

**Group Theory Methods in Physics**  
**Prof. P. Ramadevi**  
**Department of Physics**  
**Indian Institute of Technology, Bombay**

**Lecture – 46**

**Weight vector, Root vector, Comparison between SU (2) and SU (3) algebra**

Compare SU 2 and SU 3 and slowly take you to the abstract lie algebra. Simple lie algebra, what is simple lie algebra? It does not have nontrivial invariance of algebra, ok. So, I will take you from what you know and then we get on to SU 3 compare and contrast; some bit we did in the last lecture, but today we will compare and fix the root and weight vectors. So, this was the first thing that we all write the raising and lowering operators is SU2 algebra, we do not put that plus or minus 1. So, it basically shift the state  $j m$  by unit of 1 that is the operation of this raising and lowering operators, you all know this.

Similarly, we have for SU 3, the lowest nontrivial dimension is 3 cross 3 Hermitian traceless generators, number of independent real elements is 8 and out of that you can form two diagonal generators.



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**su(3) algebra**

- Here there are **8 generators**. Just like Pauli matrices for su(2), we have Gell-Mann matrices

$$\lambda_1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \lambda_2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \lambda_3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$
$$\lambda_4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}, \lambda_5 = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix}, \lambda_6 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$
$$\lambda_7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix}, \lambda_8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}.$$

Two diagonal generators - rank is 2

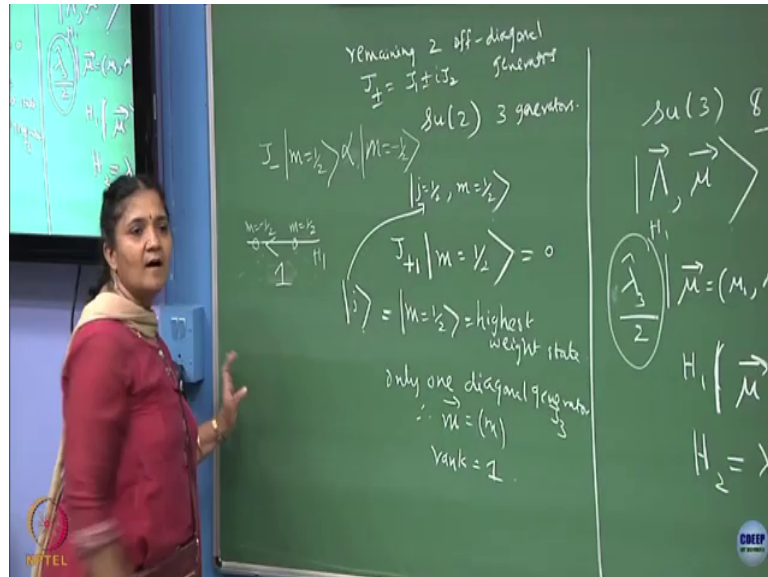


SU 2, you could form 2 cross 2 matrixes with one diagonal generator, here you can have two diagonal generator. Someone came and asked me why is this unique? So, the way to see it is that if you take this lambda 3, this part the 2 by 2 part is a sub subset sitting inside the su 3 ok. So, the whatever you did for sigma z polymatrix z component that will be sitting inside su 3 also.

The next one which you can construct should be such that that part should be same, ok. So, it is like an identity operator here, but then you have a nontrivial value so that this adds up to 0. Over all normalization is not important, but it is important in a sense that if I want to associate these numbers to some values like charge or other quantum numbers, these normalizations you can write it out later on. So, these are the eight generators of su 3 and 2 of

them are diagonal. So, any state if you want to write belonging to the su 3 irrep, then it will be a simultaneous eigenstate of lambda 3 and lambda 8, you all agree.

