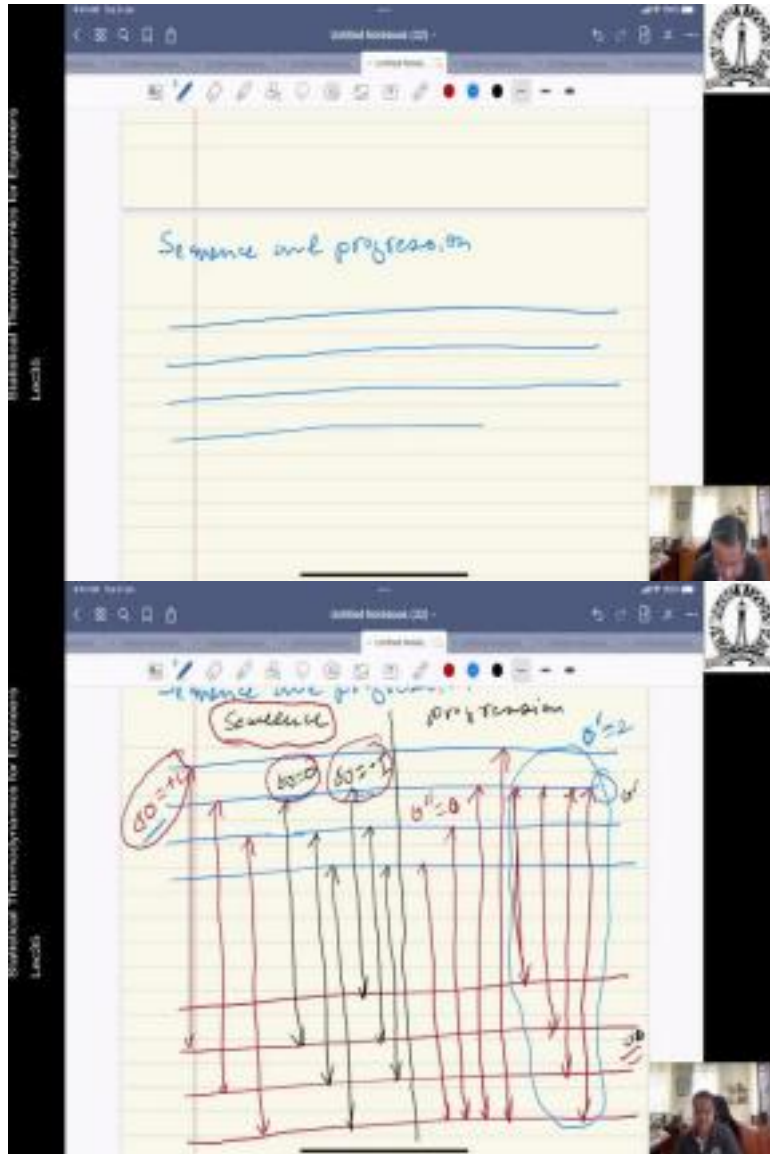


**Statistical Thermodynamics for Engineers**  
**Professor Saptarshi Basu**  
**Indian Institute of Science, Bengaluru**  
**Lecture 46**  
**Ro-vibronic spectroscopy**

(Refer Slide Time: 00:19)



Welcome to lecture number 35 of the statistical thermodynamics course. So, let us draw now, if you remember that we brought in the concept of sequence and progression. So, sequence was when the electronic spectrum containing the vibrational bands with a constant  $\Delta v$  is called a sequence whereas, with spectrum which contains regular vibrational bands with a constant  $v$  prime or  $v$  double prime is called a progression.

So, let us draw that to how this can be best represented this progression and the sequence. We already talked about sequence and progression. So, let us look at it in a little bit more details.

So, let us look at some transitions which would be relevant to. So let us draw this, is the sample example though this one. So, let us put a marker in between so, this part we are going to sequence, this is for the progression alright.

