. That would be simply first of all I will differentiate this quantity will give you minus $i\omega$ into this A will be there and then I will differentiate this. We will get $\frac{\partial A}{\partial t}$ and of course, z unit vector is

there and exponential minus $i\omega t - kz$ is there. What I am saying, that if the wave has constant amplitude then this term is 0. So, δt is simply equal to minus $i\omega$. But when the wave has a time varying amplitude then δt is not equal to this, but minus $i\omega$ plus δt has to be there means ω has to be replaced by $\omega + i\delta t$ or minus $i\omega$ has to be replaced by minus $i\omega + \delta t$.

If you compare this case with the case when there is no time variation of amplitude. Then this term is not there. The difference in the two cases is that minus $i\omega$ has to change by this or simply if I divide this equation by minus i , then this says that ω just to replace by $\omega + i\delta t$. Because if I divide this equation by minus i this becomes ω goes to $\omega + i\delta t$. And this δt operator has to operate not on the whole field, but only on the amplitude. In general your σ the conductivity is a function of ω and k . But you should do, replace ω by $\omega + i\delta t$ and similarly if you differentiate this with respect to z you can show that k has to go from $k - i\delta z$. If you can make these two changes in σ and operate this σ over the amplitude of the electric field you will get the right expression.

So, this expression that I have written here complicated expression can be deduced from \vec{J} equal to $\sigma \vec{E}$ in a simple way if I substitute σ by this and k by this let us see how do we do it.

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The image shows a whiteboard with the following handwritten equations:

$$\vec{J} = \hat{z} \sigma(\omega, k) A e^{-i(\omega t - kz)}$$

$$J_z = [\sigma(\omega + i\frac{\partial}{\partial t}, k - i\frac{\partial}{\partial z}) A] e^{-i(\omega t - kz)}$$

$$= [\sigma(\omega, k) + \frac{\partial \sigma}{\partial \omega} i\frac{\partial}{\partial t} - i\frac{\partial \sigma}{\partial k} \frac{\partial}{\partial z}] A e^{-i(\omega t - kz)}$$

$$= \sigma E_z + [i\frac{\partial \sigma}{\partial \omega} \frac{\partial A}{\partial t} - i\frac{\partial \sigma}{\partial k} \frac{\partial A}{\partial z}] e^{-i(\omega t - kz)}$$

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We begin with J is equal to σ which is a function of frequency and k in general and this direction is z into a exponential minus $i\omega t - kz$, but what I am saying that I am going to replace this ω by let me write down J_z as how much then σ which is a $\omega + i\delta t$ and k is to be replaced by $k - i\delta z$ operating over a exponential factor minus $i\omega t - kz$. Well, if this is a slow dependence as compared to ω as we have been assuming then you carry out the Taylor expansion. Any function, if there is a small increment in the variable can be expanded by using Taylor series. This expression becomes simply σ at ω and k plus there is an increment here. So, it becomes $\delta\sigma$ upon $\delta\omega$ into $i\delta t$ and because of this will be a term minus $i\delta\sigma$ upon δk into δz and these are operating over a into the exponential term. This same factor phase factor.

You may note that this is exactly same as before. Rather than going over the entire calculation and in this is valid in general. Even if there is a magnetic field using this is valid, this is a very general thing. This procedure is very similar to the procedure adopted by Schrödinger when he deduced the wave equation Schrödinger equation from de Broglie's hypothesis. Essentially, he said that ω should be replaced by an operator and you should also replace k by an operator. The same way I am suggesting here that ω can be replaced by because it is not operating over the entire electric field, it is operating this you know over the amplitude. I am replacing ω by $\omega + i\delta t$ because the time dependence of the phase term was already incorporated in deducing σ . It is only the time dependence of amplitude that was ignored. I have to replace this by ω by $\omega + i\delta t$ and so on.