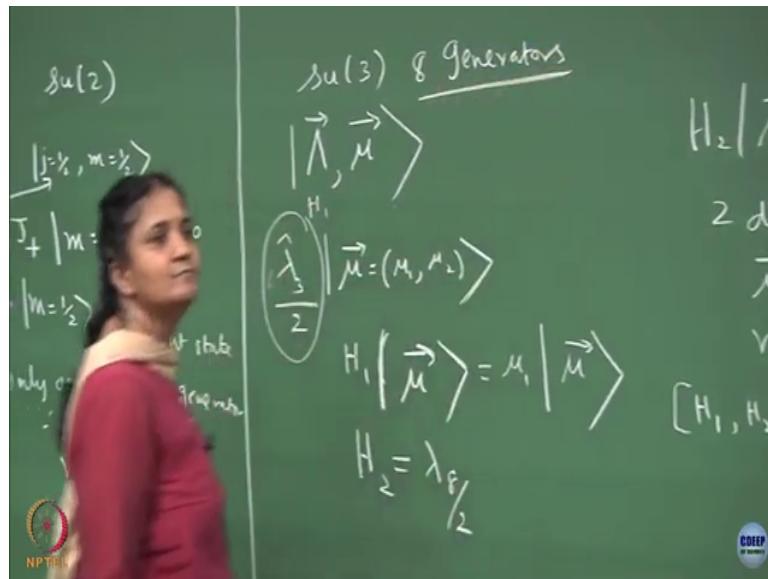
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So, su 2 su 3, you can have fundamental j equals to half, m equal to half, right. And the log of this j, I am going to call it as lambda here. I have explained what that lambda is I am going to call this as mu vector. This is also vector; we put this j equal to half in hindsight. The way we put this is that if J plus on m equal to plus half what will this be, 0 ok.

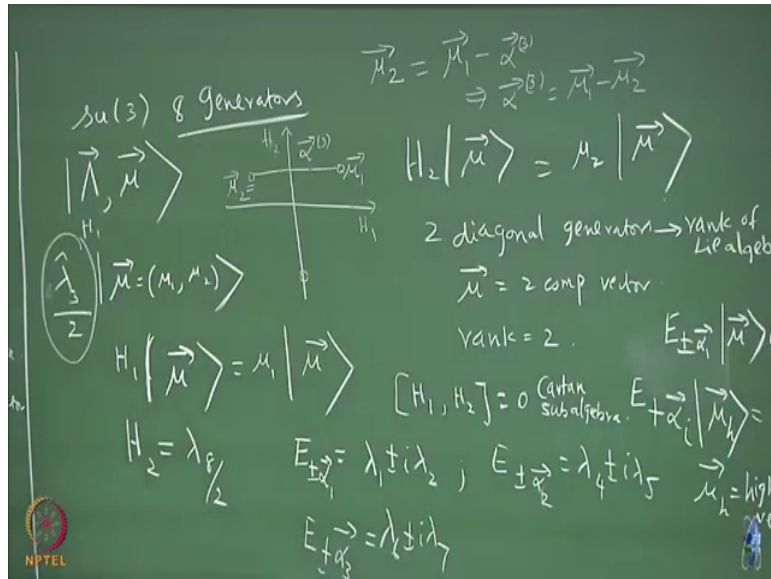
Then we call m equal to half is the highest weight state and this highest weight state is what we denote it as j clear. Anything which is any raising operator cannot take it further above that highest value is what we call it as j and we put that j value here. You do not need quadratic (Refer Time: 04:54) in group theory language to write the j value; j is the weight. If you do a raising operator j plus, it is going to be 0 ok.

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Similarly, here  $\mu$  is important.  $\mu$  has component  $\mu_2$  and you have the  $\lambda_3$  operator which is diagonal which I have called it as  $H_1$ . So,  $H_1$  if it operates on  $\mu$  vector, what will it give me? It will pull out a first component of this  $\mu$  vector as  $\mu_1$  times  $\mu$  vector.

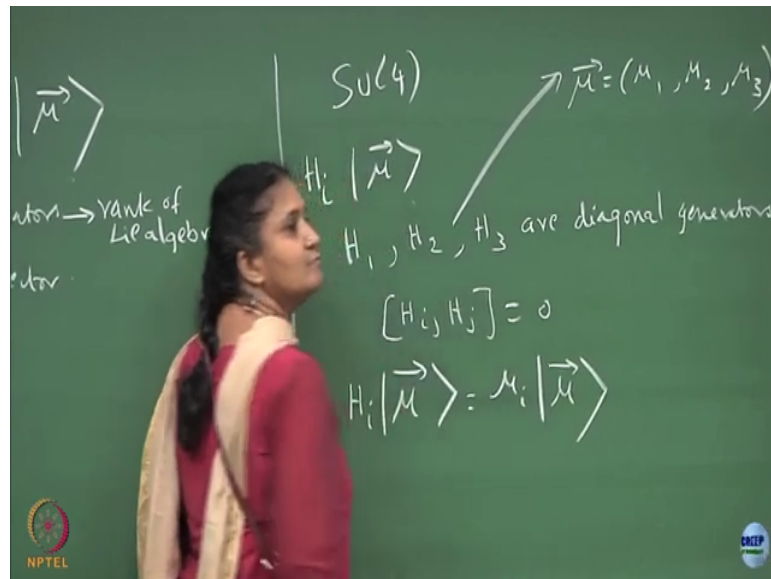
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Similarly,  $H_2$ ;  $H_2$  is what?  $H_2$  is  $\lambda_8$  by 2 and  $H_2$  on  $\mu$  vector will be  $\mu_2$  times  $\mu$  vector;  $\mu_2$  is the second component of the two vectors. So, there will be only two components for  $su(3)$ . So, here  $\mu$  is a number, it has only one component; weight vector is also just a weight, reason is only one diagonal generator, ok. So, that is why therefore,  $\mu$  vector is just one component, clear. Here,  $H_1$  and  $H_2$  are two diagonal generators two diagonal generators so, your  $\mu$  vector will be two component vector, ok.

