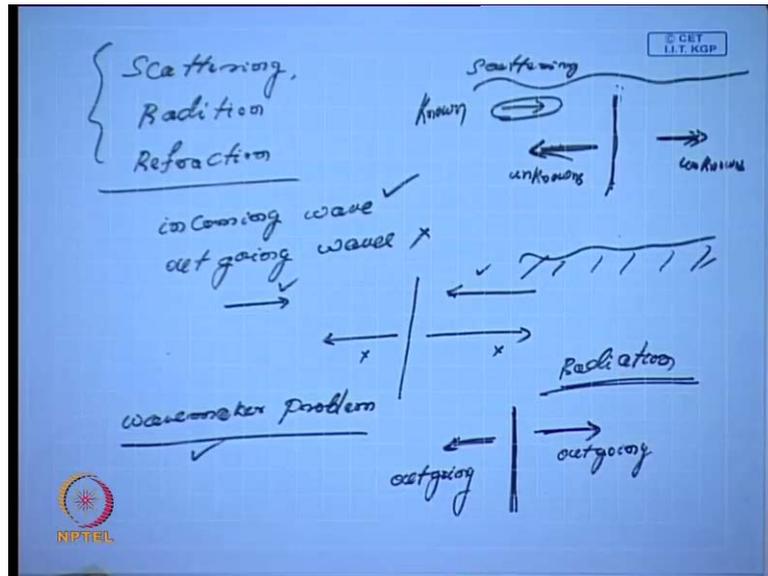


Marine Hydrodynamic
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Lecture -30
Gravity Wave Transformation and Energy Relation

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Welcome you to this lecture and marine hydrodynamics. Today will talk about wave transformation and energy relations. So, there are various transformation when it comes to waves, so here will imprecise to two things; one is the scattering and radiation, and the other one is called radiation. So, what I mean I will also if time permits only I may talk too little about to the fraction if not today, may be in the next class. So, what I mean by scattering? We what that let me talk about three waves, what I mean by in coming wave and outgoing waves. In case of a incoming waves, suppose we have a observe bar here, in case of incoming wave the wave approaches towards the observer or towards or towards the struggle bar whatever it is, it can be approaching form either direction.

On the other hand when it comes to outgoing waves, the wave that goes most away from the observe bar or the straggle, so these two are the incoming waves and these two are the outgoing waves. So, now rest on this will depend what is a scattering and then what is a radiation? In case of radiation in case of radiation what you do that this obstacle it only radiates I energy. Basically radiates the wave that means only from here initially

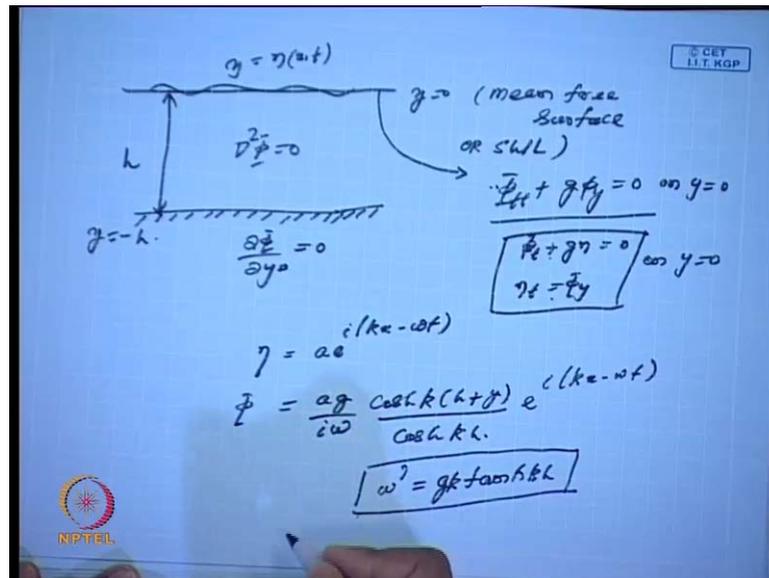
there was it is to like a car motor. What is this obstacle? It generates the wave and it oscillates or vibrates. So, when the process because of this it generates the wave and only g g radiated in both the direction, then we call this as the in call this as the outgoing there only outgoing waves exist. Here would have only outgoing waves, so that is happened in case of wave problem radiation.

So, that one process we see it here the obstacle or the structure, which are oscillates and its sometimes we call it a wave knocker problem means the wave maker problem are problem of radiation so in this what happen this generates the wave and the energy wave because of its oscillation is an is generate a wave from this wave goes away from there obstacle. So, in this case and often we call this either the radiation problem or the wave method problem like, on the other hand in case of when it comes to the scattering problem here what happened here we have a obstacle and incoming wave will be approaching to this obstacle.

So, suppose this is the bottom domain this is the free surface and because of this obstacle what will happen because there is a gap between here and here. So, it part of the energy will be reflected back and because we have gaps. So, a part of the energy will be going out in this, so there is a transformation of energy to the other side and here there is a reflection of energy. So, here there is a incoming wave in this class we know assume that the nature of the incoming wave is known and the knowing the nature of the wave obstacle we need to very quitted determine the nature of the two outgoing waves.