This is a very standard procedure in all areas of physics whenever you are dealing with waves or fields whose amplitudes vary with time and position slowly. This is a 1 important thing that I wanted to mention to you. Well, now this question of simple interpretation I think if time permits, then I would like to say something more significant from in a different way you can arrive at the same expression. Let me go a step further, let me write this expression σ e bring this a in the interior, σ into e is z rather this becomes $i\delta\sigma$ upon $\delta\omega$ into δa by δz δt minus $i\delta\sigma$ upon δk and δa by δz and obviously, in these two terms I am multiplying an exponential term minus $i\omega t - kz$. I have to interpret this. Well, I am going to make an approximation here. As we have already seen the σ that we have derived here is almost an imaginary quantity. Because the expression for σ if you look is has a i there. It is imaginary. However, if we

include collisions then it may have been small real part because collisions are a small quantity.

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$$\sigma = \sigma_r + i \sigma_i, \quad \sigma_i \gg \sigma_r$$

$$\vec{J} + \frac{\partial \vec{D}}{\partial t} = 0$$

Multiply by $\cdot \vec{E}$ and take time av.

$$\text{Re } \vec{A} \cdot \text{Re } \vec{B} = \frac{1}{2} \text{Re} [\vec{A} \cdot \vec{B}^* + \vec{A} \cdot \vec{B}]$$

↓
time av.
= 0

I can write down conductivity a sum a small real part and a very large imaginary part. where sigma i is much bigger than sigma r or magnitude bigger, but let me retain this small sigma r to have a physical interpretation. Now for electrostatic waves there is no magnetic field in the wave. Convention point in vector is 0 e cross h is 0 because h is 0. Now, let us understand the implication of J from the last Maxwell equation if curl of h is 0 then J plus delta D by delta t is 0. The last Maxwell equation becomes this if h is 0. I want to multiply this equation by dot E and take the time average. So, multiply this equation by dot E and let us see and take time average. Let us see what I get, please remember J is always a real quantity though I have written in a complex way, but whatever expression I have written real part of that expression is J.

So, J is a physical quantity and electric field also is a physical quantity though written in terms of a complex expression just takes the real part. So, I am going to use this identity that real part of A into real part of B is equal to half real part of A dot B take complex conjugate of any quantity either A or B plus a dot b. This is the identity that we had proved earlier and you can verify even now. Take A as small a 1 plus I a 2 and B as small b 1 plus I b 2 where a 1 b 1 a 2 b 2 are all real you can verify this easily. Now electric field changes with time exponential minus i omega t minus k z first time dependence. Similarly, J also has the similar time

dependence. So, if I replace this a by J and b by e, but I get this term will have 2 omega frequency variation and time average will be 0 and this term will have the exponential phase terms will cancel because this is a complex conjugate. When I multiply this equation by dot E only these terms survives this goes this goes away and let us see what the consequence is.

So, time average of this is 0. Let me just mention this. This is time average is 0, you can forget this term is 0 rather time average is not exactly 0 and how about this term. This is the only thing I have to retain a dot b star by 2 and take the real part. I will write down term by term J has three terms and this may also have two terms. Let me simplify these, I think I want to be slow on this because this derivation is very fundamental to plasma physics. I want you to follow it with greater **(audio not clear)**.

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$$\begin{aligned} \frac{\text{Re}[\vec{J} \cdot \vec{E}^*]}{2} &= \frac{1}{2} \text{Re} \left[\sigma |A|^2 + i \frac{\partial \sigma}{\partial \omega} \frac{\partial A}{\partial t} A^* \right. \\ &\quad \left. - i \frac{\partial \sigma}{\partial k} \frac{\partial A}{\partial z} A^* \right], \quad \sigma \approx i \sigma_r \\ &= \frac{1}{2} \sigma_r |A|^2 - \frac{\partial \sigma_r}{\partial \omega} \frac{1}{4} \frac{\partial}{\partial t} A A^* \\ &\quad + \frac{\partial \sigma_r}{\partial k} \frac{1}{4} \frac{\partial}{\partial z} A A^* \\ \text{Re} [A^* \frac{\partial A}{\partial t}] &\equiv \frac{1}{2} [A^* \frac{\partial A}{\partial t} + A \frac{\partial A^*}{\partial t}] \end{aligned}$$

Let me write down J dot E a star and I take the real part of this quantity and divide by 2. This is what I am going to evaluate if I do this the first term J is equal to sigma e. It becomes real part half of sigma into e dot e star which is the same thing as modulus of A square which is a real quantity. This real part means real part of sigma will survive. Second term was plus i delta sigma by delta omega into delta A by delta t.