And these two diagonal generators are sometimes we call this as rank of lie algebra. So, the rank is 2, number of diagonal generators what we call it as rank for  $su(3)$  it is 2. So, here rank is 1 and here rank is 2.

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Now tell me how you will extend it for su, you can do su 4, right. So, you will have a mu vector, you will have H 1, H 2 and H 3 are diagonal generators. So, any commutator of H i with H j will be 0. Here it is H 1 H 2 is 0, then any i and j which is 1 2 and 3 they will commute. So, mu vector if you operate H i, H i is either 1, 2 or 3 mu vector is how many component, 3 component here.

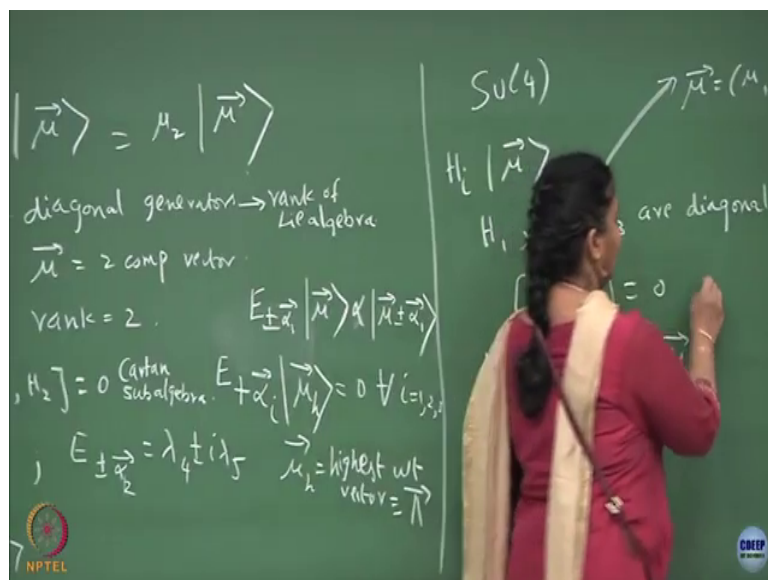
So, this implies mu vector will have mu 1, mu 2, mu 3 clear. So, this will tell you that H i on mu vector should be mu i times mu vector, clear. I have still not said about the lambda, it should be kind of clear to you just like j plus annihilates here. You should have the remaining generators, here there are 8 generators right, here there are 3 generators. One diagonal generators is gone, the two other generators there you write a complex conjugation of each other, right.

We write here the remaining two generators, 2 off-diagonal generators. What are they?  $J_1$  plus or minus  $J_2$  which is  $J_1$  plus or minus  $i J_2$ ; up to some normalization it is going to be  $J_1$  plus Hermitian conjugate is  $J_2$  minus. This is your diagonal generator we call it as  $J_3$ . This last column the first column, you all know. I am now making an extrapolation to the second which is the first non-trivial one and then generalizing into arbitrary  $su_n$  which you will have  $n - 1$  diagonal generators and how the weight becomes a weight vector. With how many components,  $n - 1$  component; is that clear, ok.

So, then out of this is what you did this to show that it is 0. So, here you have two diagonal generators, the remaining off-diagonal are 6 of them, right. So, let me call them as  $\lambda_1$  plus or minus  $i \lambda_2$  to be  $E$  plus or minus, I will explain what this means  $\alpha_1$  vector. You will have  $E$  plus or minus  $\alpha_2$  vector maybe which is  $\lambda_4$  plus or minus  $i \lambda_5$ . I will explain what this  $\alpha_1$  is,  $\alpha_1$  will be like your here technically I should put a plus 1. By that what I mean is that it will shift the  $m$  by  $m + 1$  units.