So, this is known nature of the incoming wave is known. And here we are interested the these are all knows only what happens here we only assume that the nature of the structure is known. So, we need to know what are the nature of the outgoing waves? So in that case we call this as a scattering.

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So, with this physical understanding of a scattering problem and (()) problem. Let us see how will look in to this problems in practice in the case of what waves, so here we all are going today the basic equations. Suppose I have the wave because of, let me think of finite water depth and this is my main T surface y is equal to 0 and this is y is equal to eta x T I. Consider a 2 dimensional problem flow domain, so my web profile is a 1 dimensional in nature y b comes function of eta eta is a function of x and T and this is my surface valuation and this is my surface valuation and this is mean free surface.

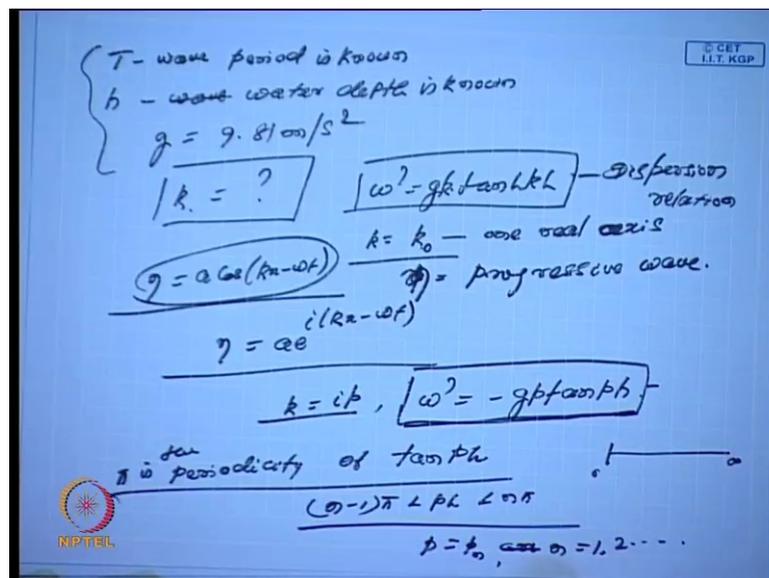
We all known the fluid domain del square pi 0 and on a main free surface that is on y is equal to 0 which I called it as the main free surface. Sometimes we call it also s W L still water label are still water label on the still water label we have pi T t plus g pi by is equal to 0 pi T t plus e pi by is equal 0. Here we have this is water depth is finite so this distance is uniform depth this I call it as s here del by del y will be 0. This is 0 and this is the line is equal to minus s this assume the depth is uniform and again this condition is satisfied this is non main free surface pi is equal to 0.

And we all know that pi T plus zeta e 0 and the surface and also eta T is pi (()). And these 2 condition are satisfied are y is equal to 0 from which we obtained this T surface boundary condition this is the clean rise dynamic condition this is the clean rise kinematic condition. And combining this we have already got the free surface boundary condition which is this and this is satisfied on this surface y is equal to 0. We have seen

if you start the wave η is equal to $A e^{i(kx - \omega t)}$, we have seen that the corresponding ω or we can see that corresponding velocity potential and will call this capital ω .

So, the corresponding velocity potential will be $\frac{g}{\omega} \cos(kx - \omega t)$ plus y by $\cos(ky - \omega t)$. So, here this ω satisfied that is $\omega^2 = gk \tanh kh$. This is the dispersion relation now. If we look at the dispersion relation what happens here this dispersion relation what about the behavior of the ω if I assume that I know the period of the wave assume we all we have offered several problems on this problem this from the dispersion relation.

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If I assume that T is the period T is known that is ω period is known. And I know what is water depth h the water depth sorry water depth is a known and I have g is equal to 9.81 meter per second square, then what will be my k a that I need to mirror? So, we have already seen that $\omega^2 = gk \tanh kh$ hyperbolic case we have worked out also several problem that this as $k \rightarrow 0$ a 1 areal root in the projective develops is so 1 areal root, now what will happen? So, that is gives us the corresponding wave number it $k \rightarrow 0$ is the 1 real root of this the this number k is equal to if I call it k is equal to If I call it k is equal to 0 .

And only would, so the corresponding π that I call per eta I call the progress able solution corresponding it is the progress able if a part, so for k is equal to k not for real so that means this for every depth for if once for a every time b the T . What are the, we have k naught and eta if so that eta is equal to a cause $k \times$ we have $k \times$ minus omega 2. We have a wave profile in that exist. Also wish have seen that the corresponding this are even if in the general 1 what I have consider just now eta is equal to a to the $(())$ $k \times$ minus omega T . So, this exist for a particular a real k exist that is k not indented.

What happen to if there any other root of this relation? This dispersion relation again if will, I will see that two analyze the roots of this dispersion relation, let me call this as the dispersion relation. If we look at this dispersions suppose, I put k is equal to $I p$ what happen or a k is equal to $I p$ than omega square will be minus $g p$ $10 p h$. Then now I have coming that from hyperbole function to tan function tens in function. And for we know that the tangent function is tan function is a function, which is a period city of tangent function periodicity of $\tan p h$. $p h \pi$ is the periodicity of $\tan p h$ see that is the ph that means you can always n minus into is less than $p h$ is less than in π .