If you recall I am multiplying this by E star they are on the same direction, just multiply by a star. A could be complex in general and similarly, there is another term which is minus i delta sigma by delta k delta A by delta z into A star, this is the meaning of real part of J dot E star by 2 means the time average value of J dot E is this the quantity. As, I mentioned sigma here

is largely σ_i . I replace by i times σ_i . This becomes means if I take the real part of this quantity it becomes half σ_r modulus of a square because this is real. So, real part will be σ_r only. Only real part will survive. And σ_i is because σ in these expressions is nearly equal to i times σ_i . If I substitute this σ in here and in there then, this becomes also real quantity i and i will become minus. This becomes a real quantity; similarly this also becomes a real quantity and then I have to take the real part of this quantity, which is you can easily show is equal to minus $\frac{1}{4} \frac{\Delta \sigma}{\Delta t} \sigma_i$ rather $\frac{1}{4} \frac{\Delta \sigma}{\Delta t} \sigma_i$.

And similarly this term becomes when I plus $\frac{1}{4} \frac{\Delta \sigma}{\Delta t} \sigma_i$ upon $\frac{1}{4} \frac{\Delta \sigma}{\Delta t} \sigma_i$ of a star. You, may wonder why I am getting 4 not 2. The reason is that I am getting a term here. I think one should understand this with clarity that real part of σ_i $\frac{1}{4} \frac{\Delta \sigma}{\Delta t} \sigma_i$, I have to evaluate which is the same thing as half this number a star $\frac{1}{4} \frac{\Delta \sigma}{\Delta t} \sigma_i$ plus. It is complex conjugate. Any complex numbers real part is half of that number plus it is complex conjugate. Complex conjugate of this quantity means replace this σ_i by σ_i and replace the un-starred quantity by the starred quantity. This should go away and this should be there this should not be there and this can be easily shown is equal to half $\frac{1}{4} \frac{\Delta \sigma}{\Delta t} \sigma_i$ same thing. I have used this is identity you can identical. This is a very important expression. It has three terms the one which involves the real part of conductivity it may be tiny, but it may this term may be comparable to this term because σ is σ a very slowly varying function of time.

Though σ_i may be large, but because the slow dependence of σ_i on time this term may be these terms may be comparable. Let us understand now the physical implication of the fourth Maxwell equation that I have written as $\nabla \cdot \mathbf{J} + \frac{d\rho}{dt} = 0$. $\nabla \cdot \mathbf{e}$ star i have obtained I have yet to obtain the other term $\frac{d\rho}{dt} \cdot \mathbf{e}$ star.

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The whiteboard shows the following steps:


$$\vec{J} + \frac{\partial \vec{D}}{\partial t} = 0$$

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$$\epsilon_0 (-i\omega \vec{E} + \frac{\partial A}{\partial t} e^{-i(\omega t - kz)})$$

$$\frac{1}{2} \text{Re} [\vec{E} \cdot \frac{\partial \vec{D}}{\partial t}] = \text{Re} \frac{\epsilon_0}{2} [A^* \cdot \frac{\partial A}{\partial t}] = \frac{\epsilon_0}{4} \frac{\partial |A|^2}{\partial t}$$

$$\frac{1}{2} \sigma |A|^2 + \frac{\epsilon_0}{4} (1 - \frac{\partial \sigma_i}{\partial \omega} \frac{1}{\epsilon_0}) \frac{\partial |A|^2}{\partial t} + \frac{1}{4} \frac{\partial \sigma_i}{\partial R} \frac{\partial |A|^2}{\partial z} = 0$$

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I am always keeping in focus $\vec{J} + \frac{\partial \vec{D}}{\partial t} = 0$. This term is how much \vec{D} is ϵ_0 into $\frac{\partial \vec{E}}{\partial t}$ which is equal to $-i\omega \vec{E} + \frac{\partial A}{\partial t} e^{-i(\omega t - kz)}$ and when you take $\vec{E} \cdot$ of this quantity and time average this quantity. If I take $\vec{E} \cdot$ and this becomes real, but this is imaginary, real part of this is 0. If I take half real part of $\vec{E} \cdot \frac{\partial \vec{D}}{\partial t}$ and take star here and then this term does not contribute. This is the only term that contributes and that gives you ϵ_0 is there; 2 is there and this becomes $A^* \cdot \frac{\partial A}{\partial t}$ and I take the real part of this quantity. Real part of this quantity has to be taken, which is, the same thing as ϵ_0 by 4 $\frac{\partial |A|^2}{\partial t}$. $\vec{J} \cdot$ to be obtained earlier $\vec{E} \cdot \frac{\partial \vec{D}}{\partial t}$ we have obtained now add the two terms and put them equal to 0, what do you get essentially this expression.