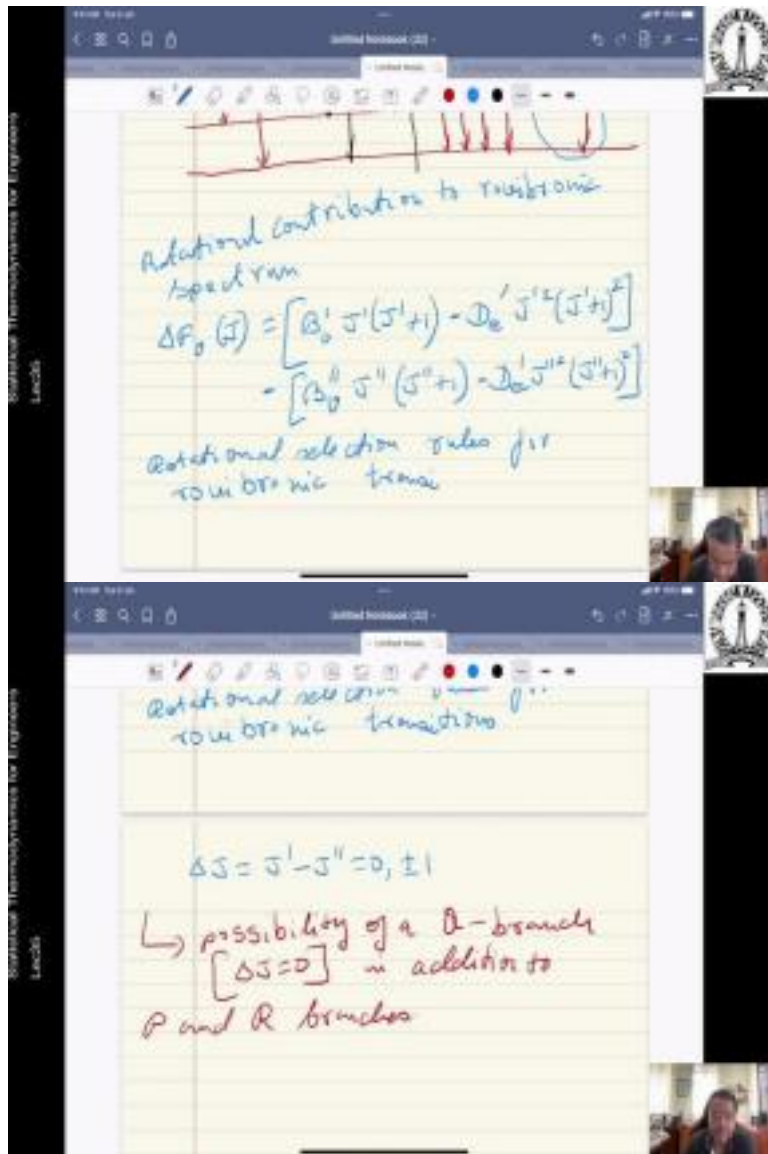
So, and this is of course  $\nu$  prime  $\nu$  double prime. So, let us look at some of the transitions now let us mark it as red (1:58). So, this is like  $\Delta \nu$  is equal to this one plus one always it is going from a second to a third vibration that is the kind of transitions that we are looking at. Now, if we so this is one family that we can look into another family now where (2:52) let us look at the thin from 2 to 2, this is of  $\Delta \nu$  is equal to 0 into  $\Delta \nu$  type of transitions this is still  $\Delta \nu$  equal to minus 1.

So, this is plus 1, this is 0, and this is minus 1. So, this is sequence so, you see that this is a constant if you look at it, it is a constant  $\Delta \nu$  so that is why it is called sequence, it is a constant  $\Delta \nu$  that you see. We will look at the transitions from here to here this is plus and or from here to there is plus and a plus and then this is the same way so you can see there is a plus one vibrational level this is 0 it is migrating from a  $\nu$  1 prime to  $\nu$  1 double prime which is equal to 0.

And so this is the way of the sequence part actually works. And this side we call it the progression. So the progression will be like it is easy to draw this is for a constant  $\nu$  prime or a  $\nu$  double prime as you can see here, where  $\nu$  double prime is equal to 0, this is the  $\nu$  double prime part. Similarly, you can have so you are not bothered about the  $\Delta \nu$  percent, you are only concerned about where the transition is taking place.

So, this part is for  $\nu$  prime is equal to 2 that is where the transitions have started from. So, this is a progression, because it has a constant  $\nu$  prime or  $\nu$  double prime. Here of course, it is a for sequence because it maintains a constant  $\Delta \nu$ . So, these are the two types of electronic spectrum which containing the vibrational bands with a constant  $\Delta \nu$  or with a constant  $\Delta \nu$  prime or  $\nu$  double prime. So, this is what is represented in this particular diagram, which we said that we are going to cover in the next in the previous class.

(Refer Slide Time: 06:25)



Now, let us rewrite the rotational contribution to the ro-vibronic spectrum. Where this  $\Delta F_0(J) = B_0' J'(J'+1) - D_0' J'^2(J'+1)^2 - [B_0'' J''(J''+1) - D_0'' J''^2(J''+1)^2]$ . So, again the single and the double as we said refer to the upper and the lower electronic states.