So, each interval I always is interval of π I can get a different $p h$, so that will give me if I look in the positive axis. So, from 0 to infinity than I will have $(())$ interval of π I will have a root. So, that will give me so in this process I can easily get it has infinitely many roots, so I call this p is equal to $p n$ are n is equal to 1 2 3 of to infinity there are infinitely many roots.

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$\omega^2 = gk \tanh kh$
 $k = k_0$ (one real root)
 $k = i p_n$ (infinitely many complex imaginary roots)
 $\phi = \eta = a e^{i(kz - \omega t)}$
 $\phi_0 = \frac{ag}{i\omega} e^{i(k_0 z - \omega t)} \frac{\cosh k_0 (h+y)}{\cosh k_0 h}$ ✓
 $\phi_n = \frac{ag}{i\omega} e^{-k_n z - i\omega t} \frac{\cos k_n (h+y)}{\cos k_n h}$ ✓
 $n = 1, 2, \dots$

So, in the process what I get that means my dispersion relation ω^2 is equal to $gk \tanh kh$ in the hyperbolic case. If this person's lesson has k is equal to k not 1 real root and the k is equal to $i p_n$ infinitely many complex roots. Because it has no real part. So, in the process what will happen to the corresponding, if I look at because I am dealing with a my original equation is a Laplace equation associated with a 2 boundary condition. So, then each suppose I have seen that my π is equal to I have seen π is equal to or if I say my η is equal to a , it is per I times $a x$ minus ωT .

Then I have seen that my corresponding π in case of the real solution in k is a real that is k not than it will be $(())$ by I ω it the power I times $k x$ minus ωT and cause hyperbolic k into s plus y because hyperbolic $k h$ that is my π naught. I call this as k naught corresponding solution. And if I look at k is equal to $i p_n$ then π_n is will be at the form $a g$ by I ω T to the power when take to the x as 0 then I can have I this $i k$ n . So, that I will give minus k_n and x minus $p_n x$ minus $\pi \omega T$ and then this will give me cause p_n cause $p_n h$. So, then what will happen to this, so this are for n is equal to 1 to $(())$. So, I can also say, so this are also one of the solutions this is one solution. So, all this last few classes on waves we have emphasized only on this solution associated π , but now see if you look into this roots infinite many roots than this also is a solution. So, it has infinitely many solution is a look at the original problem than that I will just $(())$.

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$$\frac{\partial^2 \phi}{\partial x^2} + g \phi = 0 \quad \text{on } y=0$$

$$\frac{\partial \phi}{\partial y} = 0 \quad \text{on } y=h$$

$$\phi = \frac{ag}{i\omega} \left\{ a e^{i(k_0 x - \omega t)} + b e^{-i(k_0 x - \omega t)} \right\} + \frac{ag e^{-i\omega t}}{i\omega} \sum_{n=1}^{\infty} \left\{ a_n e^{-k_n x} + b_n e^{k_n x} \right\}$$

propagating modes
 evanescent modes
 local solution
 local effect
 progressive wave soln.

$a > 0, b_n = 0$ $a < 0, a_n = 0$

If I look at the original problem that means $\frac{\partial^2 \phi}{\partial x^2} + g \phi = 0$ on $y=0$ and $\frac{\partial \phi}{\partial y} = 0$ on $y=h$. My general solution will be of this form $\phi = \frac{ag}{i\omega} \left\{ a e^{i(k_0 x - \omega t)} + b e^{-i(k_0 x - \omega t)} \right\} + \frac{ag e^{-i\omega t}}{i\omega} \sum_{n=1}^{\infty} \left\{ a_n e^{-k_n x} + b_n e^{k_n x} \right\}$. I cannot take minus $i\omega t$ plus $a g$ by $i\omega$ I can call it a_n , n is equal to 1 to infinity and into the minus $k_n x$ plus $b_n e^{k_n x}$ and this is a_n and b_n and look us a_n plus b_n and into into the for minus $i\omega t$.

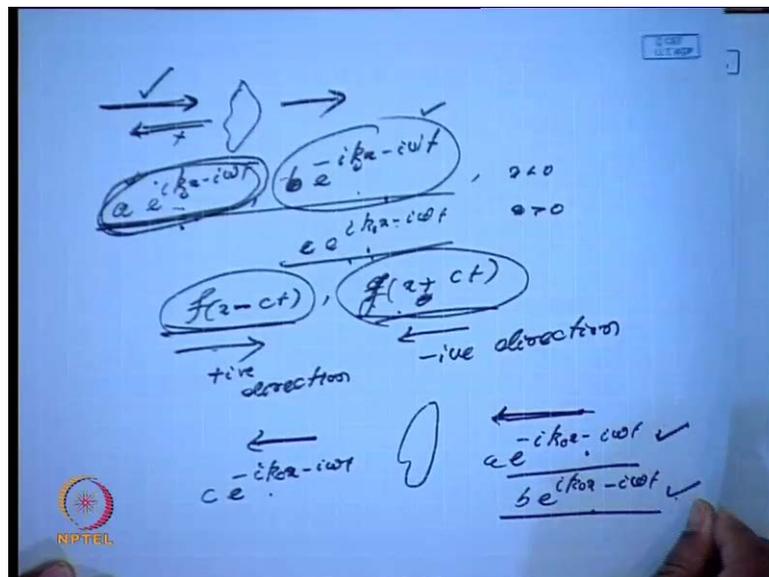
Now, this is the general for now. Question comes, what happen because I was dealing with $i\omega t$ this term as come. And again I was here I was taken minus $p_n y$ again plus $p_n x$, so this I will let us explain this terms what happen actually this part contribute to the profile the take a more solution. This parts I have the local effect, which will call the evanescent modes even says centre modes and this is the profile modes, but the progressive wave solution or call the progressive wave solution.