If you just club these terms together what you get is this half $\sigma |A|^2 + \frac{\epsilon_0}{4} (1 - \frac{\partial \sigma_i}{\partial \omega} \frac{1}{\epsilon_0}) \frac{\partial |A|^2}{\partial t} + \frac{1}{4} \frac{\partial \sigma_i}{\partial R} \frac{\partial |A|^2}{\partial z} = 0$. This has a simple interpretation because when there is conductivity real part means the current density has a component in phase with the electric field there is power dissipation. This is the power dissipation per unit volume. Then, this term should be something like rate of increase of energy density of the electric field. This factor multiplied by modulus of a square can be interpreted as energy density of

the electric field and this term should be then expressible or can be interpreted as the pointing representing divergence of pointing flux.

This equation has lot of physical implication. Let us try to appreciate in plasma physics we do not talk that much in terms of conductivity. As, we talk in terms of effective plasma permittivity, but the two are unrelated as I defined earlier.

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$$\epsilon_{eff} = 1 + i \frac{\sigma}{\omega \epsilon_0}$$

$$= 1 - \frac{\sigma_i}{\omega \epsilon_0}$$

$$\sigma_i = \omega \epsilon_0 (1 - \epsilon_{eff})$$

$$\frac{1}{2} \sigma_r |A|^2 + \frac{\partial}{\partial t} W_E + \frac{\partial}{\partial z} (W_E v_g) = 0$$

Effective plasma permittivity, we define as 1 plus i sigma upon omega epsilon 0 and sigma is almost imaginary at least in these terms that are appearing here the time derivative amplitude and space derivative amplitude in those terms. This is nearly equal to 1 minus sigma i upon omega epsilon 0, from here, I can write down sigma i is equal to omega epsilon 0 into 1 minus epsilon effective. This is a better way of expressing that equation rather than in terms of sigma. I would like to write that equation in terms of effective plasma permittivity and if you do this, the equation becomes rather simple. Half sigma r modulus of A square plus delta t of I will call a term w E plus delta z of w E into v g equal to 0. That equation can be put in this form; however, you may wonder there was delta k of epsilon effective. Where that is and what is W E. you just switch off this sigma i in that expression you will find it.

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$$W_E = \frac{1}{4} \epsilon_0 \frac{\partial}{\partial \omega} (\omega \epsilon_{eff}) |A|^2$$
 Energy stored in the electric field per unit volume

$$\partial \sigma_i / \partial k \rightarrow - \omega \epsilon_0 \frac{\partial \epsilon_{eff}}{\partial k}$$

$$\epsilon_{eff}(\omega, k) = 0$$

$$\frac{\Delta \omega}{\Delta k} = v_g \quad \frac{\partial \epsilon_{eff}}{\partial \omega} \Delta \omega + \frac{\partial \epsilon_{eff}}{\partial k} \Delta k = 0$$

If W_E turns out to be $\frac{1}{4} \epsilon_0 \frac{\partial}{\partial \omega} (\omega \epsilon_{eff}) |A|^2$. Normally, a half vector comes in here, but because we have taken time average, half comes because of time average also. This is the energy stored in the electric field. This is called energy stored in the electric field.