So, in comparison to the rovibrational case the rotational selection rule or ro-vibronic transitions this turns out to be  $\Delta J = J' - J'' = 0, \pm 1$ . So, there is a possibility now possibility of a Q branch which corresponds to  $\Delta J = 0$ , in addition to P and R branches.

(Refer Slide Time: 08:55)



transition of 0. And lastly it is forbidden, when it will be forbidden? Forbidden if this is equal to 0 for both participating electronic states.

So, some of the features therefore of the ro-vibronic transitions are, features of ro-vibronic transitions, number one strong lines correspond to the main line transitions and weak spectral lines., these corresponds to satellite transitions. So, you can see it so, these are the features of the rovibrational transition.

Now, if you study the situation over here this particular equation this represents a rotational structure within the vibronic bands which is already we have we know that if you remember what we call the band origin, so it is a rotational structure within a vibronic band, so it is a it is a rotational structure within the vibronic band.

(Refer Slide Time: 11:44)

Top Screenshot:

no. speed run

$$\Delta E_{0,0}(J) = \left[ B_0' J'(J'+1) - D_0' J'^2(J'+1)^2 \right] - \left[ B_0'' J''(J''+1) - D_0'' J''^2(J''+1)^2 \right]$$

Rotational selection rules for ro-vibronic transitions

$$\Delta J = J' - J'' = 0, \pm 1$$

↳ possibility of a Q-branch ( $\Delta J = 0$ ) in addition to

Bottom Screenshot:

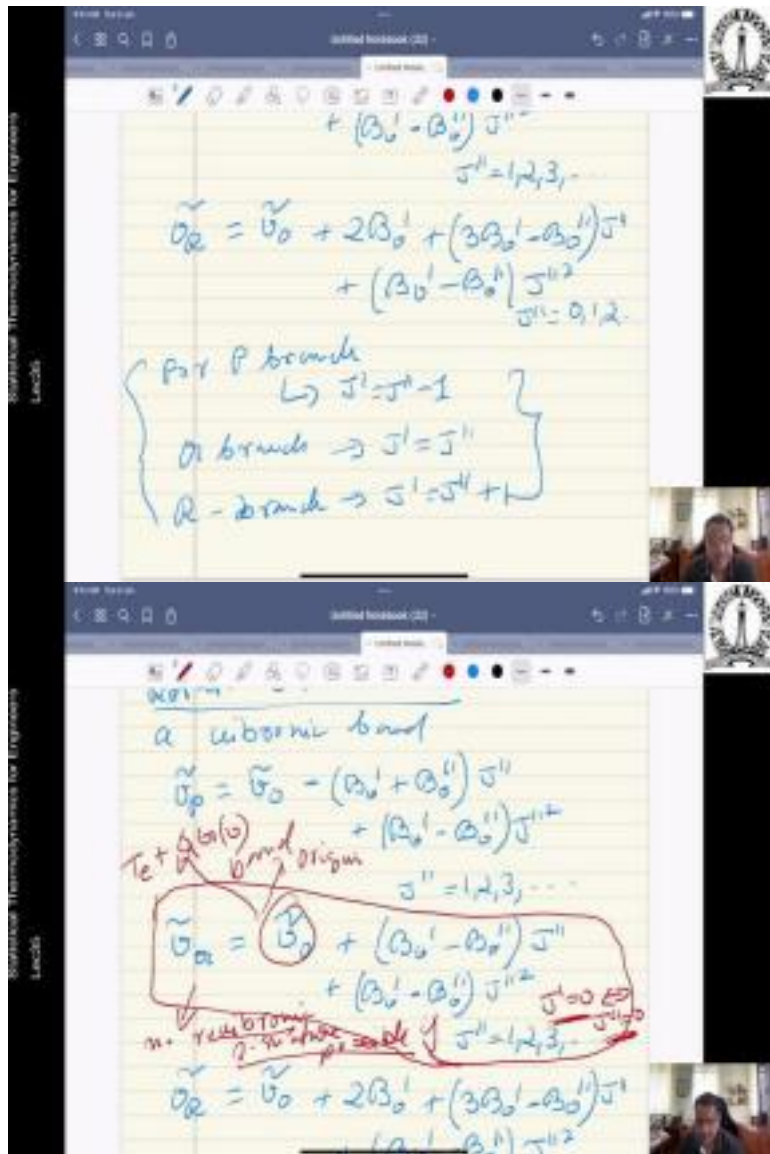
We assume  $D_0 = 0$  (negligible rotational distortion)

Rotational structure within a vibronic band

$$\tilde{G}_p = \tilde{G}_0 - (B_0' + B_0'') J'' + (B_0' - B_0'') J''^2$$

$J'' = 1, 2, 3, \dots$

$$\tilde{G}_R = \tilde{G}_0 + (B_0' - B_0'') J'$$



Now the rotational structure now, if we apply this equation and port it there we will just see what happens, let us use a new page and we assume it  $D_e$  is equal to 0 that is negligible rotational distortion rotational distortion we can now represent the three branches the rotational structure. So, the rotational structure within vibronic band is given by this equation, so this is the rotational structure.