On the other at this 7 7 modes often call as local solution or local abate, now when we say this is the other local effect what happen? Suppose, I am dealing with that my only x I am dealing with that my only x I am dealing with the x the abdomen of fluid is only positive x axis if I am dealing with x is 0. Then what will happen to this, because if x is 0 then v to the or $p_n x$ p_n are positive p_n , suppose it is $p_n x$ will tan to infinity. So, that

means I am looking for a solution, which will not behave on boundary solution of looking only for bounded solution.

So, in that case my b_n will be 0, but on the other hand if I said my I am looking at if I say I am looking at solution for x less than 0 than automatically x is less than 0. So, this part will tan to 0 where has x is less than 0 means there should be positive, so n should be unbounded x less than 0, so in that case if x is less than 0 I can call it is a n is 0. So, that will give me another set so that is a one thing than coming back to this two terms there are physically I have just be find that for a wave this are the two refer to the progression wave solution and as I have told that in case of wave.

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We have it may happen that wave by wave, which is a incoming wave and because we may have a structure or have some obstruction and a part of the wave exergual reflected. So, here we have a wave the nature of the incoming wave will be $i k x$ minus $I \omega T$ and nature of the outgoing waves some a nature of the outgoing will be some b into the or minus $I k$ call it k not x minus $I \omega T$.

So, if we have problem of a scattering problem, so we may have incident wave and also reflected wave if we are dealing with a wave. So, in that case this will contribute to the incoming wave this contributes to the outgoing wave as well as the x less than 0 is concern with the $f x$ greater than 0 than on this side the wave will be outgoing, but it is always moving in this direction.

So, in this case, so for so there are the two waves this will correspond to the incoming wave and this $(())$ will correspond to the outgoing waves for x less than 0. On the other hand if we are dealing with a scattering problem on the other hand for x greater than 0, I may have a constant c . And then my wave will be it is where $i k x - \omega t$ and that will be the type which is the outgoing wave, here always look at it is of the form $f(x - ct)$ and this is of the form $f(x + ct)$. So, this wave which is always moves in the positive direction and this wave which moves in the it gives basically this sine $(())$ minus gives the direction in which direction the wave is propagating.

So, whether is a incoming wave are outgoing wave, but this plus minus sine gives us the direction, which the wave is propagating. So, that is why here we have this one is a. Although, it is a minus sign, but still it is a incoming wave where has this is the outgoing wave both are minus sign. On the other hand here we have plus for sine, this is a minus sign similar to this what this is a outgoing wave. So, this basically determine this sine sine of this particularly $x - ct$ or $x + ct$ type. So, that determine whether the wave is progressing in the positive direction and other it is propagating always in the negative direction.

So, there are two things here whether the wave is incoming are outgoing; one as the other thing is whether the wave is propagating in the positive are negative direction that is a another things. So, when you when one deals with a physical problem one has to take this 2 aspects in the aconite which $(())$. Suppose I say that I have a wave which is approaching from this side than what should be my progressive, so I will be have a to the minus $i k x - \omega t$. If wave is if I have obstracter like this and my wave is a approaching from the right side. So, this will be my incoming wave and on the other hand when wave was approaching from this side this was becoming the incoming wave.

And here the outgoing wave will be of this type some a here some be it there I cannot accept minus $i k x - \omega t$. This will be my if the structure is there, this is the incoming wave, this is the outgoing wave, this is that is the, this is the incident wave, this is the reflecting, whereas this wave, the wave will be propagating this wave. Because approaching from this side and when it approaches from this side then this will be some of the it is nature will be a to the c into the minus $i k x - \omega t$. So, this will be the nature of the because this is this waves this two waves they propagating in the

negativity direction of x axis. And where has this wave this propagating the in the positive direction of the x axis, so anything this is propagating in the positive direction will have a minus sign of the nature will be like this anything is propagate in the negative direction the nature will be like this.

So, because of that here it is these wave is a wave propagating towards the positive direction, this is a wave which is propagating towards negative direction this is again towards positive direction this is again towards negative direction this is towards positive direction this is towards negative direction. So, respective of the fact whether it is a incoming wave for a outgoing wave the direction of propagation determine the positive or negative sign in the wave pattern. So, this two things are very clear for us now, this understanding will a look at the full solution and determine what kind of wave it will what kind of solution it will be.