Energy stored in the electric field of the wave per unit volume and what I have done regarding there was a term like $\Delta \sigma_i / \Delta k$ what did I do with that. It was replaced by ϵ_{eff} because σ_i when I wrote down in terms of ϵ_{eff} . It became is equal to $-\omega \epsilon_0 \frac{\partial \epsilon_{eff}}{\partial k}$ because they were constant. I am differentiating this expression partially (audio not clear) to k . So, ω can be treated to be constant and then $\frac{\partial \epsilon_{eff}}{\partial k}$. But one must remember one thing that ϵ_{eff} is a function of ω and k and for electrostatic wave this is always equal to 0. If I differentiate this function any increment in this function due to increase in ω in change in k would of this form $\frac{\partial \epsilon_{eff}}{\partial \omega} \Delta \omega + \frac{\partial \epsilon_{eff}}{\partial k} \Delta k = 0$ because this is 0 always at ω, k , as well as, at $\omega + \Delta \omega$ and $k + \Delta k$ also this is 0.

So, from here I can obtain the value of this expression in terms of $\frac{\partial \omega}{\partial k}$ of this quantity you may just check and then you can write down the $\mathbf{J} \cdot \mathbf{E}$ equation the way I had written and I have called $\frac{\partial \omega}{\partial k}$ or $\frac{\partial k}{\partial \omega}$ as v_g group velocity. This can be written as $\frac{\partial \omega}{\partial k}$.

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$$\frac{1}{2} \sigma_r |A|^2 + \frac{\partial W_E}{\partial t} + \frac{\partial (W_E v_g)}{\partial z} = 0$$

$\begin{array}{c} \rightarrow \\ z \end{array} \quad \begin{array}{c} | \rightarrow \\ z+dz \end{array}$

$\vec{S}_{av} = W_E v_g \hat{z}$

$v_g = \frac{v_{th}^2}{v_p}$

EM wave $\begin{array}{c} \updownarrow \\ \rightarrow \vec{k} \end{array}$
 ES $\begin{array}{c} \leftrightarrow \\ \rightarrow \vec{k} \end{array}$

Let me write that J dot E equation again, which was half sigma r A square plus delta delta t of W E plus delta delta z of W E into v g is equal 0. Now simple interpretation would say that, this is the amount of heat dissipation, the power lost by the wave per unit volume. This is the increase in energy density per unit volume. If in certain region suppose consider a region between this is z and z plus dz, if in certain volume if I multiply by the volume of this region dz each term by dz, and area I am taking unit area of cross section of this plane as unity. Then the volume element will be dz, multiply this by dz, then the first term gives you the energy dissipation per unit volume. This gives you the increase in energy per unit time in that region.

Energy is being lost here more over energy is increasing here. It means energy should be coming in here from somewhere. So, more energy should pour in and less should get out only then increase in energy density and dissipation energy is possible. This quantity must represent the energy flow. Because divergence of energy flow is represents like in pointing theorem also. This is some sort of a pointing vector, less quantity more quantity of this comes in less gets out only then there is an increase in energy density of the electric field and increase and power loss there. The energy flow density or pointing vector for the electrostatic wave would be the pointing vector would be S. I am talking of average pointing vector is equal to W E into v g, the direction of energy flow will be in the z direction in this case because my wave was taken to be in the two direction.

This is equivalent of pointing vector for electromagnetic wave this is e cross h where as for electrostatic waves for which magnetic field is not there is no magnetic perturbation time

average pointing vector is $\mathbf{W} = \mathbf{E} \times \mathbf{H}$, where W is the energy density of electric field per unit volume is energy stored in the electric field per unit volume. I have given you a derivation of two expressions one of W in a dispersive medium and second in the derivation of time average pointing vector. One thing, I would like to mention to you. For electrostatic waves the expression for conductivity was different than the expression for conductivity for electromagnetic waves. The thing was for an electromagnetic wave, if the direction of propagation is like this the electric field was in the transverse direction. So, there was no compression of charges this is EM wave.

Whereas for electrostatic wave your electric field is oscillating in the same direction as \mathbf{k} this is electrostatic wave and there is a compression rarefaction of charge and hence pressure term is significant. Here, there is no pressure term there. We had seen that in the case of electromagnetic waves conductivity does not involve any temperature, any thermal effects temperature comes through collision collisions only, but not through pressure, but here in case of electrostatic wave's temperature comes or thermal velocity comes through pressure term. It is a separate dependence. The two conductivities are different and hence epsilon effective are also different, and it will be useful for you to obtain the energy density for a plasma wave and energy density to a sound wave and also to calculate the group velocity. You can easily verify that for a plasma wave group velocity is equal to thermal velocity square upon v_{phase} .