So, once again we can find out that what is rotational structure is going to be. So, let us write these are all covered previously we just consult back the notes what we did in the last class plus square where  $J''$  is equal to 1, 2, 3, dot dot. Similar, the Q branch and this has given as  $B v' + B v'' J' J'' + J''^2$ ,  $J''$  equal to 1, 2, 3, dot dot.

Then off course we have R which is  $(J'' - 1) J''$  this  $2 B v'$  this were all not very difficult to understand plus prime double prime  $J''^2$ ,  $J''$  is equal to 0, 1. So,

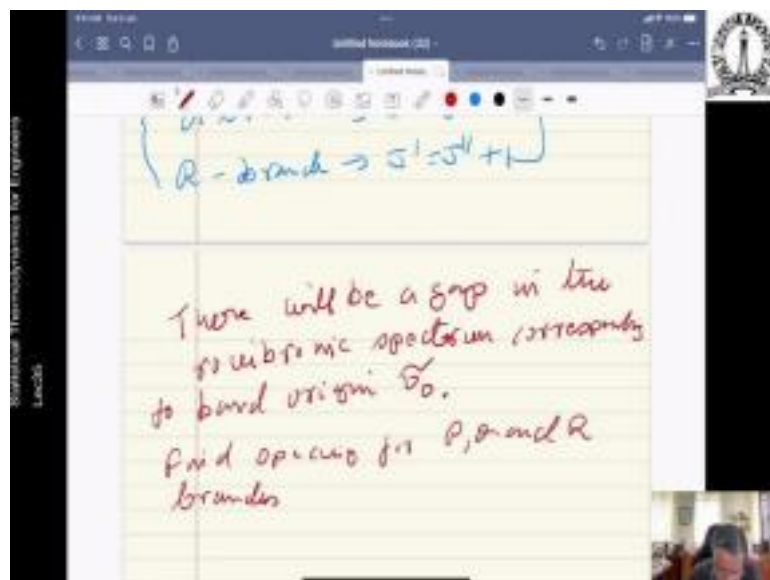
for the P is now of the P branch for P branch,  $J'$  is equal to  $J'' - 1$ . Q branch,  $J'' = J'$ , and then R branch  $J'$  is equal to  $J'' + 1$ . These are the three branches and we observe from this particular equation if you look at this particular equation, what do you see in this equation?

If you see this equation there is no ro-vibronic signature is possible for  $J'' = 0$  to  $J' = 0$ . So, here this equation suggests no ro-vibronic signature is possible because these will cancel if you look at this, this will cancel no ro-vibronic signature is possible if  $J'$  is equal to  $J''$  (16:29) to frame this exactly this.

This is actually when  $J' = 0$  to  $J'' = 0$ . So, there is no ro-vibronic signature possible if you substitute this this will just cancel out and you will be left with only the  $\nu_0$ , where  $\nu_0$  is basically the band origin,  $\nu_0$  was the band origin this was called the band origin and this was written as  $T_e + \Delta G$ .

So, no ro-vibronic signature is possible if  $J'$  is equal to 0 and  $J''$  is equal to 0. For others we have transition possible. So, in other words there will be a gap now that will appear in the ro-vibronic spectrum corresponding to the band origin at (17:37) the analogy with a similar gap with the infrared part of the spectrum.

(Refer Slide Time: 17:47)



Handwritten equations on a digital notepad:

$$\Delta \tilde{\nu}_P = (B_0' + B_0'') + \frac{(B_0'' - B_0')}{(2J'' + 1)}$$

$$\Delta \tilde{\nu}_Q = 2(J'' + 1)(B_0' - B_0'')$$

$$\Delta \tilde{\nu}_R = (3B_0' - B_0'') - \frac{(B_0'' - B_0')}{(2J'' + 1)}$$

$$\Delta \tilde{\nu}_P = (B_0' + B_0'') + \frac{(B_0'' - B_0')}{(2J'' + 1)}$$

So, there will be a gap so, this suggests there will be a gap in the ro-vibronic spectrum corresponding to band origin which is at  $\nu_0$ , this is similar to the infrared spectrum as well.