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The image shows handwritten notes on a blue background. At the top, there is a mathematical expression for the potential ϕ as a function of x and t :

$$\phi = \frac{ag}{2\omega} \left\{ \sum_{n=1}^{\infty} \left[a_n e^{i(k_n x - \omega t)} + b_n e^{-i(k_n x + \omega t)} \right] \right\} \frac{\cos(k_n y)}{\cos(k_n h)}$$

Below this, there is another expression for the potential:

$$+ \frac{ag}{i\omega} \sum_{n=1}^{\infty} \left\{ a_n e^{-k_n x - i\omega t} + b_n e^{+k_n x - i\omega t} \right\}$$

Underneath, the notes are divided into two cases:

Case 1: wave maker problem (radiation)

A diagram shows a vertical line at $x=0$ with a double-headed arrow. To the left ($x < 0$), it is noted that $a_0 = 0$ and $a_n = 0$. To the right ($x > 0$), it is noted that $b_0 = 0$ and $b_n = 0$.

Case 2: scattering problem

A diagram shows a vertical line at $x=0$ with a double-headed arrow. To the left ($x < 0$), it is noted that $a_0 = 0$ and $a_n = 0$. To the right ($x > 0$), it is noted that $b_0 = 0$ and $b_n = 0$. The region to the right is labeled "incoming wave" and "reflected wave".

Now, if I go back again I will just try it again that the general form of the velocity potential that means ϕ is I have seen that this is a $\frac{g}{\omega} \sum_{n=1}^{\infty} [a_n e^{i(k_n x - \omega t)} + b_n e^{-i(k_n x + \omega t)}] \frac{\cos(k_n y)}{\cos(k_n h)}$. So, I call this as a_0 and b_0 minus i cannot x minus i omega T , so if I take minus i comma and I will have this. So, this is I call it omega T this into, so this is the general form, so if I am looking at a case one if I look at a wave maker problem basically I look at a radiation problem (()).

So, if it is a problem of radiation than what will happen I will have only outgoing waves, so in that case if I say the wave is radiating energy will both the direction. So, for x greater than 0 my b_1 it is a positivity direction. So, b_1 will be a 0 will be 0 on the positive side because this is going sorry this is going in the this way, so this term will be so b_0 will 0.

And here this side will be b_n will be 0 that is if this is line is x is 0 and I have a wave maker problem. So, the energy is only radiated, so your b_0 will be 0 because wave is propagating only in this deduction. So, this term will contribute because of this on this term will not contribute, whereas because wave is one of the wave will exist (()). On the other hand here b_n will be, because this term will be on the right side this term will not contribute. So, b_n will be 0 on the other hand on this side for x this is for x greater than 0, but what will happen the if the wave is a generated in both the direction. In this case what will happen? This here than my a knot will be because wave is propagating the negative direction. So, a knot will be 0 and where has on from this seven side modes I will have b and a 0.

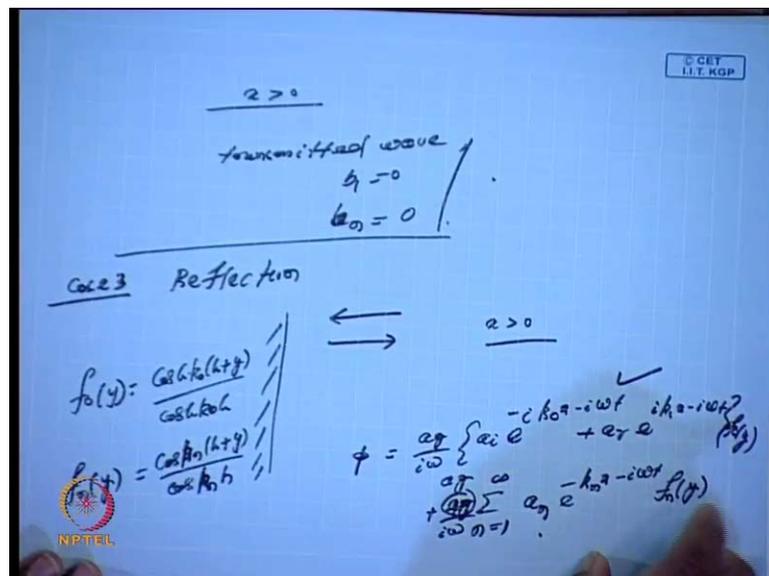
So, for wave maker problem from the general solution, this will be this is only will be and this is for x less than 0. On the other hand if I look into scattering problem case two scattering problem, in case of a scattering problem what will happen? I will have a structure, I will have one incoming wave, I will have one outgoing wave. Here I will have out going wave, so this wave is the incoming wave this is the reflected wave. This is the transmitter wave incoming wave, this is the reflected wave and this is the transmitter wave.

Here I before preceding for the right again point out high is the benefit potential, but it is the velocity potential, but finally what we are interested we are interested in knowing η x t as I act whole earlier that many situations because whole problem has been recast wave problem in π interims of the velocity potential. So, we are solving lab plus equation subject to the free surface boundary condition and the bottom boundary condition everything in terms of π . So, once we obtained π we can get η because η represents gives the surface profile surface valuation or the part on the waves π dragnets by T the coefficients associated with the basically the coefficient a_0 p_0 a_n and b_n .

All these terms they give us and understanding about what exactly happens to unplaced waves, so in many situations because all these parameters k_0 will you know because once we solve it omega squared that is dispersion will know k_0 . But what will happen in a wave propagation problem, we need to know these constants a_0, b_0, n, v, n . If I look at a radiation problem on the right side maybe we have b_0 . Because of the nature of the physical problem b_0, v, n, a_0 and here is a_0, n_0 , but we need we need other conditions to finally find what are the rest of the unknowns? Similarly, what happens in case of a scattering problem we have incoming wave we have reflected today we have a transmitted wave so if I look here we look at this problem in case of scattering problem on the left side for $x < 0$ I will have both the terms will be so a_0 will be b_0 will be.