Similarly, this  $\alpha_1$  vector will shift the weight vector by  $\alpha_1$  that is a meaning of it, is that clear? So, we will do this elaborately and what else I have left there is one more  $\alpha_3$  which will be  $\lambda_6$  plus or minus  $i \lambda_7$ . So, the remaining 6 off diagonal generators can be grouped into Hermitian pairs; Hermitian conjugate pairs in this fashion where, these  $\alpha$  ones are again two component vectors. Because, the  $\mu$  vector is two components the rank of the group is 2. So, you will have two components.

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And what its job is that if you do  $E_{\alpha_1}$  plus or minus if you operate it on  $\mu$ , it has to give you  $I$  should not say equal to there will be some proportionality constants and all. It will just rise the weight vector by  $\alpha$  units or decrease it by  $\alpha$  units depending upon whether you do a  $\lambda_1 + i\lambda_2$  or  $\lambda_1 - i\lambda_2$ , ok.

Now tell me what will be the analog of  $j$  which I call it as  $\lambda$  here. What is the requirement?  $E_{\alpha_i}$  on  $\mu$  vector if it is 0 for all  $i$  equal to 1 2 and 3. So, let me call that with a subscript  $H$  just like here  $m$  equal to half. If you do there is only one rising operator in  $su_2$  that annihilates it, but there are 3 raising operators  $\lambda_1 + i\lambda_2$ ,  $\lambda_4 + i\lambda_5$  and  $\lambda_6 + i\lambda_7$  all of them. When it hit some weight vector it gives you 0, then you call this weight as highest weight highest weight vector and that is what you are denoting it by  $\lambda$  vector.



Just like here  $m$  equal to half is what you denote it by  $j$ ; if this happens, then you call that weight vector to be lambda vector and then you put in the state to be lambda to keep track that that is the highest weight vector. These are can be lower weight vectors which are obtained by you know the lowering operators, clear.

same thing holds there I am not going to do the  $su_4$ , but I guess from these two you understand what I am trying to tell. So, this diagonal generators algebra is called as Cartan sub algebra. These things I have already explained in the slides. This is called Cartan sub algebra, and this Cartans sub algebra rank is 2. There are 2 generators; this one will be 3 because there are 3 generators right,  $i$  and  $j$  1 2 1 3 for  $su_4$  ok, is this clear now? The extrapolation of a magnetic quantum number which you are all very familiar, now gets promoted to weight vector in the lie algebra language which will be two component for  $su_3$ .

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**su(3) algebra(contd)**

- The three basis states are denoted as

$$|\Lambda\mu_1\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, |\Lambda\mu_2\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}, |\Lambda\mu_3\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

- Here  $\mu_1, \mu_2, \mu_3$  are called two-component weight vectors given by eigenvalues of  $H_1 = \lambda_3/2$ ;  $H_2 = \lambda_8/2$

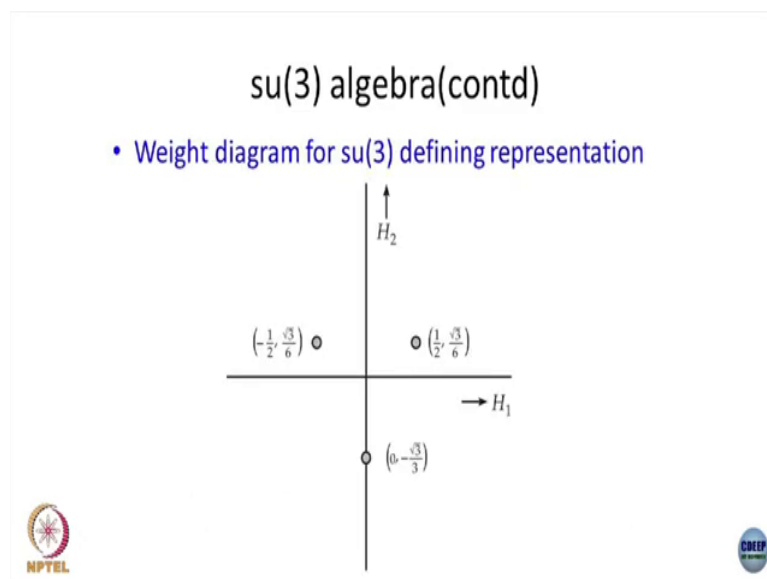
$|\Lambda\mu_1\rangle \equiv (1/2, \sqrt{3}/6); |\Lambda\mu_2\rangle \equiv (-1/2, \sqrt{3}/6); |\Lambda\mu_3\rangle \equiv (0, -\sqrt{3}/3)$

Specifically for the defining representation or the lowest non-trivial representation of  $su(3)$ , you can write these 3 basis.  $\lambda$  now I explain what that  $\lambda$  is and  $\mu$  turns out to be explicitly if I write those apply these matrixes on those two states on those basis,  $1\ 0\ 0$  if you apply. The  $H_1$  will give you a half,  $H_2$  will give you route 3 by 6, this we discussed last time.