Now, if we utilize all these equations now we find the spacing for the P, Q and R branches. So the spacing would be equal to  $B_0'$  (18:55)  $B_0''$ ,  $2J'' + 1$  therefore (19:08)  $\Delta \tilde{\nu}_Q$  equal to  $2J'' + 1$ ,  $B_0' - B_0''$  that of course where is R is  $3B_0' - B_0''$ ,  $B_0'' - B_0'$ .

So, these are the spacing's between the consecutive ro-vibronic lines within the P, Q and the R branches. So, despite that we have a similarity if we now look at the (20:01) a long time back. So, if you look at for non-electronic transmissions or non-electronic spectroscopy, which is just IR spectroscopy in that particular case. Remember we wrote  $\nu_P$  as  $B_0' + B_0'' + \frac{B_0'' - B_0'}{2J'' + 1}$ , plus  $B_0' + B_0'' + B_0'' - B_0'$ , remember that we wrote this for the P branch for the rovibrational spectroscopy.

Now, for the electronic spectroscopy, if we now look at the electronic spectroscopy now, you will find that the sequence that we wrote this was the electronics spectroscopy function, this came from the rovibrational spectroscopy. So, there is an equivalence as we can see from this situation and what you see here they look kind of very similar, similarly, you will have the same equation for the R branch as well, but, there is a catch here.

So, it looks very similar, this one and this one, similarly, you will have similarity with this one and the corresponding one in the rovibrational counterpart that we should recognize.

(Refer Slide Time: 21:36)



$$\Delta \tilde{\nu}_R = \frac{(3B_0' - 3B_0'') - (B_0'' - B_0')}{(2J'+1)}$$

$$\Delta \tilde{\nu}_P = \frac{(B_0' + B_0'') + (B_0'' - B_0')}{(2J'-1)}$$
 There is equivalence but  $B_0'$  and  $B_0''$  can differ substantially for electronic transitions compared to

infrared transitions.  
 Recall  $B_0'' > B_0' \rightarrow$  mandatory for IR spectroscopy  
 $B_0'' > B_0'$  or  $B_0'' \leq B_0'$  can apply for visible and UV spectroscopy

Now there is equivalence but there is a catch. So, we must recognize that  $B_v$  prime and  $B_v$  double prime can differ and differ substantially for electronic transitions compared to infrared

transitions. Now, recall  $B_v$  double prime has to be greater than or equal to  $B_v$  prime this is mandatory for IR spectroscopy.

Now, this is mandatory but however  $B_v$  prime greater than equal to  $B_v$  double prime greater than equal to  $B_v$  prime or  $B_v$  double prime less than or equal to  $B_v$  single prime can apply for visible and u V spectroscopy that is a rotational constant depends on the Morse potential but each specific electronic state, rotational constant depends on the Morse potential for each specific electronic state.

(Refer Slide Time: 24:56)

So, what happens with the as a result of this is that these all these features that we talked about, enforce this ro-vibronic lines lie within a band to crowd together and may eventually reverse their spectral directions also, what we call a well-defined band head. In fact, so, just putting a few points over here, these ro-vibronic lines within a band crowd together and then eventually reverse their spectral directions creating band head.

The band heads can lie within the P and the R branch, so what we can do is that, we can set we can set this to be equal to zero, this P and R if we set P is equal to 0, so we can get what we call the band head. So, for each band head, so if we set  $\Delta v_P$ ,  $\Delta v_R$  is equal to 0 you can calculate what we call a band head and that band head is given as half multiply B double prime divided by, J R double prime, double prime.

So, now, if we recall that if  $B v \text{ double prime}$  is less than equal to  $B v \text{ prime}$  then the band head which is reversal band head appears in P branch and if  $B v \text{ double prime}$  is greater than equal to  $B v \text{ prime}$  the band head appears in R branch. So, the sectoral analysis between the ro-vibronic lines of mandate use of R branch for this and P branch for other one, we just reverse it. On the other hand the band heads are inevitably the most easily observable feature in any ro-vibronic spectrum because of spectral crowding and reversal.

So, it just is very easy to identify in a particular spectrum. So, we will also show an example in the next class that how one sub spectrum actually looks like the u V spectrum, nitrogen actually. So, you can understand that why the band heads should appear in the two bands. So, in the next class we are going to look at it in a little bit more details and look at the figure for the nitrogen spectrum and see how that actually goes. So, till the next class.