So, they will be, so basically a_0 will be known b_0 unknown, whereas because this I call the incoming wave I assume the nature of the incoming wave is known. And b_0 is the nature of the outgoing wave I am teaching today associated with the outgoing wave that is the reflect that is unknown. And again on the left side x is less than 0 and out of this two because it is a less than 0 this term, so my answer should be 0 all n are 0 on the other hand if I look at the in the scattering problem.

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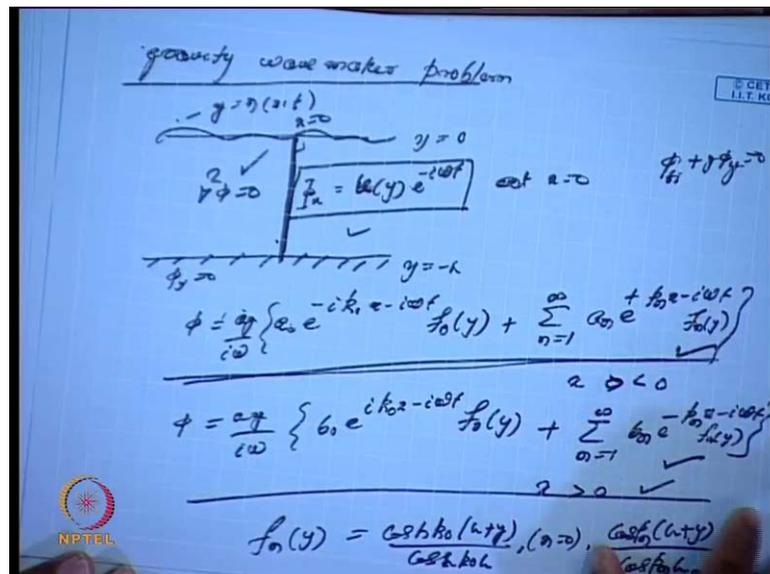
If you look at the right side for x greater than 0 in the case of the scattering problem then I will have because I have only one outgoing wave that is the transmitted wave because wave is propagating from the positive side. And then in that case my b_1 will be 0 the

nature of b_1 will be 0. On the other hand in case of divergent modes I will have x less greater than 0, so I have a_n is b_n is will be 0.

So, this is what will happen and of course in the left side constants associated constant in not be the same as the constant on the right side. They will be different than in be same I will come to a exact problem next. Now, I will go to the hard case I will come to then I will come, then I will come to 1 by 1 progress incase three suppose i case of reflection I suppose I have a (()) I have one incoming wave is coming than that will be outgoing waves.

So, in this case of, suppose I am taking that this is I am in the profile is a x is a greater than 0. So, what will be (()) now ϕ will be a g by I omega into a i , I call this as a i a to the Pera $i k x$ minus $i k x$ minus I omega T plus a $r i k$ not x minus I omega T . Then we have plus sigma 1 to infinity. I call it a n and this a n a to the minus $k_n x$ minus I omega T there is a π_0 I call it f_0 A then I call it f_1 f_0 y and here I call it f_n y an what is my f_n f_0 by is plus separable k_0 s plus y by plus high (()) k_0 is f_n y plus k_n is y by cause I use the notice on $p p n$ h. This is my general form for the replacing because I have one incident way this is the reflect today and this are the (()) modes this is a g by omega.

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This is a g now if this are the patrons so suppose so suppose I say now I will come back to a another case of a problem of a wave maker. Basically what I say laities a make problem in case of gravity a maker problem like we have seen in a wave tank. And tank

which generates wave in any hydrodynamics laboratory. One can find your wave tank particularly which generates there $j(())$ to at one end of the tank and that we apply escalate and it generates wave. So, that basically almost exist worldwide any hydrodynamics laboratory that for module testing of various physical models a wave tank available.

So, at one end of the wave tank they will be wave maker and which will be generate the waves. So suppose I say let me consider this problem like this is my free surface h by g is equal $\eta \times T$ and is this is my y is minus h y is equal 0. Suppose, there is a y maker here of this position is x is equal to 0 and this wave I have the wave maker exist here and it oscillates it makes a on the batik it makes oscillation πx is $u y$ what is a minus π ωt .

And I mean the frequency of oscillation is the same or the frequency of the waves this is geologier. Suppose, the wave maker oscillates the speed of $p i x$ is the valuental velocity in the horizontal direction on the horizontal direction it makes on oscillation of this nature and when I say πx is of this is basically on and x is equal to 0. So, at the neon position about x is equal to 0 this the $(())$ small oscillation small lumbtage oscillation with frequencies same as the wave frequency then what will happen to the wave? Then I have in the through hole fluid is a region, I have $\text{del square } \pi$ is 0 when this add or on this add if I say this is, so then what will happen for as usual my π will be I am this ferm because I am dealing with only a wave maker problem this wave maker is oscillating.