So, that will be the weight vector corresponding to  $1\ 0\ 0$ . Similarly, the weight vector  $0\ 1\ 0$ , you can that is a basis on which the matrixes operate the corresponding magnetic quantum number which (Refer Time: 19:05) promoted to weight vector will be the 2 eigenvalues of the diagonal generators. So, that value is this, ok. So, this is the way you construct the weight vectors, the 3 weight vectors corresponding to the lowest non-trivial dimension of  $su(3)$  which is called a defining representation.

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And then you can plot the weights, these are the 3 weights vector states corresponding to the 3 basis states which you are looking at. So, this is the plot of H 1, H 2 the corresponding values which are like the weight vectors associated with basis 1 0 0, this will be with basis 0 1 0 and 0 0 1 has this as a weight vectors.

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**su(3) algebra(contd)**

- Raising and lowering operators are constructed using the remaining six generators



$$E_{\pm\alpha^{(3)}} = \frac{1}{2\sqrt{2}}(\lambda_1 \pm i\lambda_2); E_{\pm\alpha^{(1)}} = \frac{1}{2\sqrt{2}}(\lambda_4 \pm i\lambda_5)$$

$$E_{\mp\alpha^{(2)}} = \frac{1}{2\sqrt{2}}(\lambda_6 \pm i\lambda_7),$$

- $E_{-\alpha^{(3)}} | \mu_1 \rangle \propto | \mu_1 - \alpha^{(3)} \rangle = | \mu_2 \rangle$  where
- $\alpha^{(3)} = \mu_1 - \mu_2$  is a **root vector**. Similarly, you can find other root vectors

$$\alpha^{(1)} = \mu_1 - \mu_3 = (1/2, \sqrt{3}/2), \alpha^{(2)} = \mu_3 - \mu_2 = (1/2, -\sqrt{3}/2)$$

- Note that  $\alpha^{(3)} = \alpha^{(1)} + \alpha^{(2)}$  is sum of two simple roots  $\alpha^{(1)}, \alpha^{(2)}$

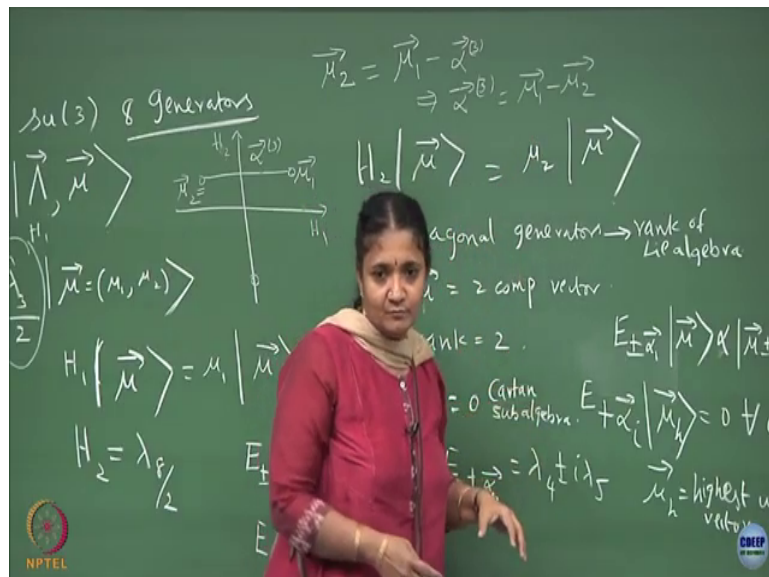
And this also I have explained here maybe the subscripts 1 2 and 3 are immaterial. So, that is just to keep track that they are 3 different two component vectors and normalizations also are some things we just uniformly followed with what I did for the su 2. So, you can see that there are 3 raising and 3 lowering operators. One more slight subtlety here, the here the lowering operator and raising operator are slightly into change here, but that does not matter.