So, my π will be a 0 in to there minus nature of the I cannot x minus I ωT , then $f 0$ by plus σn is equal to 1 to infinity I will say $z g$ by I ω this 1. This is a quantum and will be throughout $f 0 y$ plus σn is equal to 1 to infinity and this will be an in to a minus $p n x$ when the $\psi \omega T$ this is minus of the should be plus this is for $x x$ less than 0 and if I look at what happen on the right side if the wave.

The fluid is on this side this will be the form of the π and if the fluid occupies on the other side than π will be of this ferm $L g$ by I ω this will be some b not u to the Pera I cannot x and just I $\omega T f$ naught y to the σ some $b n$ plus is for x . So, this will be the nature of π for this is for x less than 0 and this will be the π by x greater than 0 now the main problem remains and this π will satisfy the free surface condition that is

the free surface condition not y is equal to 0 $\pi T T$ plus $g g \pi$ by 0 all this π are on this π both of them satisfied the lab plus equation in the fluid domain.

And also they both satisfy π by is 0 on y is equal to minus h here my $f 0$ on the fence y here. There is a $f n y f n y$ and to I have already told my $f n y$ is caused hyperbolic $k 0 s$ plus $y y$ by cause hyperbolic $k 0 h$ that is in a 0. And this is because $p n s$ plus by y because $p n h$ for all in greater than one two positive. So, this is main general, now the question comes I have one more boundary condition that is πx is $e y a$ to the minus $i \omega T$, if I substituted πx here. Suppose, I look at the positive side, let me take anyone side positive side, so what will happen to πx .

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$$p/a = 0 = \left\{ i k_0 b_0 f_0(y) + \sum_{n=1}^{\infty} k_n b_n f_n(y) \right\} e^{-i \omega t} = u(y) e^{-i \omega t}$$

$$\frac{e^{i(kx - \omega t)}}{\sin(kx - \omega t)}$$

$$i k_0 b_0 f_0(y) - \sum_{n=1}^{\infty} k_n b_n f_n(y) = u(y)$$

$$\frac{u(y) - i \sum_{n=1}^{\infty} k_n b_n f_n(y)}{b_n, n=0, 1, 2, \dots}$$

So πx at x is equal to 0 gives me I cannot will looking at the positive direction I cannot that is b naught into f not π plus because x is equal to 0 I am calculating plus sigma n is equal to 1 to infinity minus it will be minus $k n$. Then $b n f n y$ and this is nothing but $\pi x x$ is equal to 0 into into the minus $I \omega T$ that is equal to $o u y$ it to the minus $I \omega T$.

I deal with a comp less exponentials, but the in a reality the solution will be either the real are imagine a part of the comp les potential of this comp less exponential function will give me the real wave part as we say that e to the Pera $I a x$ minus ωT . The real are imagine the wave parts of the real part is sine $k x$ cause $k x$ minus ωT or sine $k x$ minus ωT for easy of mathematical simplicity. You always look in to the

complicity comp less form of the wave and in a reality when will go for real live programs will always when the I i conscious that real part of this.

So, if this will be this than what will happen my b y the if I simplify worthier I get I cannot b not f naught y minus sigma n is equal to y to infinity k n b n f n y this is going to e y. Now, I only have one question, what what happen here? My u y i know because I know the way the way (()) oscillating I assume u y is non, I looking at wave radiation problems. I know that what is the how the wave maker the structure is oscillating attach x is equal to 0, so if u is known then what will happen I have, so many onwards b not v n so I have b n n is equal to 0 1 2 this are the unknowns I need to obtain them so to obtain this unknowns.

So, what I will do I only look at the because this is kind of whether you called it a in function x function. if you look at mathematically are by it is a kind of basic simple method of separational variable and this leads to and the a standard (()) problem. So, f n y are the y again functions if I call k 0 and p n are the liagan values, then I can say that corresponding a fence will give me the liagan functions, and since it is a term level type of wonderfully problem in f n.

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Handwritten mathematical derivation on a blue background:

$$f_n'(y) + k_n^2 f_n(y) = 0$$

$$f_n''(y) - k_n^2 f_n(y) = 0$$

$$\omega^2 = g k_0 \tanh k_0 h$$

$$= -g k_n \tanh k_n h$$

$$\int_{-h}^0 \frac{f_n'(y)}{f_n(y)} dy = \begin{cases} 0 & \text{for } m \neq n \\ \frac{\sinh 2k_n h + 2k_n h}{4k_n} & \text{for } m = n = 0 \\ \frac{\sin 2k_n h + 2k_n h}{4k_n} & \text{for } m = 1, 2, \dots \end{cases}$$

$$f_n(y) = \frac{f_n(y)}{\cosh k_n h}$$

$$f_n(y) = \frac{f_n(y)}{\cosh k_n h} \quad f_n(y) = \begin{cases} \cosh k_n (a+y) & \text{for } m = 0 \\ \cos k_n (h+y) & \text{for } m = 1, 2, \dots \end{cases}$$

Because f n y will satisfy the differential equation fn double dash y minus plus p n square by sorry p n square f n y e 0 and again fn double dash y minus p 0 square f n y 0 and both satisfy the common thing of that p n satisfy. Omega square is a g and this is k naught g k

not tan hyperbely case and again it satisfy minus $g p n h$, so this is the the nature of $f n$ and that satisfy. This equation to this are all standard $(())$ program. It has to n point to point to find because it satisfy the free surface condition and also the bottom condition. We have simplified just under ordinary differential equation comes from and basically has comes to on lab plus equation the y component of the lab plus equation.

So, the because of this there also what is arithmetic criteria, so they are suppose two satisfy $(())$ multi criteria and in this case we can see that this $f n f m$ minus $s 2 0 d y$ the brute is that $f n y f m y d y$ this is 0 for z means not equal to n . And further this is equal to sine hyperbolic 2 case plus cage by 4 k . This is for m is equal to m is equal to 0 and this is sine in 2 $k h$ plus 2 $k h$ by 4 $k p p n h p n h$ by 4 $p n$ this is a m is equal to this is for 1 2 s . And here my $f m$ because hyperbolic $k 0$ is plus by this is cause hyperbole cause $p n$ into s plus $y n$ is equal to 0 1 to $(())$. So, here I am not taking the division by cause hyperbolic k naught, why because hyperbolic k not $h r$ because that is a only only a term which should be multiplied.

Here if I divided by cause of $k n h$ than another term will be here for separable square $k n h$ and here cause separable cause square $k n h$ only π I are here this is the general thing if I take $f n y$ is this then $f m f n$ will be this. And this is what a rather call it $f n \text{ bar } f n \text{ bar}$ and call it $f n \text{ bar } (()) f m y$ will be $f n \text{ bar } \pi$ by $f 0$ by y by cause hyperbolic $k 0 h$. My $f n y$ will be my 0 $f n \text{ bar}$ by $(())$ cause $k n h$. So, this will be my $f 0 f n$. So, since once this are authorial automatically $f 0 f n$, so that gives us the beauty of this that and this will be utilize.

(Refer Slide Time: 48:26)

$$f(y) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi y}{L}\right) = f(y)$$

$$b_0 = \frac{1}{L} \int_{-L}^L f(x) dx$$

$$b_n = \frac{2}{L} \int_{-L}^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

$$b_0 = \frac{2}{L} \frac{[x \cos(2kx) + 2kx^2 \sin(2kx)]}{4k}$$

$$b_n = \frac{2}{L} \frac{[x \cos(2kx) + 2kx^2 \sin(2kx)]}{4k}$$

Because my original equation, if I always b_0 of y minus sigma n is equal to 1 to infinity k n p n b n f n y is equal to my u by and if I have plate the $(())$ then I will get my b_0 as minus s 2 0 u T \cos hyperbolic k 0 s plus by s plus T d t . And b_0 will be I b_0 will be divided by divided by commas i k 0 and to I call it this c 0 0 and again b_n will be minus minus s 2 0 u T cause p n into s plus T d t divided by k n c n n a d . This c n n whether the c 0 is a 0 is I call it is sine hyperbolic to k 0 H sine hyperbolic 2 k 0 h plus 2 k 0 h .

So, an I prove will to k 0 h by plus to k 0 h by 4 k 0 and c n n and only one term. If I have here, if I will have \cos hyperbolic term a division term here, if are here if i by division by \cos p n h \tan automatically will have a square term will be \cos hyperbolic square p n h will come. Otherwise, it is remains the same this the general for, so this is my b_0 this is my b_n this is my c_0 on this. So, I know all the unknown constant b_n and what was my original equation.

(Refer Slide Time: 51:23)

Handwritten notes on a blue background showing the derivation of Havelock's expansion formula. The equation is:

$$\phi(x, y, t) = \frac{a_0 f_0(y)}{2} e^{+ik_0 x - i\omega t} + \sum_{n=1}^{\infty} \frac{b_n f_n(y)}{2} e^{-knx - i\omega t}$$

Below the equation, it is noted that b_n are known in terms of $u(y)$. The text below the equation reads: "Havelock's expansion formula" and "Gravity wave maker problem Havelock (1929) Phil. Mag. ✓". There are logos for NPTEL and CET IIT KGP in the corners.

My original expression for ϕ is $\phi(x, y, t) = \frac{a_0 f_0(y)}{2} e^{+ik_0 x - i\omega t} + \sum_{n=1}^{\infty} \frac{b_n f_n(y)}{2} e^{-knx - i\omega t}$. Here, a_0 is the amplitude of the wave maker, $f_0(y)$ is the free surface elevation, k_0 is the wavenumber, ω is the angular frequency, x is the horizontal coordinate, y is the vertical coordinate, and t is time. The term $b_n f_n(y)$ represents the higher-order modes. The b_n are known in terms of $u(y)$, the velocity potential. This is Havelock's expansion formula for the gravity wave maker problem, derived by Havelock in 1929 in the Philosophical Magazine.

We make a problem, it is called the gravity wave maker problem. In fact, this problem was first developed by Havelock in 1929. And it was probably (()) and this problem has a wide application today and they are in last line terms last more than T s. This problem has to simplify large varieties of problems in gravity wave associated today. The gravity of OA structure (()) problems. And almost this concept, which is used in the development of the wave maker classical wave maker which is used worldwide in a tank in hydrodynamics.

About with this today I will stop and tomorrow are in the next class rather will talk about how the energetic lesson is coming into picture and how the scattering problems can be handled in simple cases over complex situation will be handled rest 2 case, but accolades

the general philosophy with the radiation and scattering problem will give discussing between.

Thank you today.