This is the definition of the 3 operators which will rise by alpha 3, here it will decrease it by alpha 2, ok. So, because the E minus alpha 3 on mu 1 will reduce mu 1 by alpha 3, this state

you know that when  $m$  equal to half; if I do the  $j$  equal to with come back to the board. So, if you do  $J$  minus on  $m$  equal to half, you get it is proportional to  $m$  equal to minus half right. Here, the weight diagram is only the  $H_1$  plane, you start with the state which corresponds to  $m$  equal to half, the ladder operation took you to  $m$  equal to minus half.

These are the two states which defines the irrep of  $su_2$  it is a two-dimensional vector space and the two states whose weight values are  $m$  equal to half and  $m$  equal to minus half can be achieved by the lowering operator. You apply a lowering operator on this, you can get to this; if you do a raising operator on this, you can get to this.

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Same thing here, you have a  $H_1 H_2$  plane, there are three states here 1, 2 and a 3, ok. The raising and lowering operators should be such that if you take the weight vector  $\mu$  here and suppose this direction I call it as alpha vector 3, it has to give me. So, this is  $\mu_1$  let us say,

then this will be  $\mu_2$  and  $\mu_2$  has to be just like here. It has to decrease this by 1 unit. Similarly, this one has to decrease  $\mu_1$  vector by  $\alpha_3$  vector units, ok. So, you have to make sure that  $\mu_2$  vector has to be  $\mu_1$  vector minus  $\alpha_3$  vector. Is that clear?

So, that is the way the shift will happen, but it should give me the 3 weight states which I found by using the Gell-Mann matrixes because I know  $\mu_1$  and I know  $\mu_2$ , I can determine  $\alpha_3$ . These two points I know the difference between these two points  $\mu_2$  minus  $\mu_1$  or  $\mu_1$  minus  $\mu_2$  will give me  $\alpha_3$ . So,  $\alpha_3$  is  $\mu_1$  minus  $\mu_2$ , right. Whatever happened here is a one-dimensional weight diagram, this becomes a two-dimensional weight diagram with specifically I am looking at this state and I am doing a lowering operator.

When I do a lowering operator, I know that the weight vector has to shift by this, but I also know the weight vector for the state. It has to be one of the states because we have explicitly written down by the defining representation acting on  $1\ 0\ 0$ ,  $0\ 1\ 0$  and  $0\ 0\ 1$  and from there you can deduce what is  $\alpha_3$ .  $\alpha_3$  will be the difference between two weight vectors, similarly for  $\alpha_1$  and  $\alpha_2$ , ok.

So, this is what I am showing it on the screen again on the slides;  $E$  minus  $\alpha_3$  should be on  $\mu_1$  should be proportional to  $\mu_1$  minus  $\alpha_3$  and it should be equal to  $\mu_2$  were  $\alpha_3$  is  $\mu_1$  minus  $\mu_2$  and this is what we call it as a root vector. The difference between the 2 states; this is one state and the state the difference between any 2 states the weight vectors corresponding to that is what we call it as a root vector ok. So,  $\alpha_1$  so, this way you can find out  $\alpha_1$  is  $\mu_1$  minus  $\mu_3$ , what is that direction? This will correspond to  $\alpha_1$  and there is also one more direction which will correspond to  $\alpha_2$ .

In this case, you had only one direction, you had only one direction and the difference is only plus 1 plus or minus 1, ok. So, if I am doing it from here it is plus 1, here if I am doing it is minus 1. So, I should even not even put the sign, I should just say  $\alpha$  vector is plus or minus 1. Here  $\alpha_3$  will be the difference  $\mu_1$  minus  $\mu_2$  would be plus or minus

depending on the direction,  $\alpha_2$  is for this,  $\alpha_3$  is for this and they come naturally here by combining the generators, clear.

So, you will have how many root vectors? Root vectors will be the number of off-diagonal generators, right. Number of off diagonal generators is 2 here and that 2 should be Hermitian conjugate of each other; one is raising by plus one another one is reducing by plus 1, right. So, it is  $J + 1$ , is 1 root and you can have a negative of it which you can call it as a negative root  $J - 1$  is another root that will be one positive root and one negative root for  $\mathfrak{su}_2$ .

For  $\mathfrak{su}_3$  how many roots are there?

Student: 6.

6 are there, but 3 are positive and 3 are negative, ok. So, how will you find in general also? Total number of generators minus the rank of the algebra and then that should always be an even number and you have to divide it by 2 to say that those are the number of positive roots root vectors, clear.